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Evidence for tWZ production in proton-proton collisions at $\sqrt{s} = 13\,\text{TeV}$ in multilepton final states

The CMS Collaboration*

Abstract

The first evidence for the standard model production of a top quark in association with a W boson and a Z boson is reported. The measurement is performed in multilepton final states, where the Z boson is reconstructed via its decays to electron or muon pairs and the W boson decays either to leptons or hadrons. The analysed data were recorded by the CMS experiment at the CERN LHC in 2016–2018 in proton-proton collisions at $\sqrt{s} = 13\,\text{TeV}$, and correspond to an integrated luminosity of $138\,\text{fb}^{-1}$. The measured cross section is $354 \pm 54\,\text{(stat)} \pm 95\,\text{(syst)}\,\text{fb}$, and corresponds to a statistical significance of 3.4 standard deviations.

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1 Introduction

Measurements of rare processes are important at the CERN LHC, as these processes are tests of the validity of the standard model (SM) of particle physics, particularly in unexplored regions of the kinematic phase space. The electroweak production of the top quark in association with a W boson and a Z boson, i.e., the tWZ process, has unique features that make it suitable to probe several interactions in the SM [1, 2]. Its production cross section is calculated to be 136^{+9}_{-8} fb [3, 4] at next-to-leading order (NLO) in quantum chromodynamics (QCD). The main background process to tWZ is the production of a top quark pair in association with a Z boson ($\mathrm{t\bar{t}Z}$), which has a cross section of 859^{+76}_{-84} fb at NLO (QCD and electroweak) and next-to-next-to-leading logarithmic accuracy [5]. Some diagrams of this process exhibit interference with the tWZ process at NLO in QCD. The production of the two processes differs in principle in the number of resonant top quarks, which results in an additional jet originating from a b quark (b quark jet) in the final state of $\mathrm{t\bar{t}Z}$, thus posing experimental challenges for the tWZ measurement.

This Letter reports the first experimental evidence for the tWZ process. The measurement is performed using events with multiple leptons (electrons or muons), where a higher signal purity is expected compared to final states with hadronically decaying W or Z bosons. The top quark decays almost exclusively to a b quark and a W boson. In the multilepton final state, the Z boson and at least one of the two W bosons decay to leptons, yielding a final state with three or four leptons. Figure 1 (left) shows a Feynman diagram for the tWZ production at NLO accuracy in QCD. In Fig. 1 (right), a Feynman diagram of the most challenging background is shown, yielding the same final state as that of the signal.

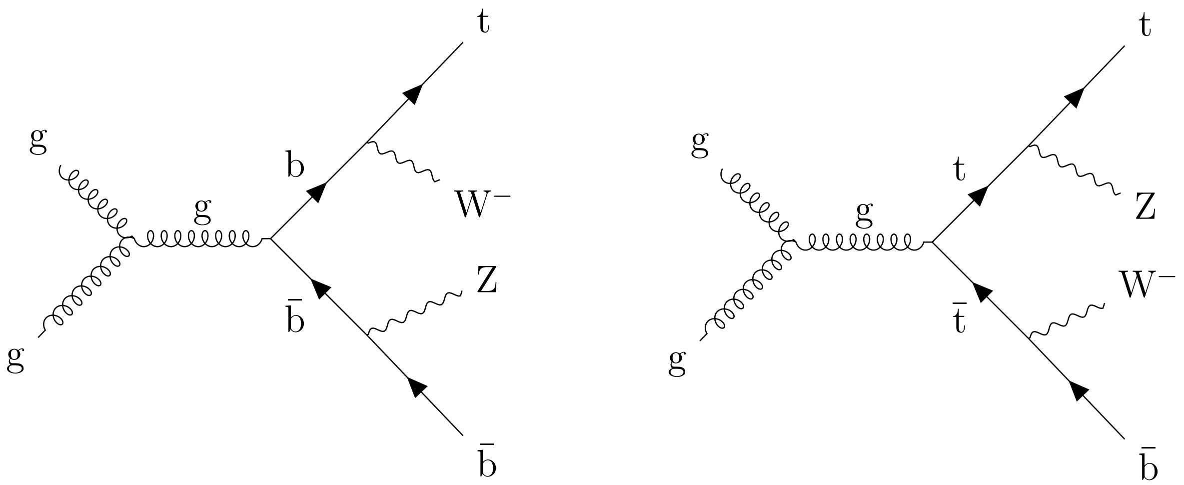

Figure 1: Feynman diagrams for the tWZ production at NLO accuracy in QCD (left), and for the $\mathrm{t\bar{t}Z}$ production arising when simulating the tWZ process at NLO accuracy in QCD (right). The two diagrams share the same final state, and are different in the number of resonant top quarks.

The analysed data set was recorded in proton-proton collisions at a center-of-mass energy of 13 TeV, collected by the CMS experiment at the LHC in 2016–2018, and corresponds to an integrated luminosity of $138\,\mathrm{fb}^{-1}$. This measurement approaches several unique challenges, from the simulation of the signal, where the interference with resonant $\mathrm{t\bar{t}Z}$ diagrams needs to be accounted for, to the use of multivariate analysis (MVA) techniques in order to achieve a more powerful discrimination between the signal and various background processes. Moreover, the analysis is carried out in two different regions of the phase space, characterized by either the top quark being almost at rest, or having a large transverse momentum p_T, referred to as low- and high-p_T (or boosted) region, respectively. The low-p_T (boosted) region contains all events with a top quark $p_T < (>) 270\,\text{GeV}$. The addition of the boosted region brings an enhanced sensitivity to new phenomena in the tWZ process, for instance in the context of SM effective field theory [2, 6]. In this Letter, we provide the inclusive measurement of the SM tWZ cross section, setting the scene for the forthcoming analyses aimed at exploring the physics beyond the SM in this signature.

Tabulated results are provided in the HEPData record for this analysis [7].

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity coverage provided by the barrel and endcap detectors. Muons are measured in gas-ionization detectors embedded in the steel flux-return yoke outside the solenoid. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Refs. [8, 9].

3 Data and simulated event samples

Events of interest are selected using a two-tiered trigger system. The first level (L1), composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of 4 μs [10]. The second level, known as the high-level trigger (HLT), consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage [11]. Collision events containing up to three leptons at HLT are selected, with an efficiency close to 100% [12], for further steps of the analysis.

Simulated Monte Carlo (MC) samples are produced with a consistent modeling of the detector operation conditions for each data-taking period. The signal sample is generated at NLO accuracy in QCD, in the five-flavor scheme (5FS), where bottom quarks are considered as sea quarks of the proton. The MADGRAPH5_aMC@NLO v2.6.5 [3, 4] MC event generator is used for the hard scattering process, while parton shower (PS) and hadronization are modeled with PYTHIA v8.240 [13]. The parton distribution function (PDF) set is NNPDF3.1 [14] at next-to-NLO (NNLO) in 5FS, and the factorization and renormalization scales are configured dynamically. At leading order (LO) accuracy and in the 5FS, the tWZ process can be easily identified through the tree-level partonic process $\text{gb} \rightarrow \text{tWZ}$, and the simulation is straightforward. However, at NLO, real emissions of the type $\text{gg} \rightarrow \text{tWZb}$ arise that may feature an additional resonant top (anti-)quark in the intermediate state. The process therefore overlaps with $\text{t}\bar{\text{t}}\text{Z}$ and with $\text{t}\bar{\text{t}}$ produced with an off-shell Z boson, at LO [2]. A schematic illustration of a $\text{t}\bar{\text{t}}\text{Z}$ diagram that appears in the generation of the signal sample at NLO in QCD is displayed in Fig. 1 (right). Such overlap is removed by means of MADSTR v1.0.0 [15], which implements diagram removal (DR) and diagram subtraction (DS) schemes. In practice, two options are used for the generation of the samples: DR1, which drops the squared resonant term from the squared amplitude of the process together with the interference term between the resonant and nonresonant contributions; and DR2, which drops the resonant term only. The DR1 scheme is used for the nominal samples, whereas DR2 is employed to assess systematic uncertainties in the signal model. The variation introduced by DS is consistently smaller than the one introduced by DR2 [16]. In order to increase the number of simulated events in the boosted region, more signal samples are generated with the top quark $p_T > 250\,\text{GeV}$. To overcome the known issue that the DR and DS scheme are no longer effective at high momenta, an additional requirement on b quarks is imposed at the generator level to retain events only from direct tWZ production [2, 15]. In practice, events with two b quarks that originate from the top quark and have $p_T > 30\,\text{GeV}$ and pseudorapidity $|\eta| < 2.4$ are removed, as they are likely to originate from the $t\bar{t}Z$ process. Signal samples at low and high momenta are checked to have no overlapping events. All signal samples are generated with the Z boson decaying to charged leptons, and with at least one of the two W bosons decaying to a lepton and a neutrino. Finally, the Z boson is simulated to be on shell. Since two of the final state leptons are identified as decay products of the Z boson in the event selection, contributions from virtual photons are not relevant in this measurement.

The MADGRAPH5_aMC@NLO MC event generator is used at NLO accuracy in QCD to model a number of background processes including $t\bar{t}Z$, the $t\bar{t}$ production in association with a W boson ($t\bar{t}W$), the single top quark production in association with a Z boson (tZq), the associated production of a W boson and a Z boson (WZ) and of two W bosons (WW), the simultaneous production of three electroweak gauge bosons (VVV, where V is either a W boson or a Z boson), and the production of a photon in association with a W or Z boson ($V\gamma$). All the samples are generated using the 5FS, apart from the tZq sample, generated in the 4FS, in which only up, down, charm, and strange quarks are considered as sea quarks of the proton. Other background processes, namely W boson production in association with jets, $t\bar{t}$ production in association with a photon ($t\bar{t}\gamma$), and with two electroweak gauge bosons ($t\bar{t}VV$) are simulated at LO accuracy with MADGRAPH5_aMC@NLO [3]. The MCFM generator v.7.0.1 [17] is used for the event simulation at LO for the gluon-initiated ZZ production ($gg \to ZZ$), while $q\bar{q} \to ZZ$ and single top production in association with a W boson are simulated with POWHEG v2 [18–21] at NLO accuracy in QCD. All events are processed with PYTHIA [13] to simulate PS and hadronization. The underlying event is modeled via the CP5 tune [22] and the NNPDF3.1 at NNLO [14] set is used for the modeling of PDFs. The matching of the matrix elements to the PS is performed using the MLM scheme [23] (FxFx merging scheme [4]) for LO (NLO) samples that are generated using MADGRAPH5_aMC@NLO. For all processes, the detector response is simulated using a detailed description of the CMS apparatus, based on the GEANT4 toolkit [24]. Additional proton-proton interactions in the same and/or neighboring bunch crossings (pileup) are simulated with PYTHIA 8. All background processes are normalized according to the measured integrated luminosity and to their most accurate theoretical cross sections (see e.g. Refs. [5, 25–28]). An important source of background is the contribution of the so-called "nonprompt" leptons from hadron decays, or from misidentification of jets or hadrons in the reconstruction process, as opposed to the prompt leptons originating from the W and Z boson decays. The processes that mostly contribute with nonprompt leptons in this measurement are Drell–Yan and $t\bar{t}$ in the dileptonic channel. This background cannot be accurately described by simulation and is estimated from data.

4 Event reconstruction and selection

The particle-flow algorithm [29] is used to reconstruct and identify individual particles in the event, combining information from different parts of the detector. The primary vertex (PV) is taken to be the vertex corresponding to the hardest scattering in the event, identified using the tracking information alone, as described in Section 9.4.1 of Ref. [30]. The energy of electrons is determined from a combination of the electron momentum at the PV, the energy of the corresponding ECAL cluster, and the energy sum of all bremsstrahlung photons spatially compatible with originating from the electron track. The energy of muons is obtained from the curvature of the corresponding track. The energy of charged hadrons is evaluated via a combination of their momentum measured in the tracker and the matching of the ECAL and HCAL energy deposits, corrected for the response function of the calorimeters to hadronic showers. Finally, the energy of neutral hadrons is obtained from the corresponding corrected ECAL and HCAL energies.

Reconstructed electrons (muons) are required to have $p_{\mathrm{T}} > 10\,\mathrm{GeV}$ and $|\eta| < 2.5$ (2.4). In order to increase the purity of prompt leptons, an MVA-based discriminant is used [31, 32] to identify the leptons that enter the final selection. Leptons that fail this requirement are subjected to additional criteria and used in the estimation of the nonprompt-lepton contribution from control regions (CRs) in data.

Jets are reconstructed by clustering the particle-flow candidates using the anti-k_{T} algorithm [33, 34] with a distance parameter of 0.4. They are required to have $p_{\mathrm{T}} > 25\,\mathrm{GeV}$, $|\eta| < 2.5$, and to be separated from selected leptons by a distance $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} > 0.4$. Here, $\Delta\eta$ and $\Delta\phi$ are the separation in pseudorapidity and azimuthal angle between the jet and the lepton, respectively. The products of hadronically decaying top quarks with p_{T} more than about twice the top quark mass, $m_{\mathrm{t}} = 172.5\,\mathrm{GeV}$, are collimated and identified as a single large-R jet. Such jets are reconstructed using the anti-k_{T} algorithm with a distance parameter of 0.8, and are employed in the analysis of the boosted region. They are required to have $p_{\mathrm{T}} > 300\,\mathrm{GeV}$ and $|\eta| < 2.4$. The magnitude of the missing transverse momentum vector $\vec{p}_{\mathrm{T}}^{\text{miss}}$, defined as the projection onto the transverse plane relative to the beam axis, of the negative vector sum of the momenta of all particle-flow candidates, is referred to as $p_{\mathrm{T}}^{\text{miss}}$ [35]. Corrections to the jet energies are propagated to $\vec{p}_{\mathrm{T}}^{\text{miss}}$.

The DEEPJET algorithm [36–38] is used to identify b quark jets. Jets are considered b tagged if they pass a requirement on the DEEPJET score that provides a misidentification rate of 1% for jets originating from u, d, or s quarks and gluons. Jets that are b tagged are referred to as "b jets" in the following.

Events in the low-p_{T} region are required to have two leptons with $p_{\mathrm{T}} > 20\,\mathrm{GeV}$ and one with $p_{\mathrm{T}} > 25\,\mathrm{GeV}$. Two of the leptons must have opposite electric charge and same flavor (OCSF). The invariant mass of the OCSF lepton pair must lie within $m_{\mathrm{Z}} \pm 15\,\mathrm{GeV}$, where m_{Z} is the Z boson mass equal to 91.2 GeV [39]. Selected events should contain at least two jets of which at least one is b tagged. To increase the sensitivity of the analysis, the low-p_{T} region is further split into $\mathrm{SR}_{3\ell,3\mathrm{j}}$ and $\mathrm{SR}_{3\ell,2\mathrm{j}}$ categories. Events with at least three jets fall into the $\mathrm{SR}_{3\ell,3\mathrm{j}}$ category whereas $\mathrm{SR}_{3\ell,2\mathrm{j}}$ contains events with exactly two jets. The $\mathrm{SR}_{3\ell,3\mathrm{j}}$ category has the largest contribution of signal with 77 signal events expected over 1252 background events, of which more than half are expected to be $\mathrm{t\bar{t}Z}$. The expected signal yield in the $\mathrm{SR}_{3\ell,2\mathrm{j}}$ category is 28 to compare with 828 background events. Background in this category is mostly from nonprompt leptons and from WZ in association with light-flavor or gluon jets (WZ+jets). The $\mathrm{t\bar{t}Z}$ process contributes to less than 15% of the background in $\mathrm{SR}_{3\ell,2\mathrm{j}}$. Another category in the low-p_{T} region targets the fully-leptonic decay of the tWZ signal. Events in this category, $\mathrm{SR}_{4\ell}$, must have at least one b-tagged jet and four leptons, three of which respect the same p_{T} requirement of $\mathrm{SR}_{3\ell,3\mathrm{j}}$ and a fourth one with $p_{\mathrm{T}} > 10\,\mathrm{GeV}$. While two out of four leptons are required to be OCSF, with an invariant mass within $m_{\mathrm{Z}} \pm 15\,\mathrm{GeV}$, the other two are required to fail at least one of these criteria. In this region, 16 signal events are expected over 162 background events, of which more than 70% is $\mathrm{t\bar{t}Z}$.

The high-p_{T} region is divided in two categories, $\mathrm{SR}_{\mathrm{Had}}^{\mathrm{Boosted}}$ and $\mathrm{SR}_{\mathrm{Lep}}^{\mathrm{Boosted}}$, respectively targeting hadronically and leptonically decaying top quarks. The selection in both categories follows that of $\mathrm{SR}_{3\ell,3j}$ with a number of additional requirements. The $\mathrm{SR}^{\mathrm{Boosted}}_{\mathrm{Had}}$ category is required to contain a large-R jet, with a soft-drop [40] mass between 105 and 210 GeV, in the vicinity of $\Delta R = 0.8$ around a b-tagged jet. In the $\mathrm{SR}^{\mathrm{Boosted}}_{\mathrm{Lep}}$ category, the lepton that is not compatible with the hypothesis of originating from the Z boson (the third lepton) is required to have $p_{\mathrm{T}} >$ 30 GeV, and to lie in a cone of $\Delta R = 2$ around a b-tagged jet with $p_{\mathrm{T}} > 200\,\mathrm{GeV}$. A fully connected deep neural network (DNN) [41] is trained on simulated $\mathrm{t\bar{t}}$ events, with one top quark decaying into a lepton, with the aim of tagging the lepton from the high-p_{T} top quark decay, using TensorFlow v2.8.0 [42]. Each third lepton is assigned a score according to its compatibility with the hypothesis of originating from the high-p_{T} top quark. Variables related to the third lepton, the b jet, as well as a combination of the two, are fed into the DNN, which is then validated on simulated $\mathrm{t\bar{t}}$ events with both top quarks decaying into leptons. Overall in the high-p_{T} region, 5 signal and 81 background events are expected.

In addition to regions enriched in signal events, two regions are defined in data to control and validate the ZZ and WZ backgrounds, in order to ensure a correct modelling. Events containing exactly four leptons, with the two pairs being compatible with the hypothesis of originating from a Z boson, are collected in the region enriched in ZZ, denoted by $\mathrm{CR}_{\mathrm{ZZ}}$. The region enriched in WZ+jets events, $\mathrm{CR}_{\mathrm{WZ}}$, is built with events containing exactly three leptons, two of them originating from the Z boson, with $p_{\mathrm{T}}^{\mathrm{miss}} > 50\,\mathrm{GeV}$, and no b-tagged jets. The expected yields for the signal and background processes, together with the observed number of events in each region, are reported in Table 1 and Table 2.

Table 1: Expected yields for signal and background processes and observed number of events in the signal regions.

	$\mathrm{SR}_{3\ell,3j}$	$\mathrm{SR}_{3\ell,2j}$	$\mathrm{SR}_{4\ell}$	$\mathrm{SR}^{\mathrm{Boosted}}$
tWZ signal	77.47 ± 0.12	28.19 ± 0.07	15.98 ± 0.06	5.44 ± 0.02
$\mathrm{t\bar{t}Z}$	657.9 ± 1.6	122.76 ± 0.61	113.86 ± 0.64	59.03 ± 0.50
Nonprompt leptons	139 ± 42	170 ± 51	1.02 ± 0.31	1.94 ± 0.58
tZq	86.45 ± 0.78	108.69 ± 0.71	0.29 ± 0.04	4.37 ± 0.17
ZZ	22.7 ± 2.4	60.6 ± 4.1	20.0 ± 2.3	0.30 ± 0.29
WZ	166.4 ± 3.3	227.8 ± 4.0	0.59 ± 0.19	6.84 ± 0.66
VV(V)	15.51 ± 0.11	10.55 ± 0.09	1.35 ± 0.03	0.64 ± 0.02
$\mathrm{t(\bar{t})X}$	108.30 ± 0.99	49.4 ± 1.2	17.32 ± 0.34	6.26 ± 0.19
$\mathrm{X}\gamma$	54.1 ± 2.6	78.3 ± 3.7	6.92 ± 0.95	1.08 ± 0.31
Total backgrounds	1249 ± 42	822 ± 51	159.9 ± 2.6	80.8 ± 1.1
Data	1463	849	180	77

5 Background estimation

The contribution of nonprompt leptons is estimated from data in a similar way to that of Ref. [12]. Event weights, derived from simulated QCD multijet samples, are assigned to collision events that are selected similarly to signal with modified (looser) lepton requirements. To validate the method, a data control region enriched in the nonprompt contribution is defined by a trilepton selection similar to that of signal, except that the OCSF pair is required to be out

Table 2: Expected yields for signal and background processes and observed number of events in the control regions.

	CR_{WZ}	CR_{ZZ}
tWZ signal	31.96 ± 0.08	2.39 ± 0.02
$t\bar{t}Z$	112.41 ± 0.73	14.44 ± 0.23
Nonprompt leptons	1450 ± 430	23.0 ± 7.0
tZq	74.84 ± 0.67	0.05 ± 0.01
ZZ	597 ± 12	2202 ± 22
WZ	10610 ± 25	0.68 ± 0.16
VV(V)	166 ± 14	16.52 ± 0.07
$t(\bar{t})X$	39.4 ± 1.8	1.08 ± 0.07
$X\gamma$	519 ± 11	2.53 ± 0.60
Total backgrounds	13520 ± 430	2028 ± 23
Data	12743	2352

of the Z boson mass window. Events should also have at least one b-tagged jet. The selection provides a good amount of nonprompt leptons. Considering key distributions such as the p_T of the trailing lepton, the nonprompt estimation is found to be compatible with data within 30%.

6 Statistical analysis

In order to increase the discrimination power between signal and backgrounds, DNN classifiers are designed and implemented via TensorFlow in $SR_{3\ell,3j}$ and $SR_{3\ell,2j}$. A fully connected DNN is trained on simulated events in $SR_{3\ell,3j}$ in the form of a multiclass classifier. Events are given a score according to their compatibility with tWZ, $t\bar{t}Z$, and other background events, and are therefore classified in three categories. The following observables yield discrimination between the processes and are fed into the classifier: the number of b jets, the p_T and η of jets and leptons, the maximum ΔR between lepton pair combinations, the maximum ΔR and invariant mass between jet pair combinations, the maximum η of all jets in the event, the maximum p_T of the lepton-jet combinations, ${p_T}_{\ell j}^{max}$, and the invariant mass of the system of leptons, jets, and $\vec{p}_T^{miss}$, m_{sys}. The last two variables are shown in Fig. 2 (upper row).

To further distinguish the tWZ signal from $t\bar{t}Z$ production, the following kinematic observables specific to the top quark and W boson are added as inputs to the classifier: the mass, p_T, η, and ϕ of the top quark and hadronically decaying W boson, as well as the p_T, η, and ϕ of the leptonically decaying W boson. This requires the reconstruction of the signal event hypothesis, taking into account the origin of the third lepton. Two top quark decay hypotheses are checked per event using $\chi^2_{t,had}$ and $\chi^2_{t,lep}$ variables,

$$\chi^2_{t,had} = \left(\frac{m_{3j} - m_t}{\sigma_{t,had}}\right)^2 \quad \text{and} \quad \chi^2_{t,lep} = \left(\frac{m_{\ell\nu j} - m_t}{\sigma_{t,lep}}\right)^2 + \left(\frac{m_{jj} - m_W}{\sigma_{W,had}}\right)^2. \tag{1}$$

In Eq. (1), m_{3j} and m_{jj} are the tri- and di-jet invariant masses, respectively, expected to be compatible with the mass of a top quark and a W boson that decay hadronically. The invariant

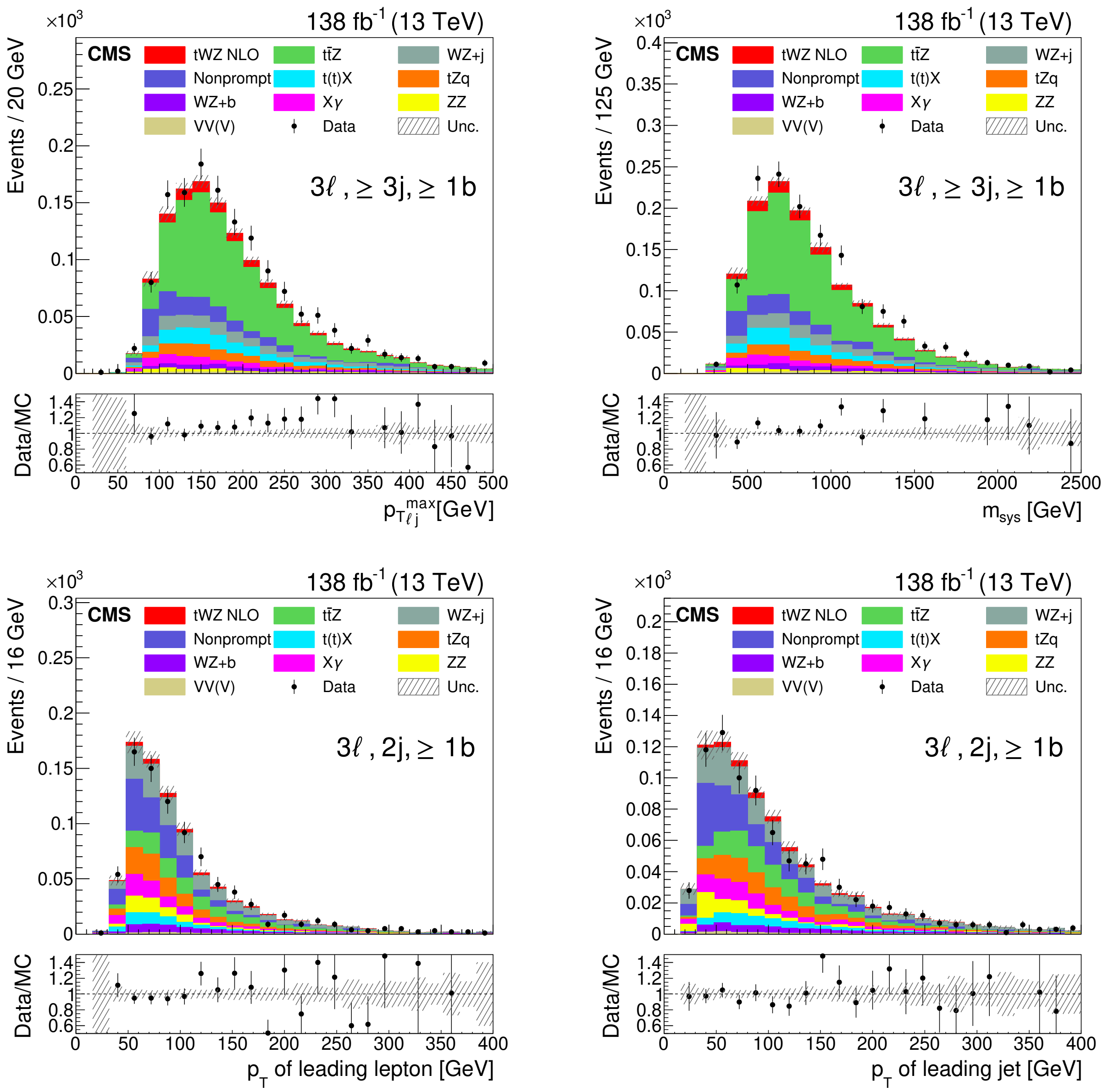

Figure 2: Pre-fit examples of the input features for the two DNN trainings: the $p_{\mathrm{T}\ell\mathrm{j}}^{\mathrm{max}}$ (upper left) and m_{sys} (upper right) observables in $\mathrm{SR}_{3\ell,3\mathrm{j}}$, as well as the p_{T} of the leading lepton (lower left) and leading jet (lower right) in $\mathrm{SR}_{3\ell,2\mathrm{j}}$. The VV(V) group in the legend denotes the VVV, WW, and W in association with jets backgrounds. The dashed band shows the total uncertainty (statistical and systematic) before the fit.

mass $m_{\ell\nu\mathrm{j}}$ of the lepton not coming from the Z boson, the neutrino, and a jet is expected to be close to the mass of the leptonically decaying top quark. The variables $\sigma_{\mathrm{t,had}}$, $\sigma_{\mathrm{t,lep}}$, and $\sigma_{\mathrm{W,had}}$ represent the mass resolutions and are estimated from simulation. The W boson mass is set to $m_{\mathrm{W}} = 80.4\,\mathrm{GeV}$ [39] and the m_{t} is set to 172.5 GeV, following the settings in the simulation. The neutrino four-momentum in the W boson decay is estimated as described in Ref. [43]. The lepton-jet combination yielding the smaller χ^2 value is selected, and the kinematic properties of the top quark and (if it exists) the hadronically decaying W boson are used in the DNN training.

Another DNN is trained as a binary classifier on simulated events in $\mathrm{SR}_{3\ell,2\mathrm{j}}$, with the aim of discriminating the signal from all other background sources, especially nonprompt leptons and WZ+jets. A broad range of observables are investigated, and those that provide the best discrimination between signal and background processes are selected: the p_T and η of the leptons and jets, the number of b jets, the p_T and invariant mass of the jet pair, the m_{sys}, and the p_T of the system of leptons and jets. Figure 2 (lower row) shows the p_T of the leading lepton and leading jet as examples of input features.

The tWZ cross section is measured via the maximization of a binned likelihood function in a fit to data. This function combines the Poisson probabilities for the observed yields given the predicted signal and background estimates in each bin, and incorporates as nuisance parameters all sources of systematic uncertainties that may affect the number of observed signal or background events. The fit returns the measured value of the signal strength, defined by the ratio of the measured cross section and the expected one, together with its uncertainty.

Inputs to the binned likelihood function are distributions and number of events, depending on the region and category. The $SR_{3\ell,3j}$ category is split into two parts: one containing events with exactly one b jet where the distribution of the tWZ output node score is taken from, and the other containing the rest and providing the distribution of the $t\bar{t}Z$ output node score to the fit. This is to enhance the sensitivity of the fit to differences between tWZ and $t\bar{t}Z$, considering the additional b jet in the $t\bar{t}Z$ background. From $SR_{3\ell,2j}$, the distribution of the tWZ output node score contributes to the fit whereas the b jet multiplicity distribution is taken from $SR_{4\ell}$. Given the limited number of events in $SR_{Had}^{Boosted}$ and $SR_{Lep}^{Boosted}$, the two SRs are merged and contribute to the signal extraction fit via event counting. The number of events in CR_{ZZ} and CR_{WZ} are also included in the fit in order to constrain the normalization of the ZZ and WZ+jets processes. To account for the different b-quark content in the SR and CR_{WZ}, the WZ background is split into WZ+b and WZ+jets based on the generator-level information. The tWZ and $t\bar{t}Z$ output node scores from the multiclass classifier in $SR_{3\ell,3j}$ and the score of the binary classifier in $SR_{3\ell,2j}$ are shown in Fig. 3 after the fit (upper and lower rows, respectively).

7 Systematic uncertainties

Several sources of systematic uncertainty are considered. We report the uncertainties in the b tagging, lepton, and trigger efficiency corrections, the distribution of pileup events in simulation, the jet energy correction (JEC) and resolution (JER), and the unclustered missing energy. The uncertainties arising from our imperfect knowledge encoded in the MC event generator parameters and settings are also included in the fit. They are systematic effects associated with the matrix element renormalization and factorization scales, the PS simulation, and the PDFs (the latter being estimated by reweighting the simulation using the corresponding variations in the NNPDF sets). Furthermore, variations in templates of the tWZ signal arising from the difference between the DR1 and DR2 schemes are included as uncertainties associated with the signal modeling. Theoretical uncertainties in the cross sections of the background processes are also taken into account. In particular, a normalization uncertainty of 15% is estimated for the $t\bar{t}Z$ cross section, which covers the relative difference between the measurement of the cross section and its theoretical prediction [12, 44]. Theoretical uncertainties of 11 (10)% are applied to tZq (ZZ, WZ+jets, and triboson and processes in association with a photon), while a 20% uncertainty is estimated for the WZ+b background [12]. All background processes, except nonprompt leptons, are estimated using simulation where uncertainties in the normalization and modeling are considered. The uncertainties in the integrated luminosity collected in each data-taking period (2016, 2017 and 2018) are calculated to be 1.2–2.5% [45–47], corresponding to an uncertainty of 1.6% for the whole 2016–2018 data-taking period. Statistical uncertainties

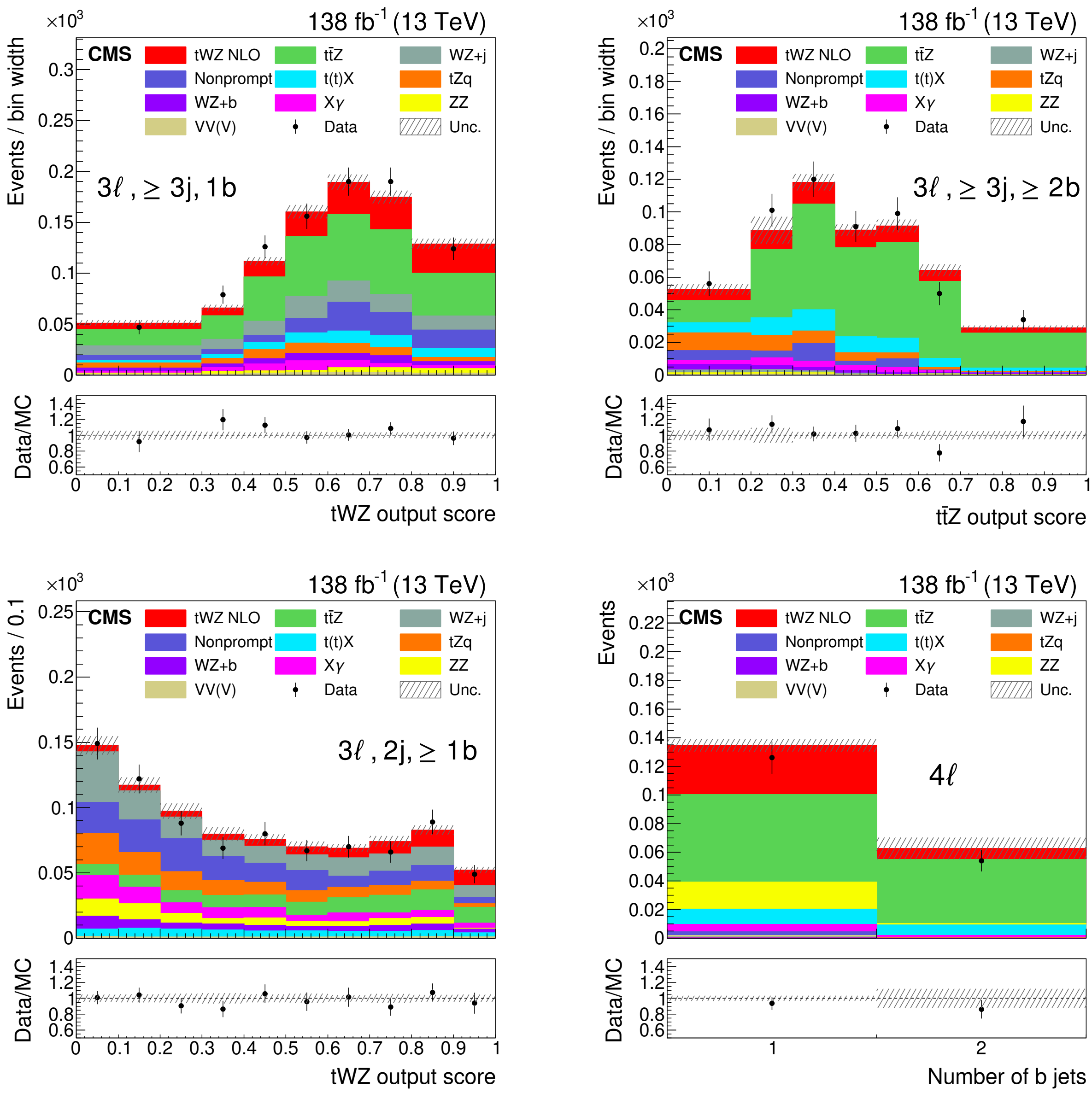

Figure 3: Score of the tWZ output node from the multiclass classifier in $SR_{3\ell,3j}$ for events with exactly one b jet (upper left), and of the $t\bar{t}Z$ output node in $SR_{3\ell,3j}$ for events with more than one b jet (upper right); score of the tWZ output node of the binary classifier in $SR_{3\ell,2j}$ (lower left), and the b jet multiplicity in $SR_{4\ell}$ (lower right). The VV(V) group in the legend denotes the VVV, WW, and W in association with jets backgrounds. The dashed band shows the total uncertainty (statistical and systematic) after the fit.

from the nonprompt-lepton background estimation, together with a 30% normalization uncertainty in order to cover with potential residual mismodeling, are treated in a bin-wise manner. Finally, uncertainties arising from the finite size of simulated samples are also accounted for following the Barlow–Beeston method [48]. In the assessment of the systematic uncertainties, the DNNs are reevaluated on modified inputs to propagate the uncertainties to the templates.

8 Results

The fit provides evidence for the tWZ process with a statistical significance of 3.4 standard deviations, for an expected one of 1.4. The signal strength is measured to be $2.6 \pm 0.4\,(\text{stat}) \pm 0.7\,(\text{syst})$, and the cross section is found to be $354 \pm 54\,(\text{stat}) \pm 95\,(\text{syst})$ fb. This result is two standard deviations above the SM expectations. The dominant systematic uncertainty is the $t\bar{t}Z$ normalization with an impact on the final measurement of 18%, constrained to the level of 10% in the fit. Fixing the $t\bar{t}Z$ cross section to the measured value in Ref. [44], the statistical significance of the signal stays above three standard deviations.

Coming next are the uncertainties in the normalization of the backgrounds other than $t\bar{t}Z$ and the experimental sources of uncertainty, with an impact of 12% and 10%, respectively. The uncertainty in the signal modeling has an impact of 5%. The region that yields the highest sensitivity to the final measurement is the $SR_{3\ell,3j}$, followed by the $SR_{4\ell}$ and the $SR_{3\ell,2j}$. The boosted region, which includes only a small number of events, contributes with a statistical significance of less than 0.2 standard deviations. Overall, results in all categories across all data-taking periods are found to be compatible within one standard deviation.

The reason for the leading impact of the $t\bar{t}Z$ normalization uncertainty lies in the very similar nature of the tWZ and $t\bar{t}Z$ processes. To further investigate the interconnection, we additionally performed a simultaneous fit of the tWZ and $t\bar{t}Z$ signal strengths. The two-dimensional likelihood scan is displayed in Fig. 4, and shows that the two parameters are anticorrelated, with the $t\bar{t}Z$ signal strength being compatible with the SM expectation within one standard deviation.

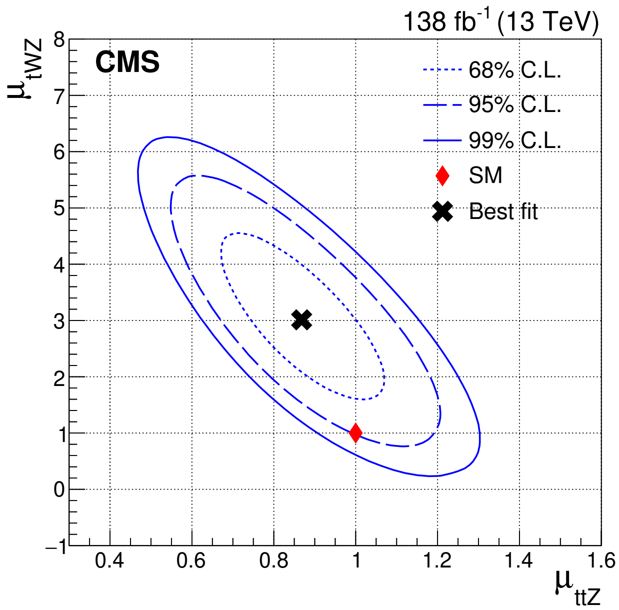

Figure 4: Two-dimensional likelihood scan of the signal strenghts of tWZ (μ_{tWZ}) vs. $t\bar{t}Z$ ($\mu_{t\bar{t}Z}$). The blue lines show the 68%, 95%, and 99% confidence level contours. The black cross represents the best-fit value, while the red diamond the SM expectation.

9 Summary

A measurement of the tWZ process is performed in final states with three or four leptons, using the data collected by the CMS experiment at the LHC in 2016–2018, at a center-of-mass energy of 13 TeV, corresponding to an integrated luminosity of 138 fb^{-1}. To assess backgrounds and establish the signal, the measurement heavily relies on multivariate techniques and the data are exploited in multiple regions and categories. The challenge of the tWZ signal overlapping with the $t\bar{t}Z$ background is overcome using the latest advancements in the tWZ modeling. We find the signal to have an observed statistical significance of 3.4 standard deviations, corresponding to a measured cross section of 354 ± 54 (stat) ± 95 (syst) fb that is two standard deviations above the standard model prediction. This is the first evidence of the tWZ process.

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[5]Also at Universidade Estadual de Campinas, Campinas, Brazil
[6]Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil
[7]Also at UFMS, Nova Andradina, Brazil
[8]Also at Nanjing Normal University, Nanjing, China
[9]Now at Henan Normal University, Xinxiang, China
[10]Now at The University of Iowa, Iowa City, Iowa, USA
[11]Also at University of Chinese Academy of Sciences, Beijing, China
[12]Also at China Center of Advanced Science and Technology, Beijing, China
[13]Also at University of Chinese Academy of Sciences, Beijing, China
[14]Also at China Spallation Neutron Source, Guangdong, China
[15]Also at Université Libre de Bruxelles, Bruxelles, Belgium
[16]Also at an institute or an international laboratory covered by a cooperation agreement with CERN
[17]Now at British University in Egypt, Cairo, Egypt
[18]Now at Cairo University, Cairo, Egypt
[19]Also at Birla Institute of Technology, Mesra, Mesra, India
[20]Also at Purdue University, West Lafayette, Indiana, USA
[21]Also at Université de Haute Alsace, Mulhouse, France
[22]Also at Department of Physics, Tsinghua University, Beijing, China
[23]Also at Tbilisi State University, Tbilisi, Georgia
[24]Also at The University of the State of Amazonas, Manaus, Brazil
[25]Also at Erzincan Binali Yildirim University, Erzincan, Turkey
[26]Also at University of Hamburg, Hamburg, Germany
[27]Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany
[28]Also at Isfahan University of Technology, Isfahan, Iran
[29]Also at Bergische University Wuppertal (BUW), Wuppertal, Germany
[30]Also at Brandenburg University of Technology, Cottbus, Germany
[31]Also at Forschungszentrum Jülich, Juelich, Germany
[32]Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland
[33]Also at Institute of Physics, University of Debrecen, Debrecen, Hungary
[34]Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary
[35]Now at Universitatea Babes-Bolyai - Facultatea de Fizica, Cluj-Napoca, Romania
[36]Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt
[37]Also at HUN-REN Wigner Research Centre for Physics, Budapest, Hungary
[38]Also at Faculty of Informatics, University of Debrecen, Debrecen, Hungary
[39]Also at Punjab Agricultural University, Ludhiana, India
[40]Also at University of Hyderabad, Hyderabad, India
[41]Also at University of Visva-Bharati, Santiniketan, India
[42]Also at Indian Institute of Science (IISc), Bangalore, India
[43]Also at IIT Bhubaneswar, Bhubaneswar, India
[44]Also at Institute of Physics, Bhubaneswar, India
[45]Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany
[46]Also at Department of Physics, Isfahan University of Technology, Isfahan, Iran
[47]Also at Sharif University of Technology, Tehran, Iran
[48]Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran

[49]Also at Helwan University, Cairo, Egypt
[50]Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy
[51]Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy
[52]Also at Università degli Studi Guglielmo Marconi, Roma, Italy
[53]Also at Scuola Superiore Meridionale, Università di Napoli 'Federico II', Napoli, Italy
[54]Also at Fermi National Accelerator Laboratory, Batavia, Illinois, USA
[55]Also at Università di Napoli 'Federico II', Napoli, Italy
[56]Also at Ain Shams University, Cairo, Egypt
[57]Also at Consiglio Nazionale delle Ricerche - Istituto Officina dei Materiali, Perugia, Italy
[58]Also at Riga Technical University, Riga, Latvia
[59]Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia
[60]Also at Consejo Nacional de Ciencia y Tecnología, Mexico City, Mexico
[61]Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka
[62]Also at Saegis Campus, Nugegoda, Sri Lanka
[63]Also at INFN Sezione di Pavia, Università di Pavia, Pavia, Italy
[64]Also at National and Kapodistrian University of Athens, Athens, Greece
[65]Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland
[66]Also at University of Vienna Faculty of Computer Science, Vienna, Austria
[67]Also at Universität Zürich, Zurich, Switzerland
[68]Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria
[69]Also at Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France
[70]Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey
[71]Also at Konya Technical University, Konya, Turkey
[72]Also at Izmir Bakircay University, Izmir, Turkey
[73]Also at Adiyaman University, Adiyaman, Turkey
[74]Also at Bozok Universitetesi Rektörlügü, Yozgat, Turkey
[75]Also at Marmara University, Istanbul, Turkey
[76]Also at Milli Savunma University, Istanbul, Turkey
[77]Also at Kafkas University, Kars, Turkey
[78]Now at stanbul Okan University, Istanbul, Turkey
[79]Also at Hacettepe University, Ankara, Turkey
[80]Also at Istanbul University - Cerrahpasa, Faculty of Engineering, Istanbul, Turkey
[81]Also at Yildiz Technical University, Istanbul, Turkey
[82]Also at Vrije Universiteit Brussel, Brussel, Belgium
[83]Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom
[84]Also at University of Bristol, Bristol, United Kingdom
[85]Also at IPPP Durham University, Durham, United Kingdom
[86]Also at Monash University, Faculty of Science, Clayton, Australia
[87]Now at an institute or an international laboratory covered by a cooperation agreement with CERN
[88]Also at Università di Torino, Torino, Italy
[89]Also at Bethel University, St. Paul, Minnesota, USA
[90]Also at Karamanoğlu Mehmetbey University, Karaman, Turkey
[91]Also at California Institute of Technology, Pasadena, California, USA

[92]Also at United States Naval Academy, Annapolis, Maryland, USA
[93]Also at Bingol University, Bingol, Turkey
[94]Also at Georgian Technical University, Tbilisi, Georgia
[95]Also at Sinop University, Sinop, Turkey
[96]Also at Erciyes University, Kayseri, Turkey
[97]Also at Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Bucharest, Romania
[98]Also at Texas A&M University at Qatar, Doha, Qatar
[99]Also at Kyungpook National University, Daegu, Korea
[100]Also at another institute or international laboratory covered by a cooperation agreement with CERN
[101]Also at Universiteit Antwerpen, Antwerpen, Belgium
[102]Also at Yerevan Physics Institute, Yerevan, Armenia
[103]Also at Northeastern University, Boston, Massachusetts, USA
[104]Also at Imperial College, London, United Kingdom
[105]Also at Institute of Nuclear Physics of the Uzbekistan Academy of Sciences, Tashkent, Uzbekistan