

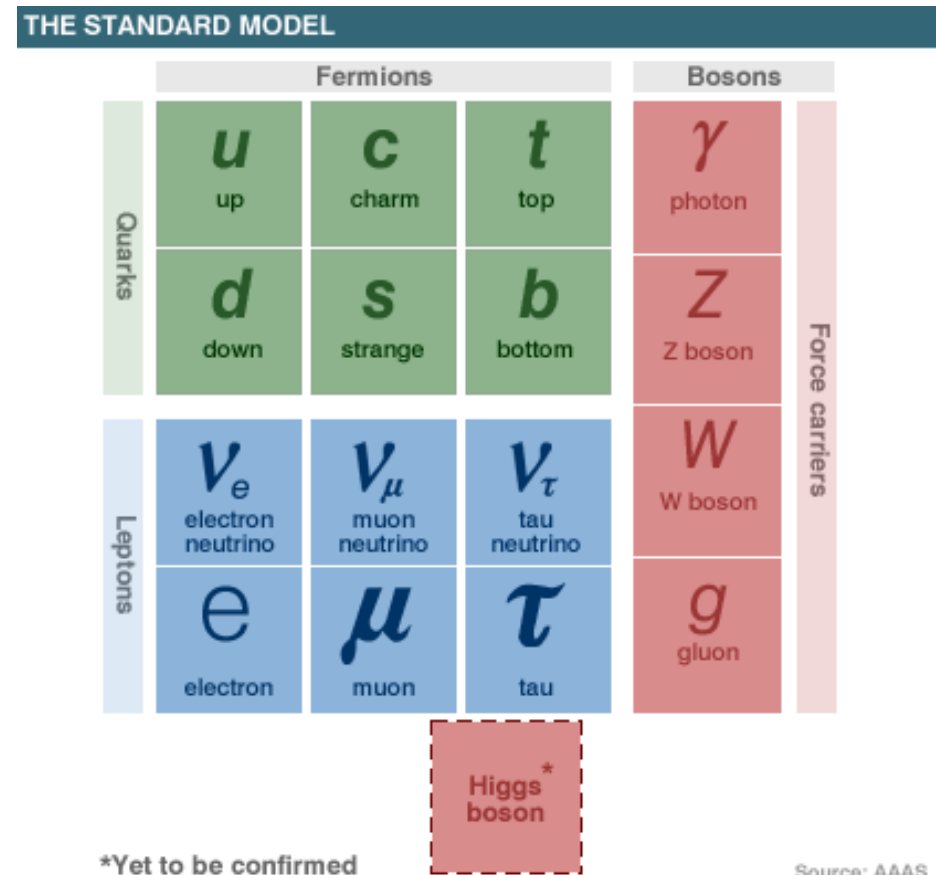
# Searching for axions and ALPs from string theory.

**Andreas Ringwald (DESY)**

PASCOS 2012,  
Merida, Mexico, 3-8 June 2012

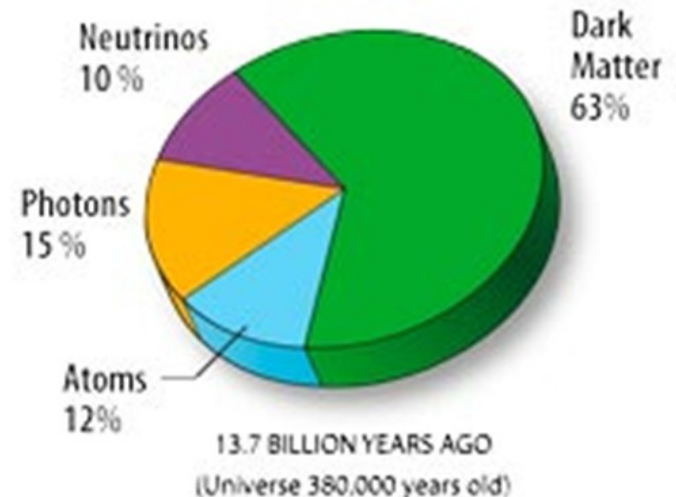
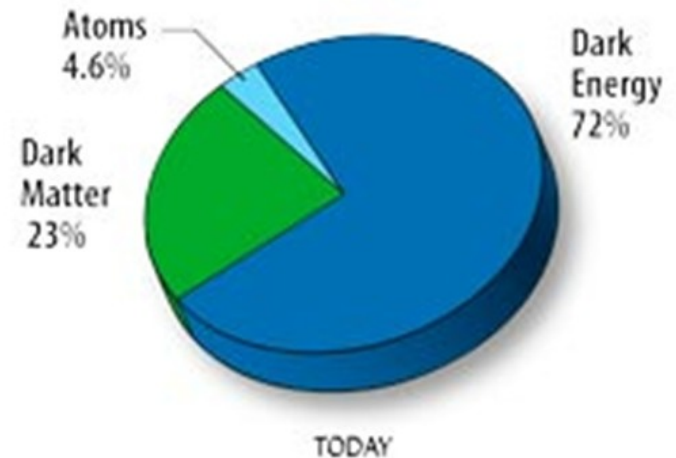
# Particles beyond the Standard Model?!

- Standard Model of Particle Physics: Fundamental description of known matter particles and gauge forces



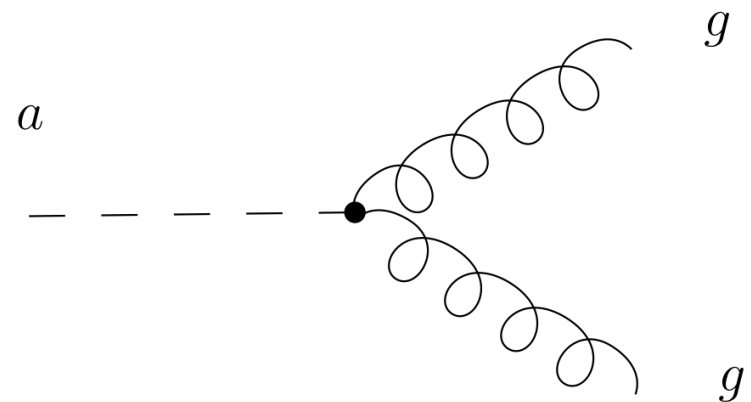
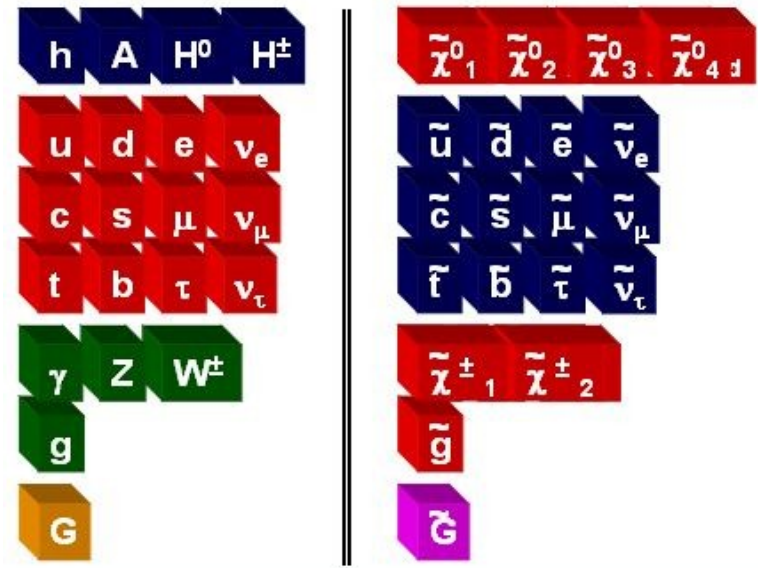
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- > Standard Model of Particle Physics: Fundamental description of known matter particles and gauge forces
- > Standard Model of Cosmology: only about 5 % of energy content of present universe consists of known particles
- > Extensions of the Standard Model of Particle Physics: several good motivated candidates for constituents of dark matter
  - SUSY: Neutralino, Gravitino
  - Peccei Quinn: Axion



# Peccei-Quinn extension of the Standard Model

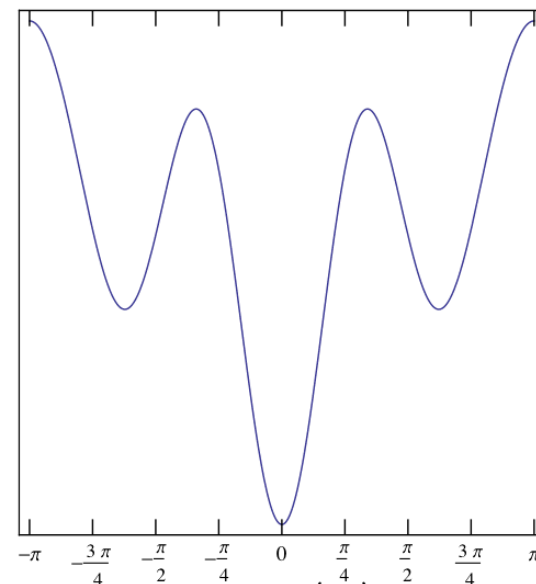
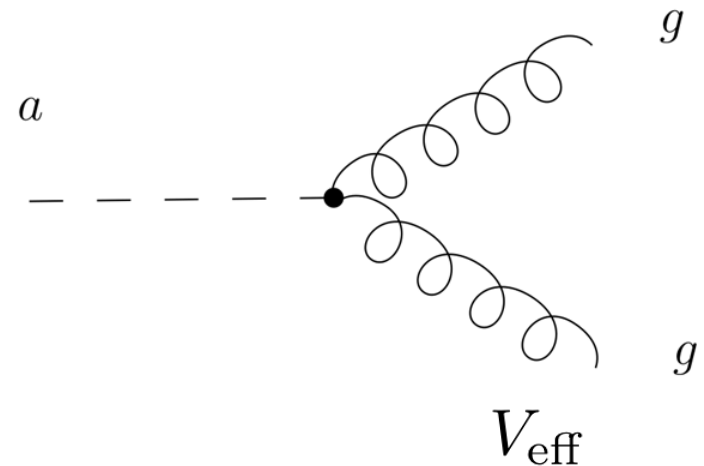
## > Motivation:

Explanation of unnatural smallness,  $\theta < 10^{-10}$ , of CP-violating topological term in QCD Lagrangian.

$$\mathcal{L}_{\text{CP-viol.}} = \frac{\alpha_s}{4\pi} \theta \text{tr} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

## > Axion field $\theta \rightarrow a(x)/f_a$

- $\langle a \rangle = 0$



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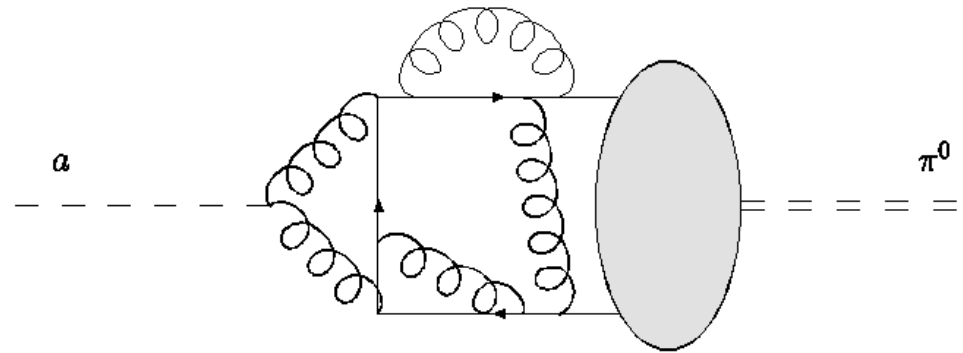
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$$m_a = \frac{m_\pi f_\pi}{f_a} \frac{\sqrt{m_u m_d}}{m_u + m_d} \simeq 6 \text{ meV} \times \left( \frac{10^9 \text{ GeV}}{f_a} \right)$$



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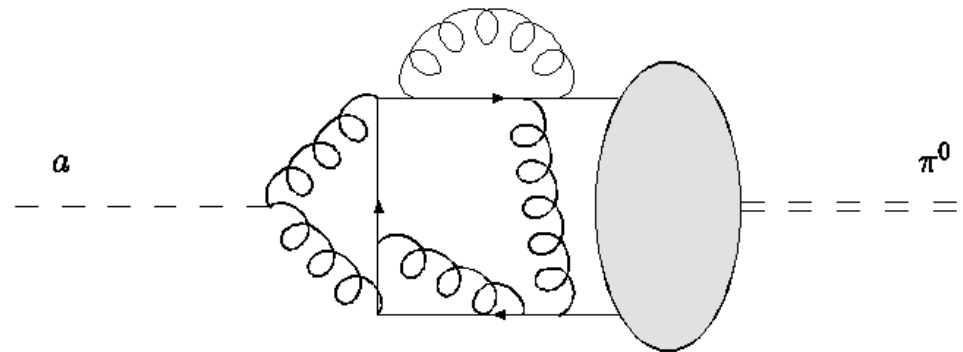
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- extremely weak interactions with Standard Model particles

$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma} a \vec{E} \cdot \vec{B},$$

$$g_{a\gamma} \simeq \frac{\alpha}{2\pi f_a} \sim 10^{-12} \text{ GeV}^{-1} \left( \frac{10^9 \text{ GeV}}{f_a} \right)$$

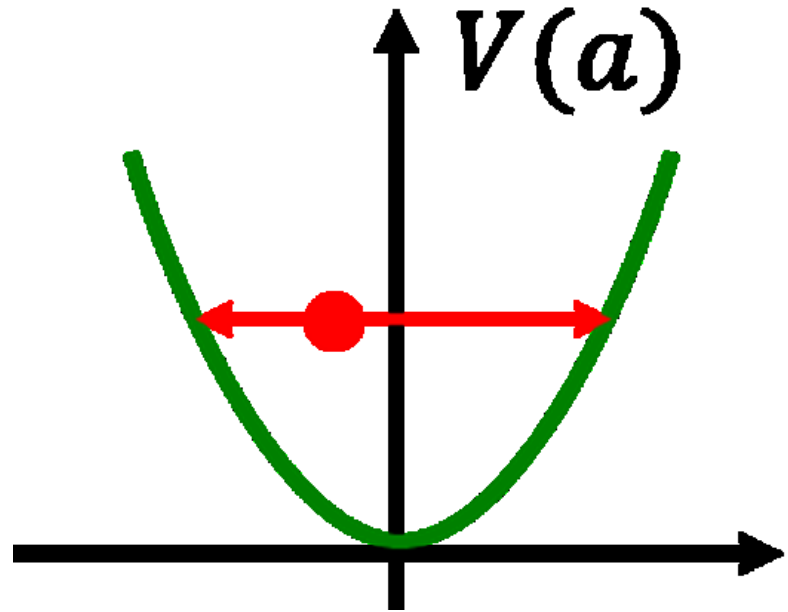


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- > Axion field  $\theta \rightarrow a(x)/f_a$ 
  - $\langle a \rangle = 0$
  - Axion: ultralight particle, cf. Weinberg '78; Wilczek '78
  - extremely weak interactions with Standard Model particles
- > Welcome side effect:
  - Axion: candidate for dark matter: created non-thermally via misalignment mechanism in form of coherent oscillations of axion field



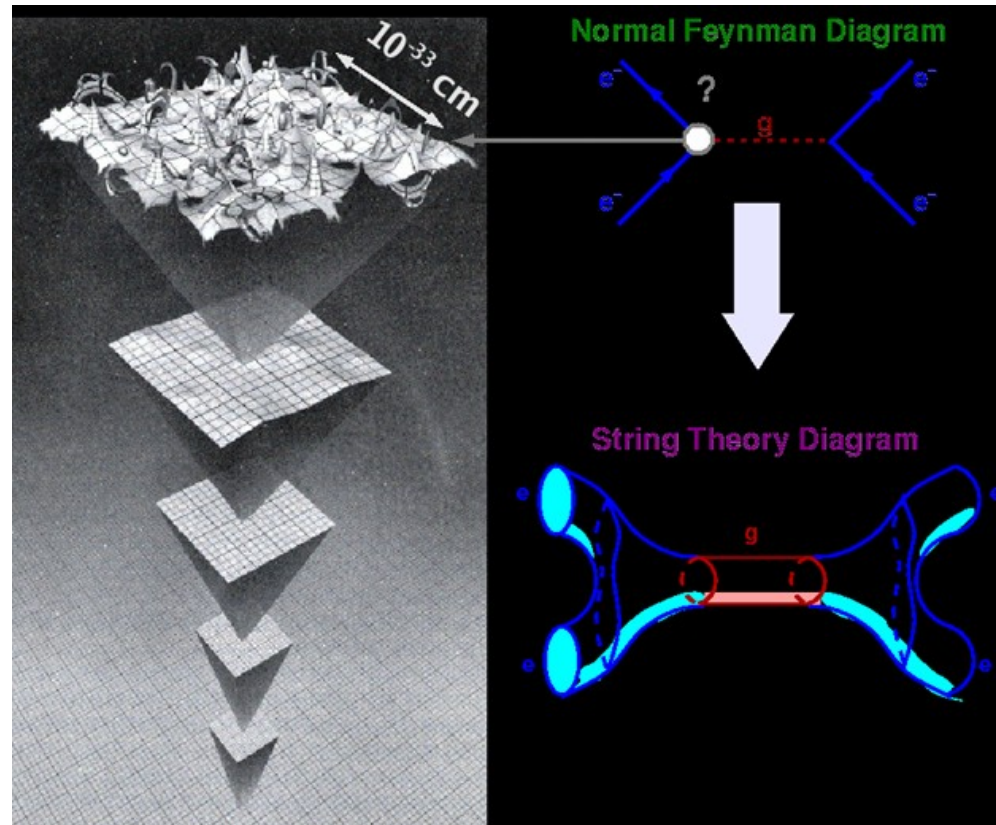
$$\Omega_a h^2 \approx 0.7 \left( \frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \left( \frac{\theta_i}{\pi} \right)^2$$



# Coexistence: SUSY and PQ extension in string theory

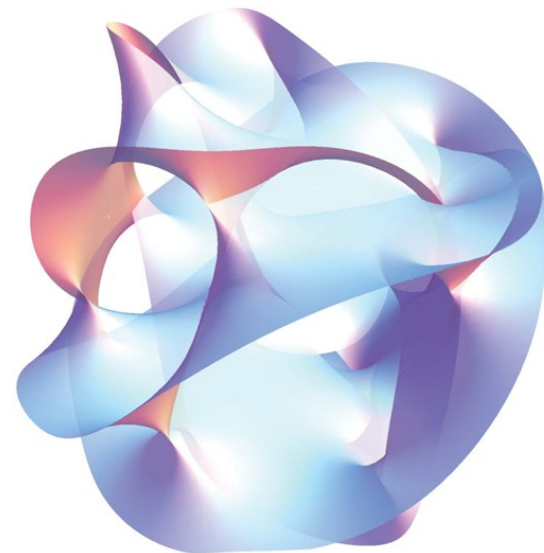
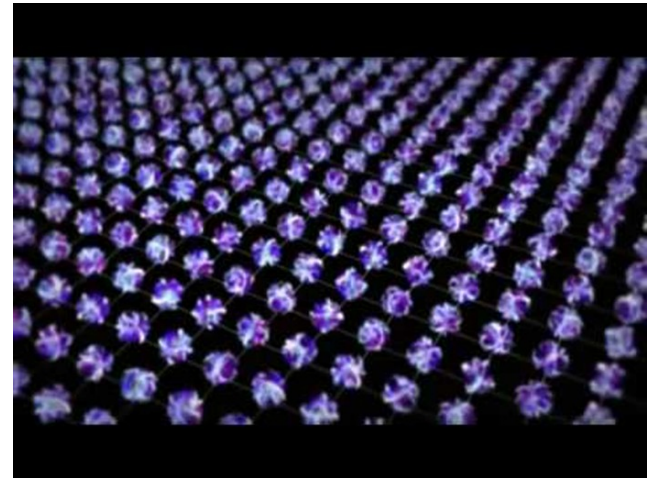
➤ Particularly strongly motivated extensions of Standard Model based on string theory:

- Unification of all forces
- Quantum gravity



# Coexistence: SUSY and PQ extension in string theory

- > Particularly strongly motivated extensions of Standard Model based on string theory:
  - Unification of all forces
  - Quantum gravity
- > Spectrum of low-energy effective theory in (3+1)-dimensions is supersymmetric and possibly contains several kinds of very weakly interacting slim particles (WISPs): **Axion, ALPs (Axion-Like Particles)**
  - if the compact space comprised of the 6 extra dimensions has certain geometrical and topological properties (Calabi-Yau; several cycles)



# Axions in string theory

- > String compactifications generically contain pseudo-scalar fields with axionic coupling to gauge fields and anomalous global shift symmetry

$$a_i F \tilde{F} \qquad a_i \rightarrow a_i + \epsilon$$

- > These axion and axion-like particle (ALP) candidates arise in string compactifications as KK zero modes of antisymmetric tensor fields:  
cf. Witten '84

heterotic string :  $B_2$

IIB string :  $C_2, C_4$



# Axions in IIB string theory

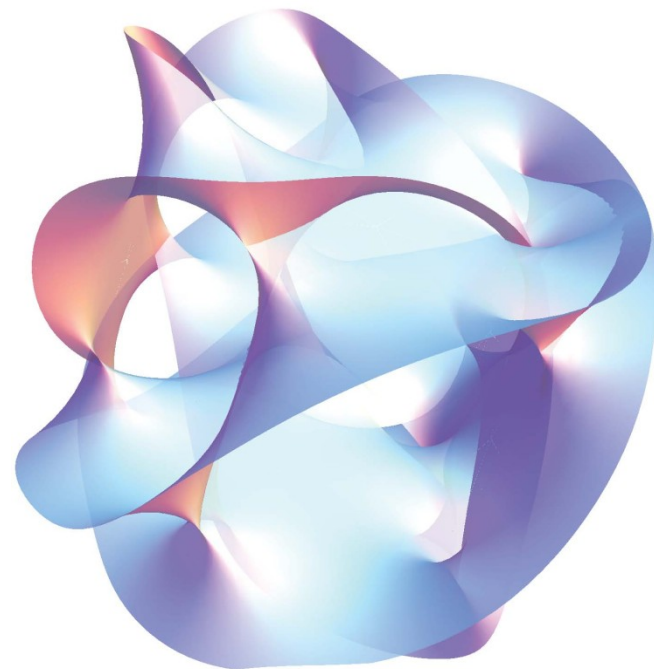
- Concentrate on IIB case (moduli stabilisation best understood):  
Realisation of brane-world scenarios in string theory

- KK reduction (expansion in harmonic forms):

$$C_2 = c^a(x)\omega_a, \quad a = 1, \dots, h_-^{1,1}$$

$$C_4 = c_\alpha(x)\tilde{\omega}^\alpha + \dots, \quad \alpha = 1, \dots, h_+^{1,1}$$

- Number of axionic fields determined by topology of CY orientifold: number of topologically non-equivalent 2-cycles or 4-cycles



# Axions in IIB string theory

➤ Number of cycles generically  $O(100)$ :

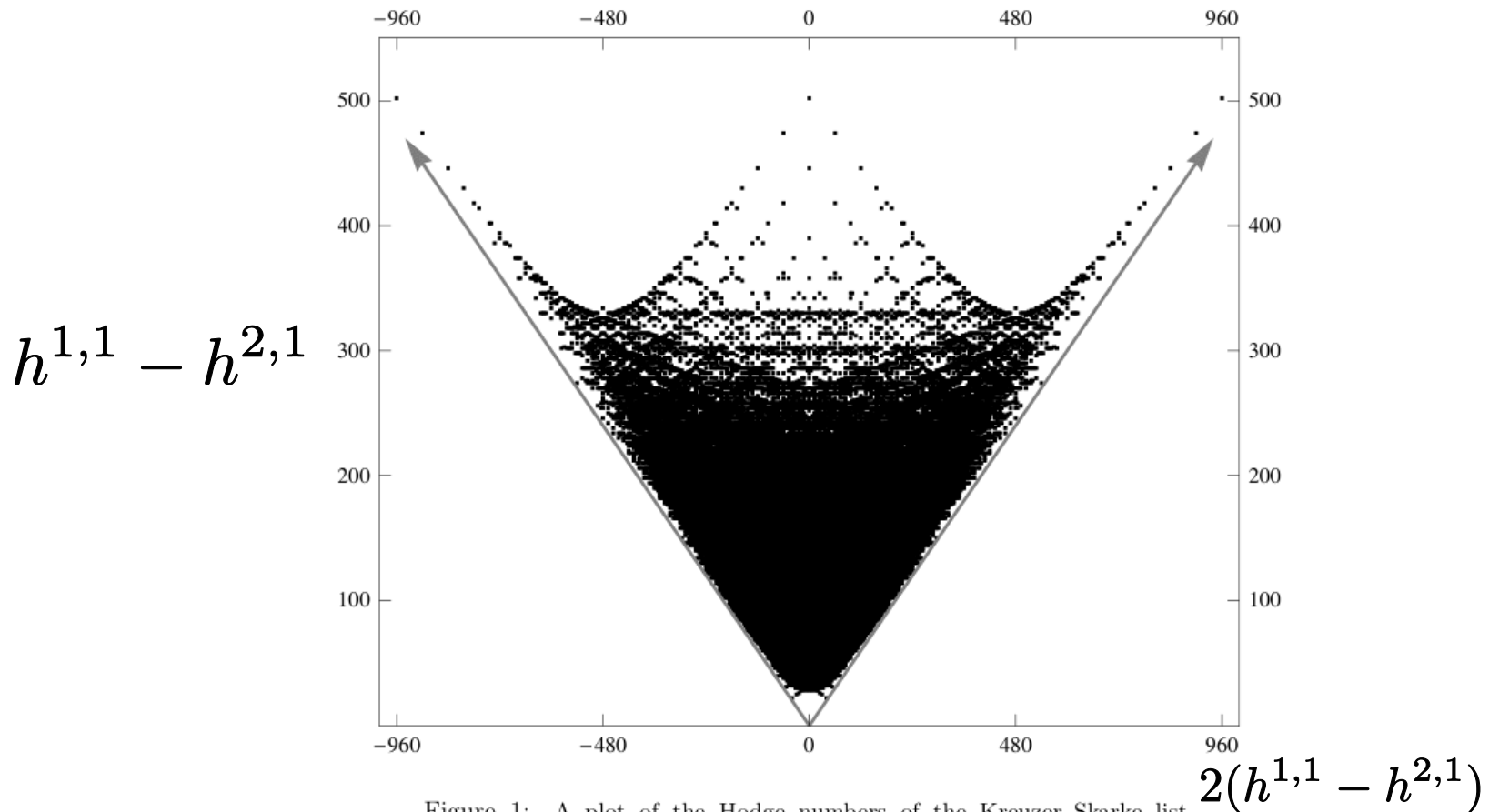
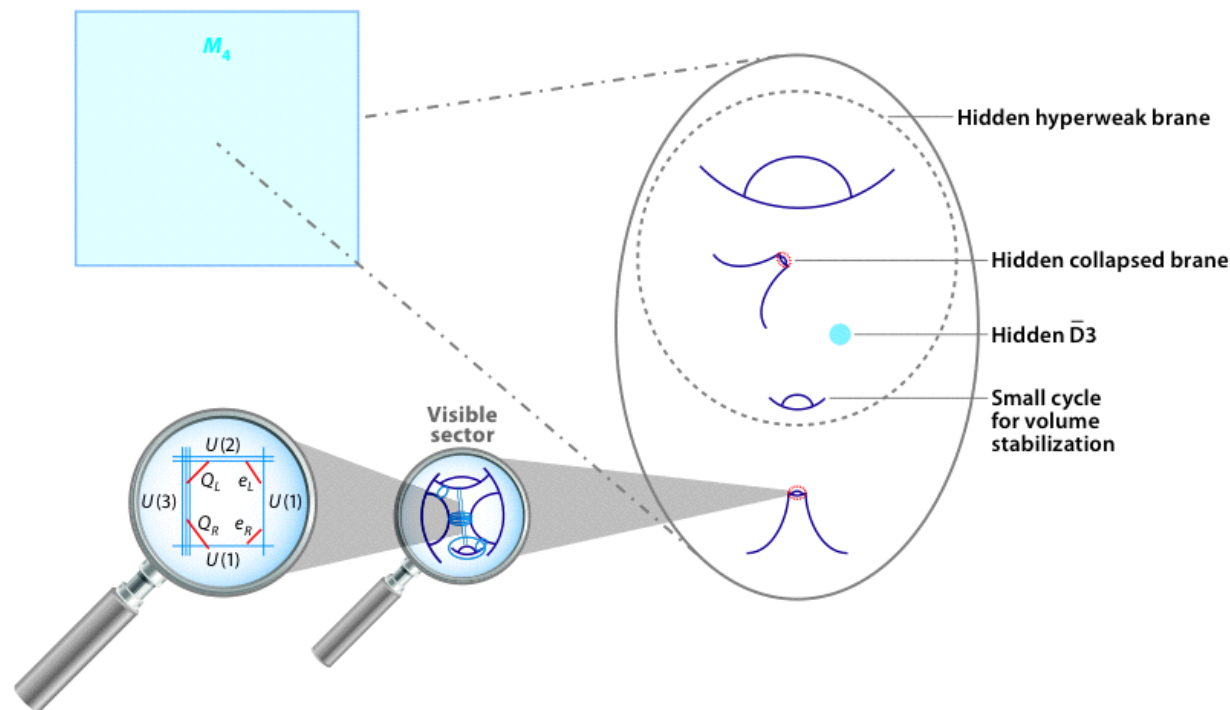


Figure 1: A plot of the Hodge numbers of the Kreuzer-Skarke list.  $\chi = 2(h^{1,1} - h^{2,1})$  is plotted horizontally and  $h^{1,1} + h^{2,1}$  is plotted vertically. The oblique axes bound the region  $h^{1,1} \geq 0$ ,  $h^{2,1} \geq 0$ .

# Axions in IIB string theory

- Cycles can be wrapped by space-time filling D-branes



- Each of these branes gives rise to a gauge theory at low energy
  - Visible sector gauge theory realized by stacks of D7 branes wrapping small 4-cycles
  - Hidden sectors, in particular hidden photons, realized by branes wrapping cycles not intersecting visible sector branes

# Axions in IIB string theory

- Each axionic field comes along with a real scalar field – saxion – which is real part of lowest component of chiral superfield,

$$T_\alpha = \tau_\alpha + i c_\alpha$$

- $\tau_\alpha$  ... Kähler modulus measuring the volume of 4-cycle  $\alpha$
- 4D EFT from KK reduction of D-brane action

$$S_p = \frac{-2\pi}{(2\pi\sqrt{\alpha'})^{p+1}} \left( \int_\Sigma d^{p+1}\xi e^{-\phi} \sqrt{\det(g + B + 2\pi\alpha'F)} + i \int_\Sigma e^{B+2\pi\alpha'F} \wedge \sum_q C_q \right).$$

- $T_\alpha$  is gauge kinetic function for theory on D7-brane:
  - volume measures gauge coupling,  $\tau_\alpha \sim g^{-2}$
  - $c_\alpha$  has axionic coupling,  $\sim c_\alpha F \wedge F$



# Axions in IIB string theory

- 4D effective field theory, cf. [Jockers, Louis '05](#)

$$S \supset -dc_\alpha \frac{\mathcal{K}_{\alpha\beta}}{8} \wedge \star dc_\beta - \frac{r^{i\alpha} \tau_\alpha}{4\pi M_P} (F_i \wedge \star F_i) + \frac{r^{i\alpha} c_\alpha}{4\pi M_P} (F \wedge F),$$

with  $\mathcal{K}_{\alpha\beta} \equiv \frac{\partial^2 K}{\partial \tau_\alpha \partial \tau_\beta}$ ,  $K = -2 \ln \mathcal{V}$ ,  $r^{i\alpha} \equiv \ell_s^{-4} \int_{D_i} \tilde{\omega}^\alpha$

- Decay constants and coupling to gauge bosons via canonical normalization of axion and gauge kinetic terms and matching to:

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu a_i \partial^\mu a_i - \frac{g_3^2}{32\pi^2} \left( \theta_0 + C_{i33} \frac{a_i}{f_{a_i}} \right) F_{3,\mu\nu}^b \tilde{F}_3^{b,\mu\nu}$$
$$- \frac{g_2^2}{32\pi^2} C_{iWW} \frac{a_i}{f_{a_i}} F_{W,\mu\nu}^b \tilde{F}_W^{b,\mu\nu} - \frac{g_Y^2}{32\pi^2} C_{iYY} \frac{a_i}{f_{a_i}} F_{Y,\mu\nu} \tilde{F}_Y^{\mu\nu}$$





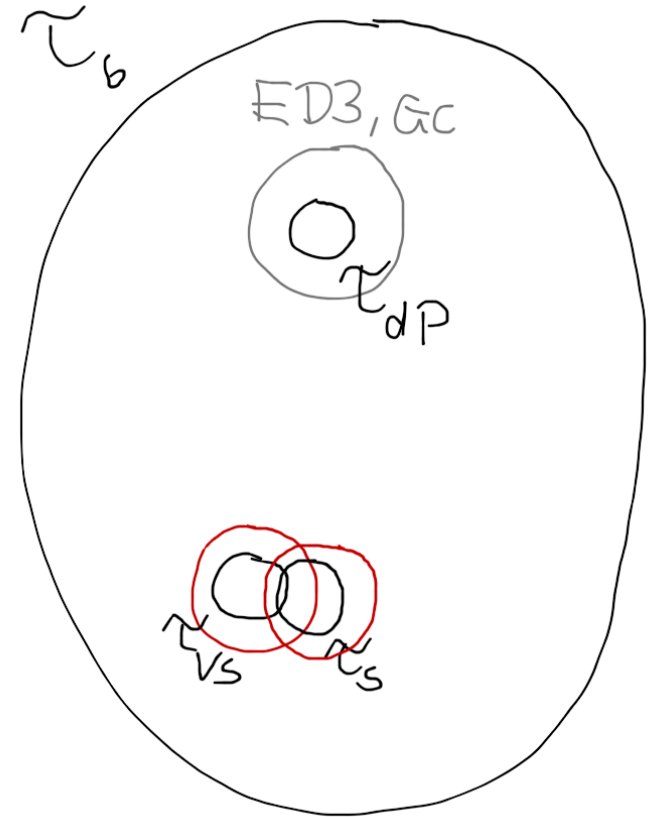
- > An axiverse - QCD axion plus possibly many ultra-light ALPs whose mass spectrum is logarithmically hierarchical – may naturally arise from strings, cf. [Arvanitaki et al. `09](#)
- > Challenges to obtain an axiverse:
  - Only axions which are not projected out by orientifold projection appear in LE EFT
  - Only axions which do not get too heavy by Kähler moduli stabilisation can be candidates for QCD axion and other light ALPs, cf. [Conlon `06](#)
  - Only axions which are not eaten by Stückelberg mechanism to give masses to brane-localized anomalous U(1) gauge bosons will appear in LE EFT
- > [Acharya, Bobkov, Kumar `11](#): Moduli stabilisation via single non-perturbative correction to superpotential (cf. [Bobkov, Braun, Kumar, Raby `10](#)) fixes all Kähler moduli plus one axion combination: axiverse with  $W_0 \ll 1$  and  $f_a \sim 10^{16}$  GeV
- > [Cicoli, Goodsell, AR, 1206.0819](#): Moduli stabilisation of the so-called LARGE Volume Scenario (LVS) which exploits both perturbative and non-perturbative effects to fix Kähler moduli gives rise to an axiverse with  $W_0 \sim 1$  and  $f_a \sim 10^{10}$  GeV

## > Strategy to fix moduli in LVS: (cf. [Cicoli, Mayrhofer, Valandro '11](#))

- Exploit del Pezzo four-cycle supporting single non-perturbative effect; dP modulus fixed at small size, dP saxion and axion heavy; interplay with leading order  $\alpha'$  correction yields exponentially large CY volume
- Visible sector with chirality built by wrapping magnetised D7-branes around rigid but not del Pezzo four cycles; D-term conditions stabilise d combinations of Kähler moduli, leaving
  - $n_{\text{ax}} \equiv h^{1,1} - 1 - d \geq 1$  flat directions; latter fixed by pert. corrections
- LVS requires  $n_{\text{ax}} \geq 2$ : one of the remaining cycles should be small to obtain correct value of  $g_{\text{VS}}^2 \sim 1/\tau_{\text{VS}}$ , while there should be at least one further which can be large; latter fixed at

$$\mathcal{V} \sim \tau_b^{3/2} \sim W_0 e^{2\pi\tau_{\text{dP}}}$$

## > More on this: talk by [Michele Cicoli](#)



➤ Mass scales for  $g_s \sim 0.1, W_0 \sim 1, \mathcal{V} \sim 10^{14}$ :

$$M_s \sim \frac{M_P}{\sqrt{4\pi\mathcal{V}}} \sim 10^{10} \text{ GeV}$$

$$m_{\tau_s} \sim \frac{M_P}{\mathcal{V}^{1/2}} \sim 10^{10} \text{ GeV}$$

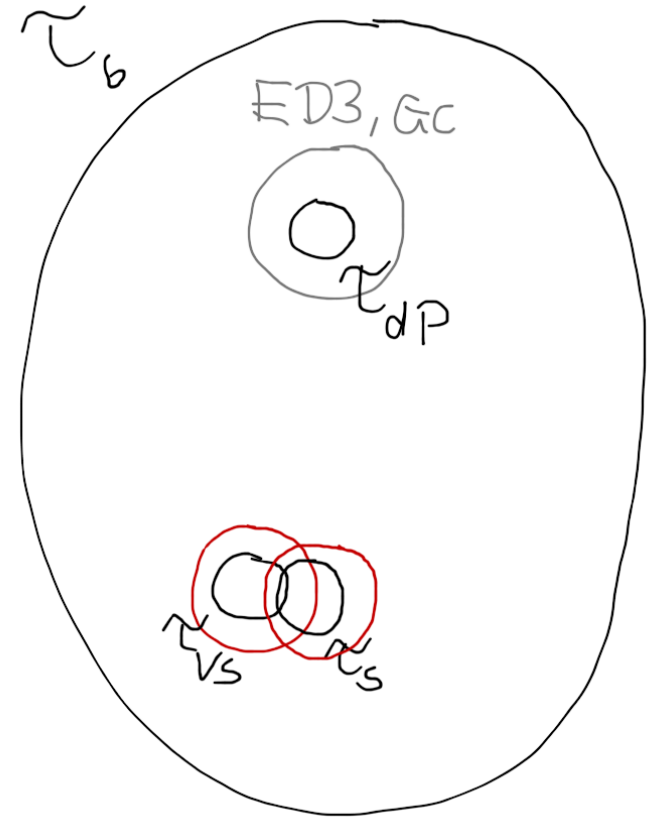
$$m_{\tau_{\text{dP}}} \sim \frac{M_P}{\mathcal{V}} \ln \mathcal{V} \sim 30 \text{ TeV}$$

$$m_{3/2} \sim \sqrt{g_s/(4\pi)} W_0 \frac{M_P}{\mathcal{V}} \sim 1 \text{ TeV}$$

$$m_{\tau_{\text{VS}}} \sim \alpha_{\text{VS}} m_{3/2} \sim 40 \text{ GeV}$$

$$m_{\tau_b} \sim \frac{M_P}{\mathcal{V}^{3/2}} \sim 0.1 \text{ MeV}$$

- No cosmological moduli problem since  $\tau_b$  diluted by entropy production due to decay of  $\tau_{\text{VS}}$  reheating universe to  $\mathcal{O}(\text{GeV})$



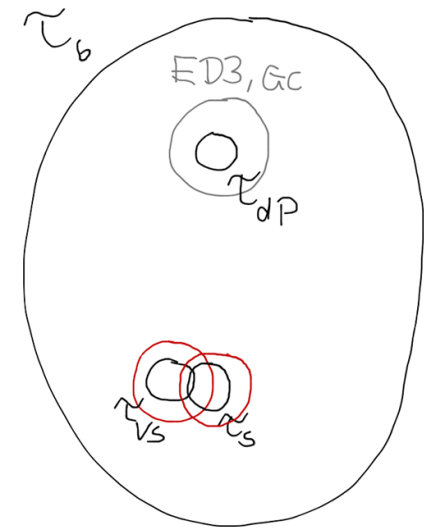
➤ Scaling of axion decay constants and couplings:

$$f_{a_b} = \frac{\sqrt{3}}{4\pi} \frac{M_P}{\tau_b} \simeq \frac{M_P}{4\pi \mathcal{V}^{2/3}} \simeq \frac{M_{\text{KK}}^{10\text{D}}}{4\pi}, \quad f_{a_s} = \frac{1}{\sqrt{6} (2\tau_s)^{1/4}} \frac{M_P}{4\pi \sqrt{\mathcal{V}}} \simeq \frac{M_s}{\sqrt{4\pi} \tau_s^{1/4}},$$

$$C_{bbb} \simeq g_b^{-2} \frac{f_{a_b}}{M_P} \simeq \mathcal{O}(1), \quad C_{sbb} \simeq g_b^{-2} \frac{f_{a_s} \tau_s^{3/4}}{\mathcal{V}^{1/2} M_P} \simeq \mathcal{O}(\epsilon) \simeq \mathcal{O}(\mathcal{V}^{-1/3}),$$

$$C_{bss} \simeq g_s^{-2} \frac{f_{a_b}}{M_P} \simeq \mathcal{O}(\epsilon^2) \simeq \mathcal{O}(\mathcal{V}^{-2/3}), \quad C_{sss} \simeq g_s^{-2} \frac{f_{a_s}}{\tau_s^{3/4} M_s} \simeq \mathcal{O}(1).$$

- $a_{\text{VS}}$  : QCD axion with  $f_{a_{\text{VS}}} \sim f_{a_s} \sim 10^{10} \text{ GeV}$
- $a_b$  : essentially massless ALP with  $f_{a_b} \sim 10^8 \text{ GeV}$   
 , but nearly decoupled,  $C_{bss} \sim 10^{-10}$



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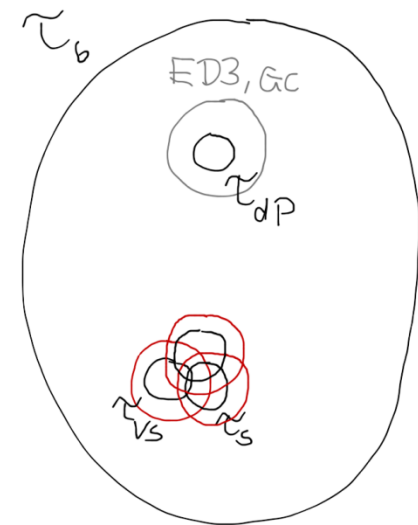
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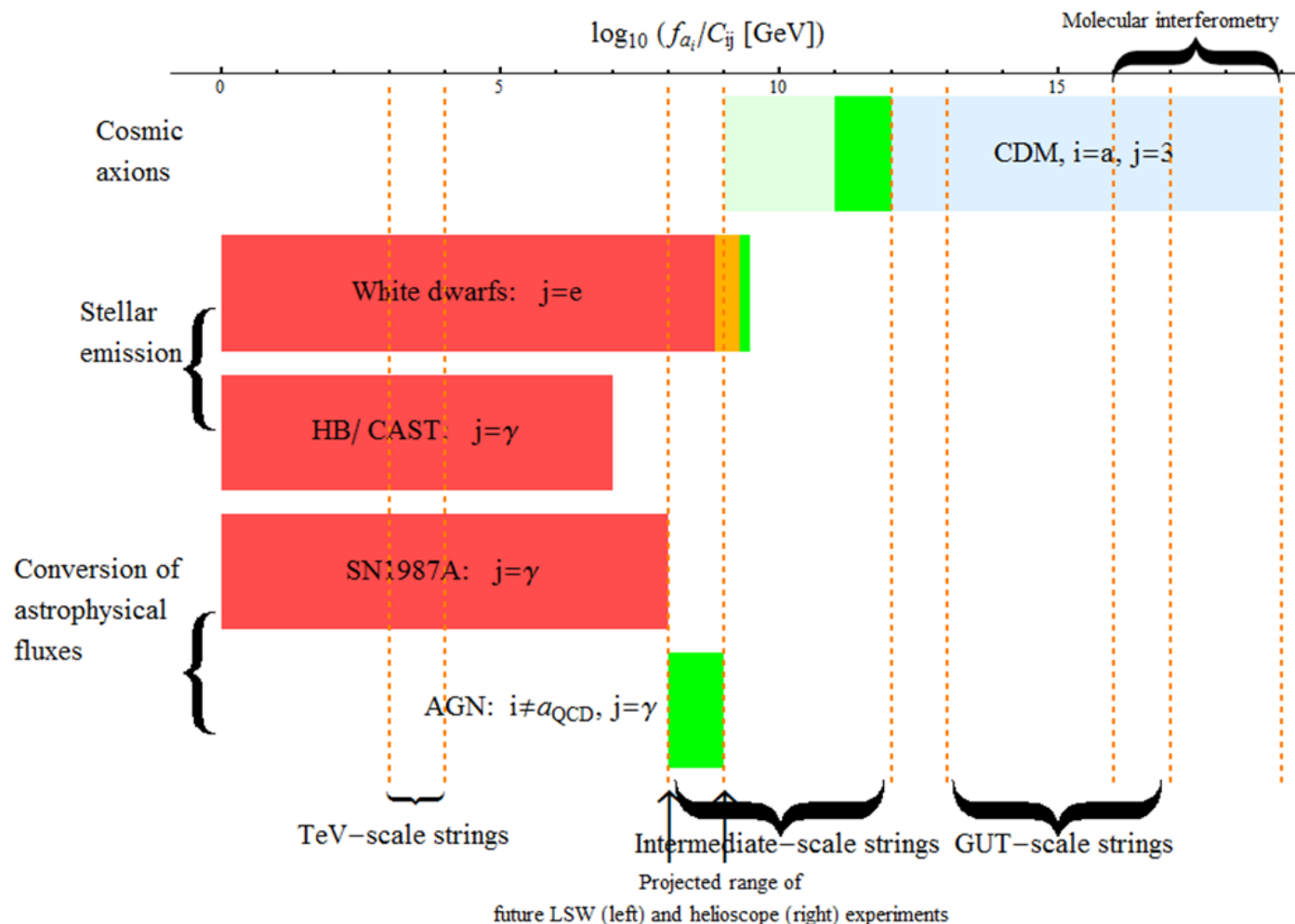
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- $a_b$  : essentially massless ALP with  $f_{a_b} \sim 10^8 \text{ GeV}$ , but nearly decoupled,  $C_{bss} \sim 10^{-10}$
- Possibly more ultralight,  $m_{a_b} \ll m_a$ , ALPs with

$$f_{a_s} \sim 10^{10} \text{ GeV}, C_{sss} \sim 1$$



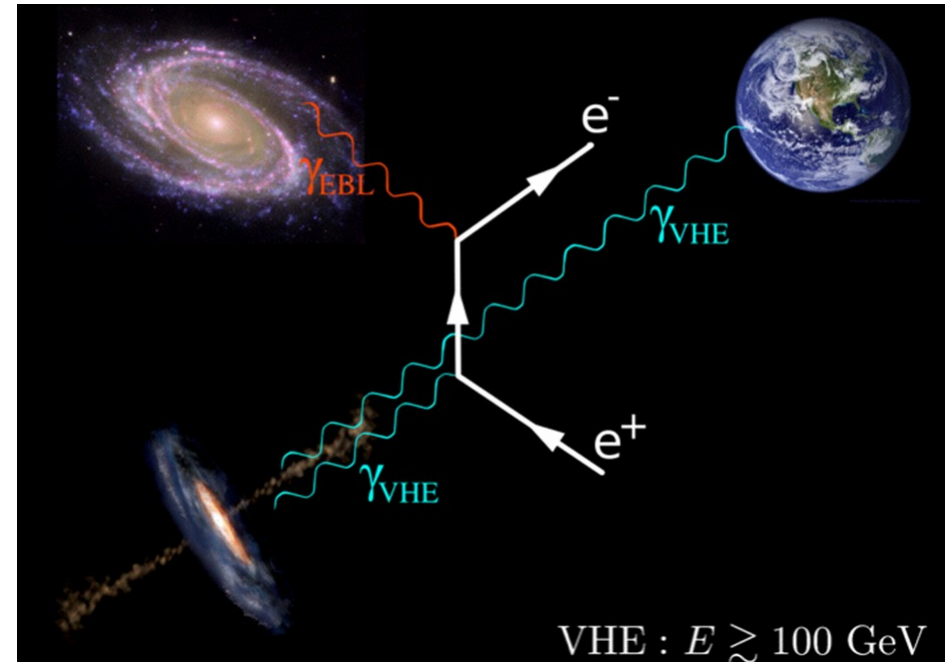
# Axion and ALPs with intermediate scale decay constant?

- Current limits and possible hints from astrophysics and cosmology:



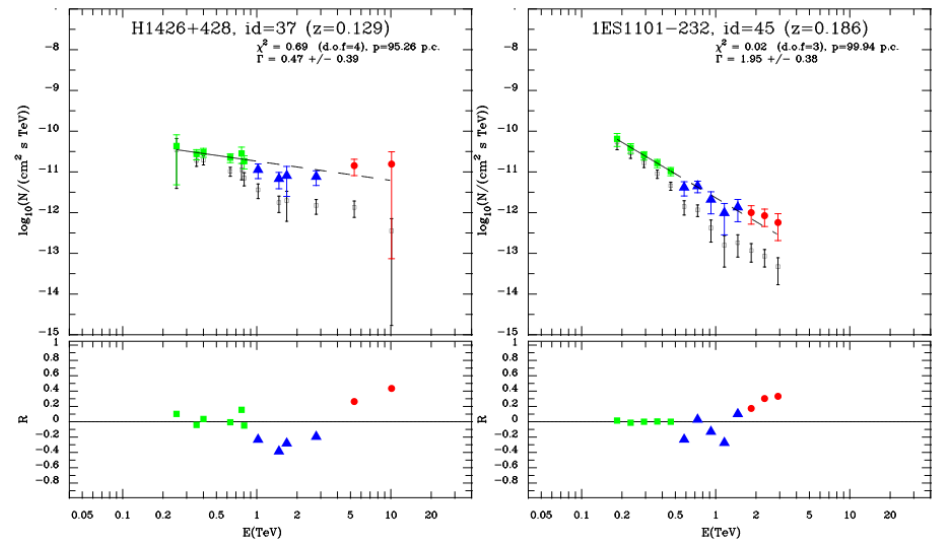
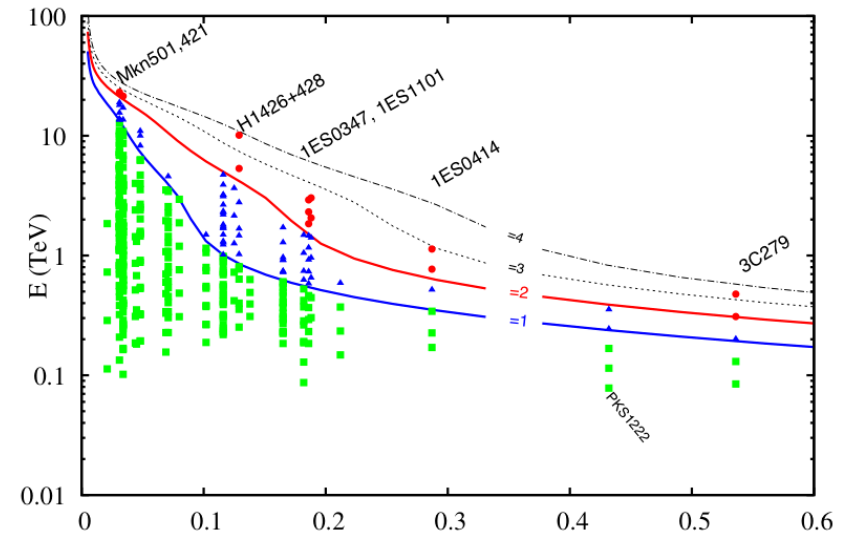
# Hints for axion or ALP effects in propagation of TeV photons through universe?

- TeV photons of distant Active Galactic Nuclei (AGN) should feature absorption breaks due to electron-positron pairproduction on the Extragalactic Background Light (EBL)



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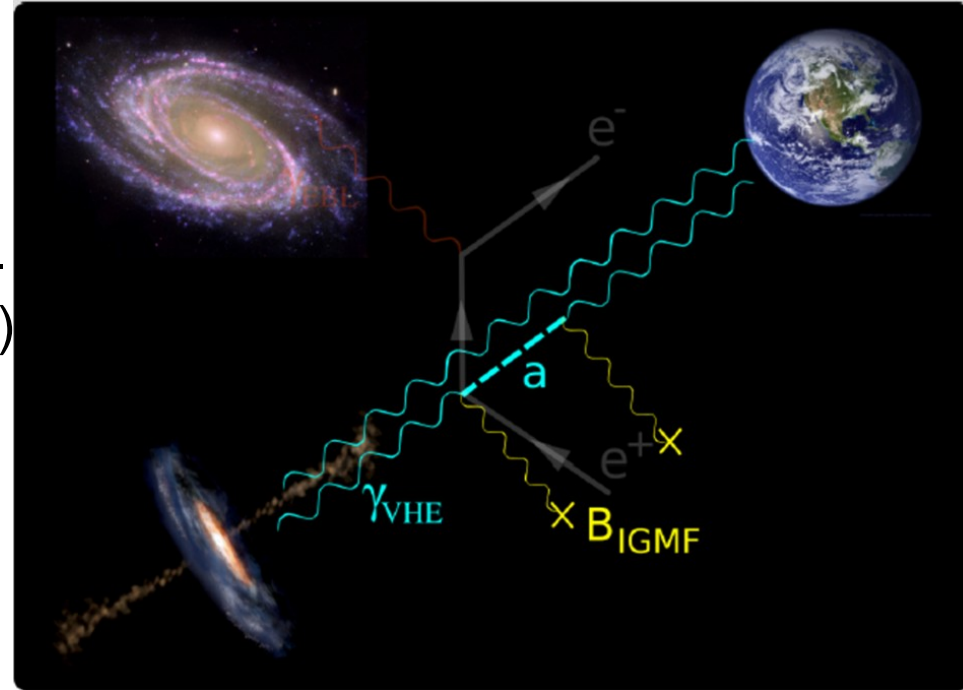
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- Not the case: (e.g. [Horns, Meyer `12](#))
  - 50 spectra (HESS, MAGIC, Veritas), assumption: minimal EBL; absorption ruled out by more than 4 sigma





# Hints for axion or ALP effects in propagation of TeV photons through universe?

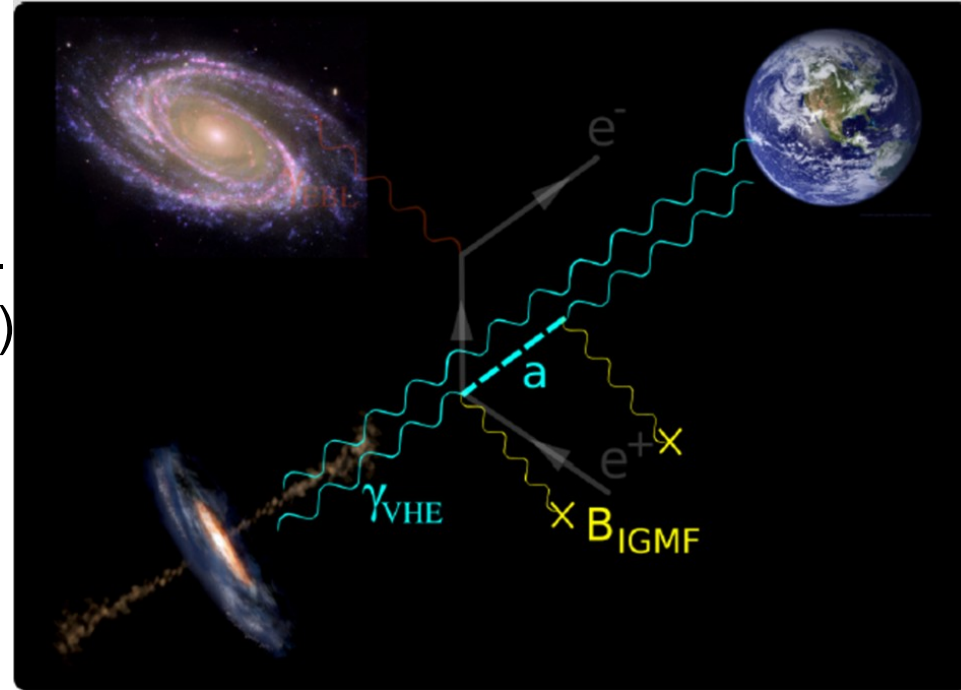
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- > Possible explanation: Photon  $\leftrightarrow$  ALP conversion in astrophysical magnetic fields ([Roncadelli et al.](#), [Sanchez-Conde et al.](#))



$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left( \frac{m_a^2}{4\omega} L_B \right)$$

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- Has to be an ALP: too light for a QCD axion with such a decay constant



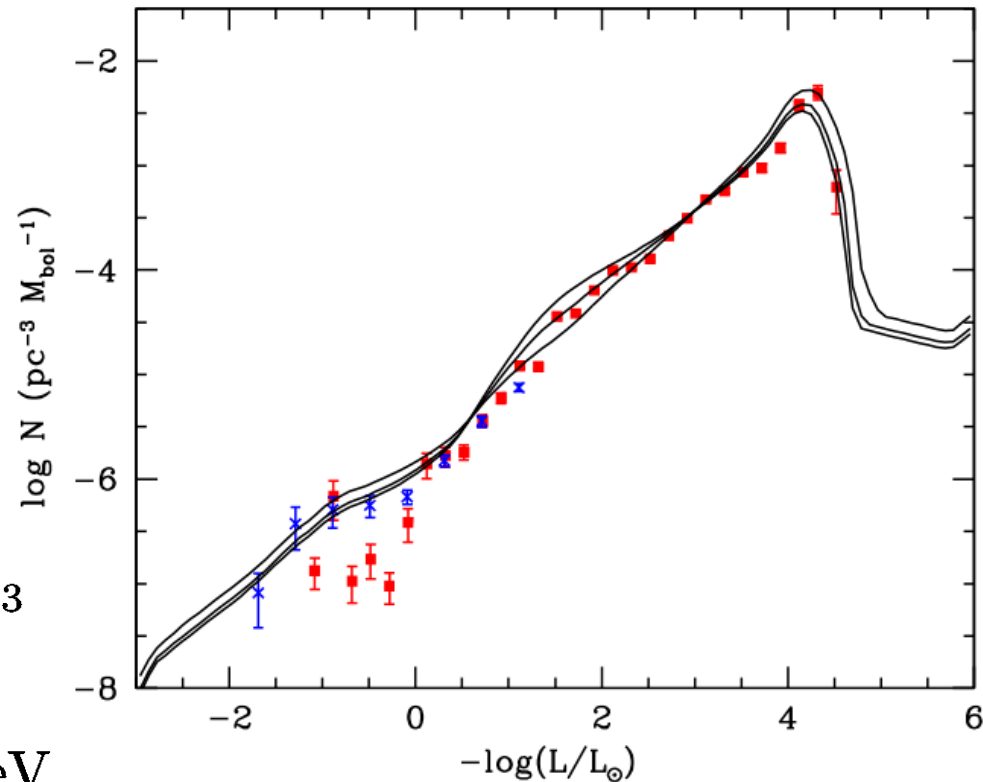
$$g_{i\gamma} \equiv \frac{\alpha C_{i\gamma}}{2\pi f_{a_i}} = 10^{-12} \div 10^{-11} \text{ GeV}^{-1}$$
$$\rightarrow \frac{f_{a_i}}{C_{i\gamma}} \simeq 10^8 \div 10^9 \text{ GeV}$$
$$\text{for } m_{a_i} \lesssim \text{neV}$$

# Hint for axion/ALP-production in white dwarfs?

- Non-standard energie-loss mechanism in white dwarfs, cf. [Isern et al.](#)
- Compatible with axion or ALP production in electron-bremsstrahlung

$$g_{ie} \equiv \frac{C_{ie} m_e}{f_{a_i}} = (2.0 \div 7.0) \times 10^{-13}$$
$$\rightarrow \frac{f_{a_i}}{C_{i,e}} \simeq (0.7 \div 2.6) \times 10^9 \text{ GeV}$$

for  $m_{a_i} \lesssim \text{keV}$



# Can we explain these hints within IIB axiverse?

- Anomalous transparency of universe and anomalous energy loss of white dwarfs could be explained by

$$C_{i\gamma}/C_{ie} \simeq 10, \quad f_{a_i}/C_{i\gamma} \simeq 10^8 \text{ GeV}, \quad m_{\text{ALP}} \lesssim 10^{-9} \div 10^{-10} \text{ eV}.$$

- Model where visible sector build from intersecting branes in geometric regime

$$\frac{C_{i\gamma}}{C_{ie}} \sim \frac{8\pi\tau_*}{3}, \quad \frac{f_{a_i}}{C_{i\gamma}} = \frac{1}{8\pi N_{i\gamma}\tau_*^{1/4}} \frac{M_P}{\sqrt{\mathcal{V}}} = \frac{1}{8\pi N_{i\gamma}\tau_*^{1/4}} \sqrt{\frac{g_s M_P m_{3/2}}{W_0}}.$$

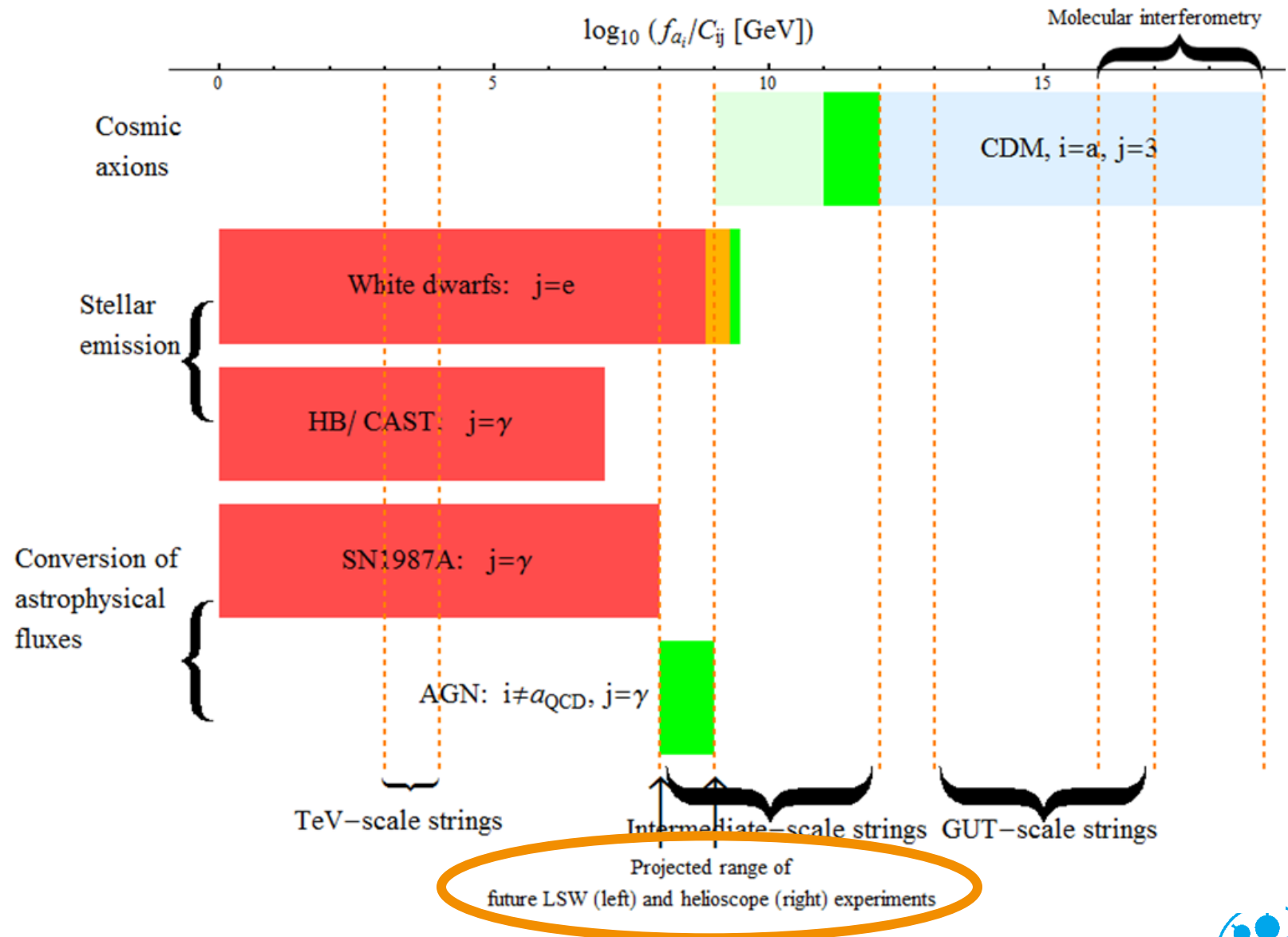
- Yields required values for  $m_{3/2} = 10 \text{ TeV}$ ,  $g_s \simeq 0.1$  and  $W_0 \sim 10$
- ALP mass could be generated by single Kähler potential instanton, e.g.  $m_{\text{ALP}} \sim m_{3/2} e^{-\pi\tau_*} \sim 10^{-10} \text{ eV}$  requires

$$\tau_* \sim \frac{1}{\pi} \ln \left( \frac{g_s m_{3/2}}{m_{\text{ALP}}} \right) \sim 16.$$

- Astrophysical hints compatible with intermediate string-scale scenario

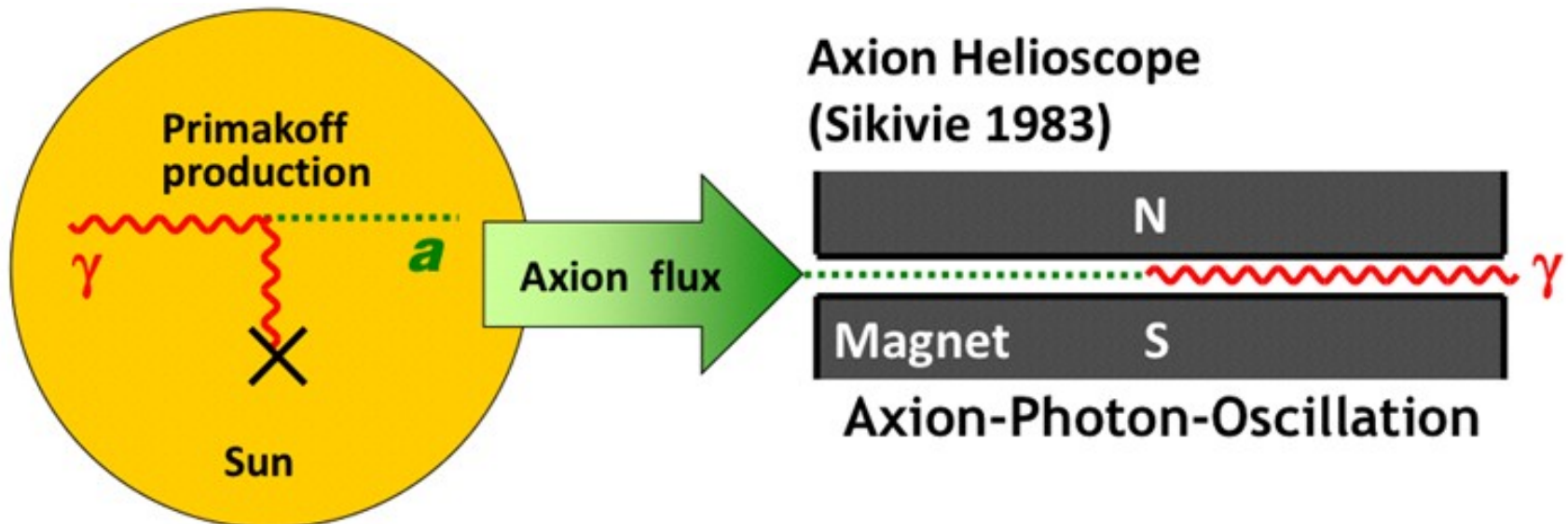


# Laboratory probes of axion and ALPs with intermediate scale decay constant?



# Search for solar axions and ALPs

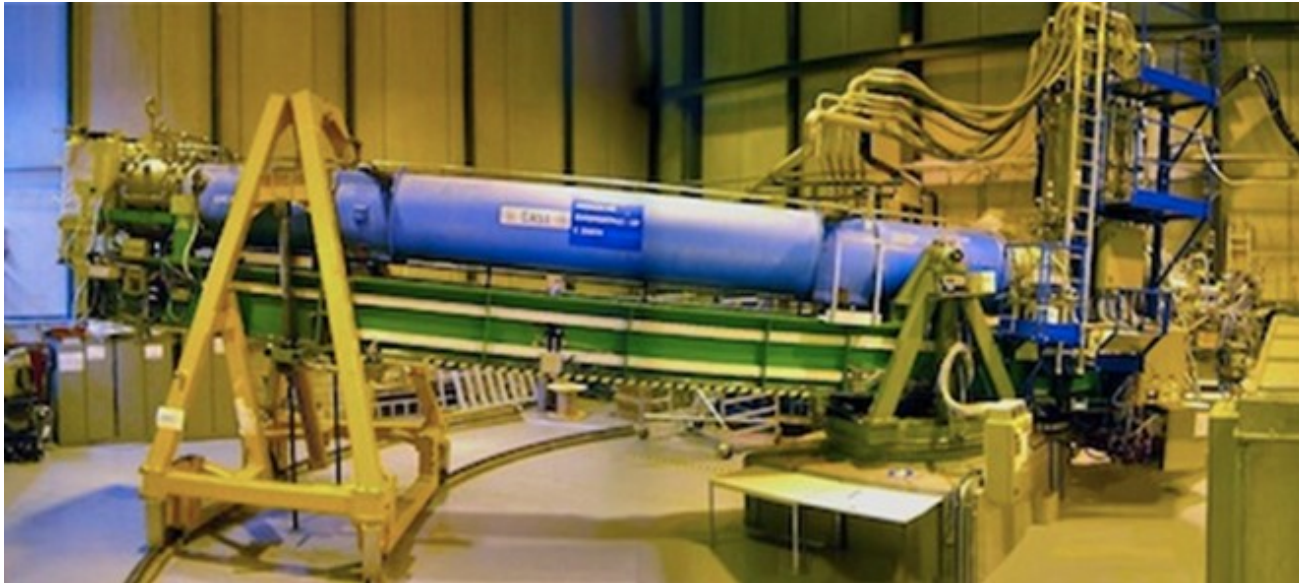
- Sun strong source of axions and ALPs
- Helioscope searches for axions/ALPs and HPs



$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left( \frac{m_a^2}{4\omega} L_B \right)$$

# Search for solar axions and ALPs

- Sun strong source of axions and ALPs
- Helioscope searches for axions/ALPs and HPs

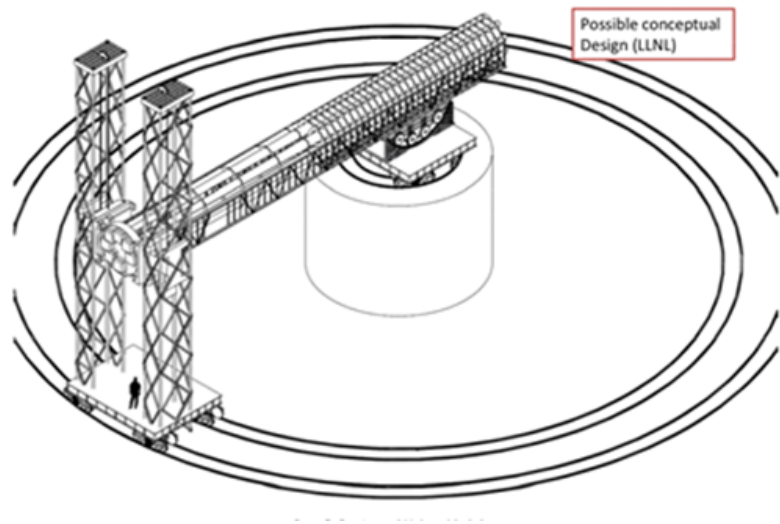
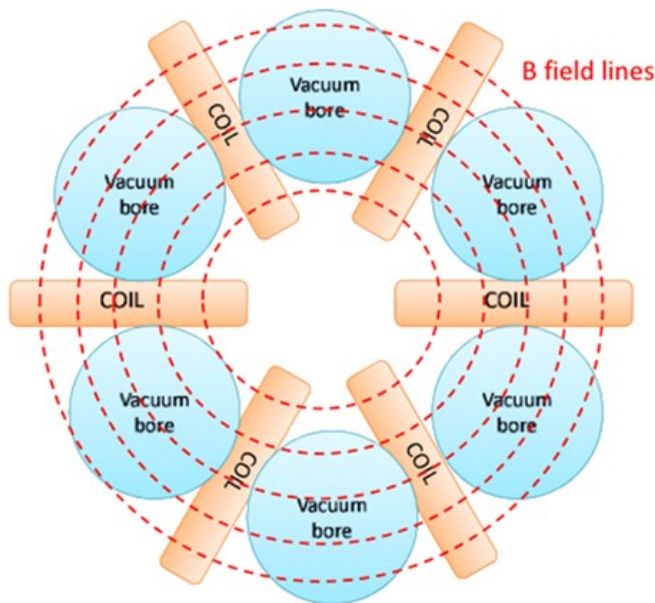


- CAST ... CERN Axion Solar Telescope



# Search for solar axions and ALPs

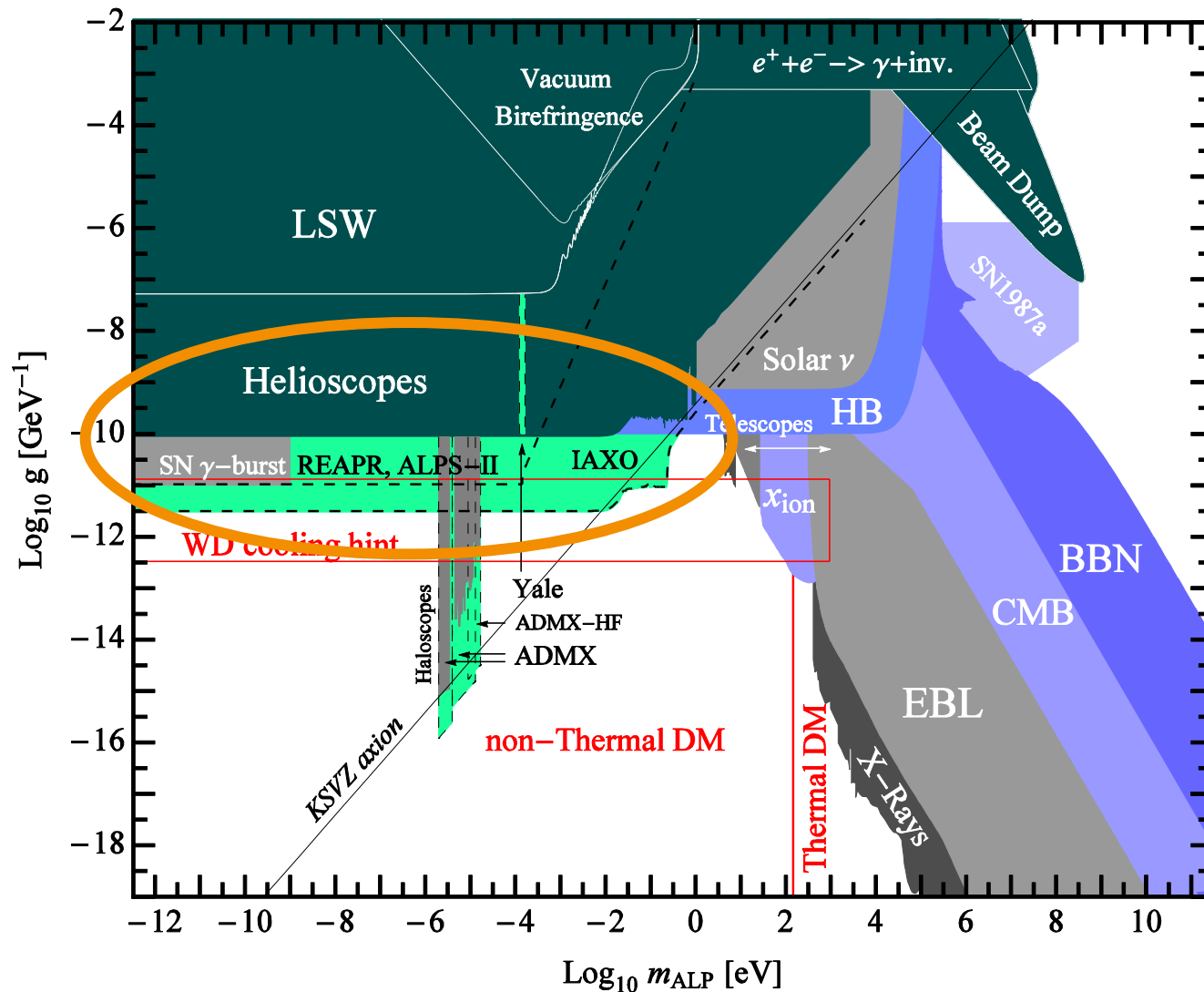
- Sun strong source of axions and ALPs
- Helioscope searches for axions/ALPs and HPs



- CAST ... CERN Axion Solar Telescope
- IAXO ... International Axion Observatory (under investigation)



# Search for solar axions and ALPs



# Search for ALPs via light-shining-through-walls

- ALPs can pass walls
- Light-shining-through-walls experiments: (here ALPS (@DESY)):

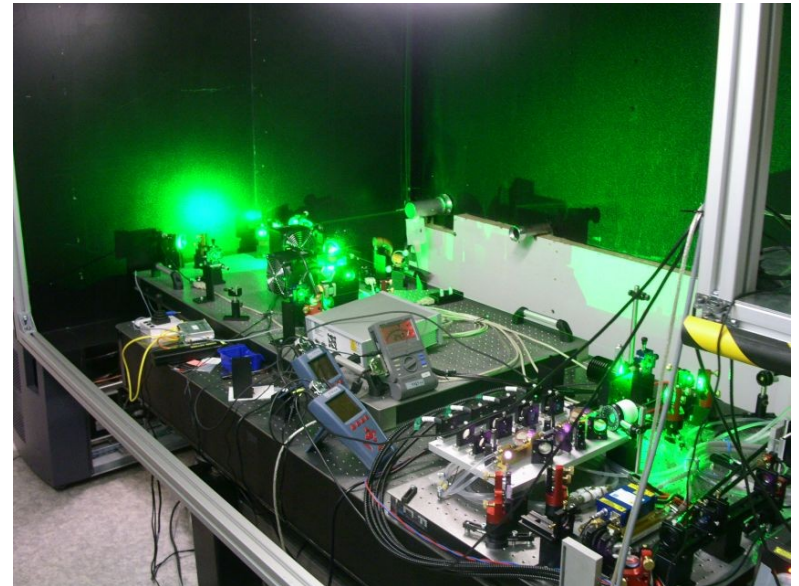


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# Search for ALPs via light-shining-through-walls

## > ALPS:

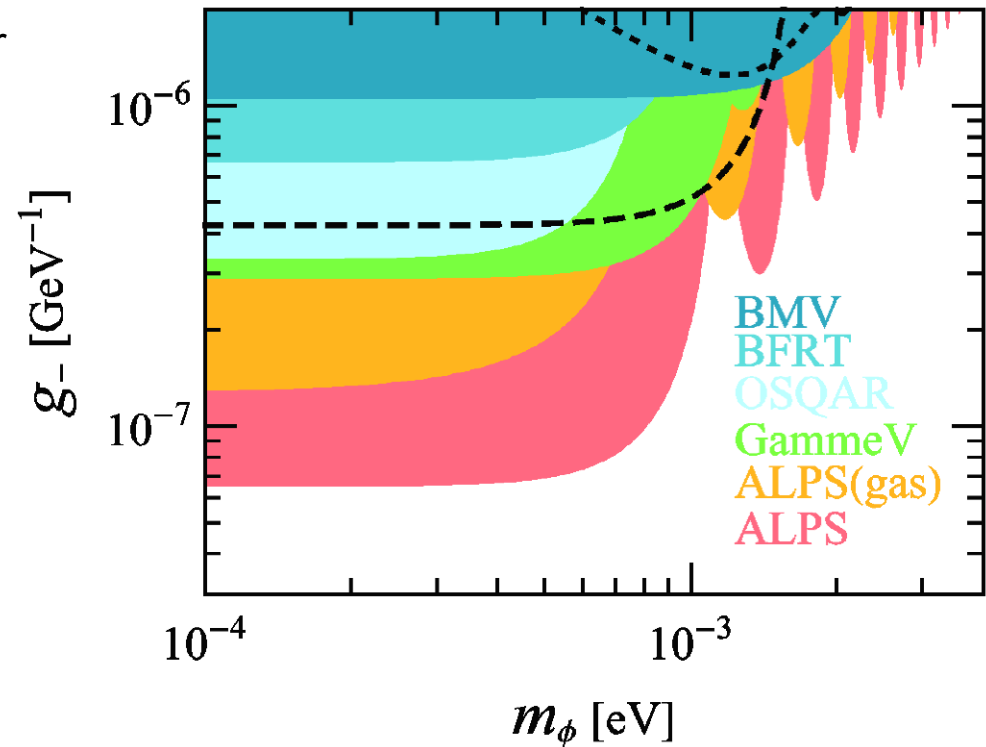
- HERA dipole (8.4 m, 5 T)
- Primary laser: enhanced LIGO laser (1064 nm, 35 W)
- Frequency doubled: 523 nm
- 300-fold power build-up in cavity



# Search for ALPs via light-shining-through-walls

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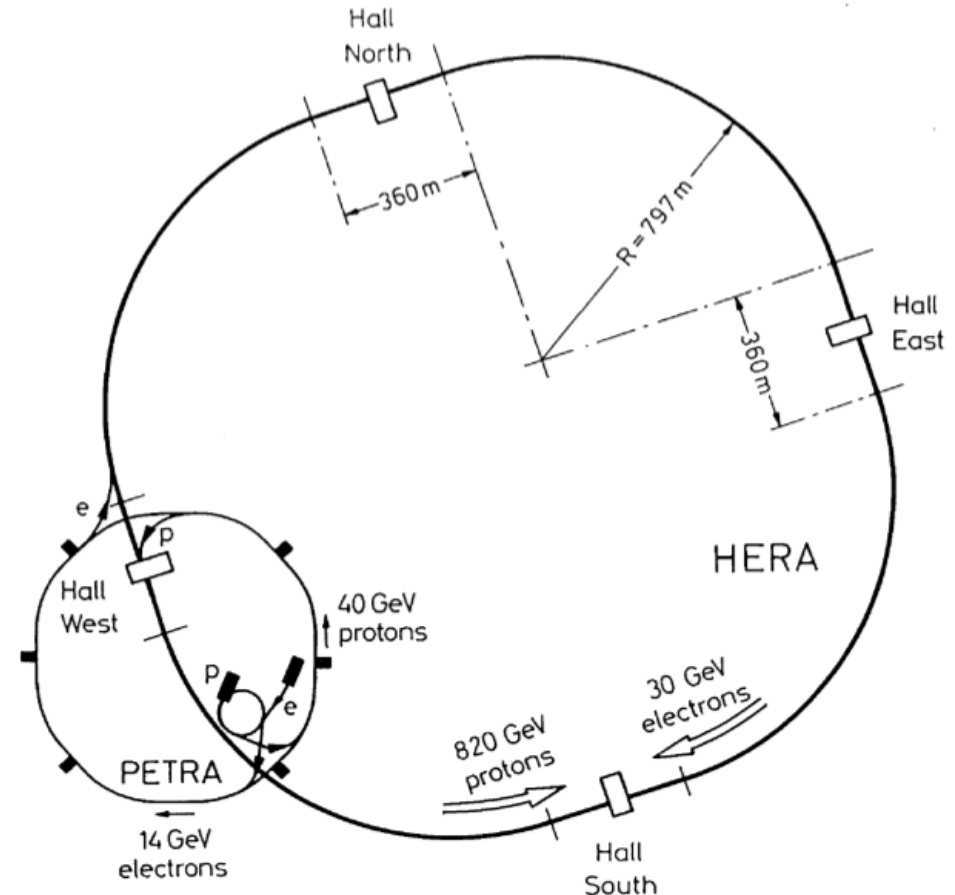
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## > ALPS-II plans (2016+):

- 12 + 12 HERA dipoles
- Increased power build-up ( $\sim 5000$ )
- Cavity also on regeneration part



# Search for ALPs via light-shining-through-walls

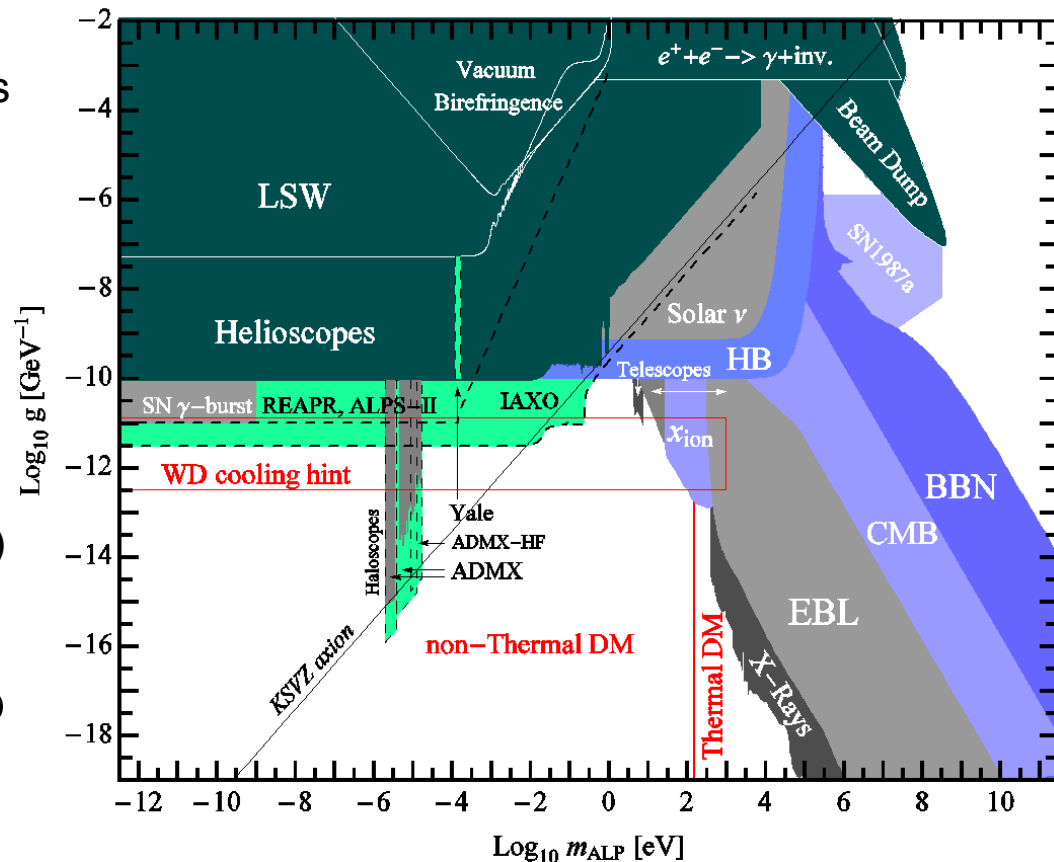
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## > ALPS-II plans (2016+):

- 12 + 12 HERA dipoles
- Increased power build-up (~5000)
- Cavity also on regeneration part

## > Similar plans also at Fermilab (REAPR)



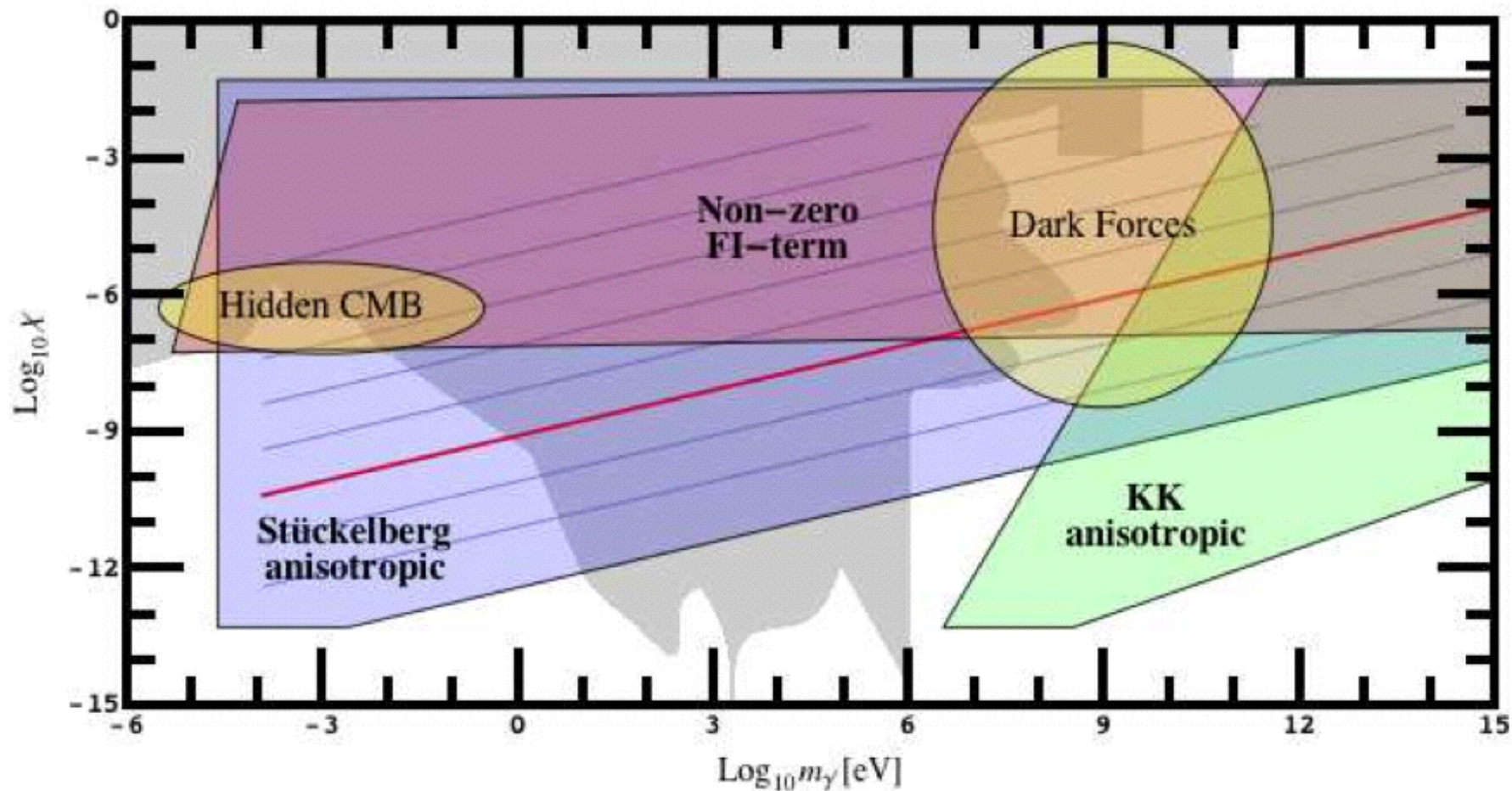
# Summary and conclusions

- String phenomenology holds the promise of an axiverse – the QCD axion plus a (possibly large) number of further ultralight axion-like particles, possibly populating each decade of mass down to the Hubble scale
- Promise fulfilled in LARGE Volume Scenario of IIB string compactifications
- Models that exhibit a QCD axion with an intermediate-scale decay constant and additional even lighter axion-like particles having the same decay constant and coupling to the photon can explain astrophysical anomalies and be tested in the next generation of helioscopes and light-shining-through-walls experiments
- Cosmology of LVS axiverse still to be investigated in detail



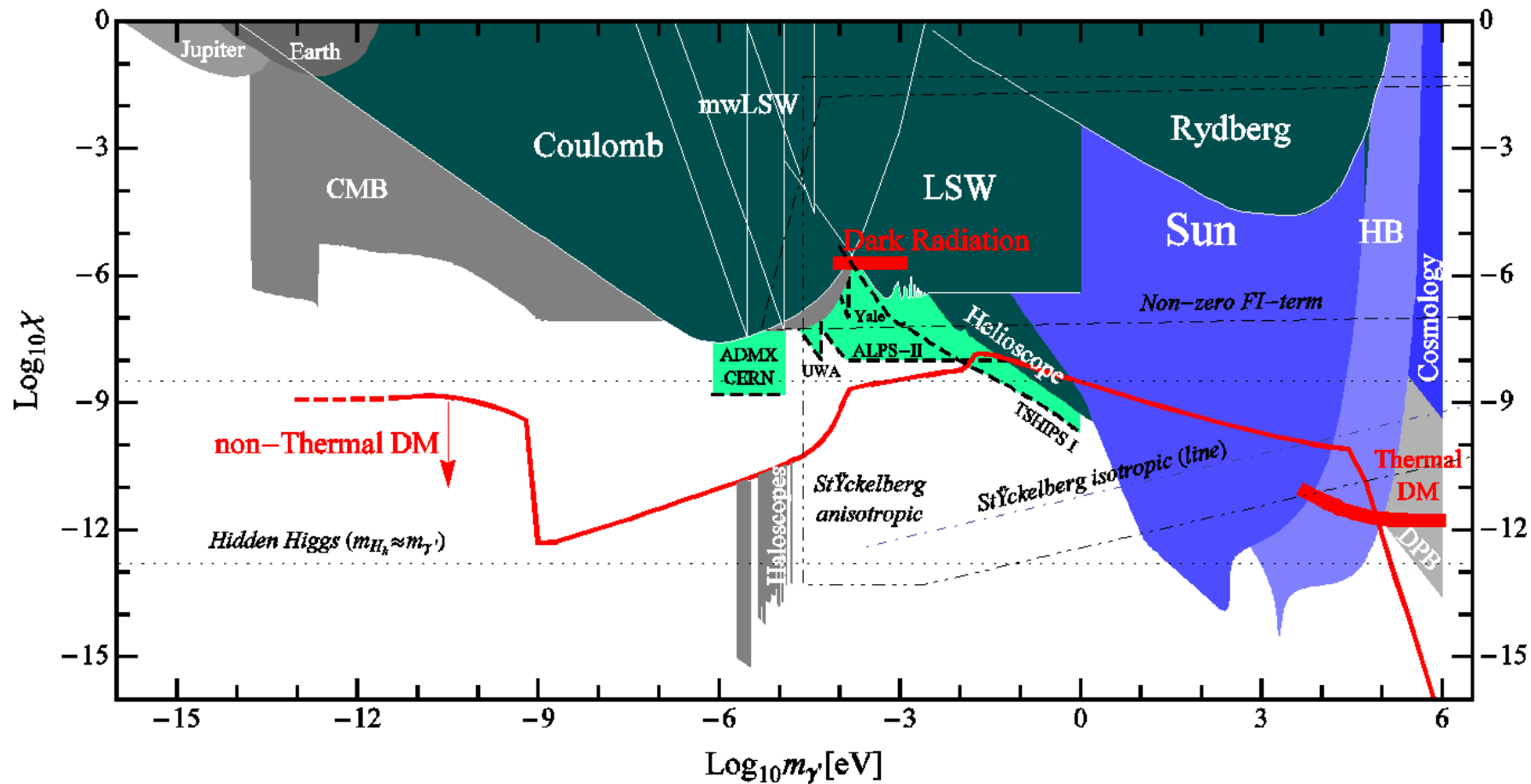


# Backup: Hidden photons

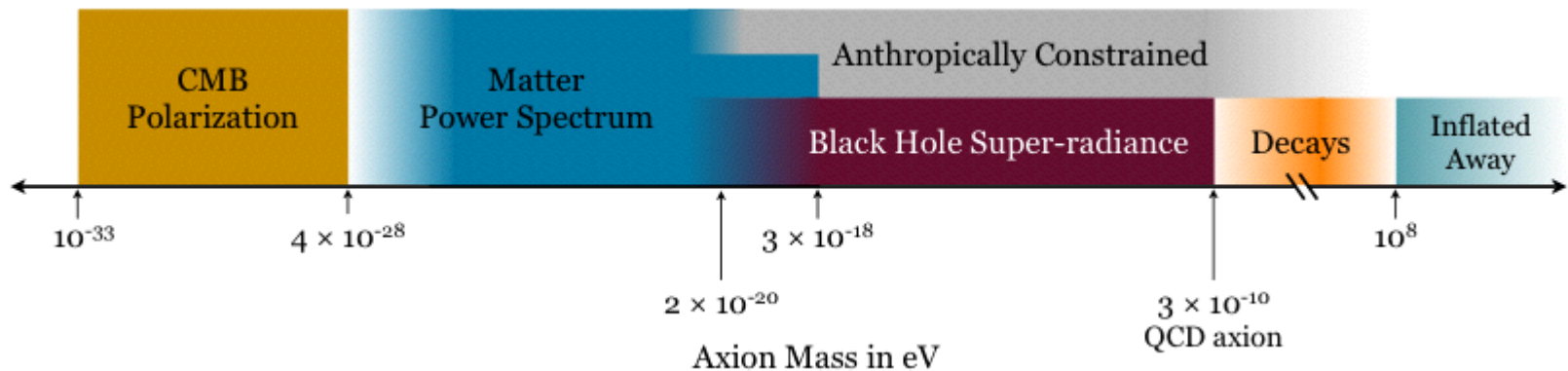




# Backup: Hidden photons



# Backup: More probes of ultralight axions and ALPs



cf. Arvanitaki, Dimopoulos, Dubovsky, Kaloper, March-Russell '09

