

# Probing Neutrino Dark Energy with Extremely High-Energy Cosmic Neutrinos

(astro-ph/0606316, accepted at JCAP, A. Ringwald, L.Schrempp)

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DESY Hamburg

- 1 Motivation
- 2 The Mass Varying Neutrino (MaVaN) Scenario
- 3 Signatures of Ultra-Energetic MaVaNs in the Sky?

# What is the nature of Dark Energy?

## Neutrino Dark Energy (Mass Varying Neutrinos)

[Fardon, Nelson, Weiner '03]

Idea of varying neutrino masses in other contexts

[Kawasaki, Murayama, Yanagida '92, Stephenson et al '97]

- Attractive scalar force between Big Bang relic neutrinos (the analog of the Cosmic Microwave Background (CMB) photons) → **smooth** background, can form a **negative pressure** fluid
- → acts as a form of Dark Energy → accelerated expansion
- → neutrino mass  $m_\nu$  becomes a function of neutrino energy density  $\rho_\nu(z)$  which evolves on cosmological timescales (here parametrized in terms of cosmic redshift  $z$ )

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# Mass Varying Neutrino (MaVaN) Scenario

## The non-SM neutrino interaction mediated by a scalar field

A concrete realization as preferred in the literature which implements the 'seesaw mechanism'

→ indirect coupling of neutrinos to the mediator of the scalar force, → light active neutrino masses

[FNW '03, FNW '05, Peccei '04, Barbieri, Hall, Oliver, Strumia '04, Takahashi, Tanimoto '05/'06...]

## Introduce the 'Dark Sector' (→ Standard Model singlets)

- A light scalar field, the acceleron  $\mathcal{A}$ , where  $H \sim 10^{-33} \text{eV} \ll m_{\mathcal{A}} < 10^{-4} \text{eV}$
- A scalar potential  $V_0(\mathcal{A})$
- Right-handed/sterile neutrinos  $N$
- Yukawa couplings  $\kappa \mathcal{A} N N \rightarrow M_N(\mathcal{A}) = \kappa \mathcal{A}$

## 'Our Sector' contains:

- Active neutrinos  $\nu$
- The Higgs field  $H$
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# 'Mini-Seesaw'-MaVaN model

## Light active neutrino masses $m_\nu$

- Scalar force of the 'Dark Sector' transmitted to the active neutrino sector via the **seesaw mechanism**
- Low energy effective theory for  $M_N(\mathcal{A}) \gg m_D$  (after integrating out the Higgs and the  $N$ )  $\rightarrow \mathcal{L} \supset \frac{m_D^2}{M_N(\mathcal{A})} \nu^2 + h.c. + V_0(\mathcal{A})$

[seesaw:Yanagida '79,Gell-Mann,Ramond '79,Minkowski' 80,Mohapatra,Senjanovic' 80]

Active neutrinos interact through a non-SM scalar force!

VEV of the accelaron,  $\mathcal{A}$ , generates light neutrino masses

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# Dynamics of the MaVaN Scenario

## Complex interplay between the neutrinos and the accelaron:

- Neutrino energy density  $\rho_\nu$  is a function of neutrino mass  $m_\nu(\mathcal{A})$   
 $\longrightarrow \rho_\nu(m_\nu(\mathcal{A}))$
- $\longrightarrow \rho_\nu(m_\nu(\mathcal{A}))$  can stabilize  $\mathcal{A}$  by contributing to its effective potential  $V_{\text{eff}}(\mathcal{A}) = \rho_\nu(m_\nu(\mathcal{A})) + V_0(\mathcal{A})$ , the total energy density of the coupled system
- Take cosmic expansion into account:  $\rho_\nu(\mathbf{z}, m_\nu(\mathcal{A})) \sim (1 + \mathbf{z})^3$
- $V_{\text{eff}}(\rho_\nu(\mathbf{z}, m(\mathcal{A})))$  **evolves** on cosmological time scales!
- For a homogeneous configuration, in the adiabatic limit  $H^2 \ll V''_{\text{eff}}(\mathcal{A}) = m_{\mathcal{A}}^2$ , the equilibrium value  $\mathcal{A}$  tracks the minimum of  $V_{\text{eff}}(\mathcal{A}, \mathbf{z})$

### Crucial effect:

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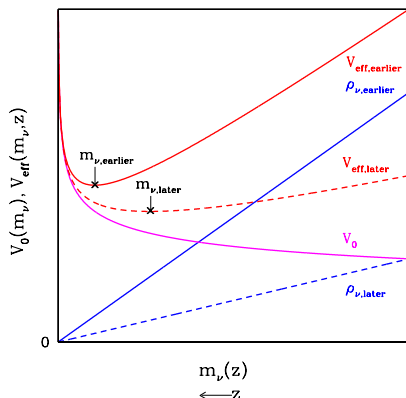
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- $\rightarrow \frac{\partial V_{\text{eff}}(\mathcal{A})}{\partial \mathcal{A}} = \frac{\partial m_\nu}{\partial \mathcal{A}} \frac{\partial V_{\text{eff}}(m_\nu)}{\partial m_\nu} = 0$
- Neutrino mass variation determined by  $V'_{\text{eff}}(m_\nu, \mathbf{z}) = 0 = \rho'(\mathbf{z}) + V'_0(m_\nu)$
- $\rightarrow$  in the non-relativistic limit,  
 $\rho_\nu(\mathbf{z}) = m_\nu n_\nu(\mathbf{z})$  ( $n_\nu(\mathbf{z}) \equiv$  neutrino number density)

- Neutrino acceleration fluid has dynamical Eq. of State

$$\omega(\mathbf{z}) + 1 = - \frac{m_\nu(\mathbf{z}) V'_0(m_\nu)}{V_{\text{eff}}(m_\nu(\mathbf{z}))}$$



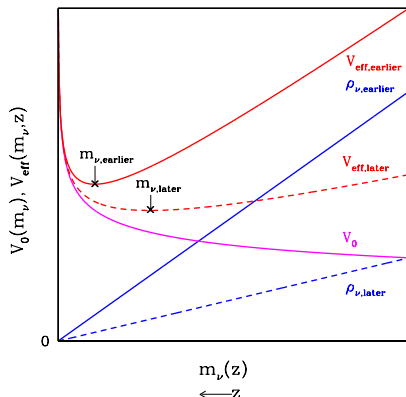
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## Possible Instabilities?

- Phenomenon similar to gravitational instabilities of cold dark matter possible for MaVaNs?
- Attractive scalar force mediated by the acceleron  $\mathcal{A}$
- For relativistic neutrinos, random motions prevent collapse
- However, when neutrinos turn non-relativistic formation of 'neutrino nuggets'?

Note: Outcome of instability inherently non-linear process

... but if nuggets really form, neutrino fluid redshifts similar to cold dark matter with  $\omega \sim 0 \leadsto -1 \rightarrow$  no acceleration

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## Ways out of catastrophic instability!

- Relativistic neutrino responsible for dark energy  $\rightarrow$

$$m_{\nu_1} \leq T_{\nu_0} = 1.69 \times 10^{-4} \text{ eV}$$

[Afshordi,Zaldarriaga,Kohri '05, Fardon,Nelson,Weiner '05]

- Certain constraints on the scalar acceleration potential  $V_0(\mathcal{A})$  can lead to stable MaVaN models also in the highly-non-relativistic regime

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# A Hybrid Model of MaVaNs

## Acceleration mechanism known from inflation

- Two light scalar fields  $\rightarrow$  one is kept in a metastable minimum due to the large value of the other
- Energy density stored in the false minimum drives accelerated expansion

[Linde '93]

## Supersymmetrization of MaVaN model naturally sets the stage!

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# MaVaN mass phenomenology

## MaVaN mass evolution with redshift, $m_\nu(z)$

- **Hybrid Model**  $\rightarrow$   
Microscopic origin for  
quadratic acceleration potential  
 $V_0(\mathcal{A}) \sim \mathcal{A}^2 \rightarrow V_0(m_\nu) \sim m_\nu^{-2}$

$V_0$  fixes mass evolution  $m_\nu(z)$ !

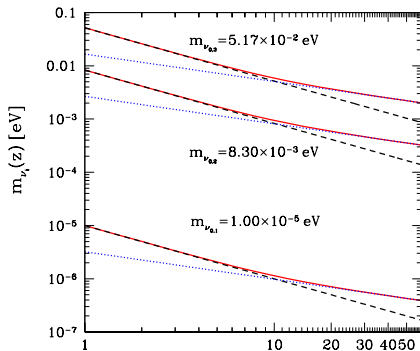
Recall:

$$V'_{\text{eff}}(m_\nu) = 0 = \rho'(m_\nu) + V'_0(m_\nu)$$

- Mass behaviour in the  
low-redshift regime  
approximated by a simple  
power law

$$m_{\nu_i} \approx m_{\nu_{0,i}} (1+z)^{-1},$$
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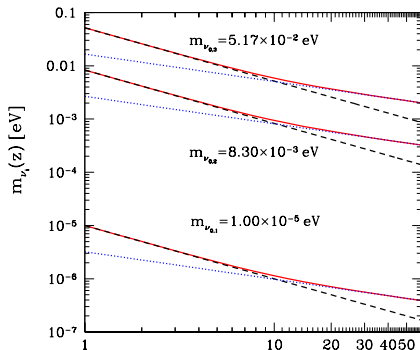
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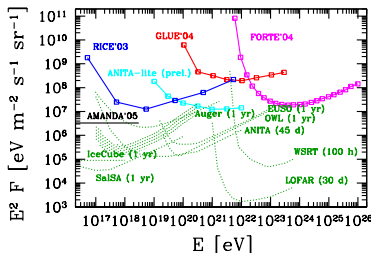
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# Extremely High Energy Cosmic Neutrinos (EHEC $\nu$ 's)

## EHEC $\nu$ 's in reach!

- Existing and planned observatories have a combined sensitivity of  
 $10^7 \text{ GeV} < E < 10^{17} \text{ GeV}$
- Appreciable event samples
- Fluxes to be detected at earth contain valuable information on  
particle physics and cosmology



## Absorption of EHEC $\nu$ 's by the Cosmic $\nu$ Background (C $\nu$ B)

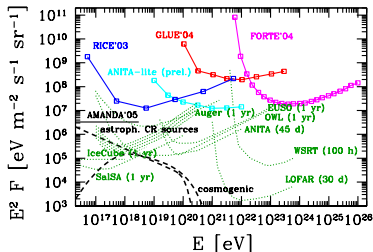
Most direct evidence for the existence of the C $\nu$ B so far  $\rightarrow$   
confirmation of standard cosmology back to  $\approx 1 \text{ s}$  after the Big Bang!

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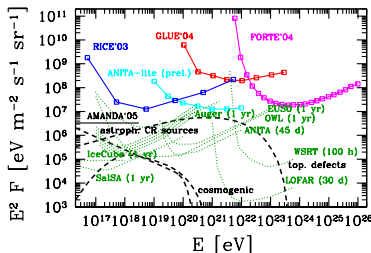
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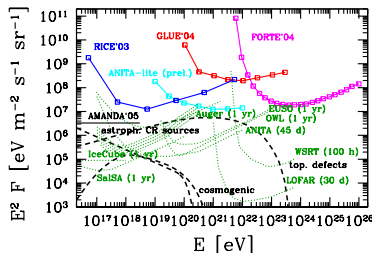
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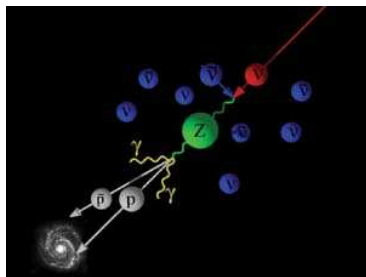
## The process $\nu_{\text{EHEC}} \bar{\nu}_{\text{relic}} \rightarrow Z$

- At the resonance energies

$$E_{\nu_i}^{\text{res}} = \frac{M_Z^2}{2m_{\nu_i}} \simeq 4 \times 10^{21} \text{ eV} \left( \frac{\text{eV}}{m_{\nu_i}} \right)$$

EHEC  $\nu$ 's can annihilate with relic  $\bar{\nu}_i$ 's (and vice versa) into  $Z$  bosons

- Exceptional loss of transparency of the C $\nu$ B for EHEC  $\nu$ 's  $\rightarrow$  **absorption dips** in the EHEC  $\nu$  spectra
- Annihilation process sensitive to  $m_{\nu_i}$  and possible time variation  $m_{\nu_i}(z)$



Relic Neutrino Absorption Spectroscopy as a Test of Neutrino Dark Energy (MaVaNs)!

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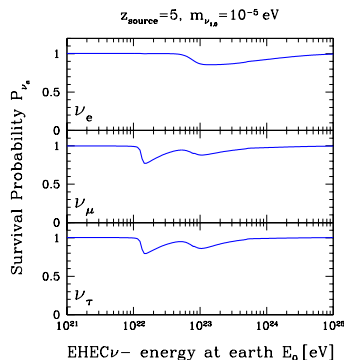
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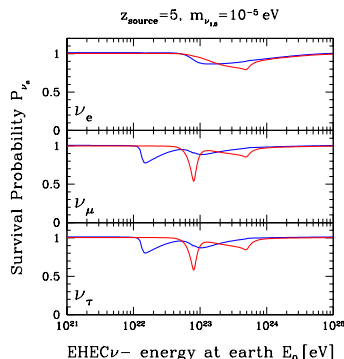
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# Relic Neutrino Absorption Spectroscopy

## Energy loss due to cosmic redshift!

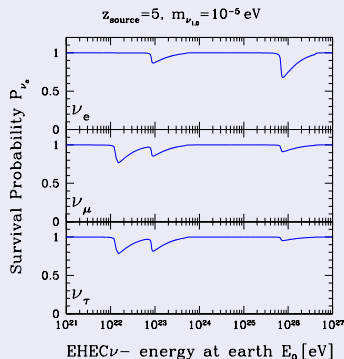
Source injects EHEC $\nu$  at  $z_s > 0$  with energy  $E \rightarrow$

EHEC $\nu$  arrives at earth at  $z = 0$  with energy  $E_0 = \frac{E}{1+z_s} < E$

## Interpretation of absorption dips for $m_\nu = \text{const.}$

As a first step, **assume** the relic neutrinos to be at rest,  $p_{\text{relic}} = 0$ , to disentangle different influences on the Survival Probability

- Why are the dips in the Survival Probability not sharply peaked at  $E \sim E_{\nu_i}^{\text{res}}$ ?
- Consider effect of cosmic redshift  $\sim (1+z)^{-1}$  on  $E \rightarrow$  Survival Probability  $P_{\nu_\alpha}$  reduced in the interval  $\frac{E_{\nu_i}^{\text{res}}}{1+z_s} \leq E_0 \leq E_{\nu_i}^{\text{res}} \rightarrow$  **redshift distortion**



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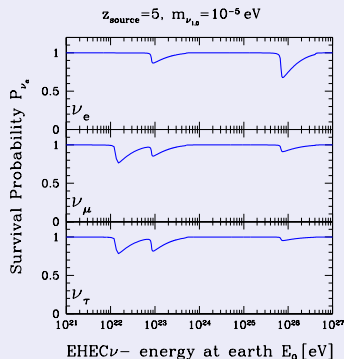
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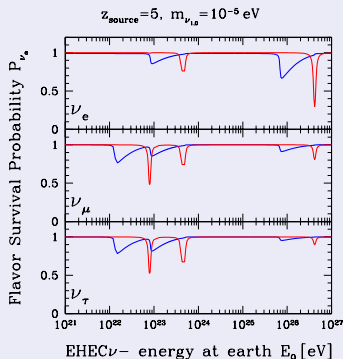
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? How does the mass variation  $m_\nu(z)$  affect the absorption features?

1) Recall:  $m_{\nu_i}(z) = m_{\nu_{0,i}} (1+z)^{-1}$  in the **low redshift** regime

2) Recall:  $E_{\nu_i}^{\text{res}} = \frac{M_Z^2}{2m_{\nu_i}} \rightarrow$  employ 1)

- $\rightarrow E_{\nu_i}^{\text{res}}(z) = E_{\nu_{0,i}}^{\text{res}} (1+z)$
- Mass induced  $z$  dependence **compensates** the effect of cosmic redshift  $\rightarrow$  all annihilations at  $0 \leq z \leq z_s$  contribute to the absorption peak at  $E_{\nu_{0,i}}^{\text{res}} = \frac{E_{\nu_i}^{\text{res}}(z)}{1+z}$



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EHEC $\nu$  arrives at earth at  $z = 0$  with energy  $E_0 = \frac{E}{1+z_s} < E$

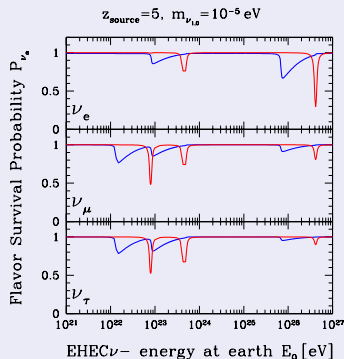
## Interpretation of absorption dips for $m_\nu(z) \neq \text{const.}$

? How does the mass variation  $m_\nu(z)$  affect the absorption features?

1) Recall:  $m_{\nu_i}(z) = m_{\nu_{0,i}} (1+z)^{-1}$  in the **low redshift** regime

2) Recall:  $E_{\nu_i}^{\text{res}} = \frac{M_Z^2}{2m_{\nu_i}} \rightarrow$  employ 1)

- $\rightarrow E_{\nu_i}^{\text{res}}(z) = E_{\nu_{0,i}}^{\text{res}} (1+z)$
- Mass induced  $z$  dependence **compensates** the effect of cosmic redshift  $\rightarrow$  all annihilations at  $0 \leq z \leq z_s$  contribute to the absorption peak at  $E_{\nu_{0,i}}^{\text{res}} = \frac{E_{\nu_i}^{\text{res}}(z)}{1+z}$



# Relic Neutrino Absorption Spectroscopy

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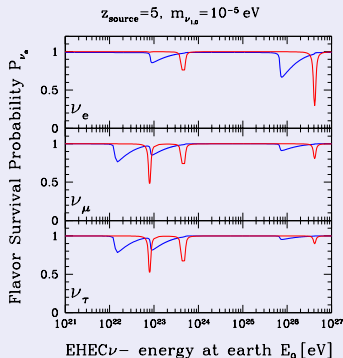
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# Relic Neutrino Absorption Spectroscopy

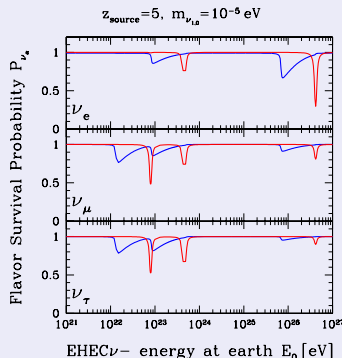
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# Relic Neutrino Absorption Spectroscopy

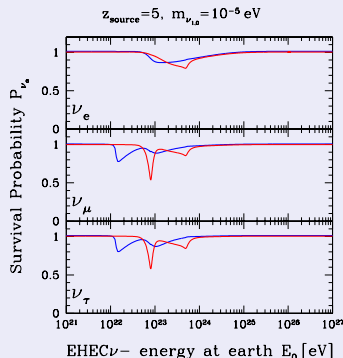
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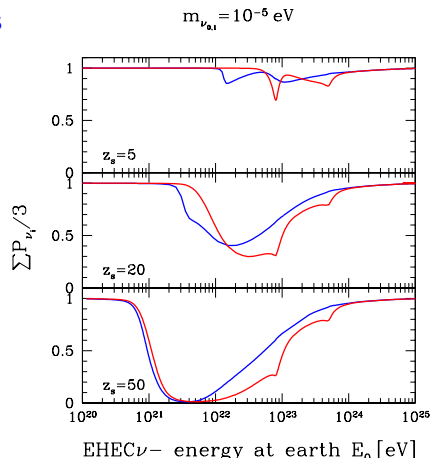


# Relic Neutrino Absorption Spectroscopy

## Sum of Survival Probabilities

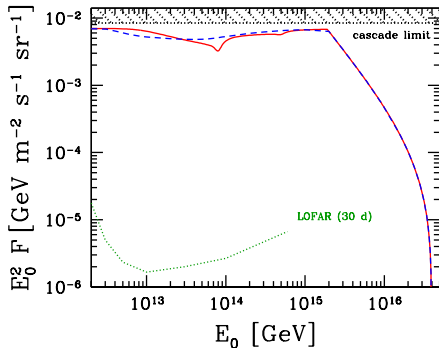
$$\frac{1}{3} \sum P_{\nu_i}$$

- What happens if the EHEC $\nu$  source is farther away ( $\rightarrow z_s$  larger)?
- $\rightarrow$  Dips get deeper and thermal broadening is more pronounced
- MaVaN absorption features (red) clearly shifted to higher energies
- MaVaN absorption features (red) clearly deeper  $\rightarrow$  better to detect

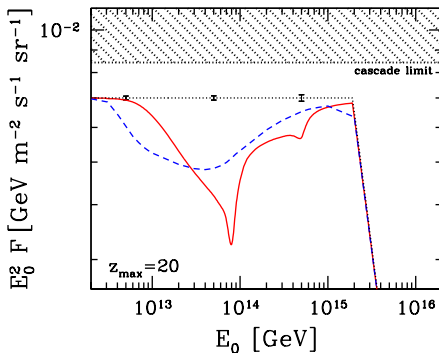


# Relic Neutrino Absorption Spectroscopy

Flavor Summed EHEC $\nu$  Flux  $\sum F_{\nu_\alpha}$  to be observed at Earth  
for an Astrophysical Source at  $z_s = 20$



[LOFAR sensitivity: O. Scholten et al. '06]

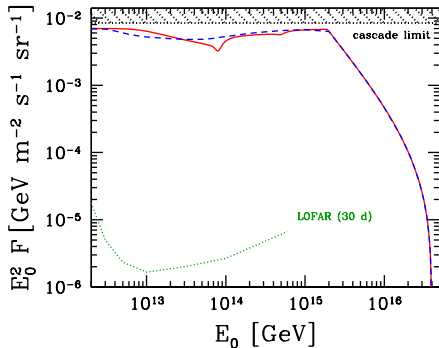


## Outlook

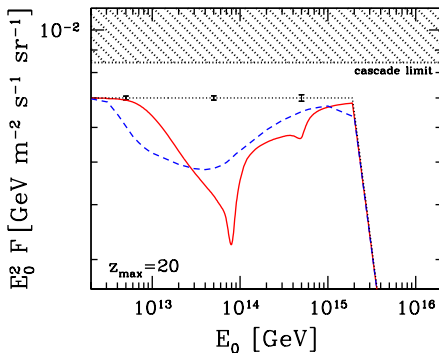
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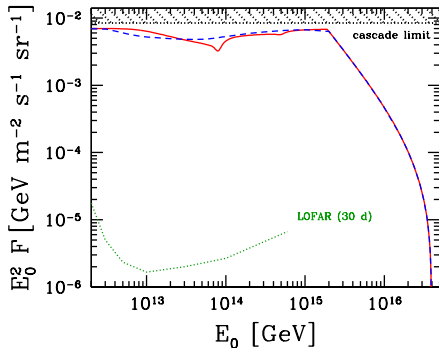


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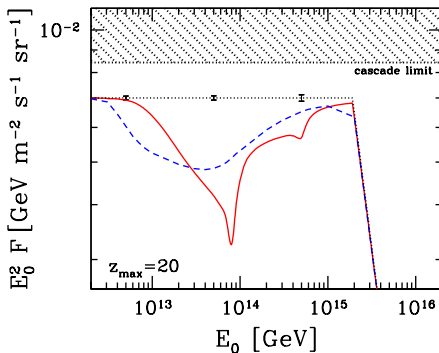
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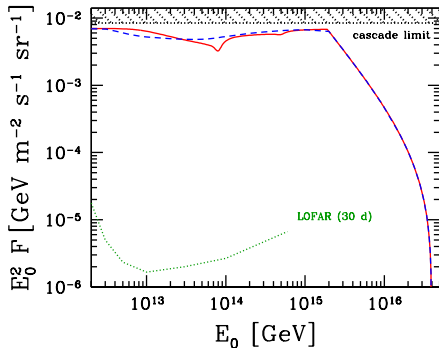


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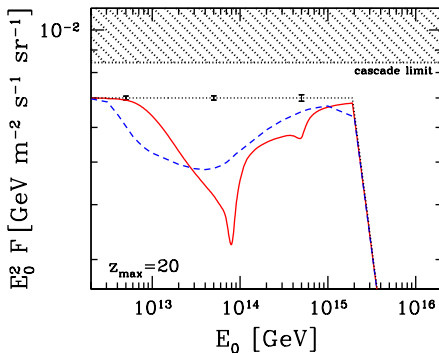
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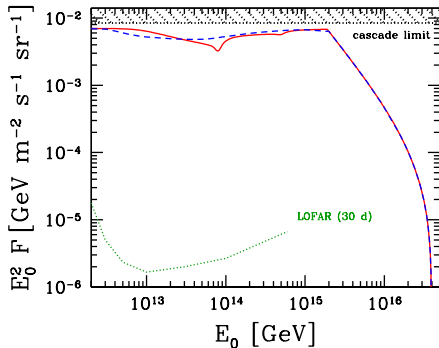


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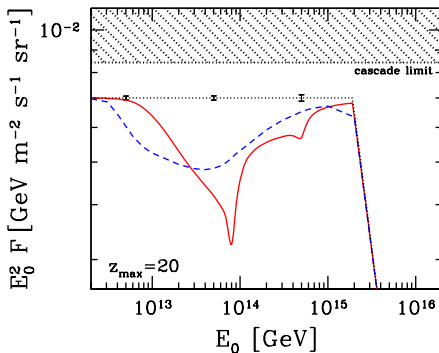
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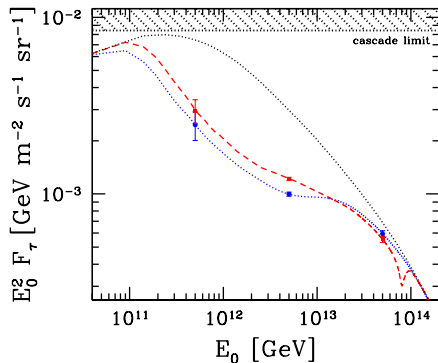
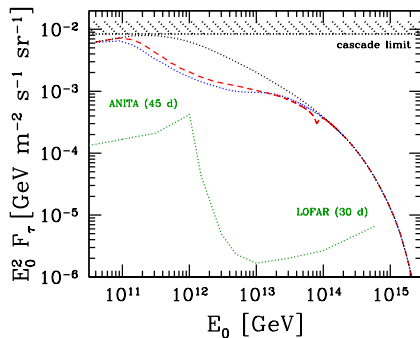
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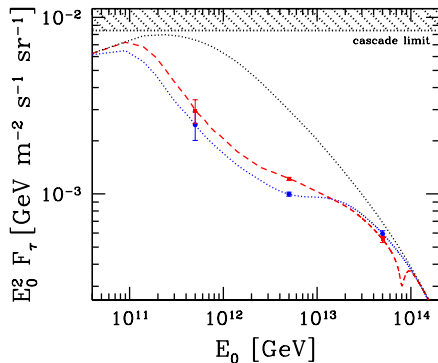
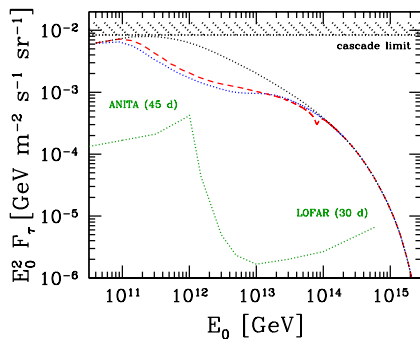
# Relic Neutrino Absorption Spectroscopy

Flux  $F_{\nu_\tau}$  to be observed at Earth for a Topological Defect Source



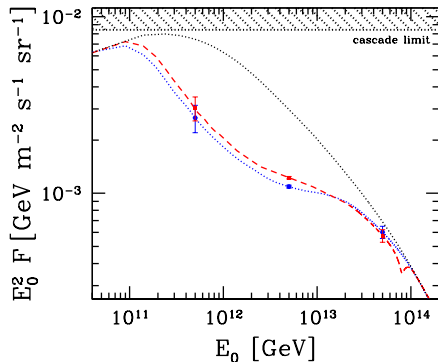
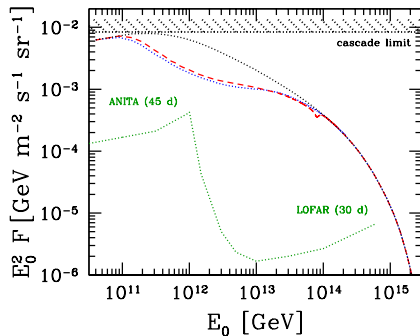
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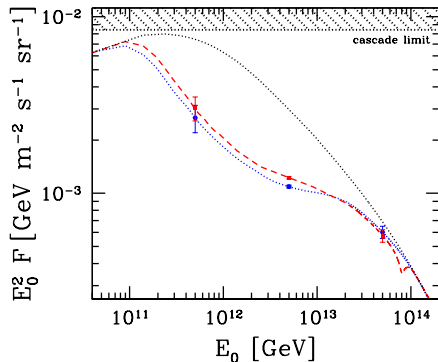
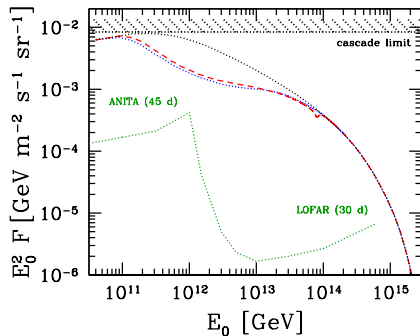
# Relic Neutrino Absorption Spectroscopy

Total EHEC $\nu$  Flux  $\sum F_{\nu_\alpha}$  to be observed at Earth for a TD Source



# Relic Neutrino Absorption Spectroscopy

Total EHEC $\nu$  Flux  $\sum F_{\nu_\alpha}$  to be observed at Earth for a TD Source



## Light active neutrino masses $m_\nu$ :

- Allow tiny coupling  $\lambda H N \nu$ ,  $\mathcal{O}(10^{-12})$ , linking ‘Our Sector’ to the ‘Dark Sector’
- At scales  $\ll 100$  GeV  $\rightarrow \mathcal{L} \supset m_D \nu N + M_N(\mathcal{A}) N N + h.c. + V_0(\mathcal{A})$ , with  $m_D = \lambda \langle H \rangle < 1$  eV
- Low energy effective theory  $\rightarrow \mathcal{L} \supset \frac{m_D^2}{M_N(\mathcal{A})} \nu^2 + h.c. + V_0(\mathcal{A})$  for  $M_N(\mathcal{A}) \gg m_D$  (after integrating out  $N$ )

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