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First exclusive reconstruction of the B^{*+} , B^{*0} , and B_s^{*0} mesons and precise measurement of their masses

The CMS Collaboration*

Abstract

Using proton-proton collision data collected by the CMS experiment at $\sqrt{s} = 13$ TeV in 2016–2018, corresponding to an integrated luminosity of 140 fb^{-1} , the first full reconstruction of the three vector B meson states, B^{*+} , B^{*0} , and B_s^{*0} , is performed. The mass differences between the excited mesons and their corresponding ground states are measured to be $m(B^{*+}) - m(B^+) = 45.277 \pm 0.039 \pm 0.027 \text{ MeV}$, $m(B^{*0}) - m(B^0) = 45.471 \pm 0.056 \pm 0.028 \text{ MeV}$, and $m(B_s^{*0}) - m(B_s^0) = 49.407 \pm 0.132 \pm 0.041 \text{ MeV}$, where the first uncertainties are statistical and the second are systematic. These results improve on the precision of previous measurements by an order of magnitude.

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While the ground states of beauty mesons B^+ , B^0 , and B_s^0 have been extensively studied, the corresponding excited vector mesons B^{*+} , B^{*0} , and B_s^{*0} are much less explored. The main challenge in the reconstruction of these excited states stems from the low-energy (40–50 MeV) photons emitted in $B^* \rightarrow B\gamma$ decays, with the energy below the detection thresholds in most experiments. Throughout this Letter, B^* corresponds to one of the B^{*+} , B^{*0} , or B_s^{*0} mesons, and B refers to one of the B^+ , B^0 , or B_s^0 mesons; the inclusion of charge-conjugate states is implied. About 30 years ago, the four LEP experiments combined b jets with photons reconstructed from their conversions into e^+e^- pairs to measure the flavor-averaged mass difference $m(B^*) - m(B)$ [1–4], also known as the hyperfine splitting. Studies of the P-wave B_s^0 mesons decaying into B^+K^- and $B^{*+}K^-$ by the LHCb and CMS experiments [5, 6], as well as to $B^0K_S^0$ and $B^{*0}K_S^0$ by CMS [6], provided the measurements of the mass differences $m(B^{*+}) - m(B^+)$ and $m(B^{*0}) - m(B^0)$ without reconstructing low-energy photons. Measurements at the dedicated B physics experiments CLEO and Belle, at the $\Upsilon(10860)$ energy, allowed indirect estimates of the B_s^{*0} meson mass [7, 8]; however, these two results are in tension with each other. Measurements of the properties of excited B mesons provide an important input to our understanding of the quantum chromodynamics (QCD) mechanisms responsible for quark dynamics and the formation of hadrons. In particular, precise measurements of the hyperfine splitting in the B meson system are an important input for quark models of hadrons [9–11].

In this Letter, the first exclusive reconstruction of the three excited vector mesons B^{*+} , B^{*0} , and B_s^{*0} , is presented. The analysis is performed using proton-proton (pp) collision data recorded by the CMS experiment during LHC Run 2 in 2016–2018 at $\sqrt{s} = 13$ TeV, corresponding to an integrated luminosity of 140 fb^{-1} [12–14]. Low-energy photons are reconstructed through their conversions into e^+e^- pairs in the beam pipe and detector material. Ground-state B mesons are reconstructed using the decays $B^+ \rightarrow \psi K^+$, $B^0 \rightarrow \psi K^{*0}$, and $B_s^0 \rightarrow \psi\phi$, where ψ stands for the J/ψ and $\psi(2S)$ mesons detected using their decays to $\mu^+\mu^-$, and K^{*0} and ϕ refer to the $K^*(892)^0$ and $\phi(1020)$ mesons, reconstructed through their decays to $K^+\pi^-$ and K^+K^- , respectively. Tabulated results are provided in the HEPData record for this analysis [15].

The CMS apparatus [16, 17] is a multipurpose, nearly hermetic detector, designed to trigger on [18–20] and identify electrons, muons, photons, and charged and neutral hadrons [21–23]. A global “particle-flow” algorithm [24] aims to reconstruct all individual particles in an event, combining information provided by the all-silicon tracker and by the crystal electromagnetic and brass-scintillator hadron calorimeters, operating inside a 3.8 T superconducting solenoid, with data from the gas-ionization muon detectors embedded in the flux-return yoke outside the solenoid. The procedure followed for aligning the detector is described in Ref. [25]. An improved pixel detector was installed before the 2017 data taking began [26].

Events of interest are selected using a two-tiered trigger system [18]. The first level (L1), composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of $4 \mu\text{s}$ [19]. The second level, known as the high-level trigger (HLT), consists of a farm of computing processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage. The events used in this analysis were selected at L1 by requiring the presence of at least two muons, and at the HLT by requiring two opposite-sign (OS) muons, with various pseudorapidity (η) and transverse momentum (p_T) thresholds, compatible with being produced in dimuon decays of either the J/ψ or $\psi(2S)$ meson.

The offline reconstruction requires two OS muons that form a vertex with fit χ^2 probability $P_{\text{vtx}} > 1\%$, each satisfying $p_T > 3.5 \text{ GeV}$ and $|\eta| < 2.5$ requirements, and passing a special

selection optimized for low- p_T muons [27]. The dimuon invariant mass $m(\mu^+\mu^-)$ must be within 2.9–3.4 GeV for the J/ψ decay channels, and 3.4–3.9 GeV for the $\psi(2S)$ decays. To build B candidates, the selected ψ candidate is combined with one or two OS high-purity tracks [23] with $p_T > 1$ GeV that do not match the selected muons. For the $B^+ \rightarrow \psi K^+$ and $B_s^0 \rightarrow \psi\phi$ selection, one or both tracks are assigned the charged kaon mass hypothesis, with an additional requirement of $1.01 < m(K^+K^-) < 1.03$ GeV for the latter. For the $B^0 \rightarrow \psi K^*(892)^0$ selection, the $K^+\pi^-$ or $K^-\pi^+$ mass assignment is made based on which combination has an invariant mass closer to the world-average K^{*0} meson mass of $M_{K^{*0}}^{\text{PDG}} = 895.55$ MeV [28], followed by the requirement $0.81 < m(K^\pm\pi^\mp) < 0.97$ GeV. In addition, a $\phi \rightarrow K^+K^-$ veto is applied in the $B^0 \rightarrow \psi K^*(892)^0$ channel to ensure that the selected $B_s^0 \rightarrow B_s^0\gamma$ candidates do not contribute to the $B^0\gamma$ sample.

The B^+ (B^0 and B_s^0) candidates are obtained via a 3-track (4-track) vertex fit of the selected muons and track(s), which constrains the dimuon invariant mass to the world-average J/ψ or $\psi(2S)$ mass [28]. From all reconstructed pp collision points, the primary vertex (PV) is chosen as the one with the smallest B pointing angle, defined as the angle between the B candidate momentum and the vector connecting the PV and the reconstructed B candidate vertex. Furthermore, if any of the 3 or 4 tracks used in the B candidate reconstruction is included in the fit of the chosen PV, they are removed and the PV position is refitted. The B candidates are required to have $p_T(B) > 10$ GeV, $P_{\text{vtx}}(B) > 1\%$, transverse vertex displacement from PV $L_{xy}(B)$ larger than 3 times its uncertainty, two-dimensional pointing angle $\cos[\vec{L}_{xy}(B), \vec{p}_T(B)] > 0.999$, and invariant mass within twice the average mass resolution of the corresponding B meson world-average mass M_B^{PDG} [28].

Photon conversion candidates are formed from two OS tracks with a small angular separation, a small distance of closest approach, a conversion vertex at least 1.5 cm away from the beam axis, and a χ^2 probability above 0.05% for a kinematic fit that constrains the two tracks to originate from a common vertex with an invariant mass of zero. More details on the selection criteria are given in Ref. [29]. Only photon candidates with $p_T > 300$ MeV and $|\eta| < 1.5$ are kept, since conversions at more forward pseudorapidity are reconstructed with a significantly worse resolution. To improve the resolution, the $B\gamma$ invariant mass is evaluated using the expression $M(B\gamma) = m(B\gamma) - m(B) + M_B^{\text{PDG}}$, where the B meson and the conversion photon four-momenta are taken from the kinematic vertex fit that constrains the B meson, the photon candidate, and all tracks forming the chosen PV, to originate from the same vertex [30–32].

The momenta of the e^+ and e^- tracks measured by the tracker tend to be lower than the true momenta because the electrons and positrons lose energy while traversing the tracker, which results in a measured photon energy lower on average than the true photon energy [29, 33]. A photon energy scale (PES) correction is performed using $\pi^0 \rightarrow \gamma\gamma$ decays reconstructed from two converted photons in several ranges of photon p_T and $|\eta|$. The energy correction factor in each range is given by the ratio between the nominal π^0 mass [28] and the mean π^0 mass in data obtained from fitting the $\gamma\gamma$ invariant mass distribution. The PES correction is 1.002–1.04 in the kinematic range relevant for this analysis, i.e., $p_T(\gamma) < 1.5$ GeV. Appendix A contains example fits to the $\gamma\gamma$ invariant mass distribution before and after the PES correction.

The PYTHIA 8.240 package [35] with the CP5 underlying event tune [36] is used to simulate the B^* meson production. The EVTGEN 1.6.0 [37] program models the decays $B^* \rightarrow B\gamma$, $B^+ \rightarrow \psi K^+$, $B^0 \rightarrow \psi K^*(892)^0$, and $B_s^0 \rightarrow \psi\phi$ followed by the $\psi \rightarrow \mu^+\mu^-$ decays [28]. Final-state photon radiation is described using PHOTOS 3.61 [38, 39]. The generated Monte Carlo (MC) events are then passed to a detailed GEANT4-based simulation [40] of the CMS detector. The simulated events are then processed by the same reconstruction algorithms as used for the

collision data. The simulation includes the effects of multiple pp interactions in the same or nearby bunch crossings (pileup) with a multiplicity distribution matching that observed in data. The simulated p_T and $|\eta|$ distributions of all three B^* mesons were adjusted using the corresponding B^{*+} meson distributions measured in data.

The $B\gamma$ mass resolution is found to depend strongly on $|\eta(\gamma)|$, ranging from about 0.7 MeV at $|\eta(\gamma)|$ close to zero to about 2 MeV at $|\eta(\gamma)|$ around 1.5. To improve the accuracy of the B^* meson mass measurements, a simultaneous fit of the $B\gamma$ mass distributions in several categories, defined by ranges of $|\eta(\gamma)|$, is performed. Pseudo-experiments were used to obtain a set of optimal $|\eta(\gamma)|$ ranges for each meson flavor. The resulting simultaneous fit structure includes: three $|\eta(\gamma)|$ ranges with boundaries at 0, 0.6, 1.0, and 1.4 for the B^{*+} signal; two $|\eta(\gamma)|$ ranges with boundaries at 0, 0.9, and 1.5 for the B^{*0} signal, and two $|\eta(\gamma)|$ ranges with boundaries at 0, 0.8, and 1.4 for the B_s^{*0} signal.

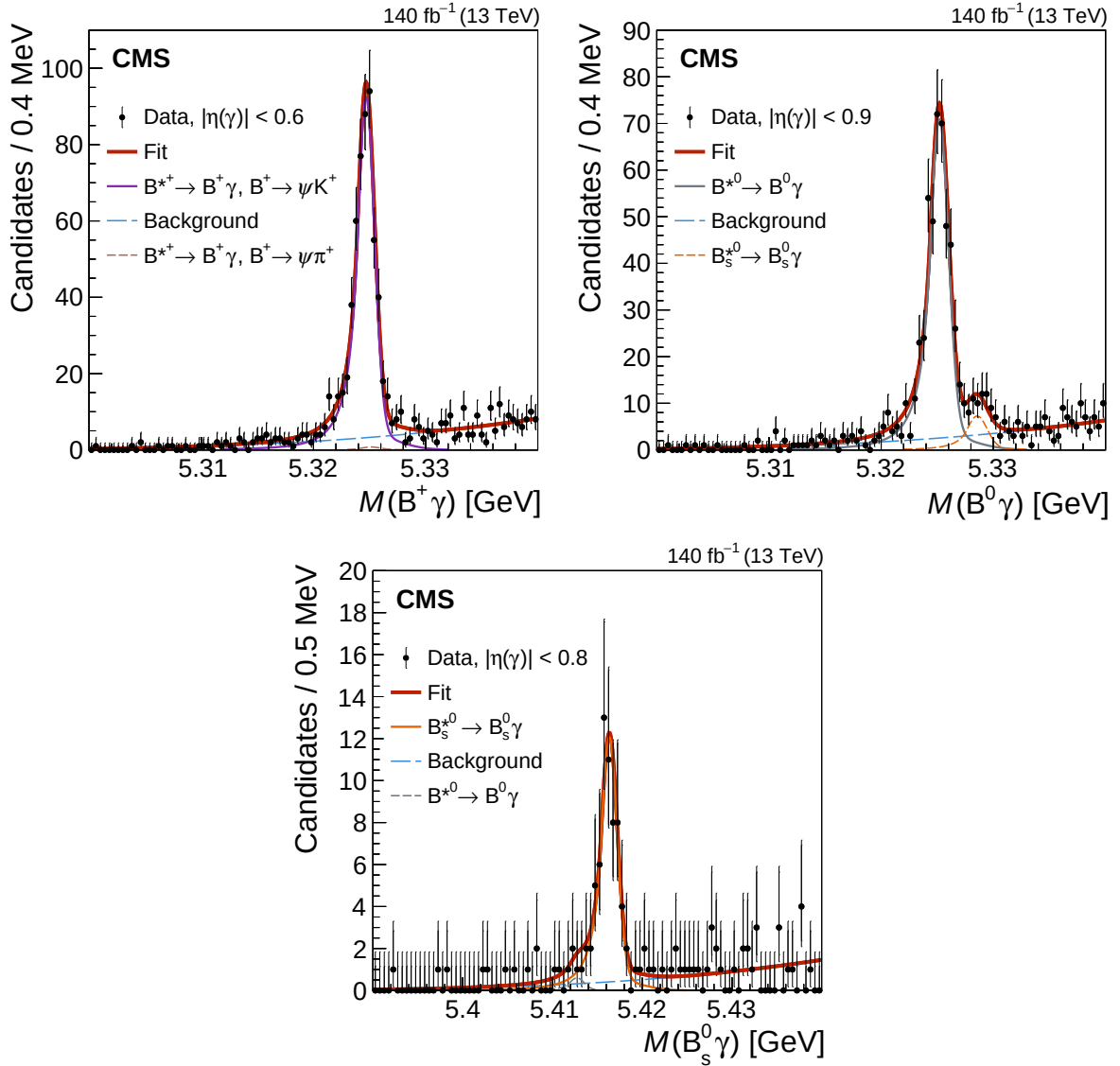


Figure 1: The measured distributions (filled circles) of the $B^+\gamma$ (upper left), $B^0\gamma$ (upper right), and $B_s^0\gamma$ (lower) invariant mass in the lowest $|\eta(\gamma)|$ range, which has the best invariant mass resolution. The simultaneous fit projections are shown with thick red lines, and the other components are described in the legends. The error bars represent statistical uncertainties in data.

The measured invariant mass distributions of the selected $B^+\gamma$, $B^0\gamma$, and $B_s^0\gamma$ candidates are shown in Fig. 1 for the three categories with the best invariant mass resolution and in Appendix B for the four remaining ones. A clear signal is seen in each of them. A simultaneous extended unbinned maximum likelihood fit to the seven distributions is performed. Each signal is described by the sum of two Crystal Ball (CB) [41] functions, with common mean and tail parameters n and α . The three mean values, which are free parameters of the fit, correspond to the measured masses of the B^{*+} , B^{*0} , and B_s^{*0} mesons. The parameters n , α , and the relative contribution of the two CB functions (f) are fixed to the ones obtained from simulation. The widths of the CB functions (σ_1 , σ_2) are multiplied by a common scale factor, which is a free parameter, to account for a possible data-simulation difference in the mass resolution. For the B^{*0} (B_s^{*0}) signals, the yield in the high $|\eta(\gamma)|$ range is set to be 41% (49%) of the yield in the low $|\eta(\gamma)|$ range, as found in simulation. No such constraint is applied to the B^{*+} signal since it has been used to correct the p_T and $|\eta|$ distributions in simulation. The peaking contribution from the Cabibbo-suppressed $B^+ \rightarrow \psi\pi^+$ decay in the $B^+\gamma$ mass spectrum is also modeled with the sum of two CB functions, with the shape parameters fixed to those found in simulation and the yield fixed relative to that for the B^{*+} signal, taking into account their relative branching fractions [28] and reconstruction efficiencies. In the $B^0\gamma$ mass distribution, the peaking contribution from $B_s^{*0} \rightarrow B_s^0\gamma \rightarrow \psi K^+ K^- \gamma$ decays (with one of the B_s^0 daughter kaons reconstructed as a pion candidate from the $B^0 \rightarrow \psi K^+ \pi^-$ decay) is described by the sum of two CB functions with the shape parameters of the B_s^{*0} signal and the mass shifted by -87.78 MeV with respect to it, as found in the simulation. The normalization of this component is left free in the fit. In the $B_s^0\gamma$ mass distribution, the contribution from the B^{*0} decays is described by the sum of two CB functions with the shape parameters of the B^{*0} signal and the mass shifted by $+87.54$ MeV with respect to it. Its normalization is fixed to 0.7% of the B^{*0} signal, according to the studies based on simulated event samples. Finally, the background in all distributions is described with a standard threshold function $(m - m_0)^\beta$, where m_0 is specific to each meson flavor, set to the corresponding ground-state meson mass, and β is a free parameter, shared across the seven categories.

The fit returns the mass differences $\Delta m(B^{*+}) \equiv m(B^{*+}) - m(B^+) = 45.277 \pm 0.039$ MeV, $\Delta m(B^{*0}) \equiv m(B^{*0}) - m(B^0) = 45.471 \pm 0.056$ MeV, and $\Delta m(B_s^{*0}) \equiv m(B_s^{*0}) - m(B_s^0) = 49.41 \pm 0.13$ MeV, where the uncertainties are statistical only.

The measurement uncertainties are dominated by the statistical component. The leading systematic uncertainties are discussed below. To evaluate the systematic uncertainties related to the choice of the fit model, several alternative functions are tested, with different n or α in the CB functions, the sum of CB and Student's t distributions, a CB function plus a Gaussian function, or a CB function plus a Novosibirsk [42] function. In order to evaluate the uncertainties reflecting the simulated signal shape, an alternative fit is performed, no longer fixing the shape parameters but imposing that they are common to the (simultaneously fitted) measured and simulated distributions. The nominal fit fixes the yield ratios between the $|\eta(\gamma)|$ categories for B^{*0} and B_s^{*0} signals to the values expected from simulation; alternatively, these constraints are removed. Variations of the $B^+ \rightarrow J/\psi\pi^+$ contribution produce negligible changes in the results. The peak positions and yields of the B_s^{*0} to $B^0\gamma$ and B^{*0} to $B_s^0\gamma$ cross-feed contributions are varied. The alternative shape for the combinatorial background has seven independent β parameters, one for each $|\eta(\gamma)|$ category. For each of these sets of variations, the largest deviation in the measured quantities, with respect to the results of the baseline fit, is taken as the corresponding systematic uncertainty.

A cross-check is performed by splitting the data into subsamples matching each of the three data-taking years (2016, 2017, and 2018) and measuring the corresponding B^{*+} mass values.

This cross-check also provides a good estimate of the uncertainty due to the detector alignment, since the pixel detector was replaced between 2016 and 2017. The results are found to be in good agreement with each other and the baseline result.

The statistical uncertainty from the PES correction is evaluated by making alternative PES corrections, with a momentum scale factor in each $|\eta(\gamma)|$ and $p_T(\gamma)$ range distributed as a Gaussian distribution, with mean equal to the baseline correction scale factor and RMS equal to its uncertainty. The simultaneous fit to the data is repeated each time, producing a set of $\Delta m(B^*)$ measurements. The RMS values of these distributions are taken as the systematic uncertainties in the measured mass differences. The systematic uncertainty from the PES correction is evaluated by comparing the weighted average of the π^0 mass values in the $|\eta(\gamma)|$ and $p_T(\gamma)$ ranges relevant for the analysis with the world-average value of the π^0 mass and scaling the difference by the ratio of $\Delta m(B^*)$ to the π^0 mass.

All the systematic uncertainties described above are summarized in Table 1, together with the total systematic uncertainty, calculated as the sum in quadrature of the individual components. Systematic uncertainties in the various ratios or differences between the three reported mass differences are evaluated in the same way and are given in Appendix C.

Table 1: Systematic uncertainties in the measured mass differences, in keV.

Source	$m(B^{*+}) - m(B^+)$	$m(B^{*0}) - m(B^0)$	$m(B_s^{*0}) - m(B_s^0)$
Signal model	4	8	21
Signal shape parameters	15	15	18
Yield ratios of $ \eta(\gamma) $ ranges	1	2	10
Background shape	2	<1	7
Cross-feed $B_s^{*0} \leftrightarrow B^{*0}$	<1	1	10
Photon energy scale, statistical	12	14	16
Photon energy scale, systematic	18	18	19
Total	27	28	41

All the parameters measured in this analysis are presented in Table 2. The mass differences (entries 1–3 in Table 2) are consistent with, and an order of magnitude more precise than, the current world-average values [28]. The measured mass difference $m(B_s^{*0}) - m(B_s^0)$ is in agreement with the Belle result [8], and about 2 standard deviations away from the CLEO result [7]. The masses of the vector states (rows 4–6) are determined by adding the measured difference between the excited and ground states to the world-average values of the ground-state masses [28].

In addition, the measurements of the differences between the vector state masses (rows 7–10 in Table 2), are reported, using the known mass differences between the B meson ground states [28]: $m(B^0) - m(B^+) = 0.31 \pm 0.05 \text{ MeV}$ and $m(B_s^0) - [m(B^0) + m(B^+)]/2 = 87.37 \pm 0.12 \text{ MeV}$, assuming their uncertainties to be uncorrelated. The result for $m(B^{*0}) - m(B^{*+})$ is consistent with the only previous measurement that used P-wave B_s^0 mesons [6], and is about 3 times more precise, while the other mass differences between vector states are measured for the first time. The differences and ratios between the mass differences are also measured for the first time (entries 11–18 in Table 2, which reflect the correlations between different uncertainties).

The measurements of $\Delta m(B_s^{*0})$, $\Delta m(B^{*+})$, and $\Delta m(B^{*0})$ agree with the predictions from lattice QCD [43–45] for $\Delta m(B_s^{*0})$ and $\Delta m(B^{*+,0}) \equiv m(B^{*+,0}) - m(B^{+,0})$ (where $B^{+,0}$ corresponds to either B^+ or B^0 since the predictions do not distinguish between the two), while the uncer-

Table 2: Results of the measurement. The first uncertainties are statistical, the second are systematic, and the third, where present, are related to the masses or mass differences of the ground states taken from Ref. [28].

	Parameter	Value
1	$\Delta m(\mathbf{B}^{*+}) \equiv m(\mathbf{B}^{*+}) - m(\mathbf{B}^+)$	$45.277 \pm 0.039 \pm 0.027 \text{ MeV}$
2	$\Delta m(\mathbf{B}^{*0}) \equiv m(\mathbf{B}^{*0}) - m(\mathbf{B}^0)$	$45.471 \pm 0.056 \pm 0.028 \text{ MeV}$
3	$\Delta m(\mathbf{B}_s^{*0}) \equiv m(\mathbf{B}_s^{*0}) - m(\mathbf{B}_s^0)$	$49.407 \pm 0.132 \pm 0.041 \text{ MeV}$
4	$m(\mathbf{B}^{*+})$	$5324.69 \pm 0.04 \pm 0.03 \pm 0.07 \text{ MeV}$
5	$m(\mathbf{B}^{*0})$	$5325.19 \pm 0.06 \pm 0.03 \pm 0.08 \text{ MeV}$
6	$m(\mathbf{B}_s^{*0})$	$5416.34 \pm 0.13 \pm 0.04 \pm 0.10 \text{ MeV}$
7	$m(\mathbf{B}^{*0}) - m(\mathbf{B}^{*+})$	$0.50 \pm 0.07 \pm 0.01 \pm 0.05 \text{ MeV}$
8	$m(\mathbf{B}_s^{*0}) - m(\mathbf{B}^{*+})$	$91.66 \pm 0.14 \pm 0.03 \pm 0.12 \text{ MeV}$
9	$m(\mathbf{B}_s^{*0}) - m(\mathbf{B}^{*0})$	$91.15 \pm 0.14 \pm 0.03 \pm 0.12 \text{ MeV}$
10	$m(\mathbf{B}_s^{*0}) - \frac{1}{2} [m(\mathbf{B}^{*0}) + m(\mathbf{B}^{*+})]$	$91.40 \pm 0.13 \pm 0.03 \pm 0.12 \text{ MeV}$
11	$\Delta m(\mathbf{B}^{*0}) - \Delta m(\mathbf{B}^{*+})$	$0.19 \pm 0.07 \pm 0.01 \text{ MeV}$
12	$\Delta m(\mathbf{B}_s^{*0}) - \Delta m(\mathbf{B}^{*+})$	$4.13 \pm 0.14 \pm 0.03 \text{ MeV}$
13	$\Delta m(\mathbf{B}_s^{*0}) - \Delta m(\mathbf{B}^{*0})$	$3.94 \pm 0.14 \pm 0.03 \text{ MeV}$
14	$\Delta m(\mathbf{B}_s^{*0}) - \frac{1}{2} [\Delta m(\mathbf{B}^{*0}) + \Delta m(\mathbf{B}^{*+})]$	$4.03 \pm 0.13 \pm 0.03 \text{ MeV}$
15	$\Delta m(\mathbf{B}^{*0}) / \Delta m(\mathbf{B}^{*+})$	$1.0043 \pm 0.0015 \pm 0.0002$
16	$\Delta m(\mathbf{B}_s^{*0}) / \Delta m(\mathbf{B}^{*+})$	$1.0912 \pm 0.0031 \pm 0.0007$
17	$\Delta m(\mathbf{B}_s^{*0}) / \Delta m(\mathbf{B}^{*0})$	$1.0866 \pm 0.0031 \pm 0.0007$
18	$2 \Delta m(\mathbf{B}_s^{*0}) / [\Delta m(\mathbf{B}^{*+}) + \Delta m(\mathbf{B}^{*0})]$	$1.0889 \pm 0.0030 \pm 0.0007$

tainties in the predictions are 1–2 orders of magnitude larger than those of the measurement. Phenomenological models provide more precise predictions including $m(\mathbf{B}^{*0}) - m(\mathbf{B}^{*+}) = 0.6 \pm 0.2$ [10] and $0.31 \pm 0.07 \text{ MeV}$ [11], close to our measurement (entry 7 in Table 2). Reference [10] also predicts hyperfine splittings $\Delta m(\mathbf{B}^{*+,0}) = 45.69 \pm 0.02 \text{ MeV}$ (above both values 1 and 2 in Table 2) and $\Delta m(\mathbf{B}_s^{*0}) = 46.73 \pm 0.08 \text{ MeV}$ (below our measurement in line 3 of Table 2). Experimental systematic uncertainties are reduced by taking the ratios of the mass differences, as shown in entries 16–18 in Table 2. In addition, the theoretical uncertainties in the lattice QCD predictions from Ref. [44] are also reduced, giving a prediction of $\Delta m(\mathbf{B}_s^{*0}) / \Delta m(\mathbf{B}^{*+,0}) = 1.007 \pm 0.034$, which is about 2.5 standard deviations below our results and with an uncertainty that remains an order of magnitude larger.

In summary, three vector B meson states, $\mathbf{B}^{*+}$, $\mathbf{B}^{*0}$, and $\mathbf{B}_s^{*0}$, have been fully reconstructed in exclusive final states for the first time by detecting low-energy photons through their conversion to e^+e^- pairs in the beam pipe and detector material. The data sample of $\sqrt{s} = 13 \text{ TeV}$ proton-proton collisions was collected by the CMS experiment and corresponds to an integrated luminosity of 140 fb^{-1} . The measurements benefit from a new photon energy scale calibration method that uses $\pi^0 \rightarrow \gamma\gamma$ decays in which both photons convert to e^+e^- pairs. The masses of the three vector states are measured with respect to the corresponding ground states to be $m(\mathbf{B}^{*+}) - m(\mathbf{B}^+) = 45.277 \pm 0.039 \pm 0.027 \text{ MeV}$, $m(\mathbf{B}^{*0}) - m(\mathbf{B}^0) = 45.471 \pm 0.056 \pm 0.028 \text{ MeV}$, and $m(\mathbf{B}_s^{*0}) - m(\mathbf{B}_s^0) = 49.407 \pm 0.132 \pm 0.041 \text{ MeV}$, where the first uncertainties are statistical and the second are systematic. A number of difference and ratio measurements between the reported masses and mass differences are also provided, where, due to cancellations in the computations of these quantities, both experimental and theoretical uncertainties are lower. These results are either the first measurement or an order of magnitude more precise than previous measurements. The new measurements are more precise than the available theoretical

predictions and provide an important input to our understanding of heavy-quark systems.

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A Example distributions of the $\pi^0 \rightarrow \gamma\gamma$ invariant mass

The distributions of the $\gamma\gamma$ invariant mass from two converted photons in 2018 data, showing the $\pi^0 \rightarrow \gamma\gamma$ peak in the kinematic range $0.8 < p_T(\gamma) < 0.9 \text{ GeV}$, $0.8 < |\eta(\gamma)| < 0.9$, are presented in Fig. A.1 before (left) and after (right) the PES correction. The distributions are fit with a sum of two Crystal Ball functions with a common mean for the signal and a polynomial multiplied by a power law function for the background.

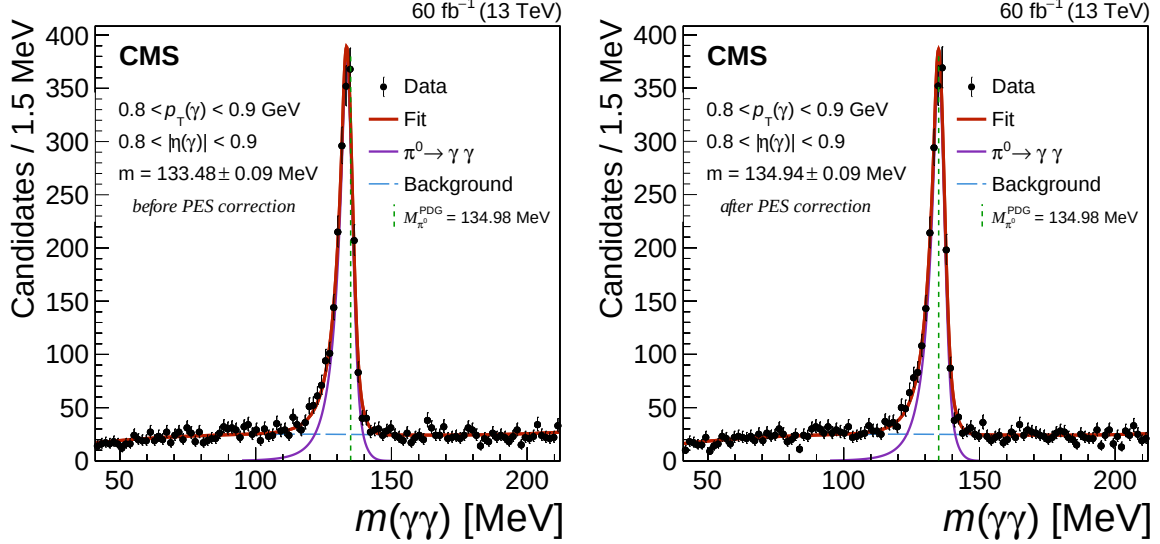


Figure A.1: The measured distributions of the diphoton invariant mass near the π^0 mass with the fit results overlaid, before (left) and after (right) the PES correction, in the kinematic range $0.8 < p_T(\gamma) < 0.9 \text{ GeV}$, $0.8 < |\eta(\gamma)| < 0.9$. The green vertical line shows the world-average value of the π^0 mass of 134.98 MeV [28].

B Fit projections in extra pseudorapidity ranges

The measured invariant mass distributions of the $B^+\gamma$, $B^0\gamma$, and $B_s^0\gamma$ candidates in the four additional categories of the simultaneous fit are shown in Fig. B.1.

C Systematic uncertainties in the differences and ratios of the mass differences

The systematic uncertainties for the differences and ratios of the mass differences (entries 11–18 in the tabulated list of results) are evaluated in the same manner as for the $\Delta m(B^*)$ uncertainties, with the variations to the baseline model applied simultaneously to all channels to account for correlations. The resulting uncertainties are shown in Tables C.1 and C.2. The experimental statistical and systematic uncertainties for lines 7–10 of the tabulated list of results are identical to those for lines 11–14 (Table C.1). As for the other measurements, the statistical uncertainties dominate.

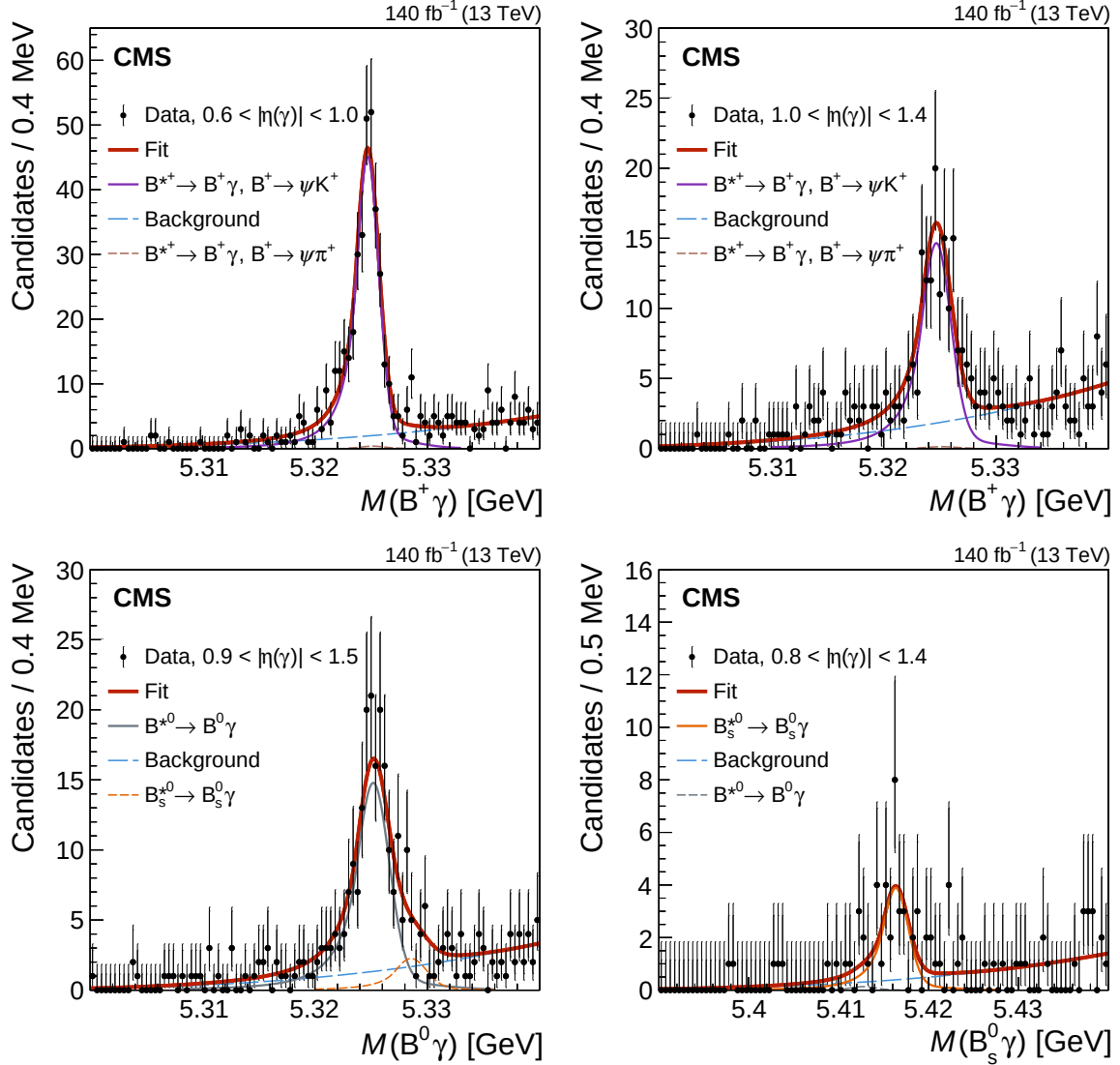


Figure B.1: The measured distributions (filled circles) of the $B^+\gamma$ (upper), $B^0\gamma$ (lower left), and $B_s^0\gamma$ (lower right) invariant mass corresponding to the most forward one or two $|\eta(\gamma)|$ regions, together with the simultaneous fit projections (thick red lines). The error bars represent statistical uncertainties in data.

Table C.1: Values, statistical, and systematic uncertainties, in keV, of the measured differences of the three mass differences.

Source	$\Delta m(B^{*0}) - \Delta m(B^{*+})$	$\Delta m(B_s^{*0}) - \Delta m(B^{*+})$	$\Delta m(B_s^{*0}) - \Delta m(B^{*0})$	$\Delta m(B_s^{*0}) - \frac{\Delta m(B^{*0}) + \Delta m(B^{*+})}{2}$
Measured value	194	4130	3936	4033
Statistical uncertainty	68	138	139	134
Signal model	4	23	23	23
Signal shape parameters	1	3	3	3
Yield ratios of $ \eta(\gamma) $ ranges	3	11	7	9
Background shape	3	9	6	8
Cross-feed $B_s^{*0} \leftrightarrow B^{*0}$	1	10	12	11
Photon energy scale, statistical	4	11	11	11
Photon energy scale, systematic	<1	2	2	2
Total systematic uncertainty	7	31	30	30

Table C.2: Values, statistical, and systematic uncertainties of the measured ratios of the three mass differences.

Source	$\Delta m(\text{B}^{*0})/\Delta m(\text{B}^{*+})$	$\Delta m(\text{B}_s^{*0})/\Delta m(\text{B}^{*+})$	$\Delta m(\text{B}_s^{*0})/\Delta m(\text{B}^{*0})$	$\frac{2\Delta m(\text{B}_s^{*0})}{\Delta m(\text{B}^{*+})+\Delta m(\text{B}^{*0})}$
Measured value	1.00428	1.09122	1.08656	1.08888
Statistical uncertainty	0.00151	0.00306	0.00309	0.00297
Signal model	0.00009	0.00050	0.00052	0.00052
Signal shape parameters	0.00002	0.00003	0.00004	0.00003
Yield ratios of $ \eta(\gamma) $ ranges	0.00008	0.00023	0.00016	0.00020
Background shape	0.00005	0.00020	0.00014	0.00017
Cross-feed $\text{B}_s^{*0} \leftrightarrow \text{B}^{*0}$	0.00003	0.00023	0.00025	0.00015
Photon energy scale	0.00009	0.00025	0.00024	0.00024
Total systematic uncertainty	0.00016	0.00067	0.00066	0.00065

D The CMS Collaboration

Yerevan Physics Institute, Yerevan, Armenia

A. Hayrapetyan, V. Makarenko , A. Tumasyan¹ 

Institut für Hochenergiephysik, Vienna, Austria

W. Adam , J.W. Andrejkovic, L. Benato , T. Bergauer , M. Dragicevic , C. Giordano, P.S. Hussain , M. Jeitler² , N. Krammer , A. Li , D. Liko , M. Matthewman, I. Mikulec , J. Schieck² , R. Schöfbeck² , D. Schwarz , M. Shooshtari, M. Sonawane , W. Waltenberger , C.-E. Wulz² 

Universiteit Antwerpen, Antwerpen, Belgium

T. Janssen , H. Kwon , D. Ocampo Henao , T. Van Laer , P. Van Mechelen 

Vrije Universiteit Brussel, Brussel, Belgium

J. Bierkens , N. Breugelmans, J. D'Hondt , S. Dansana , A. De Moor , M. Delcourt , F. Heyen, Y. Hong , P. Kashko , S. Lowette , I. Makarenko , D. Müller , J. Song , S. Tavernier , M. Tytgat³ , G.P. Van Onsem , S. Van Putte , D. Vannerom 

Université Libre de Bruxelles, Bruxelles, Belgium

B. Bilin , B. Clerbaux , A.K. Das, I. De Bruyn , G. De Lentdecker , H. Evard , L. Favart , P. Gianneios , A. Khalilzadeh, F.A. Khan , A. Malara , M.A. Shahzad, L. Thomas , M. Vanden Bemden , C. Vander Velde , P. Vanlaer , F. Zhang 

Ghent University, Ghent, Belgium

M. De Coen , D. Dobur , G. Gokbulut , J. Knolle , L. Lambrecht , D. Marckx , K. Skovpen , N. Van Den Bossche , J. van der Linden , J. Vandenbroeck , L. Wezenbeek 

Université Catholique de Louvain, Louvain-la-Neuve, Belgium

S. Bein , A. Benecke , A. Bethani , G. Bruno , A. Cappati , J. De Favereau De Jeneret , C. Delaere , A. Giammanco , A.O. Guzel , V. Lemaitre, J. Lidrych , P. Malek , P. Mastrapasqua , S. Turkcapar 

Centro Brasileiro de Pesquisas Fisicas, Rio de Janeiro, Brazil

G.A. Alves , M. Barroso Ferreira Filho , E. Coelho , C. Hensel , T. Menezes De Oliveira , C. Mora Herrera⁴ , P. Rebello Teles , M. Soeiro , E.J. Tonelli Manganote⁵ , A. Vilela Pereira⁴ 

Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

W.L. Aldá Júnior , H. Brandao Malbouisson , W. Carvalho , J. Chinellato⁶ , M. Costa Reis , E.M. Da Costa , G.G. Da Silveira⁷ , D. De Jesus Damiao , S. Fonseca De Souza , R. Gomes De Souza , S. S. Jesus , T. Laux Kuhn⁷ , M. Macedo , K. Mota Amarilo , L. Mundim , H. Nogima , J.P. Pinheiro , A. Santoro , A. Sznajder , M. Thiel , F. Torres Da Silva De Araujo⁸ 

Universidade Estadual Paulista, Universidade Federal do ABC, São Paulo, Brazil

C.A. Bernardes⁷ , F. Damas , T.R. Fernandez Perez Tomei , E.M. Gregores , B. Lopes Da Costa , I. Maietto Silverio , P.G. Mercadante , S.F. Novaes , B. Orzari , Sandra S. Padula , V. Scheurer

Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria

A. Aleksandrov , G. Antchev , P. Danev, R. Hadjiiska , P. Iaydjiev , M. Shopova , G. Sultanov 

University of Sofia, Sofia, BulgariaA. Dimitrov , L. Litov , B. Pavlov , P. Petkov , A. Petrov **Instituto De Alta Investigación, Universidad de Tarapacá, Casilla 7 D, Arica, Chile**S. Keshri , D. Laroze , S. Thakur **Universidad Tecnica Federico Santa Maria, Valparaiso, Chile**W. Brooks **Beihang University, Beijing, China**T. Cheng , T. Javaid , L. Wang , L. Yuan **Department of Physics, Tsinghua University, Beijing, China**Z. Hu , Z. Liang, J. Liu, X. Wang **Institute of High Energy Physics, Beijing, China**G.M. Chen⁹ , H.S. Chen⁹ , M. Chen⁹ , Y. Chen , Q. Hou , X. Hou, F. Iemmi , C.H. Jiang, A. Kapoor¹⁰ , H. Liao , G. Liu , Z.-A. Liu¹¹ , J.N. Song¹¹, S. Song, J. Tao , C. Wang⁹, J. Wang , H. Zhang , J. Zhao **State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China**A. Agapitos , Y. Ban , A. Carvalho Antunes De Oliveira , S. Deng , B. Guo, Q. Guo, C. Jiang , A. Levin , C. Li , Q. Li , Y. Mao, S. Qian, S.J. Qian , X. Qin, X. Sun , D. Wang , J. Wang, H. Yang, M. Zhang, Y. Zhao, C. Zhou **State Key Laboratory of Nuclear Physics and Technology, Institute of Quantum Matter, South China Normal University, Guangzhou, China**S. Yang **Sun Yat-Sen University, Guangzhou, China**Z. You **University of Science and Technology of China, Hefei, China**K. Jaffel , N. Lu **Nanjing Normal University, Nanjing, China**G. Bauer^{12,13}, B. Li¹⁴, H. Wang , K. Yi¹⁵ , J. Zhang **Institute of Modern Physics and Key Laboratory of Nuclear Physics and Ion-beam Application (MOE) - Fudan University, Shanghai, China**

Y. Li

Zhejiang University, Hangzhou, Zhejiang, ChinaZ. Lin , C. Lu , M. Xiao¹⁶ **Universidad de Los Andes, Bogota, Colombia**C. Avila , D.A. Barbosa Trujillo , A. Cabrera , C. Florez , J. Fraga , J.A. Reyes Vega**Universidad de Antioquia, Medellin, Colombia**C. Rendón , M. Rodriguez , A.A. Ruales Barbosa , J.D. Ruiz Alvarez **University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia**N. Godinovic , D. Lelas , A. Sculac **University of Split, Faculty of Science, Split, Croatia**M. Kovac , A. Petkovic , T. Sculac 

Institute Rudjer Boskovic, Zagreb, Croatia

P. Bargassa , V. Brigljevic , B.K. Chitroda , D. Ferencek , K. Jakovcic, A. Starodumov , T. Susa 

University of Cyprus, Nicosia, Cyprus

A. Attikis , K. Christoforou , A. Hadjiagapiou, C. Leonidou , C. Nicolaou, L. Paizanos , F. Ptochos , P.A. Razis , H. Rykaczewski, H. Saka , A. Stepennov 

Charles University, Prague, Czech Republic

M. Finger[†] , M. Finger Jr. 

Escuela Politecnica Nacional, Quito, Ecuador

E. Ayala 

Universidad San Francisco de Quito, Quito, Ecuador

E. Carrera Jarrin 

Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt

A.A. Abdelalim^{17,18} , R. Aly^{19,17} 

Center for High Energy Physics (CHEP-FU), Fayoum University, El-Fayoum, Egypt

M. Abdullah Al-Mashad , A. Hussein, H. Mohammed 

National Institute of Chemical Physics and Biophysics, Tallinn, Estonia

K. Ehataht , M. Kadastik, T. Lange , C. Nielsen , J. Pata , M. Raidal , N. Seeba , L. Tani 

Department of Physics, University of Helsinki, Helsinki, Finland

A. Milieva , K. Osterberg , M. Voutilainen 

Helsinki Institute of Physics, Helsinki, Finland

N. Bin Norjoharuddeen , E. Brücken , F. Garcia , P. Inkaew , K.T.S. Kallonen , R. Kumar Verma , T. Lampén , K. Lassila-Perini , B. Lehtela , S. Lehti , T. Lindén , N.R. Mancilla Xinto , M. Myllymäki , M.m. Rantanen , S. Saariokari , N.T. Toikka , J. Tuominiemi 

Lappeenranta-Lahti University of Technology, Lappeenranta, Finland

H. Kirschenmann , P. Luukka , H. Petrow 

IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France

M. Besancon , F. Couderc , M. Dejardin , D. Denegri, P. Devouge, J.L. Faure , F. Ferri , P. Gagne, S. Ganjour , P. Gras , G. Hamel de Monchenault , M. Kumar , V. Lohezic , J. Malcles , F. Orlandi , L. Portales , S. Ronchi , M.Ö. Sahin , A. Savoy-Navarro²⁰ , P. Simkina , M. Titov , M. Tornago 

Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France

F. Beaudette , G. Boldrini , P. Busson , C. Charlot , M. Chiusi , T.D. Cuisset , O. Davignon , A. De Wit , T. Debnath , I.T. Ehle , B.A. Fontana Santos Alves , S. Ghosh , A. Gilbert , R. Granier de Cassagnac , L. Kalipoliti , M. Manoni , M. Nguyen , S. Obraztsov , C. Ochando , R. Salerno , J.B. Sauvan , Y. Sirois , G. Sokmen, L. Urda Gómez , A. Zabi , A. Zghiche 

Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France

J.-L. Agram²¹ , J. Andrea , D. Bloch , J.-M. Brom , E.C. Chabert , C. Collard 

G. Coulon, S. Falke , U. Goerlach , R. Haeberle , A.-C. Le Bihan , M. Meena , O. Poncet , G. Saha , P. Vaucelle 

Centre de Calcul de l'Institut National de Physique Nucleaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France

A. Di Florio 

Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France

D. Amram, S. Beauceron , B. Blancon , G. Boudoul , N. Chanon , D. Contardo , P. Depasse , H. El Mamouni, J. Fay , S. Gascon , M. Gouzevitch , C. Greenberg , G. Grenier , B. Ille , E. Jourdhuy, I.B. Laktineh, M. Lethuillier , B. Massoteau, L. Mirabito, A. Purohit , M. Vander Donckt , J. Xiao 

Georgian Technical University, Tbilisi, Georgia

I. Lomidze , T. Toriashvili²² , Z. Tsamalaidze²³ 

RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany

V. Botta , S. Consuegra Rodríguez , L. Feld , K. Klein , M. Lipinski , D. Meuser , P. Natland , V. Oppenländer, A. Pauls , D. Pérez Adán , N. Röwert , M. Teroerde 

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

C. Daumann, S. Diekmann , A. Dodonova , N. Eich , D. Eliseev , F. Engelke , J. Erdmann , M. Erdmann , B. Fischer , T. Hebbeker , K. Hoepfner , F. Ivone , A. Jung , N. Kumar , M.y. Lee , F. Mausolf , M. Merschmeyer , A. Meyer , F. Nowotny, A. Pozdnyakov , W. Redjeb , H. Reithler , U. Sarkar , V. Sarkisovi , A. Schmidt , C. Seth, A. Sharma , J.L. Spah , V. Vaulin, S. Zaleski

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

M.R. Beckers , C. Dziwok , G. Flügge , N. Hoeflich , T. Kress , A. Nowack , O. Pooth , A. Stahl , A. Zotz 

Deutsches Elektronen-Synchrotron, Hamburg, Germany

H. Aarup Petersen , A. Abel, M. Aldaya Martin , J. Alimena , S. Amoroso, Y. An , I. Andreev , J. Bach , S. Baxter , M. Bayatmakou , H. Becerril Gonzalez , O. Behnke , A. Belvedere , F. Blekman²⁴ , K. Borras²⁵ , A. Campbell , S. Chatterjee , L.X. Coll Saravia , G. Eckerlin, D. Eckstein , E. Gallo²⁴ , A. Geiser , V. Guglielmi , M. Guthoff , A. Hinzmann , L. Jeppe , M. Kasemann , C. Kleinwort , R. Kogler , M. Komm , D. Krücker , W. Lange, D. Leyva Pernia , K.-Y. Lin , K. Lipka²⁶ , W. Lohmann²⁷ , J. Malvaso , R. Mankel , I.-A. Melzer-Pellmann , M. Mendizabal Morentin , A.B. Meyer , G. Milella , K. Moral Figueroa , A. Mussgiller , L.P. Nair , J. Niedziela , A. Nürnberg , J. Park , E. Ranken , A. Raspereza , D. Rastorguev , L. Rygaard, M. Scham^{28,25} , S. Schnake²⁵ , P. Schütze , C. Schwanenberger²⁴ , D. Selivanova , K. Sharke , M. Shchedrolosiev , D. Stafford , M. Torkian, F. Vazzoler , A. Ventura Barroso , R. Walsh , D. Wang , Q. Wang , K. Wichmann, L. Wiens²⁵ , C. Wissing , Y. Yang , S. Zakharov, A. Zimmermann Castro Santos 

University of Hamburg, Hamburg, Germany

A.R. Alves Andrade , M. Antonello , S. Bollweg, M. Bonanomi , K. El Morabit , Y. Fischer , M. Frahm, E. Garutti , A. Grohsjean , A.A. Guvenli , J. Haller , D. Hundhausen, G. Kasieczka , P. Keicher , R. Klanner , W. Korcari , T. Kramer , C.c. Kuo, F. Labe , J. Lange , A. Lobanov , L. Moureaux , A. Nigamova , K. Nikolopoulos , A. Paasch , K.J. Pena Rodriguez , N. Prouvost, T. Quadfasel , B. Raciti , M. Rieger , D. Savoie , P. Schleper , M. Schröder , J. Schwandt 

M. Sommerhalder , H. Stadie , G. Steinbrück , R. Ward , B. Wiederspan, M. Wolf 

Karlsruher Institut fuer Technologie, Karlsruhe, Germany

S. Brommer , E. Butz , Y.M. Chen , T. Chwalek , A. Dierlamm , G.G. Dincer , U. Elicabuk, N. Faltermann , M. Giffels , A. Gottmann , F. Hartmann²⁹ , R. Hofsaess , M. Horzela , F. Hummer , U. Husemann , J. Kieseler , M. Klute , R. Kunnilan Muhammed Rafeek, O. Lavoryk , J.M. Lawhorn , A. Lintuluoto , S. Maier , M. Mormile , Th. Müller , E. Pfeffer , M. Presilla , G. Quast , K. Rabbertz , B. Regnery , R. Schmieder, N. Shadskiy , I. Shvetsov , H.J. Simonis , L. Sowa , L. Stockmeier, K. Tauqeer, M. Toms , B. Topko , N. Trevisani , C. Verstege , T. Voigtländer , R.F. Von Cube , J. Von Den Driesch, M. Wassmer , R. Wolf , W.D. Zeuner , X. Zuo 

Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece

G. Anagnostou , G. Daskalakis , A. Kyriakis 

National and Kapodistrian University of Athens, Athens, Greece

G. Melachroinos, Z. Painesis , I. Paraskevas , N. Saoulidou , K. Theofilatos , E. Tziaferi , E. Tzovara , K. Vellidis , I. Zisopoulos 

National Technical University of Athens, Athens, Greece

T. Chatzistavrou, G. Karapostoli , K. Kousouris , E. Siamarkou, G. Tsipolitis 

University of Ioánnina, Ioánnina, Greece

I. Bestintzanos, I. Evangelou , C. Foudas, P. Katsoulis, P. Kokkas , P.G. Kosmoglou Kioseoglou , N. Manthos , I. Papadopoulos , J. Strologas 

HUN-REN Wigner Research Centre for Physics, Budapest, Hungary

D. Druzhkin , C. Hajdu , D. Horvath^{30,31} , K. Márton, A.J. Rádl³² , F. Sikler , V. Veszpremi 

MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary

M. Csanád , K. Farkas , A. Fehérkuti³³ , M.M.A. Gadallah³⁴ , Á. Kadlecik , M. León Coello , G. Pásztor , G.I. Veres 

Faculty of Informatics, University of Debrecen, Debrecen, Hungary

B. Ujvari , G. Zilizi 

HUN-REN ATOMKI - Institute of Nuclear Research, Debrecen, Hungary

G. Bencze, S. Czellar, J. Molnar, Z. Szillasi

Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary

T. Csorgo³³ , F. Nemes³³ , T. Novak , I. Szanyi³⁵ 

Panjab University, Chandigarh, India

S. Bansal , S.B. Beri, V. Bhatnagar , G. Chaudhary , S. Chauhan , N. Dhingra³⁶ , A. Kaur , A. Kaur , H. Kaur , M. Kaur , S. Kumar , T. Sheokand, J.B. Singh , A. Singla 

University of Delhi, Delhi, India

A. Bhardwaj , A. Chhetri , B.C. Choudhary , A. Kumar , A. Kumar , M. Naimuddin , S. Phor , K. Ranjan , M.K. Saini

University of Hyderabad, Hyderabad, India

S. Acharya³⁷ , B. Gomber³⁷ , B. Sahu³⁷ 

Indian Institute of Technology Kanpur, Kanpur, India

S. Mukherjee 

Saha Institute of Nuclear Physics, HBNI, Kolkata, India

S. Baradia , S. Bhattacharya , S. Das Gupta, S. Dutta , S. Dutta, S. Sarkar

Indian Institute of Technology Madras, Madras, India

M.M. Ameen , P.K. Behera , S. Chatterjee , G. Dash , A. Dattamunsi, P. Jana , P. Kalbhor , S. Kamble , J.R. Komaragiri³⁸ , T. Mishra , P.R. Pujahari , A.K. Sikdar , R.K. Singh , P. Verma , S. Verma , A. Vijay 

IISER Mohali, India, Mohali, India

B.K. Sirasva

Tata Institute of Fundamental Research-A, Mumbai, India

L. Bhatt, S. Dugad , G.B. Mohanty , M. Shelake , P. Suryadevara

Tata Institute of Fundamental Research-B, Mumbai, India

A. Bala , S. Banerjee , S. Barman³⁹ , R.M. Chatterjee, M. Guchait , Sh. Jain , A. Jaiswal, B.M. Joshi , S. Kumar , M. Maity³⁹, G. Majumder , K. Mazumdar , S. Parolia , R. Saxena , A. Thachayath 

National Institute of Science Education and Research, An OCC of Homi Bhabha National Institute, Bhubaneswar, Odisha, India

S. Bahinipati⁴⁰ , D. Maity⁴¹ , P. Mal , K. Naskar⁴¹ , A. Nayak⁴¹ , S. Nayak, K. Pal , R. Raturi, P. Sadangi, S.K. Swain , S. Varghese⁴¹ , D. Vats⁴¹ 

Indian Institute of Science Education and Research (IISER), Pune, India

A. Alpana , S. Dube , P. Hazarika , B. Kansal , A. Laha , R. Sharma , S. Sharma , K.Y. Vaish 

Indian Institute of Technology Hyderabad, Telangana, India

S. Ghosh 

Isfahan University of Technology, Isfahan, Iran

H. Bakhshiansohi⁴² , A. Jafari⁴³ , V. Sedighzadeh Dalavi , M. Zeinali⁴⁴ 

Institute for Research in Fundamental Sciences (IPM), Tehran, Iran

S. Bashiri , S. Chenarani⁴⁵ , S.M. Etesami , Y. Hosseini , M. Khakzad , E. Khazaie , M. Mohammadi Najafabadi , S. Tizchang⁴⁶ 

University College Dublin, Dublin, Ireland

M. Felcini , M. Grunewald 

INFN Sezione di Bari^a, Università di Bari^b, Politecnico di Bari^c, Bari, Italy

M. Abbrescia^{a,b} , M. Barbieri^{a,b}, M. Buonsante^{a,b} , A. Colaleo^{a,b} , D. Creanza^{a,c} , N. De Filippis^{a,c} , M. De Palma^{a,b} , W. Elmetenawee^{a,b,17} , N. Ferrara^{a,c} , L. Fiore^a , L. Longo^a , M. Louka^{a,b} , G. Maggi^{a,c} , M. Maggi^a , I. Margjeka^a , V. Mastrapasqua^{a,b} , S. My^{a,b} , F. Nenna^{a,b} , S. Nuzzo^{a,b} , A. Pellicchia^{a,b} , A. Pompili^{a,b} , G. Pugliese^{a,c} , R. Radogna^{a,b} , D. Ramos^a , A. Ranieri^a , L. Silvestris^a , F.M. Simone^{a,c} , Ü. Sözbilir^a , A. Stamerra^{a,b} , D. Troiano^{a,b} , R. Venditti^{a,b} , P. Verwilligen^a , A. Zaza^{a,b} 

INFN Sezione di Bologna^a, Università di Bologna^b, Bologna, Italy

G. Abbiendi^a , C. Battilana^{a,b} , D. Bonacorsi^{a,b} , P. Capiluppi^{a,b} , F.R. Cavallo^a , M. Cuffiani^{a,b} , G.M. Dallavalle^a , T. Diotallevi^{a,b} , F. Fabbri^a , A. Fanfani^{a,b} , D. Fasanella^a , P. Giacomelli^a , C. Grandi^a , S. Lo Meo^{a,47} , M. Lorusso^{a,b} , L. Lunerti^a , S. Marcellini^a , G. Masetti^a , F.L. Navarra^{a,b} , G. Paggi^{a,b} , A. Perrotta^a , F. Primavera^{a,b} , A.M. Rossi^{a,b} , S. Rossi Tisbeni^{a,b} , T. Rovelli^{a,b} , G.P. Siroli^{a,b}

INFN Sezione di Catania^a, Università di Catania^b, Catania, Italy

S. Costa^{a,b,48} , A. Di Mattia^a , A. Lapertosa^a , R. Potenza^{a,b} , A. Tricomi^{a,b,48} 

INFN Sezione di Firenze^a, Università di Firenze^b, Firenze, Italy

J. Altork^{a,b} , P. Assiouras^a , G. Barbagli^a , G. Bardelli^a , M. Bartolini^{a,b} , A. Calandri^{a,b} , B. Camaiani^{a,b} , A. Cassese^a , R. Ceccarelli^a , V. Ciulli^{a,b} , C. Civinini^a , R. D'Alessandro^{a,b} , L. Damenti^{a,b} , E. Focardi^{a,b} , T. Kello^a , G. Latino^{a,b} , P. Lenzi^{a,b} , M. Lizzo^a , M. Meschini^a , S. Paoletti^a , A. Papanastassiou^{a,b} , G. Sguazzoni^a , L. Viliani^a

INFN Laboratori Nazionali di Frascati, Frascati, Italy

L. Benussi , S. Bianco , S. Meola⁴⁹ , D. Piccolo 

INFN Sezione di Genova^a, Università di Genova^b, Genova, Italy

M. Alves Gallo Pereira^a , F. Ferro^a , E. Robutti^a , S. Tosi^{a,b} 

INFN Sezione di Milano-Bicocca^a, Università di Milano-Bicocca^b, Milano, Italy

A. Benaglia^a , F. Brivio^a , V. Camagni^{a,b} , F. Cetorelli^{a,b} , F. De Guio^{a,b} , M.E. Dinardo^{a,b} , P. Dini^a , S. Gennai^a , R. Gerosa^{a,b} , A. Ghezzi^{a,b} , P. Govoni^{a,b} , L. Guzzi^a , M.R. Kim^a , G. Lavizzari^{a,b} , M.T. Lucchini^{a,b} , M. Malberti^a , S. Malvezzi^a , A. Massironi^a , D. Menasce^a , L. Moroni^a , M. Paganoni^{a,b} , S. Palluotto^{a,b} , D. Pedrini^a , A. Perego^{a,b} , G. Pizzati^{a,b} , S. Ragazzi^{a,b} , T. Tabarelli de Fatis^{a,b}

INFN Sezione di Napoli^a, Università di Napoli 'Federico II'^b, Napoli, Italy; Università della Basilicata^c, Potenza, Italy; Scuola Superiore Meridionale (SSM)^d, Napoli, Italy

S. Buontempo^a , C. Di Fraia^{a,b} , F. Fabozzi^{a,c} , L. Favilla^{a,d} , A.O.M. Iorio^{a,b} , L. Lista^{a,b,50} , P. Paolucci^{a,29} , B. Rossi^a 

INFN Sezione di Padova^a, Università di Padova^b, Padova, Italy; Università degli Studi di Cagliari^c, Cagliari, Italy

P. Azzi^a , N. Bacchetta^{a,51} , D. Bisello^{a,b} , P. Bortignon^{a,c} , G. Bortolato^{a,b} , A.C.M. Bulla^{a,c} , R. Carlin^{a,b} , T. Dorigo^{a,52} , F. Gasparini^{a,b} , S. Giorgetti^a , M. Gulmini^{a,53} , E. Lusiani^a , M. Margoni^{a,b} , A.T. Meneguzzo^{a,b} , J. Pazzini^{a,b} , P. Ronchese^{a,b} , R. Rossin^{a,b} , F. Simonetto^{a,b} , M. Tosi^{a,b} , A. Triossi^{a,b} , S. Ventura^a , M. Zanetti^{a,b} , P. Zotto^{a,b} , A. Zucchetta^{a,b} , G. Zumerle^{a,b}

INFN Sezione di Pavia^a, Università di Pavia^b, Pavia, Italy

A. Braghieri^a , S. Calzaferri^a , P. Montagna^{a,b} , M. Pelliccioni^a , V. Re^a , C. Riccardi^{a,b} , P. Salvini^a , I. Vai^{a,b} , P. Vitulo^{a,b} 

INFN Sezione di Perugia^a, Università di Perugia^b, Perugia, Italy

S. Ajmal^{a,b} , M.E. Ascioti^{a,b} , G.M. Bilei^a , C. Carrivale^{a,b} , D. Ciangottini^{a,b} , L. Della Penna^{a,b} , L. Fanò^{a,b} , V. Mariani^{a,b} , M. Menichelli^a , F. Moscatelli^{a,54} , A. Rossi^{a,b} , A. Santocchia^{a,b} , D. Spiga^a , T. Tedeschi^{a,b} 

INFN Sezione di Pisa^a, Università di Pisa^b, Scuola Normale Superiore di Pisa^c, Pisa, Italy;

Università di Siena^d, Siena, Italy

C. Aimè^{a,b} , C.A. Alexe^{a,c} , P. Asenov^{a,b} , P. Azzurri^a , G. Bagliesi^a , R. Bhattacharya^a , L. Bianchini^{a,b} , T. Boccali^a , E. Bossini^a , D. Bruschini^{a,c} , L. Calligaris^{a,b} , R. Castaldi^a , F. Cattafesta^{a,c} , M.A. Ciocci^{a,d} , M. Cipriani^{a,b} , R. Dell'Orso^a , S. Donato^{a,b} , R. Forti^{a,b} , A. Giassi^a , F. Ligabue^{a,c} , A.C. Marini^{a,b} , D. Matos Figueiredo^a , A. Messineo^{a,b} , S. Mishra^a , V.K. Muraleedharan Nair Bindhu^{a,b} , S. Nandan^a , F. Palla^a , M. Riggirello^{a,c} , A. Rizzi^{a,b} , G. Rolandi^{a,c} , S. Roy Chowdhury^{a,55} , T. Sarkar^a , A. Scribano^a , P. Solanki^{a,b} , P. Spagnolo^a , F. Tenchini^{a,b} , R. Tenchini^a , G. Tonelli^{a,b} , N. Turini^{a,d} , F. Vaselli^{a,c} , A. Venturi^a , P.G. Verdini^a 

INFN Sezione di Roma^a, Sapienza Università di Roma^b, Roma, Italy

P. Akrap^{a,b} , C. Basile^{a,b} , S.C. Behera^a , F. Cavallari^a , L. Cunqueiro Mendez^{a,b} , F. De Ruggi^{a,b} , D. Del Re^{a,b} , E. Di Marco^a , M. Diemoz^a , F. Errico^a , L. Frosina^{a,b} , R. Gargiulo^{a,b} , B. Harikrishnan^{a,b} , F. Lombardi^{a,b} , E. Longo^{a,b} , L. Martikainen^{a,b} , J. Mijuskovic^{a,b} , G. Organtini^{a,b} , N. Palmeri^{a,b} , R. Paramatti^{a,b} , C. Quaranta^{a,b} , S. Rahatlou^{a,b} , C. Rovelli^a , F. Santanastasio^{a,b} , L. Soffi^a , V. Vladimirov^{a,b}

INFN Sezione di Torino^a, Università di Torino^b, Torino, Italy; Università del Piemonte Orientale^c, Novara, Italy

N. Amapane^{a,b} , R. Arcidiacono^{a,c} , S. Argiro^{a,b} , M. Arneodo^{a,c} , N. Bartosik^{a,c} , R. Bellan^{a,b} , A. Bellora^{a,b} , C. Biino^a , C. Borca^{a,b} , N. Cartiglia^a , M. Costa^{a,b} , R. Covarelli^{a,b} , N. Demaria^a , L. Finco^a , M. Grippo^{a,b} , B. Kiani^{a,b} , L. Lanteri^{a,b} , F. Legger^a , F. Luongo^{a,b} , C. Mariotti^a , S. Maselli^a , A. Mecca^{a,b} , L. Menzio^{a,b} , P. Meridiani^a , E. Migliore^{a,b} , M. Monteno^a , M.M. Obertino^{a,b} , G. Ortona^a , L. Pacher^{a,b} , N. Pastrone^a , M. Ruspà^{a,c} , F. Siviero^{a,b} , V. Sola^{a,b} , A. Solano^{a,b} , A. Staiano^a , C. Tarricone^{a,b} , D. Trocino^a , G. Umoret^{a,b} , E. Vlasov^{a,b} , R. White^{a,b} 

INFN Sezione di Trieste^a, Università di Trieste^b, Trieste, Italy

J. Babbar^{a,b} , S. Belforte^a , V. Candelise^{a,b} , M. Casarsa^a , F. Cossutti^a , K. De Leo^a , G. Della Ricca^{a,b} , R. Delli Gatti^{a,b} 

Kyungpook National University, Daegu, Korea

S. Dogra , J. Hong , J. Kim, T. Kim , D. Lee, H. Lee , J. Lee, S.W. Lee , C.S. Moon , Y.D. Oh , S. Sekmen , B. Tae, Y.C. Yang 

Department of Mathematics and Physics - GWNu, Gangneung, Korea

M.S. Kim 

Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea

G. Bak , P. Gwak , H. Kim , D.H. Moon , J. Seo 

Hanyang University, Seoul, Korea

E. Asilar , F. Carnevali , J. Choi⁵⁶ , T.J. Kim , Y. Ryou

Korea University, Seoul, Korea

S. Ha , S. Han, B. Hong , J. Kim , K. Lee, K.S. Lee , S. Lee , J. Yoo 

Kyung Hee University, Department of Physics, Seoul, Korea

J. Goh , J. Shin , S. Yang 

Sejong University, Seoul, Korea

Y. Kang , H. S. Kim , Y. Kim , S. Lee

Seoul National University, Seoul, Korea

J. Almond, J.H. Bhyun, J. Choi , J. Choi, W. Jun , H. Kim , J. Kim , T. Kim, Y. Kim, Y.W. Kim , S. Ko , H. Lee , J. Lee , J. Lee , B.H. Oh , S.B. Oh , J. Shin , U.K. Yang, I. Yoon 

University of Seoul, Seoul, Korea

W. Jang , D.Y. Kang, D. Kim , S. Kim , B. Ko, J.S.H. Lee , Y. Lee , I.C. Park , Y. Roh, I.J. Watson 

Yonsei University, Department of Physics, Seoul, Korea

G. Cho, K. Hwang , B. Kim , S. Kim, K. Lee , H.D. Yoo 

Sungkyunkwan University, Suwon, Korea

M. Choi , Y. Lee , I. Yu 

College of Engineering and Technology, American University of the Middle East (AUM), Dasman, Kuwait

T. Beyrouthy , Y. Gharbia 

Kuwait University - College of Science - Department of Physics, Safat, Kuwait

F. Alazemi 

Riga Technical University, Riga, Latvia

K. Dreimanis , O.M. Eberlins , A. Gaile , C. Munoz Diaz , D. Osite , G. Pikurs , R. Plese , A. Potrebko , M. Seidel , D. Sidiropoulos Kontos 

University of Latvia (LU), Riga, Latvia

N.R. Strautnieks 

Vilnius University, Vilnius, Lithuania

M. Ambrozas , A. Juodagalvis , S. Nargelas , A. Rinkevicius , G. Tamulaitis 

National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia

I. Yusuff⁵⁷ , Z. Zolkapli

Universidad de Sonora (UNISON), Hermosillo, Mexico

J.F. Benitez , A. Castaneda Hernandez , A. Cota Rodriguez , L.E. Cuevas Picos, H.A. Encinas Acosta, L.G. Gallegos Maríñez, J.A. Murillo Quijada , L. Valencia Palomo 

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

G. Ayala , H. Castilla-Valdez , H. Crotte Ledesma , R. Lopez-Fernandez , J. Mejia Guisao , R. Reyes-Almanza , A. Sánchez Hernández 

Universidad Iberoamericana, Mexico City, Mexico

C. Oropeza Barrera , D.L. Ramirez Guadarrama, M. Ramírez García 

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

I. Bautista , F.E. Neri Huerta , I. Pedraza , H.A. Salazar Ibarguen , C. Uribe Estrada 

University of Montenegro, Podgorica, Montenegro

I. Bubanja , N. Raicevic 

University of Canterbury, Christchurch, New Zealand

P.H. Butler 

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

A. Ahmad , M.I. Asghar , A. Awais , M.I.M. Awan, W.A. Khan 

AGH University of Krakow, Krakow, Poland

V. Avati, L. Forthomme , L. Grzanka , M. Malawski , K. Piotrkowski 

National Centre for Nuclear Research, Swierk, Poland

M. Bluj , M. Górski , M. Kazana , M. Szleper , P. Zalewski 

Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland

K. Bunkowski , K. Doroba , A. Kalinowski , M. Konecki , J. Krolikowski , A. Muhammad 

Warsaw University of Technology, Warsaw, Poland

P. Fokow , K. Pozniak , W. Zabolotny 

Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal

M. Araujo , D. Bastos , C. Beirão Da Cruz E Silva , A. Boletti , M. Bozzo , T. Camporesi , G. Da Molin , M. Gallinaro , J. Hollar , N. Leonardo , G.B. Marozzo , A. Petrilli , M. Pisano , J. Seixas , J. Varela , J.W. Wulff 

Faculty of Physics, University of Belgrade, Belgrade, Serbia

P. Adzic , L. Markovic , P. Milenovic , V. Milosevic 

VINCA Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia

D. Devetak , M. Dordevic , J. Milosevic , L. Nadderd , V. Rekovic, M. Stojanovic 

Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain

M. Alcalde Martinez , J. Alcaraz Maestre , Cristina F. Bedoya , J.A. Brochero Cifuentes , Oliver M. Carretero , M. Cepeda , M. Cerrada , N. Colino , J. Cuchillo Ortega, B. De La Cruz , A. Delgado Peris , A. Escalante Del Valle , D. Fernández Del Val , J.P. Fernández Ramos , J. Flix , M.C. Fouz , M. Gonzalez Hernandez , O. Gonzalez Lopez , S. Goy Lopez , J.M. Hernandez , M.I. Josa , J. Llorente Merino , C. Martin Perez , E. Martin Viscasillas , D. Moran , C. M. Morcillo Perez , R. Paz Herrera , C. Perez Dengra , A. Pérez-Calero Yzquierdo , J. Puerta Pelayo , I. Redondo , J. Vazquez Escobar 

Universidad Autónoma de Madrid, Madrid, Spain

J.F. de Trocóniz 

Universidad de Oviedo, Instituto Universitario de Ciencias y Tecnologías Espaciales de Asturias (ICTEA), Oviedo, Spain

B. Alvarez Gonzalez , J. Ayllon Torresano , A. Cardini , J. Cuevas , J. Del Riego Badas , D. Estrada Acevedo , J. Fernandez Menendez , S. Folgueras , I. Gonzalez Caballero , P. Leguina , M. Obeso Menendez , E. Palencia Cortezon , J. Prado Pico , A. Soto Rodríguez , C. Vico Villalba , P. Vischia 

Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain

S. Blanco Fernández , I.J. Cabrillo , A. Calderon , J. Duarte Campderros , M. Fernandez , G. Gomez , C. Lasiosa García , R. Lopez Ruiz , C. Martinez Rivero , P. Martinez Ruiz del Arbol , F. Matorras , P. Matorras Cuevas , E. Navarrete Ramos , J. Piedra Gomez , C. Quintana San Emeterio , L. Scodellaro , I. Vila , R. Villar Cortabitarte , J.M. Vizan Garcia 

University of Colombo, Colombo, Sri Lanka

B. Kailasapathy⁵⁸ , D.D.C. Wickramarathna 

University of Ruhuna, Department of Physics, Matara, Sri Lanka

W.G.D. Dharmaratna⁵⁹ , K. Liyanage , N. Perera 

CERN, European Organization for Nuclear Research, Geneva, Switzerland

D. Abbaneo , C. Amendola , R. Ardino , E. Auffray , J. Baechler, D. Barney , J. Bendauid , M. Bianco , A. Bocchi , L. Borghoni , C. Botta , A. Bragagnolo , C.E. Brown , C. Caillol , G. Cerminara , P. Connor , D. d'Enterria , A. Dabrowski , A. David , A. De Roeck , M.M. Defranchis , M. Deile , M. Dobson , P.J. Fernández Manteca , W. Funk , A. Gaddi, S. Giani, D. Gigi, K. Gill , F. Glege , M. Glowacki, A. Gruber , J. Hegeman , J.K. Heikkilä , B. Huber , V. Innocente , T. James , P. Janot , O. Kaluzinska , O. Karacheban²⁷ , G. Karathanasis , S. Laurila , P. Lecoq , C. Lourenço , A.-M. Lyon , M. Magherini , L. Malgeri , M. Mannelli , A. Mehta , F. Meijers , J.A. Merlin, S. Mersi , E. Meschi , M. Migliorini , F. Monti , F. Moortgat , M. Mulders , M. Musich , I. Neutelings , S. Orfanelli, F. Pantaleo , M. Pari , G. Petrucciani , A. Pfeiffer , M. Pierini , M. Pitt , H. Qu , D. Rabadý , A. Reimers , B. Ribeiro Lopes , F. Riti , P. Rosado , M. Rovere , H. Sakulin , R. Salvatico , S. Sanchez Cruz , S. Scarfi , M. Selvaggi , A. Sharma , K. Shchelina , P. Silva , P. Sphicas⁶⁰ , A.G. Stahl Leiton , A. Steen , S. Summers , D. Treille , P. Tropea , E. Vernazza , J. Wanczyk⁶¹ , J. Wang, S. Wuchterl , M. Zarucki , P. Zehetner , P. Zexjd , G. Zevi Della Porta 

PSI Center for Neutron and Muon Sciences, Villigen, Switzerland

T. Bevilacqua⁶² , L. Caminada⁶² , W. Erdmann , R. Horisberger , Q. Ingram , H.C. Kaestli , D. Kotlinski , C. Lange , U. Langenegger , L. Noehte⁶² , T. Rohe , A. Samalan 

ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland

T.K. Aarrestad , M. Backhaus , G. Bonomelli , C. Cazzaniga , K. Datta , P. De Bryas Dexmiers D'archiacchiac⁶¹ , A. De Cosa , G. Dissertori , M. Dittmar, M. Donegà , F. Eble , K. Gedia , F. Glessgen , C. Grab , N. Härringer , T.G. Harte , W. Lustermann , M. Malucchi , R.A. Manzoni , M. Marchegiani , L. Marchese , A. Mascellani⁶¹ , F. Nessi-Tedaldi , F. Pauss , V. Perovic , B. Ristic , R. Seidita , J. Steggemann⁶¹ , A. Tarabini , D. Valsecchi , R. Wallny 

Universität Zürich, Zurich, Switzerland

C. Amsler⁶³ , P. Bärtschi , F. Bilandzija , M.F. Canelli , G. Celotto, K. Cormier , M. Huwiler , W. Jin , A. Jofrehei , B. Kilminster , T.H. Kwok , S. Leontsinis , V. Lukashenko , A. Macchiolo , F. Meng , M. Missiroli , J. Motta , P. Robmann, M. Senger , E. Shokr , F. Stäger , R. Tramontano , P. Viscone 

National Central University, Chung-Li, Taiwan

D. Bhowmik, C.M. Kuo, P.K. Rout , S. Taj , P.C. Tiwari³⁸ 

National Taiwan University (NTU), Taipei, Taiwan

L. Ceard, K.F. Chen , Z.g. Chen, A. De Iorio , W.-S. Hou , T.h. Hsu, Y.w. Kao, S. Karmakar , G. Kole , Y.y. Li , R.-S. Lu , E. Paganis , X.f. Su , J. Thomas-Wilsker , L.s. Tsai, D. Tsionou, H.y. Wu , E. Yazgan 

High Energy Physics Research Unit, Department of Physics, Faculty of Science, Chulalongkorn University, Bangkok, Thailand

C. Asawatangkuldee , N. Srimanobhas 

Tunis El Manar University, Tunis, Tunisia

Y. Maghrbi 

Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey

D. Agyel , F. Dolek , I. Dumanoglu⁶⁴ , Y. Guler⁶⁵ , E. Gurpinar Guler⁶⁵ , C. Isik ,
O. Kara , A. Kayis Topaksu , Y. Komurcu , G. Onengut , K. Ozdemir⁶⁶ , B. Tali⁶⁷ ,
U.G. Tok , E. Uslan , I.S. Zorbakir 

Middle East Technical University, Physics Department, Ankara, Turkey

M. Yalvac⁶⁸ 

Bogazici University, Istanbul, Turkey

B. Akgun , I.O. Atakisi⁶⁹ , E. Gülmez , M. Kaya⁷⁰ , O. Kaya⁷¹ , M.A. Sarkisla⁷²,
S. Tekten⁷³ 

Istanbul Technical University, Istanbul, Turkey

A. Cakir , K. Cankocak^{64,74} , S. Sen⁷⁵ 

Istanbul University, Istanbul, Turkey

O. Aydilek⁷⁶ , B. Haciosahinoglu , I. Hos⁷⁷ , B. Kaynak , S. Ozkorucuklu , O. Potok ,
H. Sert , C. Simsek , C. Zorbilmez 

Yildiz Technical University, Istanbul, Turkey

S. Cerci , B. Isildak⁷⁸ , E. Simsek , D. Sunar Cerci , T. Yetkin⁷⁹ 

Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkiv, Ukraine

A. Boyaryntsev , O. Dadazhanova, B. Grynyov 

National Science Centre, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine

L. Levchuk 

University of Bristol, Bristol, United Kingdom

J.J. Brooke , A. Bundock , F. Bury , E. Clement , D. Cussans , D. Dharmender,
H. Flacher , J. Goldstein , H.F. Heath , M.-L. Holmberg , L. Kreczko , S. Parames-
varan , L. Robertshaw, M.S. Sanjrani⁴², J. Segal, V.J. Smith 

Rutherford Appleton Laboratory, Didcot, United Kingdom

A.H. Ball, K.W. Bell , A. Belyaev⁸⁰ , C. Brew , R.M. Brown , D.J.A. Cockerill ,
A. Elliot , K.V. Ellis, J. Gajownik , K. Harder , S. Harper , J. Linacre , K. Manolopoulos,
M. Moallemi , D.M. Newbold , E. Olaiya , D. Petyt , T. Reis , A.R. Sahasransu ,
G. Salvi , T. Schuh, C.H. Shepherd-Themistocleous , I.R. Tomalin , K.C. Whalen ,
T. Williams 

Imperial College, London, United Kingdom

I. Andreou , R. Bainbridge , P. Bloch , O. Buchmuller, C.A. Carrillo Montoya ,
D. Colling , J.S. Dancu, I. Das , P. Dauncey , G. Davies , M. Della Negra , S. Fayer,
G. Fedi , G. Hall , H.R. Hoorani , A. Howard, G. Iles , C.R. Knight , P. Krueper ,
J. Langford , K.H. Law , J. León Holgado , E. Leutgeb , L. Lyons , A.-M. Magnan ,
B. Maier , S. Mallios, A. Mastronikolis , M. Mieskolainen , J. Nash⁸¹ , M. Pesaresi ,
P.B. Pradeep , B.C. Radburn-Smith , A. Richards, A. Rose , L. Russell , K. Savva ,
C. Seez , R. Shukla , A. Tapper , K. Uchida , G.P. Uttley , T. Virdee²⁹ , M. Vojinovic ,
N. Wardle , D. Winterbottom 

Brunel University, Uxbridge, United Kingdom

J.E. Cole , A. Khan, P. Kyberd , I.D. Reid 

Baylor University, Waco, Texas, USA

S. Abdullin , A. Brinkerhoff , E. Collins , M.R. Darwish , J. Dittmann , K. Hatakeyama , V. Hegde , J. Hiltbrand , B. McMaster , J. Samudio , S. Sawant , C. Sutantawibul , J. Wilson 

Bethel University, St. Paul, Minnesota, USA

J.M. Hogan⁸² 

Catholic University of America, Washington, DC, USA

R. Bartek , A. Dominguez , S. Raj , A.E. Simsek , S.S. Yu 

The University of Alabama, Tuscaloosa, Alabama, USA

B. Bam , A. Buchot Perraguin , S. Campbell, R. Chudasama , S.I. Cooper , C. Crovella , G. Fidalgo , S.V. Gleyzer , A. Khukhunaishvili , K. Matchev , E. Pearson, C.U. Perez , P. Rumerio⁸³ , E. Usai , R. Yi 

Boston University, Boston, Massachusetts, USA

S. Cholak , G. De Castro, Z. Demiragli , C. Erice , C. Fangmeier , C. Fernandez Madrazo , E. Fontanesi , J. Fulcher , F. Golf , S. Jeon , J. O'Cain, I. Reed , J. Rohlf , K. Salyer , D. Sperka , D. Spitzbart , I. Suarez , A. Tsatsos , E. Wurtz, A.G. Zecchinelli 

Brown University, Providence, Rhode Island, USA

G. Barone , G. Benelli , D. Cutts , S. Ellis , L. Gouskos , M. Hadley , U. Heintz , K.W. Ho , T. Kwon , G. Landsberg , K.T. Lau , J. Luo , S. Mondal , J. Roloff, T. Russell , S. Sagir⁸⁴ , X. Shen , M. Stamenkovic , N. Venkatasubramanian 

University of California, Davis, Davis, California, USA

S. Abbott , B. Barton , R. Breedon , H. Cai , M. Calderon De La Barca Sanchez , E. Cannaert, M. Chertok , M. Citron , J. Conway , P.T. Cox , R. Erbacher , O. Kukral , G. Mocellin , S. Ostrom , I. Salazar Segovia, W. Wei , S. Yoo 

University of California, Los Angeles, California, USA

K. Adamidis, M. Bachtis , D. Campos, R. Cousins , A. Datta , G. Flores Avila , J. Hauser , M. Ignatenko , M.A. Iqbal , T. Lam , Y.f. Lo , E. Manca , A. Nunez Del Prado , D. Saltzberg , V. Valuev 

University of California, Riverside, Riverside, California, USA

R. Clare , J.W. Gary , G. Hanson 

University of California, San Diego, La Jolla, California, USA

A. Aportela , A. Arora , J.G. Branson , S. Cittolin , S. Cooperstein , B. D'Anzi , D. Diaz , J. Duarte , L. Giannini , Y. Gu, J. Guiang , V. Krutelyov , R. Lee , J. Letts , H. Li, M. Masciovecchio , F. Mokhtar , S. Mukherjee , M. Pieri , D. Primosch, M. Quinnan , V. Sharma , M. Tadel , E. Vourliotis , F. Würthwein , A. Yagil , Z. Zhao

University of California, Santa Barbara - Department of Physics, Santa Barbara, California, USA

A. Barzdukas , L. Brennan , C. Campagnari , S. Carron Montero⁸⁵ , K. Downham , C. Grieco , M.M. Hussain, J. Incandela , M.W.K. Lai, A.J. Li , P. Masterson , J. Richman , S.N. Santpur , U. Sarica , R. Schmitz , F. Setti , J. Sheplock , D. Stuart , T.Á. Vámi , X. Yan , D. Zhang 

California Institute of Technology, Pasadena, California, USA

A. Albert , S. Bhattacharya , A. Bornheim , O. Cerri, R. Kansal , J. Mao , H.B. New-

man , G. Reales Gutiérrez, T. Sievert, M. Spiropulu , J.R. Vlimant , R.A. Wynne , S. Xie 

Carnegie Mellon University, Pittsburgh, Pennsylvania, USA

J. Alison , S. An , M. Cremonesi, V. Dutta , E.Y. Ertorer , T. Ferguson , T.A. Gómez Espinosa , A. Harilal , A. Kallil Tharayil, M. Kanemura, C. Liu , P. Meiring , T. Mudholkar , S. Murthy , P. Palit , K. Park , M. Paulini , A. Roberts , A. Sanchez , W. Terrill 

University of Colorado Boulder, Boulder, Colorado, USA

J.P. Cumalat , W.T. Ford , A. Hart , A. Hassani , S. Kwan , J. Pearkes , C. Savard , N. Schonbeck , K. Stenson , K.A. Ulmer , S.R. Wagner , N. Zipper , D. Zuolo 

Cornell University, Ithaca, New York, USA

J. Alexander , X. Chen , J. Dickinson , A. Duquette, J. Fan , X. Fan , J. Grassi , S. Hogan , P. Kotamnives , J. Monroy , G. Niendorf , M. Oshiro , J.R. Patterson , A. Ryd , J. Thom , P. Wittich , R. Zou , L. Zygala 

Fermi National Accelerator Laboratory, Batavia, Illinois, USA

M. Albrow , M. Alyari , O. Amram , G. Apollinari , A. Apresyan , L.A.T. Bauerdick , D. Berry , J. Berryhill , P.C. Bhat , K. Burkett , J.N. Butler , A. Canepa , G.B. Cerati , H.W.K. Cheung , F. Chlebana , C. Cosby , G. Cummings , I. Dutta , V.D. Elvira , J. Freeman , A. Gandrakota , Z. Gecse , L. Gray , D. Green, A. Grummer , S. Grünendahl , D. Guerrero , O. Gutsche , R.M. Harris , T.C. Herwig , J. Hirschauer , B. Jayatilaka , S. Jindariani , M. Johnson , U. Joshi , T. Klijnsma , B. Klima , K.H.M. Kwok , S. Lammel , C. Lee , D. Lincoln , R. Lipton , T. Liu , K. Maeshima , D. Mason , P. McBride , P. Merkel , S. Mrenna , S. Nahn , J. Ngadiuba , D. Noonan , S. Norberg, V. Papadimitriou , N. Pastika , K. Pedro , C. Pena⁸⁶ , C.E. Perez Lara , F. Ravera , A. Reinsvold Hall⁸⁷ , L. Ristori , M. Safdari , E. Sexton-Kennedy , N. Smith , A. Soha , L. Spiegel , S. Stoynev , J. Strait , L. Taylor , S. Tkaczyk , N.V. Tran , L. Uplegger , E.W. Vaandering , C. Wang , I. Zoi 

University of Florida, Gainesville, Florida, USA

C. Aruta , P. Avery , D. Bourilkov , P. Chang , V. Cherepanov , R.D. Field, C. Huh , E. Koenig , M. Kolosova , J. Konigsberg , A. Korytov , N. Menendez , G. Mitselmakher , K. Mohrman , A. Muthirakalayil Madhu , N. Rawal , S. Rosenzweig , V. Sulimov , Y. Takahashi , J. Wang 

Florida State University, Tallahassee, Florida, USA

T. Adams , A. Al Kadhimi , A. Askew , S. Bower , R. Goff, R. Hashmi , R.S. Kim , T. Kolberg , G. Martinez , M. Mazza , H. Prosper , P.R. Prova, M. Wulansatiti , R. Yohay 

Florida Institute of Technology, Melbourne, Florida, USA

B. Alsufyani , S. Butalla , S. Das , M. Hohlmann , M. Lavinsky, E. Yanes

University of Illinois Chicago, Chicago, Illinois, USA

M.R. Adams , N. Barnett, A. Baty , C. Bennett , R. Cavanaugh , R. Escobar Franco , O. Evdokimov , C.E. Gerber , H. Gupta , M. Hawksworth, A. Hingrajiya, D.J. Hofman , J.h. Lee , C. Mills , S. Nanda , G. Nigmatkulov , B. Ozek , T. Phan, D. Pilipovic , R. Pradhan , E. Prifti, P. Roy, T. Roy , N. Singh, M.B. Tonjes , N. Varelas , M.A. Wadud , J. Yoo 

The University of Iowa, Iowa City, Iowa, USA

M. Alhusseini , D. Blend , K. Dilsiz⁸⁸ , O.K. Köseyan , A. Mestvirishvili⁸⁹ , O. Neogi,

H. Ogul⁹⁰ , Y. Onel , A. Penzo , C. Snyder, E. Tiras⁹¹ 

Johns Hopkins University, Baltimore, Maryland, USA

B. Blumenfeld , J. Davis , A.V. Gritsan , L. Kang , S. Kyriacou , P. Maksimovic ,
M. Roguljic , S. Sekhar , M.V. Srivastav , M. Swartz 

The University of Kansas, Lawrence, Kansas, USA

A. Abreu , L.F. Alcerro Alcerro , J. Anguiano , S. Arteaga Escatel , P. Baringer ,
A. Bean , Z. Flowers , D. Grove , J. King , G. Krintiras , M. Lazarovits ,
C. Le Mahieu , J. Marquez , M. Murray , M. Nickel , S. Popescu⁹² , C. Rogan ,
C. Royon , S. Rudrabhatla , S. Sanders , C. Smith , G. Wilson 

Kansas State University, Manhattan, Kansas, USA

B. Allmond , N. Islam, A. Ivanov , K. Kaadze , Y. Maravin , J. Natoli , G.G. Reddy ,
D. Roy , G. Sorrentino 

University of Maryland, College Park, Maryland, USA

A. Baden , A. Belloni , J. Bistany-riebman, S.C. Eno , N.J. Hadley , S. Jabeen ,
R.G. Kellogg , T. Koeth , B. Kronheim, S. Lascio , P. Major , A.C. Mignerey ,
C. Palmer , C. Papageorgakis , M.M. Paranjpe, E. Popova⁹³ , A. Shevelev , L. Zhang 

Massachusetts Institute of Technology, Cambridge, Massachusetts, USA

C. Baldenegro Barrera , H. Bossi , S. Bright-Thonney , I.A. Cali , Y.c. Chen ,
P.c. Chou , M. D'Alfonso , J. Eysermans , C. Freer , G. Gomez-Ceballos , M. Gon-
charov, G. Grosso , P. Harris, D. Hoang , G.M. Innocenti , D. Kovalskiy , J. Krupa ,
L. Lavezzo , Y.-J. Lee , K. Long , C. McGinn , A. Novak , M.I. Park , C. Paus ,
C. Reissel , C. Roland , G. Roland , S. Rothman , T.a. Sheng , G.S.F. Stephans ,
D. Walter , Z. Wang , B. Wyslouch , T. J. Yang 

University of Minnesota, Minneapolis, Minnesota, USA

B. Crossman , W.J. Jackson, C. Kapsiak , M. Krohn , D. Mahon , J. Mans ,
B. Marzocchi , R. Rusack , O. Sancar , R. Saradhy , N. Strobbe 

University of Nebraska-Lincoln, Lincoln, Nebraska, USA

K. Bloom , D.R. Claes , G. Haza , J. Hossain , C. Joo , I. Kravchenko , A. Rohilla ,
J.E. Siado , W. Tabb , A. Vagnerini , A. Wightman , F. Yan 

State University of New York at Buffalo, Buffalo, New York, USA

H. Bandyopadhyay , L. Hay , H.w. Hsia , I. Iashvili , A. Kalogeropoulos ,
A. Kharchilava , A. Mandal , M. Morris , D. Nguyen , S. Rappoccio , H. Rejeb Sfar,
A. Williams , P. Young , D. Yu 

Northeastern University, Boston, Massachusetts, USA

G. Alverson , E. Barberis , J. Bonilla , B. Bylsma, M. Campana , J. Dervan ,
Y. Haddad , Y. Han , I. Israr , A. Krishna , M. Lu , N. Manganelli , R. McCarthy ,
D.M. Morse , T. Orimoto , A. Parker , L. Skinnari , C.S. Thoreson , E. Tsai ,
D. Wood 

Northwestern University, Evanston, Illinois, USA

S. Dittmer , K.A. Hahn , Y. Liu , M. McGinnis , Y. Miao , D.G. Monk ,
M.H. Schmitt , A. Taliencio , M. Velasco , J. Wang 

University of Notre Dame, Notre Dame, Indiana, USA

G. Agarwal , R. Band , R. Bucci, S. Castells , A. Das , A. Ehnis, R. Goldouzian ,
M. Hildreth , K. Hurtado Anampa , T. Ivanov , C. Jessop , A. Karneyeu , K. Lannon 

J. Lawrence , N. Loukas , L. Lutton , J. Mariano , N. Marinelli, I. Mcalister, T. McCauley , C. Mcgrady , C. Moore , Y. Musienko²³ , H. Nelson , M. Osherson , A. Piccinelli , R. Ruchti , A. Townsend , Y. Wan, M. Wayne , H. Yockey

The Ohio State University, Columbus, Ohio, USA

A. Basnet , M. Carrigan , R. De Los Santos , L.S. Durkin , C. Hill , M. Joyce , M. Nunez Ornelas , D.A. Wenzl, B.L. Winer , B. R. Yates 

Princeton University, Princeton, New Jersey, USA

H. Bouchamaoui , P. Das , G. Dezoort , P. Elmer , A. Frankenthal , M. Galli , B. Greenberg , N. Haubrich , K. Kennedy, G. Kopp , Y. Lai , D. Lange , A. Loeliger , D. Marlow , I. Ojalvo , J. Olsen , F. Simpson , D. Stickland , C. Tully 

University of Puerto Rico, Mayaguez, Puerto Rico, USA

S. Malik , R. Sharma 

Purdue University, West Lafayette, Indiana, USA

S. Chandra , R. Chawla , A. Gu , L. Gutay, M. Jones , A.W. Jung , D. Kondratyev , M. Liu , G. Negro , N. Neumeister , G. Paspalaki , S. Piperov , N.R. Saha , J.F. Schulte , F. Wang , A. Wildridge , W. Xie , Y. Yao , Y. Zhong 

Purdue University Northwest, Hammond, Indiana, USA

N. Parashar , A. Pathak , E. Shumka 

Rice University, Houston, Texas, USA

D. Acosta , A. Agrawal , C. Arbour , T. Carnahan , K.M. Ecklund , S. Freed, P. Gardner, F.J.M. Geurts , T. Huang , I. Krommydas , N. Lewis, W. Li , J. Lin , O. Miguel Colin , B.P. Padley , R. Redjimi , J. Rotter , E. Yigitbasi , Y. Zhang 

University of Rochester, Rochester, New York, USA

O. Bessidskaia Bylund, A. Bodek , P. de Barbaro[†] , R. Demina , A. Garcia-Bellido , H.S. Hare , O. Hindrichs , N. Parmar , P. Parygin⁹³ , H. Seo , R. Taus 

Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA

B. Chiarito, J.P. Chou , S.V. Clark , S. Donnelly, D. Gadkari , Y. Gershtein , E. Halkiadakis , M. Heindl , C. Houghton , D. Jaroslawski , A. Kobert , S. Konstantinou , I. Laflotte , A. Lath , J. Martins , B. Rand , J. Reichert , P. Saha , S. Salur , S. Schnetzer, S. Somalwar , R. Stone , S.A. Thayil , S. Thomas, J. Vora 

University of Tennessee, Knoxville, Tennessee, USA

D. Ally , A.G. Delannoy , S. Fiorendi , J. Harris, T. Holmes , A.R. Kanuganti , N. Karunarathna , J. Lawless, L. Lee , E. Nibigira , B. Skipworth, S. Spanier 

Texas A&M University, College Station, Texas, USA

D. Aebi , M. Ahmad , T. Akhter , K. Androsov , A. Bolshov, O. Bouhali⁹⁴ , A. Cagnotta , V. D'Amante , R. Eusebi , P. Flanagan , J. Gilmore , Y. Guo, T. Kamon , S. Luo , R. Mueller , A. Safonov 

Texas Tech University, Lubbock, Texas, USA

N. Akchurin , J. Damgov , Y. Feng , N. Gogate , Y. Kazhykarim, K. Lamichhane , S.W. Lee , C. Madrid , A. Mankel , T. Peltola , I. Volobouev 

Vanderbilt University, Nashville, Tennessee, USA

E. Appelt , Y. Chen , S. Greene, A. Gurrola , W. Johns , R. Kunnawalkam Elayavalli , A. Melo , D. Rathjens , F. Romeo , P. Sheldon , S. Tuo , J. Velkovska , J. Viinikainen 

J. Zhang

University of Virginia, Charlottesville, Virginia, USA

B. Cardwell , H. Chung , B. Cox , J. Hakala , G. Hamilton Ilha Machado, R. Hirosky ,
M. Jose, A. Ledovskoy , C. Mantilla , C. Neu , C. Ramón Álvarez 

Wayne State University, Detroit, Michigan, USA

S. Bhattacharya , P.E. Karchin 

University of Wisconsin - Madison, Madison, Wisconsin, USA

A. Aravind , S. Banerjee , K. Black , T. Bose , E. Chavez , S. Dasu , P. Everaerts ,
C. Galloni, H. He , M. Herndon , A. Herve , C.K. Koraka , S. Lomte , R. Loveless ,
A. Mallampalli , A. Mohammadi , S. Mondal, T. Nelson, G. Parida , L. Pétré ,
D. Pinna , A. Savin, V. Shang , V. Sharma , W.H. Smith , D. Teague, H.F. Tsoi ,
W. Vetens , A. Warden 

Authors affiliated with an international laboratory covered by a cooperation agreement with CERN

S. Afanasiev , V. Alexakhin , Yu. Andreev , T. Aushev , D. Budkouski , R. Chistov⁹⁵ ,
M. Danilov⁹⁵ , T. Dimova⁹⁵ , A. Ershov⁹⁵ , S. Gninenko , I. Gorbunov ,
A. Gribushin⁹⁵ , A. Kamenev , V. Karjavine , M. Kirsanov , V. Klyukhin⁹⁵ ,
O. Kodolova^{96,93} , V. Korenkov , I. Korsakov, A. Kozyrev⁹⁵ , N. Krasnikov , A. Lanev ,
A. Malakhov , V. Matveev⁹⁵ , A. Nikitenko^{97,96} , V. Palichik , V. Perelygin ,
S. Petrushanko⁹⁵ , S. Polikarpov⁹⁵ , O. Radchenko⁹⁵ , M. Savina , V. Shalaev ,
S. Shmatov , S. Shulha , Y. Skovpen⁹⁵ , K. Slizhevskiy, V. Smirnov , O. Teryaev ,
I. Tlisova⁹⁵ , A. Toropin , N. Voytishin , B.S. Yuldashev^{†98}, A. Zarubin , I. Zhizhin 

Authors affiliated with an institute formerly covered by a cooperation agreement with CERN

E. Boos , V. Bunichev , M. Dubinin⁸⁶ , V. Savrin , A. Snigirev , L. Dudko ,
K. Ivanov , V. Kim²³ , V. Murzin , V. Oreshkin , D. Sosnov 

†: Deceased

¹Also at Yerevan State University, Yerevan, Armenia

²Also at TU Wien, Vienna, Austria

³Also at Ghent University, Ghent, Belgium

⁴Also at Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil

⁵Also at FACAMP - Faculdades de Campinas, Sao Paulo, Brazil

⁶Also at Universidade Estadual de Campinas, Campinas, Brazil

⁷Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil

⁸Also at The University of the State of Amazonas, Manaus, Brazil

⁹Also at University of Chinese Academy of Sciences, Beijing, China

¹⁰Also at China Center of Advanced Science and Technology, Beijing, China

¹¹Also at University of Chinese Academy of Sciences, Beijing, China

¹²Also at School of Physics, Zhengzhou University, Zhengzhou, China

¹³Now at Henan Normal University, Xinxiang, China

¹⁴Also at University of Shanghai for Science and Technology, Shanghai, China

¹⁵Now at The University of Iowa, Iowa City, Iowa, USA

¹⁶Also at Center for High Energy Physics, Peking University, Beijing, China

¹⁷Also at Helwan University, Cairo, Egypt

¹⁸Now at Zewail City of Science and Technology, Zewail, Egypt

¹⁹Also at British University in Egypt, Cairo, Egypt

²⁰Also at Purdue University, West Lafayette, Indiana, USA

- ²¹Also at Université de Haute Alsace, Mulhouse, France
- ²²Also at Tbilisi State University, Tbilisi, Georgia
- ²³Also at an institute formerly covered by a cooperation agreement with CERN
- ²⁴Also at University of Hamburg, Hamburg, Germany
- ²⁵Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany
- ²⁶Also at Bergische University Wuppertal (BUW), Wuppertal, Germany
- ²⁷Also at Brandenburg University of Technology, Cottbus, Germany
- ²⁸Also at Forschungszentrum Jülich, Juelich, Germany
- ²⁹Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland
- ³⁰Also at HUN-REN ATOMKI - Institute of Nuclear Research, Debrecen, Hungary
- ³¹Now at Universitatea Babes-Bolyai - Facultatea de Fizica, Cluj-Napoca, Romania
- ³²Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary
- ³³Also at HUN-REN Wigner Research Centre for Physics, Budapest, Hungary
- ³⁴Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt
- ³⁵Also at The University of Kansas, Lawrence, Kansas, USA
- ³⁶Also at Punjab Agricultural University, Ludhiana, India
- ³⁷Also at University of Hyderabad, Hyderabad, India
- ³⁸Also at Indian Institute of Science (IISc), Bangalore, India
- ³⁹Also at University of Visva-Bharati, Santiniketan, India
- ⁴⁰Also at IIT Bhubaneswar, Bhubaneswar, India
- ⁴¹Also at Institute of Physics, Bhubaneswar, India
- ⁴²Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany
- ⁴³Also at Isfahan University of Technology, Isfahan, Iran
- ⁴⁴Also at Sharif University of Technology, Tehran, Iran
- ⁴⁵Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran
- ⁴⁶Also at Department of Physics, Faculty of Science, Arak University, ARAK, Iran
- ⁴⁷Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy
- ⁴⁸Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy
- ⁴⁹Also at Università degli Studi Guglielmo Marconi, Roma, Italy
- ⁵⁰Also at Scuola Superiore Meridionale, Università di Napoli 'Federico II', Napoli, Italy
- ⁵¹Also at Fermi National Accelerator Laboratory, Batavia, Illinois, USA
- ⁵²Also at Lulea University of Technology, Lulea, Sweden
- ⁵³Also at Laboratori Nazionali di Legnaro dell'INFN, Legnaro, Italy
- ⁵⁴Also at Consiglio Nazionale delle Ricerche - Istituto Officina dei Materiali, Perugia, Italy
- ⁵⁵Also at UPES - University of Petroleum and Energy Studies, Dehradun, India
- ⁵⁶Also at Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France
- ⁵⁷Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia
- ⁵⁸Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka
- ⁵⁹Also at Saegis Campus, Nugegoda, Sri Lanka
- ⁶⁰Also at National and Kapodistrian University of Athens, Athens, Greece
- ⁶¹Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland
- ⁶²Also at Universität Zürich, Zurich, Switzerland
- ⁶³Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria
- ⁶⁴Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey

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- ⁶⁵Also at Konya Technical University, Konya, Turkey
- ⁶⁶Also at Izmir Bakircay University, Izmir, Turkey
- ⁶⁷Also at Adiyaman University, Adiyaman, Turkey
- ⁶⁸Also at Bozok Universititesi Rektörlüğü, Yozgat, Turkey
- ⁶⁹Also at Istanbul Sabahattin Zaim University, Istanbul, Turkey
- ⁷⁰Also at Marmara University, Istanbul, Turkey
- ⁷¹Also at Milli Savunma University, Istanbul, Turkey
- ⁷²Also at Informatics and Information Security Research Center, Gebze/Kocaeli, Turkey
- ⁷³Also at Kafkas University, Kars, Turkey
- ⁷⁴Now at Istanbul Okan University, Istanbul, Turkey
- ⁷⁵Also at Hacettepe University, Ankara, Turkey
- ⁷⁶Also at Erzincan Binali Yildirim University, Erzincan, Turkey
- ⁷⁷Also at Istanbul University - Cerrahpasa, Faculty of Engineering, Istanbul, Turkey
- ⁷⁸Also at Yildiz Technical University, Istanbul, Turkey
- ⁷⁹Also at Istinye University, Istanbul, Turkey
- ⁸⁰Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom
- ⁸¹Also at Monash University, Faculty of Science, Clayton, Australia
- ⁸²Also at Bethel University, St. Paul, Minnesota, USA
- ⁸³Also at Università di Torino, Torino, Italy
- ⁸⁴Also at Karamanoğlu Mehmetbey University, Karaman, Turkey
- ⁸⁵Also at California Lutheran University,, Thousand Oaks, California, USA
- ⁸⁶Also at California Institute of Technology, Pasadena, California, USA
- ⁸⁷Also at United States Naval Academy, Annapolis, Maryland, USA
- ⁸⁸Also at Bingol University, Bingol, Turkey
- ⁸⁹Also at Georgian Technical University, Tbilisi, Georgia
- ⁹⁰Also at Sinop University, Sinop, Turkey
- ⁹¹Also at Erciyes University, Kayseri, Turkey
- ⁹²Also at Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Bucharest, Romania
- ⁹³Now at another institute formerly covered by a cooperation agreement with CERN
- ⁹⁴Also at Hamad Bin Khalifa University (HBKU), Doha, Qatar
- ⁹⁵Also at another institute formerly covered by a cooperation agreement with CERN
- ⁹⁶Also at Yerevan Physics Institute, Yerevan, Armenia
- ⁹⁷Also at Imperial College, London, United Kingdom
- ⁹⁸Also at Institute of Nuclear Physics of the Uzbekistan Academy of Sciences, Tashkent, Uzbekistan