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Search for low-mass resonances decaying into two jets and produced in association with a photon or a jet at \sqrt{s} = 13 TeV with the ATLAS detector

The ATLAS Collaboration

A search is performed for localized excesses in the low-mass dijet invariant mass distribution, targeting a hypothetical new particle decaying into two jets and produced in association with either a high transverse momentum photon or a jet. The search uses the full Run 2 data sample from LHC proton-proton collisions collected by the ATLAS experiment at a center-of-mass energy of 13 TeV during 2015–2018. Two variants of the search are presented for each type of initial-state radiation: one that makes no jet flavor requirements and one that requires both of the jets to have been identified as containing b-hadrons. No excess is observed relative to the Standard Model prediction, and the data are used to set upper limits on the production cross-section for a benchmark Z' model and, separately, for generic, beyond the Standard Model scenarios which might produce a Gaussian-shaped contribution to dijet invariant mass distributions. The results extend the current constraints on dijet resonances to the mass range between 200 and 650 GeV.

1 Introduction

The Standard Model of particle physics (SM) successfully describes a wide range of phenomena across many orders of magnitude of energy at colliders, astrophysical observatories, and other experiments [1]. Nevertheless it fails to account for many observed phenomena, such as the presence of dark matter (DM) [2–4] and the imbalance of matter over antimatter in the observable universe [5]. While there are a wide range of proposed models of physics beyond the SM (BSM) to explain these effects [6–17], this paper focuses on a search for a hadronically decaying resonance.

Z' bosons are hypothetical spin-1 vector bosons that are singlets under electric and color charge [1], predicted in a variety of BSM models as a potential mediator to dark matter particles [14, 17–22]. Direct Z' searches have shown to be an effective way of constraining BSM models [19–21]. Searches at the Large Hadron Collider (LHC) for Z' decays into pairs of charged leptons set strong constraints [21, 23], but may be eluded by choosing the couplings of the Z' such that it is *leptophobic*, meaning it does not couple to SM leptons. However, constraints from searches for hadronically decaying Z' resonances cannot be avoided in this way, since the same Z'-q-q decay vertex would also be responsible for Z' production at the LHC for any kind of resonance or search for topologies with a single reconstructed object recoiling against a large missing transverse energy. Additionally, the Z'-q-q vertex would be necessary for the dark matter self-annihilation cross-section that drives the relic-density for weakly interacting massive particles [2].

The ATLAS and CMS Collaborations at the LHC have published several searches that are sensitive to hadronic Z' production using a variety of strategies. Searches for excesses in the invariant mass of the two jets with the highest transverse momentum (p_T) in the event constrain Z' production at high masses [24–38]. The lower reach of these searches are limited by the bandwidth available to single- and multijet triggers, and so different strategies must be used to gain sensitivity to Z' masses below around twice the p_T threshold where these triggers become fully efficient, around 1 TeV. One strategy is to record only minimal information, but at a higher rate than the standard triggers, enabling lower trigger thresholds [31, 39]. A complementary approach, explored in this paper, relies on the production of high- p_T initial-state-radiation (ISR), such as a jet or a photon, which recoils against the Z', enabling access to lower dijet masses without trigger bias. This type of search was performed by the ATLAS and CMS Collaborations in topologies where the two Z' decay products can be resolved as individual jets [40, 41] and the case where the decay products are sufficiently boosted to be reconstructed into a single large-radius jet [42–46].

This paper explores the resolved case, where the resonance decay products are reconstructed into two jets, covering a resonance mass range of 200 GeV to 650 GeV. In total, four different channels are studied, based on the type of ISR particle and the flavor composition of the Z' decay products. The first channel, γjj , selects events in which the ISR is a photon and applies an inclusive selection on the jets that form the Z'. The second channel, γbb , focuses on the case where the ISR is a photon and both jets are identified as containing b-hadrons (b-tagged). The third and fourth channels focus on the case where the ISR is a jet, both in the inclusive case (jjj), and the case where both jets from the candidate resonance decay are b-tagged (jbb). While b-tagging requirements impose assumptions on the Z' decay products, they also provide significant background suppression, resulting in a more powerful search for some models. The cross-sections for the jjj and jbb processes are higher than for γjj and γbb . However, the jjj channel has the additional challenge of identifying which jets might correspond to the Z' decay products, and which correspond to the ISR. Furthermore, different trigger thresholds apply for recorded events triggered by a photon or a jet. Hence a different kinematic acceptance is imposed on the ISR object for the various channels. The two channels where the ISR is a photon are collectively called the *photon* channels, while the two where the ISR is a jet are referred to as the *trijet* channels.

In addition to the increased luminosity, this analysis features several improvements over the previous analysis targetting the same signature [40], including a better b-tagging algorithm that reduces the background and the inclusion of the trijet channels. The dominant backgrounds for this search are multijet and single-photon production, which are estimated from data using a functional form fit to the m_{jj} distribution. For each channel, estimates for background and signal are combined into a likelihood function, where the signal yields are controlled by a free multiplicative parameter. The signal is interpreted as either a Z' boson or a generic narrow resonance producing a Gaussian-shaped bump in the measured mass spectrum.

The paper is organized as follows. The ATLAS detector at the LHC and the data and simulated samples used in this search are described in Sections 2–3. The object and event selection for the four channels are summarized in Section 4. The methodology of the search is introduced in Section 5 and the systematic uncertainties are discussed in Section 6. The results are presented in Section 7 and conclusions are drawn in Section 8.

2 The ATLAS detector

The ATLAS detector [47] at the LHC covers nearly the entire solid angle around the collision point. It consists of an inner tracking detector surrounded by a thin superconducting solenoid, electromagnetic and hadron calorimeters, and a muon spectrometer incorporating three large superconducting air-core toroidal magnets.

The inner-detector system is immersed in a 2 T axial magnetic field and provides charged-particle tracking in the range of $|\eta| < 2.5$. The high-granularity silicon pixel detector covers the vertex region and typically provides four measurements per track, the first hit normally being in the insertable B-layer installed before Run 2 [48, 49]. It is followed by the silicon microstrip tracker, which usually provides eight measurements per track. These silicon detectors are complemented by the transition radiation tracker (TRT), which enables radially extended track reconstruction up to $|\eta| = 2.0$. The TRT also provides electron identification information based on the fraction of hits (typically 30 in total) above a higher energy-deposit threshold corresponding to transition radiation.

The calorimeter system covers the pseudorapidity range $|\eta| < 4.9$. Within the region $|\eta| < 3.2$, electromagnetic calorimetry is provided by barrel and endcap high-granularity lead/liquid-argon (LAr) calorimeters, with an additional thin LAr presampler covering $|\eta| < 1.8$ to correct for energy loss in material upstream of the calorimeters. Hadron calorimetry is provided by the steel/scintillator-tile calorimeter, segmented into three barrel structures within $|\eta| < 1.7$, and two copper/LAr hadron endcap calorimeters. The solid angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeter modules optimized for electromagnetic and hadronic energy measurements respectively.

The muon spectrometer comprises separate trigger and high-precision tracking chambers measuring the deflection of muons in a magnetic field generated by the superconducting air-core toroidal magnets. The field integral of the toroids ranges between 2.0 and 6.0 Tm across most of the detector. Three layers

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the *z*-axis along the beam pipe. The *x*-axis points from the IP to the center of the LHC ring, and the *y*-axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the *z*-axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta y)^2 + (\Delta \phi)^2}$, where $y = (1/2)[(E + p_z)/(E - p_z)]$ is the object's rapidity defined by its energy and longitudinal momentum.

of precision chambers, each consisting of layers of monitored drift tubes, cover the region $|\eta| < 2.7$, complemented by cathode-strip chambers in the forward region, where the background is highest. The muon trigger system covers the range $|\eta| < 2.4$ with resistive-plate chambers in the barrel, and thin-gap chambers in the endcap regions.

Interesting events are selected by the first-level trigger system implemented in custom hardware, followed by selections made by algorithms implemented in software in the high-level trigger [50]. The first-level trigger accepts events from the 40 MHz bunch crossings at a rate below 100 kHz, which the high-level trigger further reduces to record events to disk at about 1 kHz.

An extensive software suite [51] is used in data simulation, in the reconstruction and analysis of real and simulated data, in detector operations, and in the trigger and data acquisition systems of the experiment.

3 Data and Monte Carlo simulated samples

This analysis is performed using data from proton–proton (pp) collisions at $\sqrt{s} = 13$ TeV at the LHC, collected during 2015–2018 with the ATLAS detector. The total integrated luminosity of this data sample is 140 fb⁻¹ [52], obtained using the LUCID-2 detector [53] for the primary luminosity measurements. Due to the high instantaneous luminosity and the large total inelastic pp cross-section, there are, on average, 33.7 collisions in each bunch crossing. The additional collisions not selected as the hard-scatter process are referred to as pileup collisions. Data in this analysis are required to satisfy standard quality requirements [54].

Samples of Monte Carlo (MC) simulated signal and background events are used in this analysis for optimization, estimate of possible signal contributions, and validation of background estimation strategies.

The main benchmark model is a simplified leptophobic Z' axial-vector mediator signal, which is used to test the sensitivity to potential resonant signals. The choice of this benchmark signal is motivated by the role that dijet resonance searches play in constraining possible interactions between dark matter and SM particles. This simplified model considers an s-channel mediator with axial-vector interactions, as described in Refs. [18, 55]. Two relevant parameters of this model in the dijet search are the Z' mass, and the coupling of the Z' to quarks (g_q) , which affect the intrinsic width of the resonance. The mass of the dark matter fermion is set to be much heavier than the Z' mass, such that the decay width into dark matter is zero. In all channels, signal samples are generated for Z' masses between 200 GeV and 650 GeV, with the coupling of the mediator to non-top-quarks set to 0.2, regardless of the quark flavor. Such a choice for the coupling value is a benchmark model already used in other analyses [40]. Also, given the Z' intrinsic width for $g_q = 0.2$ is much smaller than the experimental resolution, such a sample is adequate to also model the shape of signals with slightly different or smaller g_q values. Signal samples are generated separately for a Z' decaying into u,d,s and c quark flavors and z' decaying into b-quarks, in order to enhance the signal statistical precision for the b-tagged channels.

The signal samples are simulated using the MadGraph generator [56] at leading-order (LO) in perturbative quantum chromodynamics (QCD) with the NNPDF2.3Lo parton distribution function (PDF) set [57]. The events are interfaced to Pythia8.244 [58, 59] to model the parton shower, hadronization, and underlying event, with parameter values set according to the A14 tune [60]. EvtGen [61] is used to model decays of heavy flavor hadrons. To further increase statistical precision in the γjj channel, a generator-level filter is applied on the photon p_T . Two sets of samples are produced: one set with this filter defined to be $p_{T,\gamma} \geq 130$ GeV and a second with the filter at $p_{T,\gamma} \geq 65$ GeV. For the trijet channels, a generator-level jet

filter of 370 GeV is applied on the leading- $p_T R = 0.4$ jet, which is fully efficient for the reconstructed jet p_T threshold. The generator-level filters are chosen based on the triggers used in the two different analyses, to ensure high statistical precision for each mass point without introducing kinematic biases.

While the final background estimate is performed using a data-driven approach, background multijet and photon-plus-jets samples are used to optimize the event selection and validate the background estimation strategy. For the photon channels, prompt single-photon production is simulated with the Sherpa 2.2 [62] generator. In this arrangement and thanks to the Comix [63] and OpenLoops [64–66] libraries, matrix elements for up to two partons are calculated to next-to-leading-order (NLO) accuracy, while matrix elements for up to four partons are calculated at LO in QCD. They are matched with the Sherpa parton shower [67] using the MEPS@NLO prescription [68–71] with a dynamic merging scale cut [72] of 20 GeV. Photons are required to be isolated according to a smooth-cone isolation criterion [73]. Samples are simulated using the NNPDF3.0NNLO PDF set [74], along with the dedicated set of tuned parton-shower parameters developed by the Sherpa authors.

For the trijet channels, background samples of simulated multijet processes are generated using Pythia 8.230 with the $2 \rightarrow 2$ matrix element, and additional jets are produced through the parton shower. Events are simulated using the A14 tune [60], the Lund string hadronization model and the NNPDF2.3Lo PDF set. The Pythia parton shower algorithm uses a dipole-style $p_{\rm T}$ -ordered evolution, and its renormalisation and factorisation scales are set to the geometric mean of the squared transverse masses of the outgoing particles.

The generated background and signal events are passed through a detailed detector simulation [75] based on Geant4 [76]. Additional minimum-bias interactions simulated using Pythia8 with the A3 tune [77] and the NNPDF2.3Lo PDF set [57] are overlayed to the background and signal events, to model pileup interactions. The distribution of the average number of pileup interactions in simulation is reweighted during data analysis to match that observed in Run 2 data.

4 Event reconstruction and selection

4.1 Object reconstruction

While the four different channels use different combinations of objects, the object reconstruction is standardized across all of the different channels. Three different objects are used: jets, *b*-tagged jets, and photons.

Selected events are required to contain a primary vertex with at least two associated tracks with transverse momentum $p_T > 0.5$ GeV [78]. The vertex with the highest sum of p_T^2 of the associated tracks is taken as the primary vertex.

Jets in this analysis are reconstructed from particle flow objects [79] using the anti- k_t algorithm [80] as implemented in FastJet [81], using a jet radius parameter R = 0.4. The ATLAS particle flow algorithm combines measurements from the ATLAS inner detector and calorimeter systems [82] to improve the jet energy resolution, reduce sensitivity to pileup effects, and improve jet reconstruction efficiency, especially at low jet p_T . The jet energy scale of particle flow jets is calibrated with a combination of simulation-based and *in situ* corrections [83]. Calibrated jets are required to have a $p_T > 25$ GeV and satisfy $|\eta| < 2.5$. To reduce the effects of pileup, jets with $p_T < 60$ GeV and $|\eta| < 2.4$ are required to satisfy the "tight" working point of the jet vertex tagger criteria [84], which selects jets originating from the selected primary vertex.

A set of quality criteria is also applied to reject events containing at least one jet arising from non-collision sources or detector noise [85].

Reconstructed jets are b-tagged using the DL1r multivariate algorithm [86]. Jets are considered b-tagged if they satisfy the 77% b-tag efficiency working point, as measured in inclusive $t\bar{t}$ events. The corresponding rejection factors for gluon/light-quark jets and charm-quark jets are approximately 200 and 5, respectively. This working point was found to result in optimal signal sensitivity for both the γbb and jbb channels.

Photons are reconstructed from energy deposits in the electromagnetic calorimeter [87]. Reconstructed photon candidates are required to have $p_T > 25$ GeV, and only central photons with $|\eta| < 2.5$ are considered, excluding the barrel/endcap calorimeter transition region (1.37 < $|\eta| < 1.52$). Only photons that satisfy the *tight* identification requirement are considered, and additional tight isolation requirements are imposed to further suppress contamination from jets [87].

An overlap removal procedure is applied to avoid any double-counting between the reconstructed photons and jets, by removing any jet with an angular distance $\sqrt{\Delta\eta^2 + \Delta\phi^2} < 0.4$ to any photon.

4.2 Photon channel selection

The two photon channels (γjj) and $\gamma bb)$ are required to satisfy the singel-photon trigger [88] with the lowest threshold where all events passing it are saved (unprescaled), and a single photon with $p_T > 150$ GeV is required in order for the trigger selection to be fully efficient. In addition, the event is required to have two reconstructed jets passing the selection in Section 4.1, and the two highest- p_T (leading) jets are used to reconstruct the resonance candidate. Several dijet and γ -jet kinematic variables were studied to enhance the $Z' \to q\bar{q}$ signal over the expected background, and the strongest discriminating variable was found to be the asymmetry $y^* = \frac{|y_1 - y_2|}{2}$, where y_1 and y_2 are the rapidities of the two selected jets. It is expected that a signal from Z' production produces isotropic jets while forward jets are more common in background events, causing the signal to have smaller values of y^* than the background. The shape of the y^* variable is found to be well described by the background simulation, with maximum differences of a few percent. A cut of $y^* < 0.825$ is applied to maximize the signal sensitivity across the different mass ranges. For the ybb channel, in addition to satisfying the selection for the yjj channel, the two leading jets in the event are required to satisfy the b-tagging requirements described in Section 4.1.

4.3 Trijet channel selection

In both the jjj and jbb channels, events are required to have at least three reconstructed jets, where the leading jet is required to be above 475 GeV to be fully efficient for the lowest unprescaled single jet trigger [89]. When considering only the three leading- p_T jets in a given event, there are three possible choices to assign two jets as corresponding to the resonance decay product candidates.² In the case where the resonance decays into two b-quarks, the combinatorial problem can be in principle trivially solved by identifying the two b-tagged-jets in the event. While possible, this causes challenges in background modeling, and so this trivial solution is not used. To illustrate the combinatorial problem more clearly, Figure 1 shows the particle-level dijet mass spectrum of different jet pairs for three different Z' signal samples. Each jet is assigned a number based on its p_T ordering, e.g., m_{12} is the dijet mass distribution

² While a fourth jet can contribute, it is typically not the dominant contribution, and so it is not considered, to simplify the combinatorics.

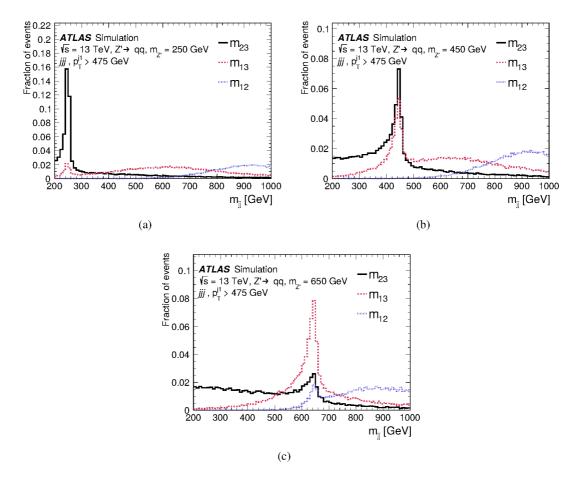


Figure 1: Distributions of particle-level dijet mass in simulated $Z' \rightarrow qq$ events corresponding to different combinations of two jets, where jets 1,2 and 3 correspond to the leading, sub-leading and third-leading jets. Signal events generated with (a) $m_{Z'} = 250$ GeV, (b) $m_{Z'} = 450$ GeV and (c) $m_{Z'} = 650$ GeV are shown.

 (m_{jj}) of the leading and subleading jets. For lower Z' mass, the m_{23} contribution clearly peaks near the Z' mass, while at higher Z' mass, contributions from m_{13} become more important.

The combinatorial challenge is further complicated by the background estimate, which relies on a functional fit to a smoothly falling m_{jj} spectrum, meaning that any additional selection to solve the combinatorial problem must also result in a smooth m_{jj} spectrum. The dominant background for these channels may be simplistically modeled as balanced high- p_T dijet events where one jet emits a softer jet, resulting in a trijet event. Since the emitting jet loses some energy to its emission, the three jets are roughly p_T ordered from low to high as the emitted jet, the emitting jet, and the non-emitting jet, it is possible to roughly map m_{23} to $m_{\text{emitted,emitting}}$, m_{13} to $m_{\text{non-emitting,emitted}}$ and m_{12} to $m_{\text{non-emitting,emitting}}$. In reality, this description of the background does not fully describe the observed background. Multiple emissions may occur, and various detector effects can spoil the p_T balance of the reconstructed jets, meaning that the assumption that m_{23} corresponds to the mass of the emitting and emitted jet may not hold in some events.

The mass spectra of the three possible pairings have different shapes, since different underlying physics contributes to each of them. This is illustrated in Figure 2(a), which shows the m_{jj} distribution for m_{12} , m_{13} , and m_{23} distributions. Any generic combination of the different pairings will result in non-smoothness in

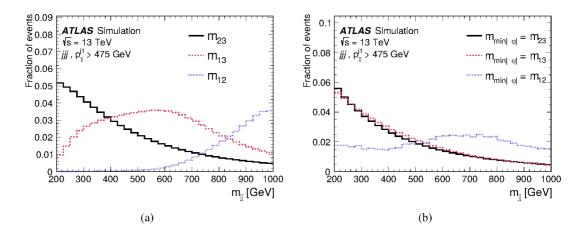


Figure 2: (a) Distributions of reconstructed dijet mass in simulated background events corresponding to different combinations of two jets, where jets 1,2 and 3 correspond to the leading, sub-leading and third-leading jets, and (b) comparison of the shape of the normalized selected dijet $m_{\rm jj}$ spectra in the case where the selected jets are the second and third leading jets ($m_{\rm min}|\Delta\phi|=m_{23}$), the first and third leading jets ($m_{\rm min}|\Delta\phi|=m_{13}$), or the first and second leading jets ($m_{\rm min}|\Delta\phi|=m_{12}$). Events are simulated with PYTHIA.

the spectrum, roughly corresponding to the turn-ons in these different pairings. Hence instead of pairing jets based solely on $p_{\rm T}$ -ordering, a different pairing method based on the kinematics of the three leading jets need to be used in the analysis.

Three similar options for choosing the dijet pairing were compared to determine which one results in the smoothest spectrum, each roughly corresponding to the mass of the emitting and emitted jet: the second-and third-leading jets (as described above), the jet pair with the minimum $|\Delta\phi|$, and the pair of jets with the lowest mass. The smoothness of the $m_{\rm jj}$ spectrum at low masses determines the fit range therefore limiting the lowest Z' mass reach covered by the search. The signal efficiencies of these three $m_{\rm jj}$ selection options are similar. However, selecting jet pairs with the minimum $|\Delta\phi|$ between them results in the smoothest low-mass spectrum in studies of multijet MC simulation.

One final aspect to consider is that in dijet-like topologies, where one of the high- p_T jets radiates a low- p_T gluon reconstructed as a jet, the gluon radiation tends to happen at small angles, while the two high- p_T jets are produced back-to-back. This means that the $|\Delta\phi|$ selection tends to identify the radiating jet and the gluon. Depending on which jet radiates the gluon, this selection can identify either the second-and third-leading jets or the first- and third-leading jets. In a small fraction of events, the $|\Delta\phi|$ selection identifies the two leading jets in the events as trijets, but where the underlying topological assumptions break down. These events, comprising around 0.2% of background events, do not produce a smooth m_{jj} spectrum, as shown in Figure 2(b). Based on this, events where the $|\Delta\phi|$ selection identifies the two leading jets are rejected to avoid biasing the fit.

To summarize, the algorithm selected to identify the Z' decay candidates takes the jet pair with the minimum $|\Delta\phi|$ between them and rejects the event if such jets correspond to the two p_T -leading jets. The pair of jets passing such requirements are referred to as the *selected jets* and, as explained above, produces the smoothest low-mass spectrum in studies of multijet MC simulation. No additional background suppression is applied, since cuts on observables such as y^* can sculpt the background distribution. For the jbb channel, in addition to passing the same jet requirements as the jjj channel, the two jets with the minimum $|\Delta\phi|$ between them are also required to be b-tagged.

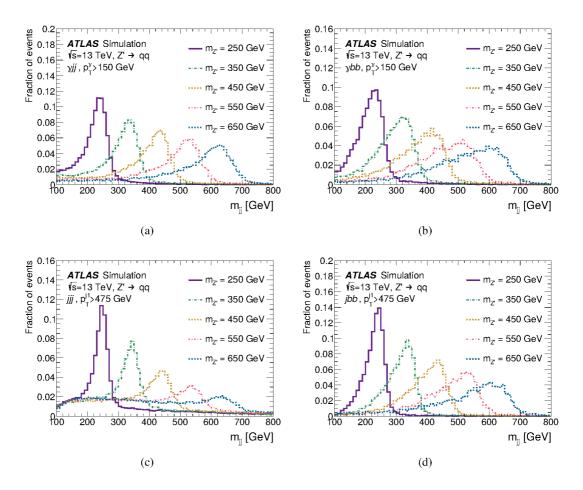


Figure 3: Distributions of reconstructed dijet mass in simulated $Z' \to qq$ events with different $m_{Z'}$ values in the (a) γjj , (b) γbb , (c) jjj, and (d) jbb channels after full event selection.

The resulting signal m_{jj} distributions after the event selection are shown for each channel in Figure 3. The jjj channel, at higher masses, has a clear tail in the m_{jj} distribution resulting from the incorrect selection of the dijet pair. This tail is less pronounced for the jbb channel, owing to the additional b-tagging requirement imposed on the selected jets. The signal efficiency for the jbb selection gets smaller at high masses, as the selected dijet pair is often not b-tagged there. Such an effect is considered less relevant than obtaining a smoothly falling m_{jj} distribution for the background at lower m_{jj} values, as the more energetic part of the m_{ij} spectrum is better studied by other analyses [31, 39].

5 Methodology

The numbers of signal and background events are estimated channel-by-channel by binned profile likelihood fits of a signal-plus-background model to the measured m_{jj} distributions. The number of selected events as a function of m_{jj} is parameterized in the fit as:

$$n_{\text{tot}}(m_{\text{ij}}) = n_S \times f_S(m_{\text{ij}}) + n_B \times f_B(m_{\text{ij}})$$
,

where n_S and n_B are the number of signal and background events, respectively, and $f_S(m_{jj})$ and $f_B(m_{jj})$ are their probability distribution functions. All the distributions entering the fit are binned finely in bins of 1 GeV width, to minimize the loss of information and the typical biases of binned likelihood fits [90].

5.1 Signal templates

The likelihood fit results are interpreted using model-independent and model-dependent strategies. For the model-independent limits, the signal probability distribution functions are taken as Gaussian distributions, with a mean equal to the resonance mass, with widths ranging from 5% to 15%. The upper end of the template width is determined based on the results of the spurious signal test and signal injection test described in Section 5.2.

For the model-dependent limits, the shape of the m_{jj} distributions from simulated samples are used directly. Signal samples are produced with a limited set of signal masses, and these templates are used as inputs to an interpolation between mass points to provide a finer signal grid using the method described in Ref. [91].

5.2 Background estimate

The selected jets in non-resonant QCD processes, which constitute the SM background for this search, result in dijet systems with smoothly-falling invariant mass distributions. In order to estimate this background in the search regions, a class of parametric functions is fit to the m_{ii} distributions:

$$f_B(x) = p_1(1-x)^{p_2}x^{p_3+p_4\ln(x)+p_5\ln^2(x)+p_6\ln^3(x)},$$

where $x = m_{jj}/\sqrt{s}$, and p_1, p_2, p_3, p_4, p_5 , and p_6 are the fitted parameters. For an *N*-parameter fit, only the first *N* parameters are allowed to vary, while the higher-*N* terms are fixed to zero. Such functions were successfully used by a wide variety of dijet and multijet resonance searches by the CDF, CMS, and ATLAS Collaborations [25, 27, 32, 34, 92–97].

The data-driven background fitting procedure employs a 5-parameter function and was validated using several cross-checks, including spurious signal tests and signal injection tests. These tests were performed using simulated samples, with a subset of checks performed on a partial data sample.

The spurious signal test evaluates whether the fitting procedure is biased such that it produces a non-zero extracted signal when fitting a data sample with no true signal. This test is performed using many pseudodata distributions which are generated from background-only fits to the simulated $m_{\rm jj}$ distribution with a 6-parameter function, which has more flexibility than the final fit function. Each distribution is tested by performing signal+background fits with a 5-parameter fit function and various signal hypotheses. Gaussian-shaped signals with widths ranging from 5 to 15% of the signal mean are tested, as well as using the signal templates from the Z' simulated samples directly. For each pseudodata distribution and signal hypothesis n_S is determined, and the median value and standard deviation of the n_S distribution are taken to be $S_{\rm spur}$ and $\sigma_{\rm spur}$ respectively. To satisfy the spurious signal requirements, $S_{\rm spur}/\sigma_{\rm spur}$ is required to be less than 0.5 for every signal hypothesis mass point. This requirement is satisfied for all channels for Z' masses between 250 GeV and 650 GeV. While the fits extend to $m_{\rm jj}$ values beyond 650 GeV, higher signal masses do not always satisfy the spurious signal tests, and are not considered further. In addition, for the jjj channel, the signal shape degrades significantly at higher signal masses, since the selected jets often do not correspond to the resonance decay products. Finally, the jbb channel passes this requirement

for resonance masses down to 200 GeV, which is possible since the m_{jj} spectrum in this channel is smooth down to lower masses than the other channels.

The signal injection test is performed to ensure that the fit is able to extract a signal component with the expected signal strength. Simulated signal models, with both Gaussian shapes and in a range of widths and signal templates, are included in the fitted background distribution with a given signal cross-section selected to be in the range $0 - 5\sigma$, where $\sigma = n_S/\sqrt{n_B}$ and the number of signal and background events are determined using a 2σ window around the injected signal peak in each test. The injected signals in this study were extracted for pseudodata distributions, with the requirement that the median of the extracted significances is within 0.5σ of the injected significance, for all the signal hypotheses and mass points individually. These tests were satisfied for all channels for the same range of resonance masses and widths as with the spurious signal tests, while they failed for Gaussian-shaped signals with widths above 15%.

The fitted m_{jj} region is determined separately for each channel. For the chosen range, the fit quality, spurious signal, and signal injection tests are must be satisfied for the MC simulation. The fitted m_{jj} region is chosen to prioritize a fit that starts as low in m_{jj} as possible while maintaining a wide enough fit range to study the relevant signal models. The lowest part of the m_{jj} spectrum does not have a smoothly falling shape because of trigger inefficiency and analysis selection effects, such as the y^* cut, and such a behavior cannot be fit by the class of functions considered. The $\gamma j j$ and $\gamma b b$ channels are found to have a smoothly falling behavior starting from $m_{jj} = 200$ GeV, while the j j j and j b b channels from $m_{jj} = 225$ GeV and $m_{jj} = 160$ GeV respectively. The upper fit ranges are defined from the edges of bins defined by the m_{jj} resolution, and are taken to be the highest among those where the background fit functions to the simulated m_{jj} distributions have a $\chi^2 p$ -value above 0.05. The procedure to determine the fit range is repeated on the measured data, to ensure the fit ranges estimated from the MC simulation are adequate. If they are not, the upper limit of the fit range is reduced by one m_{jj} resolution bin, and all the tests are repeated.

For the γbb channel, the background estimate is further validated using an ABCD method with simulated samples, following the strategy of Ref. [40]. Such method allows a test of the background estimate in the γbb channel with a higher statistical precision than available in the existing simulated sample and also serves as a test for potential biases in the m_{jj} spectrum introduced by b-tagging. The ABCD re-weighting method defines different regions based on whether events satisfy or fail to meet the y^* selection and the b-tagging requirements, providing an estimate for the b-tagged spectrum which is based on the un-tagged distribution. The ABCD-reweighted m_{jj} spectrum is found to be compatible with the one obtained by applying the γbb event selection, within the statistical precision of the re-weighted and high-statistics spectrum. Given the level of agreement, no uncertainty related to b-tagging is considered to affect the background fit. The resulting distribution is found to satisfy the spurious signal and signal injection tests described above, providing further confirmation of the fit strategy. Since this method relies on the y^* selection, the technique cannot easily be applied to the jbb channel, but the conclusions on the possible b-tagging biases from γbb apply to the jbb channel as well.

6 Systematic uncertainties

When interpreting the analysis in terms of candidate signal models, the impact of various experimental and theoretical sources of uncertainty is considered.

Luminosity. The uncertainty in the combined 2015–2018 integrated luminosity is 0.83% [52], obtained using the LUCID-2 detector [53] for the primary luminosity measurements. It is treated as a single normalization uncertainty applied as a scale factor to the signal models.

Jets. Systematic uncertainties in the R = 0.4 jet energy scale and resolution are evaluated using a series of *in situ* measurements and simulation-based techniques, thoroughly documented in Ref. [83]. The most significant source of uncertainty in the JES originates from uncertainties on the flavor of the jet. Uncertainties on the JVT are also included, but are negligible, since the uncertainties are applied only to the signal. For the photon channels, these uncertainties have a negligible impact on the efficiency, and around a 2% impact on the location of the Z' mass. For the trijet channels, the impact on the location of the Z' mass is around 1%, while the impact on the selection efficiency is around 3%.

Photons. The systematic uncertainties in the photon identification and isolation efficiencies are estimated following the prescriptions in Ref. [87]. They are evaluated by varying the correction factors for photon selection efficiencies in simulation by the corresponding uncertainties and affect the diphoton selection efficiency. The experimental uncertainties in the photon energy scale and resolution are obtained from Ref. [87], and are only applied to the γjj and γbb channels. These systematic uncertainties primarily impact the selection efficiency of the signal events, with a total impact on the efficiency of around 2%.

B-tagging. For final states with requirements on the number of b-tagged jets, additional systematic uncertainties are applied. The systematic uncertainty in the b-tagging efficiency is measured using data enriched in $t\bar{t}$ events for jet $p_T < 400$ GeV and extrapolated to higher p_T regions using a method similar to the one described in Ref. [98]. The impact of this systematic uncertainty in the efficiency of the event selection of signal events ranges from 2–3%.

Parton distribution functions. The theoretical uncertainty envelope associated with the NNPDF2.3LO PDF set is propagated through the analysis, where their impact is primarily on the normalisation of the signal events. The change in analysis selection efficiency is recalculated for each provided PDF variation, and the standard deviation of all such variations is taken as a measure of the systematic uncertainty due to PDFs. This uncertainty affects signal samples and ranges from 1–5% across the different channels, with larger effects on the photon channels than on the dijet channels.

Background modeling A systematic uncertainty to cover potential biases in modeling the background shape is accounted for using the spurious signal S_{modeling} . The value of S_{modeling} is determined as the envelope of $|S_{\text{spur}}|$ over m_{jj} and is considered to cover features of the m_{jj} background spectrum which are not modeled by the background parametrisation chosen. This is implemented in the likelihood fit as an additional signal contribution, such that

$$N_{\text{signal}}(m_{Z'}) = \sigma \times A \times BR \times \mathcal{L} + S_{\text{modeling}}(m_{Z'})\theta_{\text{modeling}},$$

where $N_{\text{signal}}(m_{Z'})$ is the number of extracted signal events at a given $m_{Z'}$, σ , A, BR, and \mathcal{L} are the signal cross-section, acceptance, signal branching ratio, and integrated luminosity, respectively; and θ_{modeling} is a nuisance parameter associated with the background modeling uncertainty. The signal acceptance is defined as the fraction of simulated events at detector level that pass the analysis selection cuts.

The spurious signal uncertainty is found to be, together with the statistical uncertainty, the leading source of uncertainty on the extracted number of signal events.

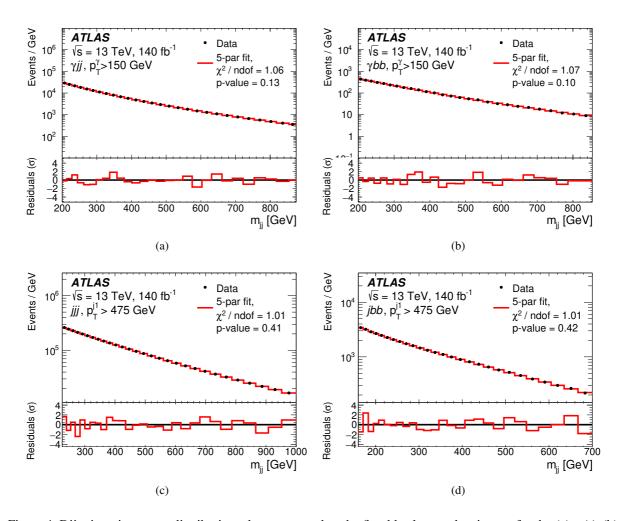


Figure 4: Dijet invariant mass distributions data compared to the fitted background estimates for the (a) γjj , (b) γbb , (c) jjj, and (d) jbb channels. Bottom panel shows the fit residuals in terms of standard deviations (σ). The distributions are shown here with the m_{jj} resolution binning.

7 Results

The dijet invariant mass distributions in data, together with the corresponding fitted background estimates, are shown in Figure 4 for all four channels. The data are well described by the 5-parameter fit function in all channels, and the global χ^2 p-value ranges from 0.10 to 0.42 for the different channels. Fit ranges are chosen such that the fits satisfy the tests explained in Section 5.2. For the photon channels, the background fit spans the m_{jj} range of 200 GeV to 875 GeV for the γjj channel and of 200 GeV to 854 GeV for the γbb channel, while for the jjj and jbb channels, the fit spans 225 GeV to 1000 GeV and 160 GeV to 700 GeV respectively.

The BumpHunter [99, 100] algorithm, as implemented in PYBumpHunter [101, 102], is used to measure the statistical significance of localized excesses of the measured data relative to the estimated background in the m_{jj} distributions, which could be due to the presence of resonant signals. This is done using mass bins with a bin width determined by the mass resolution of m_{jj} as a function of the mass, where the mass resolution is determined using a Gaussian function fit to the m_{jj} response distribution. Windows are allowed

to have a width of up to three mass-resolution bins, hence corresponding to three times the $m_{\rm jj}$ resolution, and for each scanned window BumpHunter evaluates the statistical significance of the observed difference between the data distribution and the background fit. The BumpHunter p-value is defined as the smallest observed probability for the data in a given window to deviate from the background prediction by the observed amount due to a Poissonian fluctuation of the background, using pseudoexperiments generated from the background prediction, and without considering systematic uncertainties. The most significant localized excesses identified by the BumpHunter algorithm are found at 348 GeV with a local significance of 1.9σ for the γjj channel, at 380 GeV with a local significance of 1.8σ for the γbb channel, at 404 GeV with a local significance of 1.8σ for the jbb channel.

As no significant deviation from the background expectation is observed, upper limits are set on the signal production rate as a function of the hypothesized resonance mass. The numbers of signal and background events are estimated from maximum-likelihood fits of the signal-plus-background models to the corresponding m_{jj} distributions. Systematic uncertainties described in Section 6 are included in the fits via nuisance parameters constrained by Gaussian-distributed penalty terms. Systematic uncertainties affecting the signal shape have only a small impact on the result, and the dominant uncertainties arise from the statistical and spurious signal uncertainties. The p-value is determined from a profile-likelihood-ratio-test statistic [103]. The local p-value for compatibility with the background-only hypothesis when testing a given signal hypothesis (p_0) is evaluated based on the asymptotic approximation [103]. Global significance values are computed from background-only pseudo-experiments to account for the trial factors due to scanning both the signal mass and the width hypotheses. The expected and observed 95% confidence level (CL) exclusion limits on the product of the cross-section, branching ratio, and acceptance are computed using a modified frequentist-approach [104], in an asymptotic approximation to the test-statistic distribution [103].

Figure 5 shows the 95% CL upper limits on the $\sigma \times A \times BR$ of the Z' axial-vector dark-matter mediator as a function of its mass, derived using the signal templates. Results are interpolated linearly in the log of the cross-section. Similar results are shown in Figure 6 for generic signals with Gaussian function shapes with 5%, 7%, 10%, 12%, and 15% signal widths.

Given the relation between couplings and cross-section, the limits on the Z' cross-sections can be translated to contraints on the g_q coupling, by taking into account the values of the simulated sample cross-sections and the analysis acceptance. A comparison of the limits on the g_q coupling as a function of the Z' mass for all channels is shown in Figure 7. For the flavor-inclusive channels, the cross-section combines the contributions from the signal samples where the Z' decays are restricted to light flavors with the samples where the Z' decay products decay into b-quarks. While the limits on $\sigma \times A \times BR$ are much stronger for the γbb and jbb channels than for the γjj and jjj channels, the cross-sections and branching ratios are much lower, meaning that the limits on g_q are more similar across channels. Overall, the jjj and γjj channels set the strongest limits on the g_q coupling.

The limits on the g_q coupling obtained from the γjj and jjj channels are similar in value and there is only a small overlap in the selected events across the two categories. Hence a combination of the two channels can further strengthen the constraints on g_q . No combination of the b-tagged channel is considered, as the constraints on g_q from such channels are weaker than the flavor-inclusive ones.

Events passing both the γjj and jjj selections are removed from the jjj channel, to have two statistically independent measurements. Such double-counted events amount to 0.2% of all the jjj-selected events and their impact on the jjj channel fits were found to be negligible. A combined likelihood function, which

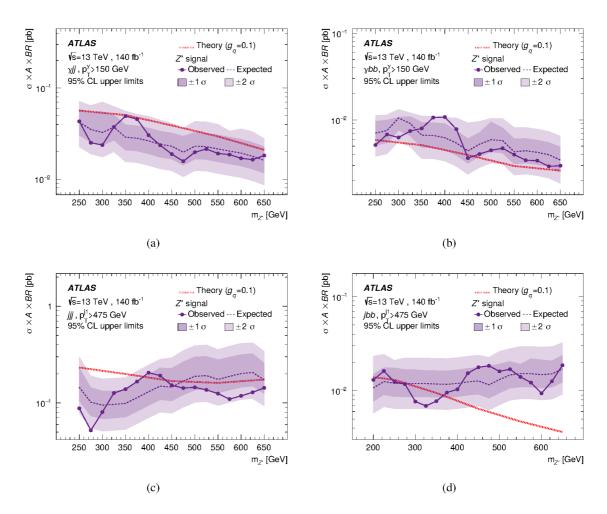


Figure 5: Expected and observed 95% CL upper limits on the product of the cross-section, acceptance and branching ratio ($\sigma \times A \times BR$) as a function of the Z' mass for the (a) γjj , (b) γbb , (c) jjj, and (d) jbb channels with the background and signal contributions modeled by a 5-parameter fit function and templates from the signal samples, respectively. Observed limits are indicated with markers and a solid line, and expected limits are indicated with a dashed line. The shaded bands around the expected limit indicates the (darker band) 1σ and (lighter band) 2σ uncertainty range. The dotted line represents the expected $\sigma \times A \times BR$ assuming $g_q = 0.1$, as obtained from the MadGraph signal samples.

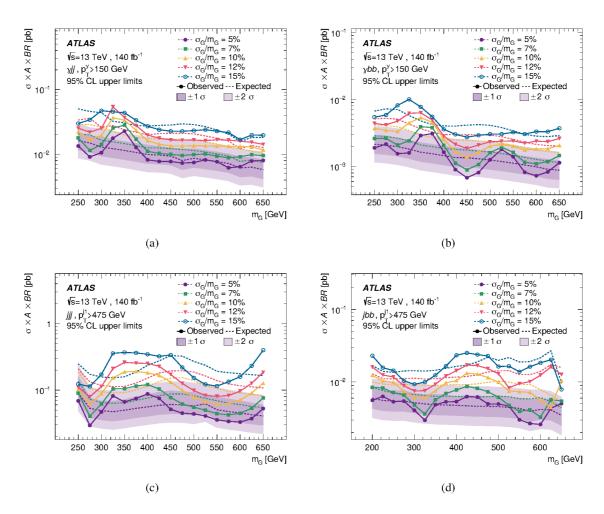


Figure 6: Expected and observed 95% CL upper limits on the product of the cross-section, acceptance and branching ratio $(\sigma \times A \times BR)$ for the (a) γjj , (b) γbb , (c) jjj, and (d) jbb channels with the background and signal contributions modeled by a 5-parameter fit function and Gaussian distributed signal templates with a mass m_G and width σ_G , respectively. Observed limits are indicated with markers and a solid line, and expected limits are indicated with a dashed line. Observed and expected limits corresponding to different choices of template widths (5%, 7%, 10%, 12% and 15%) are indicated as different sets of markers or line styles, respectively. The shaded bands around the expected limit for templates with 5% width indicate the (darker band) 1σ and (lighter band) 2σ uncertainty range.

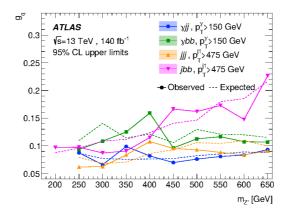


Figure 7: Expected and observed 95% CL upper limits on the g_q coupling as a function of the Z' mass on the γjj , γbb , jjj, and jbb channels with the background and signal contributions modeled by a 5-parameter fit function and templates from the signal samples, respectively. Observed limits are indicated with markers and a solid line, and expected limits are indicated with a dashed line.

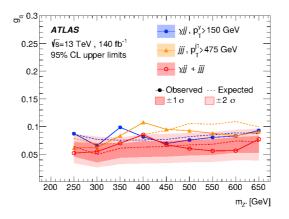


Figure 8: Expected and observed 95% CL upper limits on the g_q coupling as a function of the Z' mass on the γjj channel, the jjj channel and the combination of the two. The background and signal contributions are modeled by a 5-parameter fit function and templates from the signal samples, respectively. Observed limits are indicated with markers and a solid line, and expected limits are indicated with a dashed line. The shaded bands indicate the (darker band) 1σ and (lighter band) 2σ uncertainty range on the combined result.

is the product of the γjj and jjj likelihood functions, is used to obtain 95% CL upper limits on g_q . In the combined likelihood, the jet and luminosity systematic uncertainties are taken as correlated among the channels. Spurious signal and PDF systematic uncertainties are instead taken as uncorrelated, as the background fit functions and production mechanisms of the two processes are independent across channels. Photon systematic uncertainties only affect the γjj part of the combined likelihood. The combination improves the limits on g_q relative to the individual channels, as shown in Figure 8, with both the expected and observed limits excluding coupling values down to 0.05–0.07.

8 Conclusion

Dijet resonances with a width up to 15% of the mass, produced in association with a photon or jet, were searched for in 140 fb⁻¹ of LHC pp collisions recorded by the ATLAS experiment at $\sqrt{s} = 13$ TeV. This search expands on previous similar searches by using the full Run 2 data samples, and by including the case where the initial-state radiation is a jet. It considers the cases where no flavor requirements are placed on the resonance decay products, and the case where both of the decay products are required to be b-tagged, resulting in four channels, depending on the type of initial-state radiation and the flavor requirements. In all four channels, the observed m_{jj} distribution is well-described by a smooth functional fit without contributions from such resonances. No significant excess of events beyond the Standard Model expectation is observed, and so upper limits are set on two models: Z' axial-vector dark-matter mediators and Gaussian-shape signal contributions for resonant masses between 200 GeV and 650 GeV. Relative to already published results using a similar analysis technique, this search improves the limits on the Z'-q-q coupling g_q by up to 50%, with the most stringent limits on g_q set by the jjj channel for lower Z' candidate masses and by the γjj channel for higher Z' masses. A further combination of the jjj and γjj channels is performed, which set limits on g_q down to 0.05–0.07.

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The ATLAS Collaboration

```
G. Aad 10<sup>103</sup>, E. Aakvaag 10<sup>16</sup>, B. Abbott 10<sup>121</sup>, K. Abeling 10<sup>55</sup>, N.J. Abicht 10<sup>49</sup>, S.H. Abidi 10<sup>29</sup>,
M. Aboelela (D<sup>44</sup>, A. Aboulhorma (D<sup>35e</sup>, H. Abramowicz (D<sup>152</sup>, H. Abreu (D<sup>151</sup>, Y. Abulaiti (D<sup>118</sup>,
B.S. Acharya 69a,69b,1, A. Ackermann 63a, C. Adam Bourdarios 4, L. Adamczyk 8aa,
S.V. Addepalli <sup>©</sup><sup>26</sup>, M.J. Addison <sup>©</sup><sup>102</sup>, J. Adelman <sup>©</sup><sup>116</sup>, A. Adiguzel <sup>©</sup><sup>21c</sup>, T. Adye <sup>©</sup><sup>135</sup>, A.A. Affolder <sup>©</sup><sup>137</sup>, Y. Afik <sup>©</sup><sup>39</sup>, M.N. Agaras <sup>©</sup><sup>13</sup>, J. Agarwala <sup>©</sup><sup>73a,73b</sup>, A. Aggarwal <sup>©</sup><sup>101</sup>,
C. Agheorghiesei ©<sup>27c</sup>, A. Ahmad ©<sup>36</sup>, F. Ahmadov ©<sup>38,y</sup>, W.S. Ahmed ©<sup>105</sup>, S. Ahuja ©<sup>96</sup>, X. Ai ©<sup>62e</sup>,
G. Aielli 676a,76b, A. Aikot 6164, M. Ait Tamlihat 635e, B. Aitbenchikh 635a, I. Aizenberg 6170,
M. Akbiyik \mathbb{D}^{101}, T.P.A. Åkesson \mathbb{D}^{99}, A.V. Akimov \mathbb{D}^{37}, D. Akiyama \mathbb{D}^{169}, N.N. Akolkar \mathbb{D}^{24},
S. Aktas ©21a, K. Al Khoury ©41, G.L. Alberghi ©23b, J. Albert ©166, P. Albicocco ©53, G.L. Albouy ©60,
S. Alderweireldt <sup>52</sup>, Z.L. Alegria <sup>122</sup>, M. Aleksa <sup>36</sup>, I.N. Aleksandrov <sup>38</sup>, C. Alexa <sup>27b</sup>,
T. Alexopoulos (D10), F. Alfonsi (D23b), M. Algren (D56), M. Alhroob (D142), B. Ali (D133), H.M.J. Ali (D92),
S. Ali <sup>149</sup>, S.W. Alibocus <sup>193</sup>, M. Aliev <sup>133c</sup>, G. Alimonti <sup>171a</sup>, W. Alkakhi <sup>155</sup>, C. Allaire <sup>166</sup>, B.M.M. Allbrooke <sup>147</sup>, J.F. Allen <sup>152</sup>, C.A. Allendes Flores <sup>138f</sup>, P.P. Allport <sup>120</sup>, A. Aloisio <sup>172a,72b</sup>,
F. Alonso <sup>109</sup>, C. Alpigiani <sup>139</sup>, M. Alvarez Estevez <sup>100</sup>, A. Alvarez Fernandez <sup>101</sup>,
M. Alves Cardoso 656, M.G. Alviggi 672a,72b, M. Aly 6102, Y. Amaral Coutinho 683b, A. Ambler 6105,
C. Amelung<sup>36</sup>, M. Amerl (10)<sup>102</sup>, C.G. Ames (10)<sup>110</sup>, D. Amidei (10)<sup>107</sup>, K.J. Amirie (10)<sup>156</sup>,
S.P. Amor Dos Santos ©<sup>131a</sup>, K.R. Amos ©<sup>164</sup>, S. An<sup>84</sup>, V. Ananiev ©<sup>126</sup>, C. Anastopoulos ©<sup>140</sup>,
T. Andeen 11, J.K. Anders 136, S.Y. Andrean 147a,47b, A. Andreazza 171a,71b, S. Angelidakis 19
A. Angerami (D<sup>41</sup>,aa), A.V. Anisenkov (D<sup>37</sup>, A. Annovi (D<sup>74a</sup>, C. Antel (D<sup>56</sup>, M.T. Anthony (D<sup>140</sup>),
E. Antipov 146, M. Antonelli 53, F. Anulli 75a, M. Aoki 84, T. Aoki 154, J.A. Aparisi Pozo 164,
M.A. Aparo <sup>147</sup>, L. Aperio Bella <sup>48</sup>, C. Appelt <sup>18</sup>, A. Apyan <sup>26</sup>, S.J. Arbiol Val <sup>87</sup>,
C. Arcangeletti (D<sup>53</sup>, A.T.H. Arce (D<sup>51</sup>, E. Arena (D<sup>93</sup>, J-F. Arguin (D<sup>109</sup>, S. Argyropoulos (D<sup>54</sup>,
J.-H. Arling <sup>648</sup>, O. Arnaez <sup>64</sup>, H. Arnold <sup>6115</sup>, G. Artoni <sup>675a,75b</sup>, H. Asada <sup>6112</sup>, K. Asai <sup>6119</sup>,
S. Asai 6154, N.A. Asbah 636, K. Assamagan 629, R. Astalos 628a, K.S.V. Astrand 699, S. Atashi 6160,
R.J. Atkin (D<sup>33a</sup>, M. Atkinson<sup>163</sup>, H. Atmani<sup>35f</sup>, P.A. Atmasiddha (D<sup>129</sup>, K. Augsten (D<sup>133</sup>),
S. Auricchio (D<sup>72a,72b</sup>, A.D. Auriol (D<sup>20</sup>, V.A. Austrup (D<sup>102</sup>, G. Avolio (D<sup>36</sup>, K. Axiotis (D<sup>56</sup>),
G. Azuelos (D109,ae), D. Babal (D28b), H. Bachacou (D136), K. Bachas (D153,p), A. Bachiu (D34),
F. Backman (10<sup>47a,47b</sup>), A. Badea (10<sup>39</sup>), T.M. Baer (10<sup>107</sup>), P. Bagnaia (10<sup>75a,75b</sup>), M. Bahmani (10<sup>18</sup>),
D. Bahner (D<sup>54</sup>, K. Bai (D<sup>124</sup>, J.T. Baines (D<sup>135</sup>, L. Baines (D<sup>95</sup>, O.K. Baker (D<sup>173</sup>, E. Bakos (D<sup>15</sup>),
D. Bakshi Gupta <sup>68</sup>, V. Balakrishnan <sup>6121</sup>, R. Balasubramanian <sup>6115</sup>, E.M. Baldin <sup>637</sup>, P. Balek <sup>686</sup>,
E. Ballabene (D<sup>23b,23a</sup>, F. Balli (D<sup>136</sup>, L.M. Baltes (D<sup>63a</sup>, W.K. Balunas (D<sup>32</sup>, J. Balz (D<sup>101</sup>, E. Banas (D<sup>87</sup>,
M. Bandieramonte (D<sup>130</sup>, A. Bandyopadhyay (D<sup>24</sup>, S. Bansal (D<sup>24</sup>, L. Barak (D<sup>152</sup>, M. Barakat (D<sup>48</sup>),
E.L. Barberio (106), D. Barberis (157b,57a), M. Barbero (103), M.Z. Barel (115), K.N. Barends (133a),
T. Barillari 111, M-S. Barisits 136, T. Barklow 1144, P. Baron 1123, D.A. Baron Moreno 1102,
A. Baroncelli 62a, G. Barone 29, A.J. Barr 127, J.D. Barr 97, F. Barreiro 100,
J. Barreiro Guimarães da Costa 14a, U. Barron 152, M.G. Barros Teixeira 131a, S. Barsov 37,
F. Bartels ^{\odot}6^{3a}, R. Bartoldus ^{\odot}1^{44}, A.E. Barton ^{\odot}9^{2}, P. Bartos ^{\odot}2^{8a}, A. Basan ^{\odot}10^{1}, M. Baselga ^{\odot}4^{9},
A. Bassalat 666,b, M.J. Basso 6157a, R. Bate 6165, R.L. Bates 659, S. Batlamous 100, B. Batool 6142,
M. Battaglia 137, D. Battulga 18, M. Bauce 75a,75b, M. Bauer 36, P. Bauer 24,
L.T. Bazzano Hurrell (1030), J.B. Beacham (1051), T. Beau (10128), J.Y. Beaucamp (1091), P.H. Beauchemin (10159),
P. Bechtle (10<sup>24</sup>, H.P. Beck (10<sup>19</sup>,0), K. Becker (10<sup>168</sup>, A.J. Beddall (10<sup>82</sup>, V.A. Bednyakov (10<sup>38</sup>, C.P. Bee (10<sup>146</sup>),
L.J. Beemster \bigcirc^{15}, T.A. Beermann \bigcirc^{36}, M. Begalli \bigcirc^{83d}, M. Begel \bigcirc^{29}, A. Behera \bigcirc^{146}, J.K. Behr \bigcirc^{48},
J.F. Beirer <sup>©36</sup>, F. Beisiegel <sup>©24</sup>, M. Belfkir <sup>©117b</sup>, G. Bella <sup>©152</sup>, L. Bellagamba <sup>©23b</sup>, A. Bellerive <sup>©34</sup>,
P. Bellos 620, K. Beloborodov 637, D. Benchekroun 635a, F. Bendebba 635a, Y. Benhammou 6152,
```

```
K.C. Benkendorfer 61, L. Beresford 48, M. Beretta 53, E. Bergeaas Kuutmann 5162, N. Berger 4,
B. Bergmann 6133, J. Beringer 617a, G. Bernardi 5, C. Bernius 6144, F.U. Bernlochner 624,
F. Bernon 636,103, A. Berrocal Guardia 613, T. Berry 696, P. Berta 6134, A. Berthold 650, S. Bethke 6111,
A. Betti (D<sup>75a,75b</sup>, A.J. Bevan (D<sup>95</sup>, N.K. Bhalla (D<sup>54</sup>, M. Bhamjee (D<sup>33c</sup>, S. Bhatta (D<sup>146</sup>,
D.S. Bhattacharya 6167, P. Bhattarai 6144, K.D. Bhide 654, V.S. Bhopatkar 6122, R.M. Bianchi 6130,
G. Bianco (D<sup>23b,23a</sup>, O. Biebel (D<sup>110</sup>, R. Bielski (D<sup>124</sup>, M. Biglietti (D<sup>77a</sup>, C.S. Billingsley<sup>44</sup>, M. Bindi (D<sup>55</sup>),
A. Bingul <sup>©21b</sup>, C. Bini <sup>©75a,75b</sup>, A. Biondini <sup>©93</sup>, C.J. Birch-sykes <sup>©102</sup>, G.A. Bird <sup>©32</sup>, M. Birman <sup>©170</sup>,
M. Biros 134, S. Biryukov 147, T. Bisanz 49, E. Bisceglie 43b,43a, J.P. Biswal 135, D. Biswas 142,
I. Bloch 🕫 A. Blue 👨 U. Blumenschein 🕞 J. Blumenthal 🕫 101, V.S. Bobrovnikov 🕞 7,
M. Boehler <sup>©54</sup>, B. Boehm <sup>©167</sup>, D. Bogavac <sup>©36</sup>, A.G. Bogdanchikov <sup>©37</sup>, C. Bohm <sup>©47a</sup>,
V. Boisvert (10)96, P. Bokan (10)36, T. Bold (10)86a, M. Bomben (10)5, M. Bona (10)95, M. Boonekamp (10)136,
C.D. Booth 696, A.G. Borbély 59, I.S. Bordulev 57, H.M. Borecka-Bielska 6109, G. Borissov 692,
D. Bortoletto 127, D. Boscherini 23b, M. Bosman 13, J.D. Bossio Sola 36, K. Bouaouda 35a,
N. Bouchhar 16164, J. Boudreau 16130, E.V. Bouhova-Thacker 1692, D. Boumediene 1640,
R. Bouquet (10,57b,57a), A. Boveia (10,120), J. Boyd (10,36), D. Boye (10,29), I.R. Boyko (10,38), J. Bracinik (10,20),
N. Brahimi <sup>6</sup>, G. Brandt <sup>6</sup>, O. Brandt <sup>6</sup>, F. Braren <sup>6</sup>, B. Brau <sup>6</sup>, J.E. Brau <sup>6</sup>,
R. Brener 170, L. Brenner 115, R. Brenner 162, S. Bressler 170, D. Britton 159, D. Britzger 111,
I. Brock ^{\odot 24}, G. Brooijmans ^{\odot 41}, E. Brost ^{\odot 29}, L.M. Brown ^{\odot 166}, L.E. Bruce ^{\odot 61}, T.L. Bruckler ^{\odot 127},
P.A. Bruckman de Renstrom <sup>687</sup>, B. Brüers <sup>648</sup>, A. Bruni <sup>623b</sup>, G. Bruni <sup>623b</sup>, M. Bruschi <sup>623b</sup>,
N. Bruscino (^{15}), T. Buanes (^{16}), Q. Buat (^{139}), D. Buchin (^{111}), A.G. Buckley (^{59}), O. Bulekov (^{37}).
B.A. Bullard 144, S. Burdin 93, C.D. Burgard 49, A.M. Burger 36, B. Burghgrave 8,
O. Burlayenko <sup>654</sup>, J.T.P. Burr <sup>632</sup>, C.D. Burton <sup>611</sup>, J.C. Burzynski <sup>6143</sup>, E.L. Busch <sup>641</sup>,
V. Büscher <sup>101</sup>, P.J. Bussey <sup>159</sup>, J.M. Butler <sup>125</sup>, C.M. Buttar <sup>159</sup>, J.M. Butterworth <sup>197</sup>,
W. Buttinger 135, C.J. Buxo Vazquez 10108, A.R. Buzykaev 137, S. Cabrera Urbán 164,
L. Cadamuro 666, D. Caforio 58, H. Cai 5130, Y. Cai 144,14e, Y. Cai 14c, V.M.M. Cairo 536,
O. Cakir <sup>103a</sup>, N. Calace <sup>1036</sup>, P. Calafiura <sup>1017a</sup>, G. Calderini <sup>10128</sup>, P. Calfayan <sup>1068</sup>, G. Callea <sup>1059</sup>,
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S. Camarda 636, D. Camarero Munoz 626, P. Camarri 676a,76b, M.T. Camerlingo 72a,72b,
D. Cameron <sup>166</sup>, C. Camincher <sup>166</sup>, M. Campanelli <sup>169</sup>, A. Camplani <sup>164</sup>, V. Canale <sup>172a,72b</sup>,
A.C. Canbay (D<sup>3a</sup>, E. Canonero (D<sup>96</sup>, J. Cantero (D<sup>164</sup>, Y. Cao (D<sup>163</sup>, F. Capocasa (D<sup>26</sup>, M. Capua (D<sup>43b,43a</sup>,
A. Carbone 171a,71b, R. Cardarelli 176a, J.C.J. Cardenas 18, F. Cardillo 164, G. Carducci 1643b,43a,
T. Carli (1036), G. Carlino (1072a), J.I. Carlotto (1013), B.T. Carlson (10130,q), E.M. Carlson (10166,157a),
L. Carminati (10,71a,71b), A. Carnelli (10,136), M. Carnesale (10,75a,75b), S. Caron (10,134), E. Carquin (10,138f),
S. Carrá (D<sup>71a</sup>, G. Carratta (D<sup>23b,23a</sup>, A.M. Carroll (D<sup>124</sup>, T.M. Carter (D<sup>52</sup>, M.P. Casado (D<sup>13</sup>,i),
M. Caspar <sup>1</sup>0<sup>48</sup>, F.L. Castillo <sup>1</sup>0<sup>4</sup>, L. Castillo Garcia <sup>1</sup>13, V. Castillo Gimenez <sup>1</sup>164, N.F. Castro <sup>1</sup>131a,131e,
A. Catinaccio (D<sup>36</sup>, J.R. Catmore (D<sup>126</sup>, T. Cavaliere (D<sup>4</sup>, V. Cavaliere (D<sup>29</sup>, N. Cavalli (D<sup>23b,23a</sup>,
Y.C. Cekmecelioglu \mathbb{D}^{48}, E. Celebi \mathbb{D}^{21a}, S. Cella \mathbb{D}^{36}, F. Celli \mathbb{D}^{127}, M.S. Centonze \mathbb{D}^{70a,70b}, V. Cepaitis \mathbb{D}^{56}, K. Cerny \mathbb{D}^{123}, A.S. Cerqueira \mathbb{D}^{83a}, A. Cerri \mathbb{D}^{147}, L. Cerrito \mathbb{D}^{76a,76b}, F. Cerutti \mathbb{D}^{17a},
B. Cervato 142, A. Cervelli 23b, G. Cesarini 53, S.A. Cetin 82, D. Chakraborty 116, J. Chan 17a,
W.Y. Chan 10154, J.D. Chapman 1032, E. Chapon 10136, B. Chargeishvili 10150b, D.G. Charlton 1020,
M. Chatterjee 19, C. Chauhan 134, Y. Che 14c, S. Chekanov 6, S.V. Chekulaev 157a,
G.A. Chelkov (D<sup>38</sup>,a), A. Chen (D<sup>107</sup>, B. Chen (D<sup>152</sup>, B. Chen (D<sup>166</sup>, H. Chen (D<sup>14c</sup>, H. Chen (D<sup>29</sup>),
J. Chen 62c, J. Chen 6143, M. Chen 6127, S. Chen 6154, S.J. Chen 614c, X. Chen 662c, 136,
X. Chen (14b,ad), Y. Chen (16a, C.L. Cheng (17a), H.C. Cheng (16a, S. Cheong (17a), A. Cheplakov (17a), A.
E. Cheremushkina <sup>648</sup>, E. Cherepanova <sup>6115</sup>, R. Cherkaoui El Moursli <sup>635</sup>e, E. Cheu <sup>67</sup>, K. Cheung <sup>65</sup>,
L. Chevalier 136, V. Chiarella 53, G. Chiarelli 74a, N. Chiedde 103, G. Chiodini 70a,
A.S. Chisholm <sup>©20</sup>, A. Chitan <sup>©27b</sup>, M. Chitishvili <sup>©164</sup>, M.V. Chizhov <sup>©38</sup>, K. Choi <sup>©11</sup>, Y. Chou <sup>©139</sup>,
```

```
E.Y.S. Chow 114, K.L. Chu 170, M.C. Chu 164a, X. Chu 14a, 14e, J. Chudoba 132,
J.J. Chwastowski 687, D. Cieri 6111, K.M. Ciesla 686a, V. Cindro 694, A. Ciocio 617a, F. Cirotto 672a,72b,
Z.H. Citron 170, M. Citterio 71a, D.A. Ciubotaru<sup>27b</sup>, A. Clark 56, P.J. Clark 55, C. Clarry 156,
J.M. Clavijo Columbie <sup>1048</sup>, S.E. Clawson <sup>1048</sup>, C. Clement <sup>1047a,47b</sup>, J. Clercx <sup>1048</sup>, Y. Coadou <sup>10103</sup>,
M. Cobal 69a,69c, A. Coccaro 57b, R.F. Coelho Barrue 131a, R. Coelho Lopes De Sa 104,
S. Coelli <sup>1071a</sup>, B. Cole <sup>1041</sup>, J. Collot <sup>1060</sup>, P. Conde Muiño <sup>10131a,131g</sup>, M.P. Connell <sup>1033c</sup>,
S.H. Connell 633c, E.I. Conroy 6127, F. Conventi 672a, af, H.G. Cooke 620, A.M. Cooper-Sarkar 6127,
F.A. Corchia (D<sup>23b,23a</sup>, A. Cordeiro Oudot Choi (D<sup>128</sup>, L.D. Corpe (D<sup>40</sup>, M. Corradi (D<sup>75a,75b</sup>),
F. Corriveau 10105,w, A. Cortes-Gonzalez 1018, M.J. Costa 10164, F. Costanza 104, D. Costanzo 10140,
B.M. Cote ©120, G. Cowan ©96, K. Cranmer ©171, D. Cremonini ©23b,23a, S. Crépé-Renaudin ©60,
F. Crescioli 128, M. Cristinziani 142, M. Cristoforetti 178a, 78b, V. Croft 115, J.E. Crosby 122,
G. Crosetti 643b,43a, A. Cueto 6100, H. Cui 614a,14e, Z. Cui 67, W.R. Cunningham 659, F. Curcio 6164,
J.R. Curran 652, P. Czodrowski 636, M.M. Czurylo 636, M.J. Da Cunha Sargedas De Sousa 657b,57a,
J.V. Da Fonseca Pinto <sup>683b</sup>, C. Da Via <sup>6102</sup>, W. Dabrowski <sup>686a</sup>, T. Dado <sup>649</sup>, S. Dahbi <sup>6149</sup>,
T. Dai 6107, D. Dal Santo 619, C. Dallapiccola 6104, M. Dam 642, G. D'amen 629, V. D'Amico 6110,
J. Damp (b<sup>101</sup>, J.R. Dandoy (b<sup>34</sup>, M. Danninger (b<sup>143</sup>, V. Dao (b<sup>36</sup>, G. Darbo (b<sup>57b</sup>, S.J. Das (b<sup>29</sup>, ag),
F. Dattola <sup>648</sup>, S. D'Auria <sup>671a,71b</sup>, A. D'avanzo <sup>672a,72b</sup>, C. David <sup>633a</sup>, T. Davidek <sup>6134</sup>,
B. Davis-Purcell <sup>©34</sup>, I. Dawson <sup>©95</sup>, H.A. Day-hall <sup>©133</sup>, K. De <sup>©8</sup>, R. De Asmundis <sup>©72a</sup>,
N. De Biase 648, S. De Castro 623b,23a, N. De Groot 6114, P. de Jong 6115, H. De la Torre 6116,
A. De Maria 1014c, A. De Salvo 1075a, U. De Sanctis 1076a, 76b, F. De Santis 1070a, 70b, A. De Santo 10147,
J.B. De Vivie De Regie 60, D.V. Dedovich J. Degens 69, A.M. Deiana 64, F. Del Corso 623b,23a,
J. Del Peso 100, F. Del Rio 163a, L. Delagrange 128, F. Deliot 136, C.M. Delitzsch 149,
M. Della Pietra (D<sup>72a,72b</sup>, D. Della Volpe (D<sup>56</sup>, A. Dell'Acqua (D<sup>36</sup>, L. Dell'Asta (D<sup>71a,71b</sup>, M. Delmastro (D<sup>4</sup>,
P.A. Delsart <sup>60</sup>, S. Demers <sup>173</sup>, M. Demichev <sup>38</sup>, S.P. Denisov <sup>37</sup>, L. D'Eramo <sup>40</sup>,
D. Derendarz <sup>1087</sup>, F. Derue <sup>10128</sup>, P. Dervan <sup>1093</sup>, K. Desch <sup>1024</sup>, C. Deutsch <sup>1024</sup>, F.A. Di Bello <sup>1057b,57a</sup>,
A. Di Ciaccio (D<sup>76a,76b</sup>, L. Di Ciaccio (D<sup>4</sup>, A. Di Domenico (D<sup>75a,75b</sup>, C. Di Donato (D<sup>72a,72b</sup>,
A. Di Girolamo (D<sup>36</sup>, G. Di Gregorio (D<sup>36</sup>, A. Di Luca (D<sup>78a,78b</sup>, B. Di Micco (D<sup>77a,77b</sup>, R. Di Nardo (D<sup>77a,77b</sup>).
M. Diamantopoulou (1)34, F.A. Dias (1)115, T. Dias Do Vale (1)143, M.A. Diaz (1)138a,138b,
F.G. Diaz Capriles (D<sup>24</sup>, M. Didenko (D<sup>164</sup>, E.B. Diehl (D<sup>107</sup>, S. Díez Cornell (D<sup>48</sup>, C. Diez Pardos (D<sup>142</sup>,
C. Dimitriadi 6162,24, A. Dimitrievska 620, J. Dingfelder 24, I-M. Dinu 627b, S.J. Dittmeier 63b,
F. Dittus 636, M. Divisek 6134, F. Djama 6103, T. Djobava 6150b, C. Doglioni 6102,99, A. Dohnalova 628a,
J. Dolejsi 134, Z. Dolezal 134, K.M. Dona 39, M. Donadelli 83c, B. Dong 108, J. Donini 40,
A. D'Onofrio (D<sup>72a,72b</sup>, M. D'Onofrio (D<sup>93</sup>, J. Dopke (D<sup>135</sup>, A. Doria (D<sup>72a</sup>, N. Dos Santos Fernandes (D<sup>131a</sup>,
P. Dougan (b102), M.T. Dova (b91), A.T. Doyle (b59), M.A. Draguet (b127), E. Dreyer (b170),
I. Drivas-koulouris <sup>10</sup>, M. Drnevich <sup>118</sup>, M. Drozdova <sup>56</sup>, D. Du <sup>62a</sup>, T.A. du Pree <sup>115</sup>,
F. Dubinin 637, M. Dubovsky 628a, E. Duchovni 6170, G. Duckeck 6110, O.A. Ducu 627b, D. Duda 652,
A. Dudarev 1036, E.R. Duden 1026, M. D'uffizi 10102, L. Duflot 1066, M. Dührssen 1036, I. Duminica 1027g,
A.E. Dumitriu (D<sup>27b</sup>, M. Dunford (D<sup>63a</sup>, S. Dungs (D<sup>49</sup>, K. Dunne (D<sup>47a,47b</sup>, A. Duperrin (D<sup>103</sup>),
H. Duran Yildiz <sup>©3a</sup>, M. Düren <sup>©58</sup>, A. Durglishvili <sup>©150b</sup>, B.L. Dwyer <sup>©116</sup>, G.I. Dyckes <sup>©17a</sup>,
M. Dyndal <sup>6</sup>86a, B.S. Dziedzic <sup>6</sup>87, Z.O. Earnshaw <sup>6</sup>147, G.H. Eberwein <sup>6</sup>127, B. Eckerova <sup>6</sup>28a,
S. Eggebrecht <sup>©55</sup>, E. Egidio Purcino De Souza <sup>©128</sup>, L.F. Ehrke <sup>©56</sup>, G. Eigen <sup>©16</sup>, K. Einsweiler <sup>©17a</sup>,
T. Ekelof (162), P.A. Ekman (1999), S. El Farkh (1935b), Y. El Ghazali (1935b), H. El Jarrari (1936),
A. El Moussaouy (10) V. Ellajosyula (11) M. Ellert (11) F. Ellinghaus (11) N. Ellis (13), N. Ellis (13)
J. Elmsheuser <sup>©29</sup>, M. Elsawy <sup>©117a</sup>, M. Elsing <sup>©36</sup>, D. Emeliyanov <sup>©135</sup>, Y. Enari <sup>©154</sup>, I. Ene <sup>©17a</sup>,
S. Epari <sup>13</sup>, P.A. Erland <sup>187</sup>, M. Errenst <sup>172</sup>, M. Escalier <sup>66</sup>, C. Escobar <sup>164</sup>, E. Etzion <sup>152</sup>,
G. Evans (D<sup>131a</sup>, H. Evans (D<sup>68</sup>, L.S. Evans (D<sup>96</sup>, A. Ezhilov (D<sup>37</sup>, S. Ezzarqtouni (D<sup>35a</sup>, F. Fabbri (D<sup>23b,23a</sup>,
L. Fabbri (D<sup>23b,23a</sup>, G. Facini (D<sup>97</sup>, V. Fadeyev (D<sup>137</sup>, R.M. Fakhrutdinov (D<sup>37</sup>, D. Fakoudis (D<sup>101</sup>),
```

```
S. Falciano <sup>1075a</sup>, L.F. Falda Ulhoa Coelho <sup>1036</sup>, P.J. Falke <sup>1024</sup>, F. Fallavollita <sup>1111</sup>, J. Faltova <sup>10134</sup>,
C. Fan 6163, Y. Fan 614a, Y. Fang 614a, 14e, M. Fanti 671a, 71b, M. Faraj 669a, 69b, Z. Farazpay 698,
A. Farbin <sup>®</sup>8, A. Farilla <sup>®77a</sup>, T. Farooque <sup>®108</sup>, S.M. Farrington <sup>®52</sup>, F. Fassi <sup>®35e</sup>, D. Fassouliotis <sup>®9</sup>,
M. Faucci Giannelli (1076a,76b), W.J. Fawcett (1032), L. Fayard (1066), P. Federic (10134), P. Federicova (10132),
O.L. Fedin 637,a, M. Feickert 171, L. Feligioni 10103, D.E. Fellers 124, C. Feng 162b, M. Feng 14b,
Z. Feng 115, M.J. Fenton 160, L. Ferencz 48, R.A.M. Ferguson 92, S.I. Fernandez Luengo 138f,
P. Fernandez Martinez (13), M.J.V. Fernoux (10), J. Ferrando (19), A. Ferrari (10), P. Ferrari (11), P. Ferr
R. Ferrari <sup>1073a</sup>, D. Ferrere <sup>1056</sup>, C. Ferretti <sup>10107</sup>, F. Fiedler <sup>10101</sup>, P. Fiedler <sup>10133</sup>, A. Filipčič <sup>1094</sup>,
E.K. Filmer <sup>1</sup>, F. Filthaut <sup>114</sup>, M.C.N. Fiolhais <sup>131a,131c,c</sup>, L. Fiorini <sup>164</sup>, W.C. Fisher <sup>108</sup>,
T. Fitschen 6102, P.M. Fitzhugh 136, I. Fleck 142, P. Fleischmann 10107, T. Flick 172, M. Flores 133d,ab,
L.R. Flores Castillo 64a, L. Flores Sanz De Acedo 636, F.M. Follega 78a,78b, N. Fomin 616,
J.H. Foo 6156, A. Formica 6136, A.C. Forti 6102, E. Fortin 6136, A.W. Fortman 617a, M.G. Foti 617a,
L. Fountas (D9, j), D. Fournier (D66), H. Fox (D92), P. Francavilla (D74a,74b), S. Francescato (D61),
S. Franchellucci 656, M. Franchini 623b,23a, S. Franchino 63a, D. Francis 6, L. Franco 6114,
V. Franco Lima 636, L. Franconi 648, M. Franklin 661, G. Frattari 626, W.S. Freund 683b, Y.Y. Frid 6152,
J. Friend <sup>659</sup>, N. Fritzsche <sup>50</sup>, A. Froch <sup>54</sup>, D. Froidevaux <sup>36</sup>, J.A. Frost <sup>127</sup>, Y. Fu <sup>62a</sup>,
S. Fuenzalida Garrido (138), M. Fujimoto (103), K.Y. Fung (164), E. Furtado De Simas Filho (183),
M. Furukawa <sup>154</sup>, J. Fuster <sup>164</sup>, A. Gabrielli <sup>23b,23a</sup>, A. Gabrielli <sup>156</sup>, P. Gadow <sup>36</sup>,
G. Gagliardi (1)57b,57a, L.G. Gagnon (1)17a, S. Gaid (1)161, S. Galantzan (1)152, E.J. Gallas (1)127,
B.J. Gallop 6135, K.K. Gan 6120, S. Ganguly 6154, Y. Gao 652, F.M. Garay Walls 6138a,138b, B. Garcia 29,
C. García 6164, A. Garcia Alonso 6115, A.G. Garcia Caffaro 6173, J.E. García Navarro 6164,
M. Garcia-Sciveres <sup>17a</sup>, G.L. Gardner <sup>129</sup>, R.W. Gardner <sup>39</sup>, N. Garelli <sup>159</sup>, D. Garg <sup>80</sup>,
R.B. Garg (D144,m), J.M. Gargan<sup>52</sup>, C.A. Garner<sup>156</sup>, C.M. Garvey (D33a), P. Gaspar (D83b), V.K. Gassmann<sup>159</sup>,
G. Gaudio <sup>13</sup>, V. Gautam <sup>13</sup>, P. Gauzzi <sup>15</sup>, Sa. J. L. Gavrilenko <sup>15</sup>, A. Gavrilyuk <sup>15</sup>, C. Gay <sup>165</sup>,
G. Gaycken (D<sup>48</sup>, E.N. Gazis (D<sup>10</sup>), A.A. Geanta (D<sup>27b</sup>), C.M. Gee (D<sup>137</sup>), A. Gekow (120), C. Gemme (D<sup>57b</sup>),
M.H. Genest 60, A.D. Gentry 613, S. George 69, W.F. George 620, T. Geralis 64,
P. Gessinger-Befurt 636, M.E. Geyik 6172, M. Ghani 6168, K. Ghorbanian 695, A. Ghosal 6142,
A. Ghosh 6160, A. Ghosh 657, B. Giacobbe 6523b, S. Giagu 6575a,75b, T. Giani 65115, P. Giannetti 6574a,
A. Giannini 62a, S.M. Gibson 696, M. Gignac 6137, D.T. Gil 686b, A.K. Gilbert 686a, B.J. Gilbert 641,
D. Gillberg 634, G. Gilles 6115, L. Ginabat 6128, D.M. Gingrich 62, ae, M.P. Giordani 669a, 69c,
P.F. Giraud <sup>136</sup>, G. Giugliarelli <sup>69a,69c</sup>, D. Giugni <sup>71a</sup>, F. Giuli <sup>36</sup>, I. Gkialas <sup>9,j</sup>, L.K. Gladilin <sup>37</sup>,
C. Glasman 6, G.R. Gledhill 6, G. Glemža 6, M. Glisic 124, I. Gnesi 6, Y. Go 629,
M. Goblirsch-Kolb • 36, B. Gocke • 49, D. Godin 109, B. Gokturk • 21a, S. Goldfarb • 106, T. Golling • 56, M.G.D. Gololo • 33g, D. Golubkov • 37, J.P. Gombas • 108, A. Gomes • 131a,131b, G. Gomes Da Silva • 142,
A.J. Gomez Delegido 164, R. Gonçalo 131a,131c, L. Gonella 20, A. Gongadze 150c, F. Gonnella 20,
J.L. Gonski 1014, R.Y. González Andana 1052, S. González de la Hoz 1016, R. Gonzalez Lopez 1093,
C. Gonzalez Renteria (1017a), M.V. Gonzalez Rodrigues (1048), R. Gonzalez Suarez (10162),
S. Gonzalez-Sevilla 656, L. Goossens 636, B. Gorini 636, E. Gorini 670a,70b, A. Gorišek 694,
T.C. Gosart (D129), A.T. Goshaw (D51), M.I. Gostkin (D38), S. Goswami (D122), C.A. Gottardo (D36),
S.A. Gotz 110, M. Gouighri 35b, V. Goumarre 48, A.G. Goussiou 139, N. Govender 33c,
I. Grabowska-Bold 686a, K. Graham 634, E. Gramstad 6126, S. Grancagnolo 670a,70b, C.M. Grant 1,136,
P.M. Gravila (D<sup>27f</sup>, F.G. Gravili (D<sup>70a,70b</sup>, H.M. Gray (D<sup>17a</sup>, M. Greco (D<sup>70a,70b</sup>, C. Grefe (D<sup>24</sup>,
I.M. Gregor (D48), K.T. Greif (D160), P. Grenier (D144), S.G. Grewe 111, A.A. Grillo (D137), K. Grimm (D31),
S. Grinstein 613,8, J.-F. Grivaz 666, E. Gross 6170, J. Grosse-Knetter 655, J.C. Grundy 6127,
L. Guan 6107, C. Gubbels 6165, J.G.R. Guerrero Rojas 6164, G. Guerrieri 669a,69c, F. Guescini 6111,
R. Gugel 10101, J.A.M. Guhit 10107, A. Guida 1018, E. Guilloton 10168, S. Guindon 1036, F. Guo 1014a,14e,
J. Guo 62c, L. Guo 48, Y. Guo 107, R. Gupta 48, R. Gupta 130, S. Gurbuz 124, S.S. Gurdasani 154,
```

```
G. Gustavino 636, M. Guth 556, P. Gutierrez 5121, L.F. Gutierrez Zagazeta 5129, M. Gutsche 550,
C. Gutschow \mathbb{D}^{97}, C. Gwenlan \mathbb{D}^{127}, C.B. Gwilliam \mathbb{D}^{93}, E.S. Haaland \mathbb{D}^{126}, A. Haas \mathbb{D}^{118},
M. Habedank <sup>6</sup><sup>48</sup>, C. Haber <sup>6</sup><sup>17a</sup>, H.K. Hadavand <sup>6</sup><sup>8</sup>, A. Hadef <sup>6</sup><sup>50</sup>, S. Hadzic <sup>6</sup><sup>111</sup>, A.I. Hagan <sup>6</sup><sup>92</sup>,
J.J. Hahn (b)142, E.H. Haines (b)97, M. Haleem (b)167, J. Haley (b)122, J.J. Hall (b)140, G.D. Hallewell (b)103,
L. Halser 19, K. Hamano 166, M. Hamer 124, G.N. Hamity 152, E.J. Hampshire 196, J. Han 1626,
K. Han 62a, L. Han 14c, L. Han 62a, S. Han 617a, Y.F. Han 6156, K. Hanagaki 684, M. Hance 6137,
D.A. Hangal (D<sup>41</sup>, H. Hanif (D<sup>143</sup>, M.D. Hank (D<sup>129</sup>, J.B. Hansen (D<sup>42</sup>, P.H. Hansen (D<sup>42</sup>, K. Hara (D<sup>158</sup>),
D. Harada 656, T. Harenberg 6172, S. Harkusha 637, M.L. Harris 6104, Y.T. Harris 6127, J. Harrison 613,
N.M. Harrison (D120), P.F. Harrison 168, N.M. Hartman (D111), N.M. Hartmann (D110), Y. Hasegawa (D141),
S. Hassan 6, R. Hauser 10, C.M. Hawkes 20, R.J. Hawkings 6, Y. Hayashi 6, Y. Hayashi 6,
S. Hayashida (112), D. Hayden (108), C. Hayes (107), R.L. Hayes (115), C.P. Hays (127), J.M. Hays (1095),
H.S. Hayward (1993), F. He (1962a), M. He (1914a, 14e), Y. He (1915b), Y. He (1948), Y. He (1997), N.B. Heatley (1995),
V. Hedberg <sup>126</sup>, A.L. Heggelund <sup>126</sup>, N.D. Hehir <sup>126</sup>, C. Heidegger <sup>154</sup>, K.K. Heidegger <sup>154</sup>,
W.D. Heidorn (1081), J. Heilman (1034), S. Heim (1048), T. Heim (1017a), J.G. Heinlein (10129), J.J. Heinrich (10124),
L. Heinrich 111, ac, J. Hejbal 122, A. Held 171, S. Hellesund 16, C.M. Helling 165,
S. Hellman <sup>6</sup> <sup>47a,47b</sup>, R.C.W. Henderson <sup>92</sup>, L. Henkelmann <sup>6</sup> <sup>32</sup>, A.M. Henriques Correia <sup>36</sup>, H. Herde <sup>6</sup> <sup>99</sup>,
Y. Hernández Jiménez 6146, L.M. Herrmann 624, T. Herrmann 550, G. Herten 54, R. Hertenberger 6110,
L. Hervas (1)36, M.E. Hesping (1)101, N.P. Hessey (1)157a, M. Hidaoui (1)35b, E. Hill (1)156, S.J. Hillier (1)20,
J.R. Hinds 108, F. Hinterkeuser 124, M. Hirose 125, S. Hirose 158, D. Hirschbuehl 172,
T.G. Hitchings (102), B. Hiti (1094), J. Hobbs (1146), R. Hobincu (1027e), N. Hod (1170), M.C. Hodgkinson (10140),
B.H. Hodkinson (D<sup>127</sup>, A. Hoecker (D<sup>36</sup>, D.D. Hofer (D<sup>107</sup>, J. Hofer (D<sup>48</sup>, T. Holm (D<sup>24</sup>, M. Holzbock (D<sup>111</sup>),
L.B.A.H. Hommels (D<sup>32</sup>, B.P. Honan (D<sup>102</sup>, J. Hong (D<sup>62c</sup>, T.M. Hong (D<sup>130</sup>, B.H. Hooberman (D<sup>163</sup>,
W.H. Hopkins 6, Y. Horii 112, S. Hou 1149, A.S. Howard 194, J. Howarth 159, J. Hoya 6,
M. Hrabovsky 123, A. Hrynevich 48, T. Hryn'ova 4, P.J. Hsu 65, S.-C. Hsu 139, T. Hsu 66,
M. Hu 17a, Q. Hu 162a, S. Huang 164b, X. Huang 114a, 14e, Y. Huang 1140, Y. Huang 1101,
Y. Huang 14a, Z. Huang 1012, Z. Hubacek 133, M. Huebner 124, F. Huegging 124, T.B. Huffman 127,
C.A. Hugli \mathbb{D}^{48}, M. Huhtinen \mathbb{D}^{36}, S.K. Huiberts \mathbb{D}^{16}, R. Hulsken \mathbb{D}^{105}, N. Huseynov \mathbb{D}^{12}, J. Huston \mathbb{D}^{108},
J. Huth 661, R. Hyneman 6144, G. Iacobucci 56, G. Iakovidis 29, I. Ibragimov 142,
L. Iconomidou-Fayard 66, J.P. Iddon 536, P. Iengo 572a,72b, R. Iguchi 5154, T. Iizawa 5127,
Y. Ikegami <sup>684</sup>, N. Ilic <sup>6156</sup>, H. Imam <sup>635a</sup>, M. Ince Lezki <sup>656</sup>, T. Ingebretsen Carlson <sup>647a,47b</sup>,
G. Introzzi (10,73a,73b), M. Iodice (10,77a), V. Ippolito (10,75a,75b), R.K. Irwin (10,93), M. Ishino (10,154), W. Islam (10,171),
C. Issever (18,48), S. Istin (121a,ai), H. Ito (169), R. Iuppa (178a,78b), A. Ivina (170), J.M. Izen (1845),
P. Jain 654, K. Jakobs 554, T. Jakoubek 6170, J. Jamieson 559, K.W. Janas 686a, M. Javurkova 6104,
L. Jeanty 124, J. Jejelava 150a, P. Jenni 54, C. E. Jessiman 34, C. Jia 54, X. Jia 56, X. Jia 56,
X. Jia 614a,14e, Z. Jia 614c, C. Jiang 52, S. Jiggins 648, J. Jimenez Pena 613, S. Jin 614c, A. Jinaru 627b,
O. Jinnouchi © 155, P. Johansson © 140, K.A. Johns © 7, J.W. Johnson © 137, D.M. Jones © 147, E. Jones © 48,
P. Jones (532, R.W.L. Jones (592, T.J. Jones (593, H.L. Joos (55,36), R. Joshi (5120, J. Jovicevic (515),
X. Ju 17a, J.J. Junggeburth 17a, T. Junkermann 16a, A. Juste Rozas 17a, M.K. Juzek 18a,
S. Kabana (D<sup>138e</sup>, A. Kaczmarska (D<sup>87</sup>, M. Kado (D<sup>111</sup>, H. Kagan (D<sup>120</sup>, M. Kagan (D<sup>144</sup>, A. Kahn<sup>41</sup>,
A. Kahn <sup>129</sup>, C. Kahra <sup>101</sup>, T. Kaji <sup>154</sup>, E. Kajomovitz <sup>151</sup>, N. Kakati <sup>170</sup>, I. Kalaitzidou <sup>54</sup>,
C.W. Kalderon (D<sup>29</sup>, N.J. Kang (D<sup>137</sup>, D. Kar (D<sup>33g</sup>, K. Karava (D<sup>127</sup>, M.J. Kareem (D<sup>157b</sup>, E. Karentzos (D<sup>54</sup>,
I. Karkanias <sup>153</sup>, O. Karkout <sup>151</sup>, S.N. Karpov <sup>38</sup>, Z.M. Karpova <sup>38</sup>, V. Kartvelishvili <sup>92</sup>,
A.N. Karyukhin (1)37, E. Kasimi (1)153, J. Katzy (1)48, S. Kaur (1)34, K. Kawade (1)141, M.P. Kawale (1)121,
C. Kawamoto <sup>688</sup>, T. Kawamoto <sup>62a</sup>, E.F. Kay <sup>36</sup>, F.I. Kaya <sup>159</sup>, S. Kazakos <sup>108</sup>, V.F. Kazanin <sup>37</sup>,
Y. Ke (10)146, J.M. Keaveney (10)33a, R. Keeler (10)166, G.V. Kehris (10)61, J.S. Keller (10)34, A.S. Kelly (11)97,
J.J. Kempster (147), P.D. Kennedy (101), O. Kepka (132), B.P. Kerridge (135), S. Kersten (172),
```

```
B.P. Kerševan \bigcirc^{94}, L. Keszeghova \bigcirc^{28a}, S. Ketabchi Haghighat \bigcirc^{156}, R.A. Khan \bigcirc^{130}, A. Khanov \bigcirc^{122},
A.G. Kharlamov (^{\circ})<sup>37</sup>, T. Kharlamova (^{\circ})<sup>37</sup>, E.E. Khoda (^{\circ})<sup>139</sup>, M. Kholodenko (^{\circ})<sup>37</sup>, T.J. Khoo (^{\circ})<sup>18</sup>,
G. Khoriauli <sup>167</sup>, J. Khubua <sup>150b,*</sup>, Y.A.R. Khwaira <sup>66</sup>, B. Kibirige<sup>33g</sup>, A. Kilgallon <sup>124</sup>,
D.W. Kim (D47a,47b), Y.K. Kim (D39), N. Kimura (D97), M.K. Kingston (D55), A. Kirchhoff (D55), C. Kirfel (D24),
F. Kirfel 624, J. Kirk 6135, A.E. Kiryunin 6111, C. Kitsaki 610, O. Kivernyk 624, M. Klassen 6159,
C. Klein 634, L. Klein 6167, M.H. Klein 644, S.B. Klein 656, U. Klein 693, P. Klimek 636,
A. Klimentov <sup>©29</sup>, T. Klioutchnikova <sup>©36</sup>, P. Kluit <sup>©115</sup>, S. Kluth <sup>©111</sup>, E. Kneringer <sup>©79</sup>,
T.M. Knight 6156, A. Knue 649, R. Kobayashi 88, D. Kobylianskii 6170, S.F. Koch 6127,
M. Kocian 6144, P. Kodyš 6134, D.M. Koeck 6124, P.T. Koenig 624, T. Koffas 634, O. Kolay 650,
I. Koletsou <sup>64</sup>, T. Komarek <sup>6123</sup>, K. Köneke <sup>654</sup>, A.X.Y. Kong <sup>61</sup>, T. Kono <sup>6119</sup>, N. Konstantinidis <sup>697</sup>,
P. Kontaxakis <sup>656</sup>, B. Konya <sup>699</sup>, R. Kopeliansky <sup>641</sup>, S. Koperny <sup>686a</sup>, K. Korcyl <sup>687</sup>, K. Kordas <sup>6153,e</sup>,
A. Korn (D<sup>97</sup>, S. Korn (D<sup>55</sup>, I. Korolkov (D<sup>13</sup>, N. Korotkova (D<sup>37</sup>, B. Kortman (D<sup>115</sup>, O. Kortner (D<sup>111</sup>),
S. Kortner 111, W.H. Kostecka 116, V.V. Kostyukhin 142, A. Kotsokechagia 136, A. Kotwal 51,
A. Koulouris <sup>©36</sup>, A. Kourkoumeli-Charalampidi <sup>©73a,73b</sup>, C. Kourkoumelis <sup>©9</sup>, E. Kourlitis <sup>©111,ac</sup>,
O. Kovanda 124, R. Kowalewski 166, W. Kozanecki 136, A.S. Kozhin 137, V.A. Kramarenko 137,
G. Kramberger <sup>194</sup>, P. Kramer <sup>101</sup>, M.W. Krasny <sup>128</sup>, A. Krasznahorkay <sup>136</sup>, J.W. Kraus <sup>172</sup>,
J.A. Kremer <sup>1048</sup>, T. Kresse <sup>1050</sup>, J. Kretzschmar <sup>1093</sup>, K. Kreul <sup>1018</sup>, P. Krieger <sup>10156</sup>,
S. Krishnamurthy 10<sup>104</sup>, M. Krivos 10<sup>134</sup>, K. Krizka 10<sup>20</sup>, K. Kroeninger 10<sup>49</sup>, H. Kroha 10<sup>111</sup>, J. Kroll 10<sup>132</sup>,
J. Kroll <sup>129</sup>, K.S. Krowpman <sup>108</sup>, U. Kruchonak <sup>38</sup>, H. Krüger <sup>24</sup>, N. Krumnack M.C. Kruse <sup>51</sup>,
O. Kuchinskaia ^{37}, S. Kuday ^{3a}, S. Kuehn ^{36}, R. Kuesters ^{54}, T. Kuhl ^{48}, V. Kukhtin ^{38},
Y. Kulchitsky (1037,a), S. Kuleshov (1018d,138b), M. Kumar (1033g), N. Kumari (1048), P. Kumari (10157b)
A. Kupco (D<sup>132</sup>, T. Kupfer<sup>49</sup>, A. Kupich (D<sup>37</sup>, O. Kuprash (D<sup>54</sup>, H. Kurashige (D<sup>85</sup>, L.L. Kurchaninov (D<sup>157a</sup>),
O. Kurdysh 666, Y.A. Kurochkin 637, A. Kurova 637, M. Kuze 6155, A.K. Kvam 6104, J. Kvita 6123,
T. Kwan <sup>10105</sup>, N.G. Kyriacou <sup>10107</sup>, L.A.O. Laatu <sup>10103</sup>, C. Lacasta <sup>10164</sup>, F. Lacava <sup>1075a,75b</sup>,
H. Lacker <sup>18</sup>, D. Lacour <sup>128</sup>, N.N. Lad <sup>97</sup>, E. Ladygin <sup>38</sup>, A. Lafarge <sup>40</sup>, B. Laforge <sup>128</sup>,
T. Lagouri (10173), F.Z. Lahbabi (1035a), S. Lai (1055), I.K. Lakomiec (1086a), J.E. Lambert (10166), S. Lammers (1068),
W. Lampl <sup>107</sup>, C. Lampoudis <sup>101</sup>, A.N. Lancaster <sup>101</sup>, E. Lançon <sup>102</sup>,
U. Landgraf <sup>654</sup>, M.P.J. Landon <sup>695</sup>, V.S. Lang <sup>654</sup>, O.K.B. Langrekken <sup>6126</sup>, A.J. Lankford <sup>6160</sup>,
F. Lanni 636, K. Lantzsch 624, A. Lanza 673a, A. Lapertosa 657b,57a, J.F. Laporte 6136, T. Lari 671a,
F. Lasagni Manghi (D<sup>23b</sup>, M. Lassnig (D<sup>36</sup>, V. Latonova (D<sup>132</sup>, A. Laudrain (D<sup>101</sup>, A. Laurier (D<sup>151</sup>),
S.D. Lawlor 6140, Z. Lawrence 6102, R. Lazaridou 168, M. Lazzaroni 671a,71b, B. Le<sup>102</sup>,
E.M. Le Boulicaut <sup>151</sup>, L.T. Le Pottier <sup>17a</sup>, B. Leban <sup>23b,23a</sup>, A. Lebedev <sup>81</sup>, M. LeBlanc <sup>102</sup>,
F. Ledroit-Guillon 60, A.C.A. Lee 149, S. Lee 149, S. Lee 149, T.F. Lee 193, L.L. Leeuw 133c,
H.P. Lefebvre 696, M. Lefebvre 6166, C. Leggett 617a, G. Lehmann Miotto 636, M. Leigh 656,
W.A. Leight © 104, W. Leinonen © 114, A. Leisos © 153,r, M.A.L. Leite © 83c, C.E. Leitgeb © 18,
R. Leitner 134, K.J.C. Leney 44, T. Lenz 24, S. Leone 74a, C. Leonidopoulos 52, A. Leopold 145,
C. Leroy 109, R. Les 108, C.G. Lester 132, M. Levchenko 137, J. Levêque 14, L.J. Levinson 170,
G. Levrini (D<sup>23b,23a</sup>, M.P. Lewicki (D<sup>87</sup>, C. Lewis (D<sup>139</sup>, D.J. Lewis (D<sup>4</sup>, A. Li (D<sup>5</sup>, B. Li (D<sup>62b</sup>, C. Li<sup>62a</sup>,
C-Q. Li 111, H. Li 162a, H. Li 162b, H. Li 1614c, H. Li 1614b, H. Li 162b, J. Li 162c, K. Li 16139,
L. Li 662c, M. Li 14a, 14e, O.Y. Li 662a, S. Li 14a, 14e, S. Li 662d, 62c, d. T. Li 55, X. Li 6105, Z. Li 6127,
Z. Li 6154, Z. Li 614a,14e, S. Liang 14a,14e, Z. Liang 14a, M. Liberatore 136, B. Liberti 1576a, K. Lie 1564c,
J. Lieber Marin <sup>683</sup>e, H. Lien <sup>668</sup>, K. Lin <sup>6108</sup>, R.E. Lindley <sup>67</sup>, J.H. Lindon <sup>62</sup>, E. Lipeles <sup>6129</sup>,
A. Lipniacka 616, A. Lister 6165, J.D. Little 64, B. Liu 614a, B.X. Liu 6143, D. Liu 662d,62c,
E.H.L. Liu 620, J.B. Liu 62a, J.K.K. Liu 62a, K. Liu 62d, K. Liu 62d, M. Liu 62a, M.Y. Liu 62a,
P. Liu 624, Q. Liu 662d,139,62c, X. Liu 662a, X. Liu 662b, Y. Liu 662b, Y. Liu 662b, Y. Liu 662a,
J. Llorente Merino (143), S.L. Lloyd (195), E.M. Lobodzinska (148), P. Loch (17), T. Lohse (181),
K. Lohwasser 140, E. Loiacono 48, M. Lokajicek 132,*, J.D. Lomas 20, J.D. Long 163,
```

```
I. Longarini <sup>160</sup>, L. Longo <sup>170a,70b</sup>, R. Longo <sup>163</sup>, I. Lopez Paz <sup>167</sup>, A. Lopez Solis <sup>48</sup>,
N. Lorenzo Martinez <sup>104</sup>, A.M. Lory <sup>110</sup>, G. Löschcke Centeno <sup>147</sup>, O. Loseva <sup>137</sup>, X. Lou <sup>147</sup>a,47b,
X. Lou 14a,14e, A. Lounis 66, P.A. Love 92, G. Lu 14a,14e, M. Lu 66, S. Lu 129, Y.J. Lu 65,
H.J. Lubatti 139, C. Luci 75a,75b, F.L. Lucio Alves 14c, F. Luehring 68, I. Luise 146,
O. Lukianchuk 666, O. Lundberg 6145, B. Lund-Jensen 145, N.A. Luongo 66, M.S. Lutz 636,
A.B. Lux (D<sup>25</sup>, D. Lynn (D<sup>29</sup>, R. Lysak (D<sup>132</sup>, E. Lytken (D<sup>99</sup>, V. Lyubushkin (D<sup>38</sup>, T. Lyubushkina (D<sup>38</sup>,
M.M. Lyukova 10146, M.Firdaus M. Soberi 1052, H. Ma 1029, K. Ma 1062a, L.L. Ma 1062b, W. Ma 1062a,
Y. Ma 6122, D.M. Mac Donell 6166, G. Maccarrone 653, J.C. MacDonald 6101,
P.C. Machado De Abreu Farias <sup>63e</sup>, R. Madar <sup>40</sup>, T. Madula <sup>697</sup>, J. Maeda <sup>685</sup>, T. Maeno <sup>629</sup>,
H. Maguire (D<sup>140</sup>, V. Maiboroda (D<sup>136</sup>, A. Maio (D<sup>131a,131b,131d</sup>, K. Maj (D<sup>86a</sup>, O. Majersky (D<sup>48</sup>),
S. Majewski <sup>124</sup>, N. Makovec <sup>66</sup>, V. Maksimovic <sup>15</sup>, B. Malaescu <sup>128</sup>, Pa. Malecki <sup>87</sup>,
V.P. Maleev (D<sup>37</sup>, F. Malek (D<sup>60,n</sup>, M. Mali (D<sup>94</sup>, D. Malito (D<sup>96</sup>, U. Mallik (D<sup>80</sup>, S. Maltezos<sup>10</sup>,
S. Malyukov<sup>38</sup>, J. Mamuzic <sup>13</sup>, G. Mancini <sup>53</sup>, M.N. Mancini <sup>26</sup>, G. Manco <sup>73a,73b</sup>,
J.P. Mandalia <sup>695</sup>, I. Mandić <sup>694</sup>, L. Manhaes de Andrade Filho <sup>683a</sup>, I.M. Maniatis <sup>6170</sup>,
J. Manjarres Ramos (590, D.C. Mankad (5170), A. Mann (5110), S. Manzoni (536), L. Mao (562c),
X. Mapekula (1033c), A. Marantis (10153,r), G. Marchiori (105), M. Marcisovsky (10132), C. Marcon (1071a),
M. Marinescu (D<sup>20</sup>, S. Marium (D<sup>48</sup>, M. Marjanovic (D<sup>121</sup>, A. Markhoos (D<sup>54</sup>, M. Markovitch (D<sup>66</sup>),
E.J. Marshall <sup>1092</sup>, Z. Marshall <sup>17a</sup>, S. Marti-Garcia <sup>164</sup>, T.A. Martin <sup>168</sup>, V.J. Martin <sup>152</sup>,
B. Martin dit Latour <sup>16</sup>, L. Martinelli <sup>575a,75b</sup>, M. Martinez <sup>13,8</sup>, P. Martinez Agullo <sup>164</sup>,
V.I. Martinez Outschoorn (D<sup>104</sup>, P. Martinez Suarez (D<sup>13</sup>, S. Martin-Haugh (D<sup>135</sup>, G. Martinovicova (D<sup>134</sup>,
V.S. Martoiu (D<sup>27b</sup>, A.C. Martyniuk (D<sup>97</sup>, A. Marzin (D<sup>36</sup>, D. Mascione (D<sup>78a,78b</sup>, L. Masetti (D<sup>101</sup>),
T. Mashimo 154, J. Masik 1010, A.L. Maslennikov 137, P. Massarotti 172a,72b, P. Mastrandrea 174a,74b,
A. Mastroberardino (10<sup>43b,43a</sup>, T. Masubuchi (10<sup>154</sup>, T. Mathisen (10<sup>162</sup>, J. Matousek (10<sup>134</sup>, N. Matsuzawa (15<sup>4</sup>, J. Maurer (10<sup>27b</sup>, A.J. Maury (10<sup>66</sup>, B. Maček (10<sup>94</sup>, D.A. Maximov (10<sup>37</sup>, A.E. May (10<sup>102</sup>, R. Mazini (10<sup>149</sup>),
I. Maznas (116), M. Mazza (1010), S.M. Mazza (10137), E. Mazzeo (1071a,71b), C. Mc Ginn (1029),
J.P. Mc Gowan <sup>166</sup>, S.P. Mc Kee <sup>10107</sup>, C.C. McCracken <sup>165</sup>, E.F. McDonald <sup>10106</sup>,
A.E. McDougall (D115), J.A. Mcfayden (D147), R.P. McGovern (D129), G. Mchedlidze (D150b),
R.P. Mckenzie 633g, T.C. Mclachlan 648, D.J. Mclaughlin 697, S.J. McMahon 6135,
C.M. Mcpartland <sup>693</sup>, R.A. McPherson <sup>6166,w</sup>, S. Mehlhase <sup>6110</sup>, A. Mehta <sup>693</sup>, D. Melini <sup>6164</sup>,
B.R. Mellado Garcia (D<sup>33g</sup>, A.H. Melo (D<sup>55</sup>, F. Meloni (D<sup>48</sup>, A.M. Mendes Jacques Da Costa (D<sup>102</sup>)
H.Y. Meng 6156, L. Meng 692, S. Menke 6111, M. Mentink 636, E. Meoni 643b,43a, G. Mercado 6116,
C. Merlassino (69a,69c), L. Merola (672a,72b), C. Meroni (671a,71b), J. Metcalfe (66), A.S. Mete (66),
C. Meyer (58, J-P. Meyer (5136, R.P. Middleton (5135, L. Mijović (552, G. Mikenberg (5170,
M. Mikestikova (132), M. Mikuž (194), H. Mildner (1011), A. Milic (136), D.W. Miller (1039), E.H. Miller (10144),
L.S. Miller (1034), A. Milov (10170), D.A. Milstead (47a,47b), T. Min (14c), A.A. Minaenko (1037),
I.A. Minashvili 6150b, L. Mince 659, A.I. Mincer 6118, B. Mindur 686a, M. Mineev 638, Y. Mino 688,
L.M. Mir <sup>13</sup>, M. Miralles Lopez <sup>59</sup>, M. Mironova <sup>17a</sup>, A. Mishima <sup>154</sup>, M.C. Missio <sup>114</sup>,
A. Mitra 10168, V.A. Mitsou 10164, Y. Mitsumori 10112, O. Miu 10156, P.S. Miyagawa 1095,
T. Mkrtchyan 663a, M. Mlinarevic 697, T. Mlinarevic 697, M. Mlynarikova 636, S. Mobius 619,
P. Mogg 110, M.H. Mohamed Farook 113, A.F. Mohammed 14a,14e, S. Mohapatra 141,
G. Mokgatitswane <sup>©33g</sup>, L. Moleri <sup>©170</sup>, B. Mondal <sup>©142</sup>, S. Mondal <sup>©133</sup>, K. Mönig <sup>©48</sup>,
E. Monnier (103), L. Monsonis Romero (164), J. Montejo Berlingen (1013), M. Montella (10120),
F. Montereali (1077a,77b), F. Monticelli (1091), S. Monzani (1069a,69c), N. Morange (1066),
A.L. Moreira De Carvalho <sup>648</sup>, M. Moreno Llácer <sup>6164</sup>, C. Moreno Martinez <sup>656</sup>, P. Morettini <sup>657b</sup>,
S. Morgenstern (b<sup>36</sup>, M. Morii (b<sup>61</sup>, M. Morinaga (b<sup>154</sup>, F. Morodei (b<sup>75a,75b</sup>, L. Morvaj (b<sup>36</sup>,
P. Moschovakos <sup>©36</sup>, B. Moser <sup>©36</sup>, M. Mosidze <sup>©150b</sup>, T. Moskalets <sup>©54</sup>, P. Moskvitina <sup>©114</sup>,
J. Moss 631,k, P. Moszkowicz 686a, A. Moussa 635d, E.J.W. Moyse 6104, O. Mtintsilana 633g,
```

```
S. Muanza 10103, J. Mueller 10130, D. Muenstermann 1092, R. Müller 1019, G.A. Mullier 10162,
A.J. Mullin<sup>32</sup>, J.J. Mullin<sup>129</sup>, D.P. Mungo 156, D. Munoz Perez 164, F.J. Munoz Sanchez 1010,
M. Murin (102), W.J. Murray (108,135), M. Muškinja (1094), C. Mwewa (1029), A.G. Myagkov (1037,a),
A.J. Myers ^{\bullet 8}, G. Myers ^{\bullet 107}, M. Myska ^{\bullet 133}, B.P. Nachman ^{\bullet 17a}, O. Nackenhorst ^{\bullet 49}, K. Nagai ^{\bullet 127},
K. Nagano 684, J.L. Nagle 629, E. Nagy 6103, A.M. Nairz 636, Y. Nakahama 684, K. Nakamura 684,
K. Nakkalil <sup>©5</sup>, H. Nanjo <sup>©125</sup>, R. Narayan <sup>©44</sup>, E.A. Narayanan <sup>©113</sup>, I. Naryshkin <sup>©37</sup>, M. Naseri <sup>©34</sup>,
S. Nasri © 117b, C. Nass © 24, G. Navarro © 22a, J. Navarro-Gonzalez © 164, R. Nayak © 152, A. Nayaz © 18,
P.Y. Nechaeva 6<sup>37</sup>, S. Nechaeva 6<sup>23b,23a</sup>, F. Nechansky 6<sup>48</sup>, L. Nedic 6<sup>127</sup>, T.J. Neep 6<sup>20</sup>,
A. Negri (10<sup>73a,73b</sup>, M. Negrini (10<sup>23b</sup>, C. Nellist (10<sup>115</sup>, C. Nelson (10<sup>105</sup>, K. Nelson (10<sup>107</sup>, S. Nemecek (10<sup>132</sup>),
M. Nessi (^{\circ}36,h, M.S. Neubauer (^{\circ}163, F. Neuhaus (^{\circ}101, J. Neundorf (^{\circ}48, R. Newhouse (^{\circ}165,
P.R. Newman (D<sup>20</sup>, C.W. Ng (D<sup>130</sup>, Y.W.Y. Ng (D<sup>48</sup>, B. Ngair (D<sup>117a</sup>, H.D.N. Nguyen (D<sup>109</sup>),
R.B. Nickerson <sup>127</sup>, R. Nicolaidou <sup>136</sup>, J. Nielsen <sup>137</sup>, M. Niemeyer <sup>55</sup>, J. Niermann <sup>55</sup>,
N. Nikiforou <sup>©36</sup>, V. Nikolaenko <sup>©37,a</sup>, I. Nikolic-Audit <sup>©128</sup>, K. Nikolopoulos <sup>©20</sup>, P. Nilsson <sup>©29</sup>,
I. Ninca <sup>©48</sup>, H.R. Nindhito <sup>©56</sup>, G. Ninio <sup>©152</sup>, A. Nisati <sup>©75a</sup>, N. Nishu <sup>©2</sup>, R. Nisius <sup>©111</sup>,
J-E. Nitschke <sup>650</sup>, E.K. Nkadimeng <sup>633g</sup>, T. Nobe <sup>6154</sup>, D.L. Noel <sup>632</sup>, T. Nommensen <sup>6148</sup>,
M.B. Norfolk (10) 140, R.R.B. Norisam (10) 97, B.J. Norman (10) 4, M. Noury (10) 35a, J. Novak (10) 4, T. Novak (10) 48,
L. Novotny <sup>133</sup>, R. Novotny <sup>133</sup>, L. Nozka <sup>123</sup>, K. Ntekas <sup>160</sup>, N.M.J. Nunes De Moura Junior <sup>83b</sup>,
J. Ocariz (128), A. Ochi (185), I. Ochoa (131a), S. Oerdek (148,t), J.T. Offermann (139), A. Ogrodnik (134),
A. Oh 6102, C.C. Ohm 6145, H. Oide 684, R. Oishi 6154, M.L. Ojeda 648, Y. Okumura 6154,
L.F. Oleiro Seabra 6131a, S.A. Olivares Pino 6138d, G. Oliveira Correa 613, D. Oliveira Damazio 629,
D. Oliveira Goncalves <sup>683a</sup>, J.L. Oliver <sup>6160</sup>, Ö.O. Öncel <sup>654</sup>, A.P. O'Neill <sup>619</sup>, A. Onofre <sup>6131a,131e</sup>,
P.U.E. Onyisi 611, M.J. Oreglia 639, G.E. Orellana 691, D. Orestano 677a,77b, N. Orlando 613,
R.S. Orr 6156, V. O'Shea 659, L.M. Osojnak 6129, R. Ospanov 662a, G. Otero y Garzon 630,
H. Otono (1089), P.S. Ott (1063a), G.J. Ottino (1017a), M. Ouchrif (1035d), F. Ould-Saada (10126),
T. Ovsiannikova (D<sup>139</sup>, M. Owen (D<sup>59</sup>, R.E. Owen (D<sup>135</sup>, K.Y. Oyulmaz (D<sup>21a</sup>, V.E. Ozcan (D<sup>21a</sup>,
F. Ozturk <sup>687</sup>, N. Ozturk <sup>68</sup>, S. Ozturk <sup>682</sup>, H.A. Pacey <sup>6127</sup>, A. Pacheco Pages <sup>613</sup>,
C. Padilla Aranda 13, G. Padovano 75a,75b, S. Pagan Griso 17a, G. Palacino 68, A. Palazzo 70a,70b,
J. Pampel 624, J. Pan 6173, T. Pan 64a, D.K. Panchal 611, C.E. Pandini 6115, J.G. Panduro Vazquez 696,
H.D. Pandya 1, H. Pang 14b, P. Pani 48, G. Panizzo 69a,69c, L. Panwar 128, L. Paolozzi 56,
S. Parajuli ^{\circ} 163, A. Paramonov ^{\circ} 6, C. Paraskevopoulos ^{\circ} 53, D. Paredes Hernandez ^{\circ} 64b,
A. Pareti (D<sup>73a,73b</sup>, K.R. Park (D<sup>41</sup>, T.H. Park (D<sup>156</sup>, M.A. Parker (D<sup>32</sup>, F. Parodi (D<sup>57b,57a</sup>, E.W. Parrish (D<sup>116</sup>,
V.A. Parrish 652, J.A. Parsons 641, U. Parzefall 654, B. Pascual Dias 6109, L. Pascual Dominguez 6152,
E. Pasqualucci <sup>1075a</sup>, S. Passaggio <sup>1057b</sup>, F. Pastore <sup>1096</sup>, P. Patel <sup>1087</sup>, U.M. Patel <sup>1051</sup>, J.R. Pater <sup>10102</sup>,
T. Pauly 636, C.I. Pazos 6159, J. Pearkes 6144, M. Pedersen 6126, R. Pedro 6131a, S.V. Peleganchuk 637,
O. Penc (D<sup>36</sup>, E.A. Pender (D<sup>52</sup>, G.D. Penn (D<sup>173</sup>, K.E. Penski (D<sup>110</sup>, M. Penzin (D<sup>37</sup>, B.S. Peralva (D<sup>83d</sup>,
A.P. Pereira Peixoto 139, L. Pereira Sanchez 144, D.V. Perepelitsa 29, ag, E. Perez Codina 157a,
M. Perganti 1010, H. Pernegger 1036, O. Perrin 1040, K. Peters 1048, R.F.Y. Peters 10102, B.A. Petersen 1036,
T.C. Petersen <sup>642</sup>, E. Petit <sup>6103</sup>, V. Petousis <sup>6133</sup>, C. Petridou <sup>6153</sup>, T. Petru <sup>6134</sup>, A. Petrukhin <sup>6142</sup>,
M. Pettee 17a, N.E. Pettersson 36, A. Petukhov 37, K. Petukhova 134, R. Pezoa 186,
L. Pezzotti 636, G. Pezzullo 6173, T.M. Pham 6171, T. Pham 6106, P.W. Phillips 6135, G. Piacquadio 6146,
E. Pianori 617a, F. Piazza 6124, R. Piegaia 630, D. Pietreanu 627b, A.D. Pilkington 6102,
M. Pinamonti 69a,69c, J.L. Pinfold 2, B.C. Pinheiro Pereira 131a, A.E. Pinto Pinoargote 101,136,
L. Pintucci 69a,69c, K.M. Piper 147, A. Pirttikoski 56, D.A. Pizzi 34, L. Pizzimento 64b,
A. Pizzini 115, M.-A. Pleier 29, V. Plesanovs 4, V. Pleskot 134, E. Plotnikova 8, G. Poddar 55,
R. Poettgen (D<sup>99</sup>, L. Poggioli (D<sup>128</sup>, I. Pokharel (D<sup>55</sup>, S. Polacek (D<sup>134</sup>, G. Polesello (D<sup>73a</sup>, A. Poley (D<sup>143,157a</sup>),
A. Polini ©<sup>23b</sup>, C.S. Pollard ©<sup>168</sup>, Z.B. Pollock ©<sup>120</sup>, E. Pompa Pacchi ©<sup>75a,75b</sup>, D. Ponomarenko ©<sup>114</sup>,
L. Pontecorvo 636, S. Popa 627a, G.A. Popeneciu 627d, A. Poreba 636, D.M. Portillo Quintero 6157a,
```

```
S. Pospisil 133, M.A. Postill 140, P. Postolache 127c, K. Potamianos 168, P.A. Potepa 186a,
I.N. Potrap 638, C.J. Potter 532, H. Potti 61, J. Poveda 6164, M.E. Pozo Astigarraga 536,
A. Prades Ibanez <sup>164</sup>, J. Pretel <sup>54</sup>, D. Price <sup>102</sup>, M. Primavera <sup>70a</sup>, M.A. Principe Martin <sup>100</sup>,
R. Privara 123, T. Procter 59, M.L. Proffitt 139, N. Proklova 129, K. Prokofiev 64c, G. Proto 111,
J. Proudfoot 6, M. Przybycien 86a, W.W. Przygoda 68b, A. Psallidas 646, J.E. Puddefoot 140,
D. Pudzha 637, D. Pyatiizbyantseva 637, J. Qian 6107, D. Qichen 6102, Y. Qin 613, T. Qiu 652,
A. Quadt <sup>655</sup>, M. Queitsch-Maitland <sup>6102</sup>, G. Quetant <sup>656</sup>, R.P. Quinn <sup>6165</sup>, G. Rabanal Bolanos <sup>661</sup>,
D. Rafanoharana <sup>654</sup>, F. Ragusa <sup>671a,71b</sup>, J.L. Rainbolt <sup>639</sup>, J.A. Raine <sup>656</sup>, S. Rajagopalan <sup>629</sup>,
E. Ramakoti ©<sup>37</sup>, I.A. Ramirez-Berend ©<sup>34</sup>, K. Ran ©<sup>48,14e</sup>, N.P. Rapheeha ©<sup>33g</sup>, H. Rasheed ©<sup>27b</sup>,
V. Raskina © 128, D.F. Rassloff © 63a, A. Rastogi © 17a, S. Rave © 101, B. Ravina © 55, I. Ravinovich © 170,
M. Raymond 636, A.L. Read 6126, N.P. Readioff 6140, D.M. Rebuzzi 673a,73b, G. Redlinger 629,
A.S. Reed ©111, K. Reeves ©26, J.A. Reidelsturz ©172, D. Reikher ©152, A. Rej ©49, C. Rembser ©36,
M. Renda (10<sup>27b</sup>, M.B. Rendel 111, F. Renner (10<sup>48</sup>, A.G. Rennie (10<sup>160</sup>, A.L. Rescia (10<sup>48</sup>, S. Resconi (10<sup>71a</sup>),
M. Ressegotti 5<sup>57b,57a</sup>, S. Rettie 5<sup>36</sup>, J.G. Reyes Rivera 5<sup>108</sup>, E. Reynolds 5<sup>17a</sup>, O.L. Rezanova 5<sup>37</sup>,
P. Reznicek 134, H. Riani 35d, N. Ribaric 92, E. Ricci 78a,78b, R. Richter 111, S. Richter 47a,47b
E. Richter-Was 686, M. Ridel 128, S. Ridouani 35d, P. Rieck 118, P. Riedler 36, E.M. Riefel 47a,47b,
J.O. Rieger 115, M. Rijssenbeek 146, M. Rimoldi 36, L. Rinaldi 23b,23a, T.T. Rinn 29,
M.P. Rinnagel 10110, G. Ripellino 10162, I. Riu 1013, J.C. Rivera Vergara 10166, F. Rizatdinova 10122,
E. Rizvi <sup>195</sup>, B.R. Roberts <sup>17a</sup>, S.H. Robertson <sup>105</sup>, w, D. Robinson <sup>32</sup>, C.M. Robles Gajardo <sup>138</sup>f,
M. Robles Manzano (D<sup>101</sup>, A. Robson (D<sup>59</sup>), A. Rocchi (D<sup>76a,76b</sup>, C. Roda (D<sup>74a,74b</sup>, S. Rodriguez Bosca (D<sup>36</sup>),
Y. Rodriguez Garcia (D<sup>22a</sup>, A. Rodriguez Rodriguez (D<sup>54</sup>, A.M. Rodríguez Vera (D<sup>116</sup>, S. Roe<sup>36</sup>),
J.T. Roemer (160), A.R. Roepe-Gier (137), J. Roggel (172), O. Røhne (126), R.A. Rojas (104),
C.P.A. Roland <sup>128</sup>, J. Roloff <sup>29</sup>, A. Romaniouk <sup>37</sup>, E. Romano <sup>73a,73b</sup>, M. Romano <sup>23b</sup>,
A.C. Romero Hernandez 16163, N. Rompotis 193, L. Roos 128, S. Rosati 175a, B.J. Rosser 139,
E. Rossi (127), E. Rossi (172a,72b), L.P. Rossi (161), L. Rossini (154), R. Rosten (120), M. Rotaru (127b),
B. Rottler <sup>654</sup>, C. Rougier <sup>690</sup>, D. Rousseau <sup>666</sup>, D. Rousso <sup>648</sup>, A. Roy <sup>6163</sup>, S. Roy-Garand <sup>6156</sup>,
A. Rozanov 10103, Z.M.A. Rozario 1059, Y. Rozen 10151, A. Rubio Jimenez 10164, A.J. Ruby 1093,
V.H. Ruelas Rivera <sup>18</sup>, T.A. Ruggeri <sup>1</sup>, A. Ruggiero <sup>127</sup>, A. Ruiz-Martinez <sup>164</sup>, A. Rummler <sup>36</sup>,
Z. Rurikova <sup>©54</sup>, N.A. Rusakovich <sup>©38</sup>, H.L. Russell <sup>©166</sup>, G. Russo <sup>©75a,75b</sup>, J.P. Rutherfoord <sup>©7</sup>,
S. Rutherford Colmenares <sup>32</sup>, K. Rybacki<sup>92</sup>, M. Rybar <sup>134</sup>, E.B. Rye <sup>126</sup>, A. Ryzhov <sup>44</sup>,
J.A. Sabater Iglesias <sup>656</sup>, P. Sabatini <sup>6164</sup>, H.F-W. Sadrozinski <sup>6137</sup>, F. Safai Tehrani <sup>675a</sup>,
B. Safarzadeh Samani (135, S. Saha (1)1, M. Sahinsoy (111)1, A. Saibel (164, M. Saimpert (136)1, M. Sahinsoy (111)1, A. Saibel (164)1, M. Saimpert (1136)1, M. Saibel (114)1, 
M. Saito 154, T. Saito 154, A. Sala 71a,71b, D. Salamani 36, A. Salnikov 144, J. Salt 164, A. Salvador Salas 152, D. Salvatore 43b,43a, F. Salvatore 147, A. Salzburger 36, D. Sammel 54,
E. Sampson 692, D. Sampsonidis 6153,e, D. Sampsonidou 6124, J. Sánchez 6164,
V. Sanchez Sebastian 6164, H. Sandaker 6126, C.O. Sander 648, J.A. Sandesara 6104, M. Sandhoff 6172,
C. Sandoval (D<sup>22b</sup>, D.P.C. Sankey (D<sup>135</sup>, T. Sano (D<sup>88</sup>, A. Sansoni (D<sup>53</sup>, L. Santi (D<sup>75a,75b</sup>, C. Santoni (D<sup>40</sup>),
H. Santos © 131a,131b, A. Santra © 170, E. Sanzani © 23b,23a, K.A. Saoucha © 161, J.G. Saraiva © 131a,131d,
J. Sardain <sup>107</sup>, O. Sasaki <sup>1084</sup>, K. Sato <sup>10158</sup>, C. Sauer <sup>1058</sup>, F. Sauerburger <sup>1054</sup>, E. Sauvan <sup>104</sup>,
P. Savard 156,ae, R. Sawada 154, C. Sawyer 155, L. Sawyer 158, I. Sayago Galvan 164, C. Sbarra 152,
A. Sbrizzi (D<sup>23b,23a</sup>, T. Scanlon (D<sup>97</sup>, J. Schaarschmidt (D<sup>139</sup>, U. Schäfer (D<sup>101</sup>, A.C. Schaffer (D<sup>66,44</sup>,
D. Schaile 110, R.D. Schamberger 146, C. Scharf 18, M.M. Schefer 19, V.A. Schegelsky 37,
D. Scheirich 134, F. Schenck 18, M. Schernau 160, C. Scheulen 55, C. Schiavi 57b,57a,
M. Schioppa (D43b,43a), B. Schlag (D144,m), K.E. Schleicher (D54), S. Schlenker (D36), J. Schmeing (D172),
M.A. Schmidt ^{172}, K. Schmieden ^{101}, C. Schmitt ^{101}, N. Schmitt ^{101}, S. Schmitt ^{104},
L. Schoeffel <sup>136</sup>, A. Schoening <sup>63b</sup>, P.G. Scholer <sup>34</sup>, E. Schopf <sup>127</sup>, M. Schott <sup>101</sup>,
J. Schovancova 636, S. Schramm 656, T. Schroer 656, H-C. Schultz-Coulon 663a, M. Schumacher 654,
```

```
B.A. Schumm 137, Ph. Schune 136, A.J. Schuy 139, H.R. Schwartz 137, A. Schwartzman 144,
T.A. Schwarz 10<sup>107</sup>, Ph. Schwemling 13<sup>108</sup>, R. Schwienhorst 10<sup>108</sup>, A. Sciandra 10<sup>29</sup>, G. Sciolla 10<sup>26</sup>,
F. Scuri ©<sup>74a</sup>, C.D. Sebastiani ©<sup>93</sup>, K. Sedlaczek ©<sup>116</sup>, P. Seema ©<sup>18</sup>, S.C. Seidel ©<sup>113</sup>, A. Seiden ©<sup>137</sup>,
B.D. Seidlitz <sup>1</sup>, C. Seitz <sup>48</sup>, J.M. Seixas <sup>83b</sup>, G. Sekhniaidze <sup>72a</sup>, L. Selem <sup>60</sup>,
N. Semprini-Cesari (D<sup>23b,23a</sup>, D. Sengupta (D<sup>56</sup>, V. Senthilkumar (D<sup>164</sup>, L. Serin (D<sup>66</sup>, L. Serkin (D<sup>69a,69b</sup>),
M. Sessa © 76a,76b, H. Severini © 121, F. Sforza © 57b,57a, A. Sfyrla © 56, Q. Sha © 14a, E. Shabalina © 55,
A.H. Shah ©<sup>32</sup>, R. Shaheen ©<sup>145</sup>, J.D. Shahinian ©<sup>129</sup>, D. Shaked Renous ©<sup>170</sup>, L.Y. Shan ©<sup>14a</sup>,
M. Shapiro \bigcirc^{17a}, A. Sharma \bigcirc^{36}, A.S. Sharma \bigcirc^{165}, P. Sharma \bigcirc^{80}, P.B. Shatalov \bigcirc^{37}, K. Shaw \bigcirc^{147},
S.M. Shaw 102, A. Shcherbakova 137, Q. Shen 162c,5, D.J. Sheppard 143, P. Sherwood 197, L. Shi 197,
X. Shi 614a, C.O. Shimmin 6173, J.D. Shinner 696, I.P.J. Shipsey 6127, S. Shirabe 689,
M. Shiyakova <sup>©38,u</sup>, J. Shlomi <sup>©170</sup>, M.J. Shochet <sup>©39</sup>, J. Shojaii <sup>©106</sup>, D.R. Shope <sup>©126</sup>,
B. Shrestha (D121), S. Shrestha (D120,ah), E.M. Shrif (D33g), M.J. Shroff (D166), P. Sicho (D132),
A.M. Sickles 163, E. Sideras Haddad 33g, A.C. Sidley 115, A. Sidoti 23b, F. Siegert 50,
Dj. Sijacki <sup>15</sup>, F. Sili <sup>91</sup>, J.M. Silva <sup>52</sup>, M.V. Silva Oliveira <sup>29</sup>, S.B. Silverstein <sup>47a</sup>, S. Simion<sup>66</sup>,
R. Simoniello (1036), E.L. Simpson (1010), H. Simpson (10147), L.R. Simpson (10107), N.D. Simpson (10107), N.D
S. Simsek <sup>682</sup>, S. Sindhu <sup>655</sup>, P. Sinervo <sup>6156</sup>, S. Singh <sup>6156</sup>, S. Sinha <sup>648</sup>, S. Sinha <sup>6102</sup>,
M. Sioli (10<sup>23b,23a</sup>, I. Siral (10<sup>36</sup>, E. Sitnikova (10<sup>48</sup>, J. Sjölin (10<sup>47a,47b</sup>, A. Skaf (10<sup>55</sup>), E. Skorda (10<sup>20</sup>),
P. Skubic ©121, M. Slawinska ©87, V. Smakhtin<sup>170</sup>, B.H. Smart ©135, S.Yu. Smirnov ©37, Y. Smirnov ©37,
L.N. Smirnova (D<sup>37</sup>,a), O. Smirnova (D<sup>99</sup>), A.C. Smith (D<sup>41</sup>), D.R. Smith (D<sup>30</sup>), E.A. Smith (D<sup>39</sup>),
H.A. Smith 127, J.L. Smith 144, M. Smizanska 192, K. Smolek 133, A.A. Snesarev 137,
S.R. Snider (D156), H.L. Snoek (D115), S. Snyder (D29), R. Sobie (D166, w), A. Soffer (D152),
C.A. Solans Sanchez <sup>163</sup>, E.Yu. Soldatov <sup>163</sup>, U. Soldevila <sup>164</sup>, A.A. Solodkov <sup>163</sup>, S. Solomon <sup>164</sup>,
A. Soloshenko <sup>638</sup>, K. Solovieva <sup>54</sup>, O.V. Solovyanov <sup>640</sup>, P. Sommer <sup>636</sup>, A. Sonay <sup>613</sup>,
W.Y. Song (157b), A. Sopczak (133), A.L. Sopio (197), F. Sopkova (128b), J.D. Sorenson (113),
I.R. Sotarriva Alvarez (D155), V. Sothilingam<sup>63a</sup>, O.J. Soto Sandoval (D138c,138b), S. Sottocornola (D68),
R. Soualah <sup>161</sup>, Z. Soumaimi <sup>535e</sup>, D. South <sup>48</sup>, N. Soybelman <sup>170</sup>, S. Spagnolo <sup>70a,70b</sup>,
M. Spalla (111), D. Sperlich (154), G. Spigo (154), S. Spinali (159), D.P. Spiteri (159), M. Spousta (1514),
E.J. Staats <sup>134</sup>, R. Stamen <sup>163a</sup>, A. Stampekis <sup>120</sup>, M. Standke <sup>124</sup>, E. Stanecka <sup>187</sup>,
W. Stanek-Maslouska <sup>648</sup>, M.V. Stange <sup>50</sup>, B. Stanislaus <sup>617a</sup>, M.M. Stanitzki <sup>648</sup>, B. Stapf <sup>648</sup>,
E.A. Starchenko \mathbb{D}^{37}, G.H. Stark \mathbb{D}^{137}, J. Stark \mathbb{D}^{90}, P. Staroba \mathbb{D}^{132}, P. Starovoitov \mathbb{D}^{63a}, S. Stärz \mathbb{D}^{105},
R. Staszewski 687, G. Stavropoulos 646, J. Steentoft 6162, P. Steinberg 629, B. Stelzer 6143,157a,
H.J. Stelzer (D<sup>130</sup>, O. Stelzer-Chilton (D<sup>157a</sup>, H. Stenzel (D<sup>58</sup>, T.J. Stevenson (D<sup>147</sup>, G.A. Stewart (D<sup>36</sup>,
J.R. Stewart <sup>122</sup>, M.C. Stockton <sup>36</sup>, G. Stoicea <sup>27b</sup>, M. Stolarski <sup>131a</sup>, S. Stonjek <sup>111</sup>,
A. Straessner 650, J. Strandberg 6145, S. Strandberg 647a,47b, M. Stratmann 6172, M. Strauss 6121,
T. Strebler 103, P. Strizenec 28b, R. Ströhmer 167, D.M. Strom 124, R. Stroynowski 44,
A. Strubig (D47a,47b), S.A. Stucci (D29), B. Stugu (D16), J. Stupak (D121), N.A. Styles (D48), D. Su (D144),
S. Su 62a, W. Su 62d, X. Su 62a, D. Suchy 62a, K. Sugizaki 6154, V.V. Sulin 637, M.J. Sullivan 693,
D.M.S. Sultan © 127, L. Sultanaliyeva © 37, S. Sultansoy © 3b, T. Sumida © 88, S. Sun © 107, S. Sun © 171,
O. Sunneborn Gudnadottir 6162, N. Sur 6103, M.R. Sutton 6147, H. Suzuki 6158, M. Svatos 6132,
M. Swiatlowski (157a), T. Swirski (167), I. Sykora (128a), M. Sykora (134), T. Sykora (134), D. Ta (1610),
K. Tackmann 648,t, A. Taffard 6160, R. Tafirout 6157a, J.S. Tafoya Vargas 666, Y. Takubo 684,
M. Talby 103, A.A. Talyshev 137, K.C. Tam 164b, N.M. Tamir 152, A. Tanaka 154, J. Tanaka 154,
R. Tanaka 666, M. Tanasini 657b,57a, Z. Tao 6165, S. Tapia Araya 6138f, S. Tapprogge 6101,
A. Tarek Abouelfadl Mohamed <sup>108</sup>, S. Tarem <sup>151</sup>, K. Tariq <sup>14a</sup>, G. Tarna <sup>127b</sup>, G.F. Tartarelli <sup>171a</sup>,
M.J. Tartarin (50), P. Tas (5134), M. Tasevsky (5132), E. Tassi (543b,43a), A.C. Tate (5163), G. Tateno (5154),
Y. Tayalati 635e, V., G.N. Taylor 6106, W. Taylor 6157b, A.S. Tee 6171, R. Teixeira De Lima 6144,
P. Teixeira-Dias 696, J.J. Teoh 6156, K. Terashi 6154, J. Terron 6100, S. Terzo 613, M. Testa 653,
```

```
R.J. Teuscher (156,w), A. Thaler (179, O. Theiner (156, N. Themistokleous (152, T. Theveneaux-Pelzer (150, N. Themistokleous (152, T. Theveneaux-Pelzer (152, T. Thev
O. Thielmann \bigcirc^{172}, D.W. Thomas \bigcirc^{96}, J.P. Thomas \bigcirc^{20}, E.A. Thompson \bigcirc^{17a}, P.D. Thompson \bigcirc^{20},
E. Thomson 6 129, R.E. Thornberry 644, Y. Tian 655, V. Tikhomirov 637,a, Yu.A. Tikhonov 637,
S. Timoshenko<sup>37</sup>, D. Timoshyn <sup>134</sup>, E.X.L. Ting <sup>1</sup>, P. Tipton <sup>173</sup>, S.H. Tlou <sup>33g</sup>, K. Todome <sup>155</sup>,
S. Todorova-Nova 134, S. Todt<sup>50</sup>, L. Toffolin 69a,69c, M. Togawa 84, J. Tojo 89, S. Tokár 28a,
K. Tokushuku <sup>©84</sup>, O. Toldaiev <sup>©68</sup>, R. Tombs <sup>©32</sup>, M. Tomoto <sup>©84,112</sup>, L. Tompkins <sup>©144,m</sup>,
K.W. Topolnicki 686b, E. Torrence 6124, H. Torres 690, E. Torró Pastor 6164, M. Toscani 630,
C. Tosciri (539, M. Tost (511, D.R. Tovey (5140, A. Traeet 16, I.S. Trandafir (527b, T. Trefzger (5167,
A. Tricoli <sup>©29</sup>, I.M. Trigger <sup>©157a</sup>, S. Trincaz-Duvoid <sup>©128</sup>, D.A. Trischuk <sup>©26</sup>, B. Trocmé <sup>©60</sup>,
L. Truong 633c, M. Trzebinski 687, A. Trzupek 687, F. Tsai 6146, M. Tsai 6107, A. Tsiamis 6153,e,
P.V. Tsiareshka<sup>37</sup>, S. Tsigaridas <sup>157a</sup>, A. Tsirigotis <sup>153</sup>, V. Tsiskaridze <sup>156</sup>, E.G. Tskhadadze <sup>150a</sup>,
M. Tsopoulou 6153, Y. Tsujikawa 688, I.I. Tsukerman 637, V. Tsulaia 617a, S. Tsuno 684, K. Tsuri 6119,
D. Tsybychev 146, Y. Tu 64b, A. Tudorache 27b, V. Tudorache 27b, A.N. Tuna 61,
S. Turchikhin <sup>©</sup><sup>57b,57a</sup>, I. Turk Cakir <sup>©</sup><sup>3a</sup>, R. Turra <sup>©</sup><sup>71a</sup>, T. Turtuvshin <sup>©</sup><sup>38,x</sup>, P.M. Tuts <sup>©</sup><sup>41</sup>,
S. Tzamarias (D153,e), E. Tzovara (D101), F. Ukegawa (D158), P.A. Ulloa Poblete (D138c,138b), E.N. Umaka (D29),
G. Unal 636, A. Undrus 629, G. Unel 6160, J. Urban 628b, P. Urquijo 6106, P. Urrejola 6138a, G. Usai 68,
R. Ushioda (155), M. Usman (179), Z. Uysal (182), V. Vacek (173), B. Vachon (175), T. Vafeiadis (174),
A. Vaitkus ©97, C. Valderanis ©110, E. Valdes Santurio ©47a,47b, M. Valente ©157a, S. Valentinetti ©23b,23a,
A. Valero 164, E. Valiente Moreno 164, A. Vallier 190, J.A. Valls Ferrer 164, D.R. Van Arneman 1515,
T.R. Van Daalen 139, A. Van Der Graaf 49, P. Van Gemmeren 6, M. Van Rijnbach 126,
S. Van Stroud 697, I. Van Vulpen 6115, P. Vana 6134, M. Vanadia 76a, W. Vandelli 636,
E.R. Vandewall <sup>122</sup>, D. Vannicola <sup>152</sup>, L. Vannoli <sup>53</sup>, R. Vari <sup>75a</sup>, E.W. Varnes <sup>75a</sup>, C. Varni <sup>17b</sup>,
T. Varol 149, D. Varouchas 66, L. Varriale 164, K.E. Varvell 148, M.E. Vasile 1276, L. Vaslin 44,
G.A. Vasquez 6166, A. Vasyukov 638, R. Vavricka 101, F. Vazeille 640, T. Vazquez Schroeder 636,
J. Veatch (D31, V. Vecchio (D102, M.J. Veen (D104, I. Veliscek (D29, L.M. Veloce (D156, F. Veloso (D131a,131c,
S. Veneziano (D<sup>75a</sup>, A. Ventura (D<sup>70a,70b</sup>, S. Ventura Gonzalez (D<sup>136</sup>, A. Verbytskyi (D<sup>111</sup>),
M. Verducci (D<sup>74a,74b</sup>, C. Vergis (D<sup>95</sup>, M. Verissimo De Araujo (D<sup>83b</sup>, W. Verkerke (D<sup>115</sup>),
J.C. Vermeulen 115, C. Vernieri 144, M. Vessella 1014, M.C. Vetterli 143,ae, A. Vgenopoulos 153,e,
N. Viaux Maira 138f, T. Vickey 140, O.E. Vickey Boeriu 140, G.H.A. Viehhauser 127, L. Vigani 163b,
M. Villa (D<sup>23b,23a</sup>, M. Villaplana Perez (D<sup>164</sup>, E.M. Villhauer<sup>52</sup>, E. Vilucchi (D<sup>53</sup>, M.G. Vincter (D<sup>34</sup>,
G.S. Virdee ©<sup>20</sup>, A. Visibile<sup>115</sup>, C. Vittori ©<sup>36</sup>, I. Vivarelli ©<sup>23b,23a</sup>, E. Voevodina ©<sup>111</sup>, F. Vogel ©<sup>110</sup>,
J.C. Voigt <sup>650</sup>, P. Vokac <sup>6133</sup>, Yu. Volkotrub <sup>686b</sup>, J. Von Ahnen <sup>648</sup>, E. Von Toerne <sup>624</sup>,
B. Vormwald <sup>136</sup>, V. Vorobel <sup>134</sup>, K. Vorobev <sup>37</sup>, M. Vos <sup>164</sup>, K. Voss <sup>142</sup>, M. Vozak <sup>115</sup>,
L. Vozdecky 121, N. Vranjes 15, M. Vranjes Milosavljevic 15, M. Vreeswijk 115, N.K. Vu 62d,62c,
R. Vuillermet • 36, O. Vujinovic • 101, I. Vukotic • 39, S. Wada • 158, C. Wagner • 1. J.M. Wagner • 17a, W. Wagner • 172, S. Wahdan • 172, H. Wahlberg • 1, M. Wakida • 12, J. Walder • 135, R. Walker • 110,
W. Walkowiak (10) 142, A. Wall (10) 129, E.J. Wallin (10) 7, T. Wamorkar (10) 6, A.Z. Wang (10) 137, C. Wang (10) 101,
C. Wang 11, H. Wang 17a, J. Wang 164c, R.-J. Wang 101, R. Wang 161, R. Wang 166,
S.M. Wang 6149, S. Wang 62b, S. Wang 614a, T. Wang 62a, W.T. Wang 680, W. Wang 614a,
X. Wang 614c, X. Wang 6163, X. Wang 662c, Y. Wang 662d, Y. Wang 614c, Z. Wang 6107,
Z. Wang 62d,51,62c, Z. Wang 6107, A. Warburton 6105, R.J. Ward 620, N. Warrack 659,
S. Waterhouse 6, A.T. Watson 6, M.F. Watson 6, M.F. Watson 6, E. Watton 5, G. Watts 6, G. Watts 6, Watts 6, Watton 6, Watts 6, Watton 7, Watson 6, M.F. Watson 6, Watton 6, Watton 6, Watton 7, Watt
B.M. Waugh (D<sup>97</sup>, J.M. Webb (D<sup>54</sup>, C. Weber (D<sup>29</sup>, H.A. Weber (D<sup>18</sup>, M.S. Weber (D<sup>19</sup>, S.M. Weber (D<sup>63a</sup>,
C. Wei 62a, Y. Wei 54, A.R. Weidberg 127, E.J. Weik 118, J. Weingarten 49, M. Weirich 1101,
C. Weiser <sup>654</sup>, C.J. Wells <sup>648</sup>, T. Wenaus <sup>629</sup>, B. Wendland <sup>649</sup>, T. Wengler <sup>636</sup>, N.S. Wenke <sup>111</sup>,
N. Wermes \bigcirc^{24}, M. Wessels \bigcirc^{63a}, A.M. Wharton \bigcirc^{92}, A.S. White \bigcirc^{61}, A. White \bigcirc^{8}, M.J. White \bigcirc^{1},
D. Whiteson 160, L. Wickremasinghe 125, W. Wiedenmann 171, M. Wielers 135,
```

```
C. Wiglesworth 642, D.J. Wilbern 121, H.G. Wilkens 636, J.J.H. Wilkinson 632, D.M. Williams 641,
H.H. Williams <sup>129</sup>, S. Williams <sup>1032</sup>, S. Willocq <sup>104</sup>, B.J. Wilson <sup>102</sup>, P.J. Windischhofer <sup>1039</sup>,
F.I. Winkel 630, F. Winklmeier 6124, B.T. Winter 654, J.K. Winter 6102, M. Wittgen 144, M. Wobisch 698,
T. Wojtkowski<sup>60</sup>, Z. Wolffs <sup>115</sup>, J. Wollrath <sup>160</sup>, M.W. Wolter <sup>87</sup>, H. Wolters <sup>131a,131c</sup>, M.C. Wong <sup>137</sup>,
E.L. Woodward 641, S.D. Worm 648, B.K. Wosiek 687, K.W. Woźniak 687, S. Wozniewski 655,
K. Wraight <sup>659</sup>, C. Wu <sup>620</sup>, M. Wu <sup>614d</sup>, M. Wu <sup>6114</sup>, S.L. Wu <sup>6171</sup>, X. Wu <sup>656</sup>, Y. Wu <sup>662a</sup>, Z. Wu <sup>64</sup>,
J. Wuerzinger (D<sup>111,ac</sup>, T.R. Wyatt (D<sup>102</sup>, B.M. Wynne (D<sup>52</sup>, S. Xella (D<sup>42</sup>, L. Xia (D<sup>14c</sup>, M. Xia (D<sup>14b</sup>),
J. Xiang 64c, M. Xie 62a, X. Xie 62a, S. Xin 614a, 14e, A. Xiong 6124, J. Xiong 617a, D. Xu 614a,
H. Xu 662a, L. Xu 662a, R. Xu 6129, T. Xu 6107, Y. Xu 614b, Z. Xu 652, Z. Xu 14c, B. Yabsley 6148,
S. Yacoob 633a, Y. Yamaguchi 6155, E. Yamashita 6154, H. Yamauchi 6158, T. Yamazaki 617a,
Y. Yamazaki <sup>625</sup>, J. Yan <sup>559</sup>, Z. Yan <sup>6104</sup>, H.J. Yang <sup>62c,62d</sup>, H.T. Yang <sup>62a</sup>, S. Yang <sup>62a</sup>.
T. Yang 64c, X. Yang 36, X. Yang 14a, Y. Yang 44, Y. Yang 62a, Z. Yang 62a, W-M. Yao 17a,
H. Ye 14c, H. Ye 155, J. Ye 154a, S. Ye 1529, X. Ye 162a, Y. Yeh 157, I. Yeletskikh 1538, B.K. Yeo 151b,
M.R. Yexley \mathbb{D}^{97}, T.P. Yildirim \mathbb{D}^{127}, P. Yin \mathbb{D}^{41}, K. Yorita \mathbb{D}^{169}, S. Younas \mathbb{D}^{27b}, C.J.S. Young \mathbb{D}^{36},
C. Young 6144, C. Yu 614a,14e, Y. Yu 62a, M. Yuan 6107, R. Yuan 62d,62c, L. Yue 697,
M. Zaazoua 662a, B. Zabinski 87, E. Zaid52, Z.K. Zak 87, T. Zakareishvili 164, N. Zakharchuk 634,
S. Zambito 656, J.A. Zamora Saa 6138d,138b, J. Zang 6154, D. Zanzi 654, O. Zaplatilek 6133,
C. Zeitnitz 172, H. Zeng 14a, J.C. Zeng 163, D.T. Zenger Jr 26, O. Zenin 37, T. Ženiš 28a,
S. Zenz (5)5, S. Zerradi (5)35a, D. Zerwas (5)66, M. Zhai (5)14a,14e, D.F. Zhang (5)140, J. Zhang (5)62b,
J. Zhang 66, K. Zhang 14a,14e, L. Zhang 62a, L. Zhang 14c, P. Zhang 14a,14e, R. Zhang 171,
S. Zhang 10<sup>107</sup>, S. Zhang 10<sup>44</sup>, T. Zhang 10<sup>154</sup>, X. Zhang 10<sup>62c</sup>, X. Zhang 10<sup>62c</sup>, Y. Zhang 10<sup>62c</sup>,
Y. Zhang <sup>1097</sup>, Y. Zhang <sup>1014c</sup>, Z. Zhang <sup>1017a</sup>, Z. Zhang <sup>1066</sup>, H. Zhao <sup>10139</sup>, T. Zhao <sup>1062b</sup>, Y. Zhao <sup>10137</sup>,
Z. Zhao (62a, Z. Zhao (62a, A. Zhemchugov (638, J. Zheng (614c, K. Zheng (6163, X. Zheng (62a,
Z. Zheng (D<sup>144</sup>, D. Zhong (D<sup>163</sup>, B. Zhou (D<sup>107</sup>, H. Zhou (D<sup>7</sup>, N. Zhou (D<sup>62c</sup>, Y. Zhou (D<sup>14c</sup>, Y. Zhou<sup>7</sup>,
C.G. Zhu 62b, J. Zhu 6107, X. Zhu 62d, Y. Zhu 62c, Y. Zhu 62a, X. Zhuang 614a, K. Zhukov 637,
N.I. Zimine <sup>©38</sup>, J. Zinsser <sup>©63b</sup>, M. Ziolkowski <sup>©142</sup>, L. Živković <sup>©15</sup>, A. Zoccoli <sup>©23b,23a</sup>, K. Zoch <sup>©61</sup>.
T.G. Zorbas 140, O. Zormpa 46, W. Zou 41, L. Zwalinski 36.
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- 6611GL 1 11 1 1/4 D 1 G 1 GNDG/DIODS 01405 G
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