

Precise predictions for the trilinear Higgs self-coupling in the Standard Model and beyond

Henning Bahl,^{a,b} Johannes Braathen,^c Martin Gabelmann^{c,*} and Georg Weiglein^{c,d}

^aUniversity of Chicago, Department of Physics and Enrico Fermi Institute, 5720 South Ellis Avenue, Chicago, IL 60637 USA

^bInstitut für Theoretische Physik, Philosophenweg 16, 69120 Heidelberg, Germany

^cDeutsches Elektronen-Synchrotron DESY, Notkestr. 85, 22607 Hamburg, Germany

^dII. Institut für Theoretische Physik, Universität Hamburg, Luruper Chaussee 149, 22761 Hamburg, Germany

E-mail: bahl@thphys.uni-heidelberg.de, johannes.braathen@desy.de, martin.gabelmann@desy.de, georg.weiglein@desy.de

Deviations in the trilinear self-coupling of the Higgs boson at 125 GeV from the Standard Model (SM) prediction are a sensitive test of physics Beyond the SM (BSM). The LHC experiments searching for the simultaneous production of two Higgs bosons start to become sensitive to such deviations. Therefore, precise predictions for the trilinear Higgs self-coupling in different BSM models are required in order to be able to test them against current and future bounds. We present the new framework `anyH3`, which is a Python library that can be utilized to obtain predictions for trilinear scalar couplings up to the one-loop level in any renormalisable theory. The program makes use of the UFO format as input and is able to automatically apply a wide variety of renormalisation schemes involving minimal and non-minimal subtraction conditions. External-leg corrections are also computed automatically, and finite external momenta can be optionally taken into account. The Python library comes with convenient command-line as well as Mathematica user interfaces. We perform cross-checks using consistency conditions such as UV-finiteness and decoupling, and also by comparing against results known in the literature. As example applications, we obtain results for the trilinear self-coupling of the SM-like Higgs boson in various concrete BSM models, study the effect of external momenta as well as of different renormalisation schemes.

The European Physical Society Conference on High Energy Physics (EPS-HEP2023)
21-25 August 2023
Hamburg, Germany

*Speaker

1. Introduction

Probing the self-interactions of the Higgs boson at 125 GeV found at the LHC is among the main goals of current and future high-energy physics experiments. A useful way to describe deviations in this coupling from the SM prediction is via the coupling modifier $\kappa_\lambda = \frac{\lambda_{hhh}}{\lambda_{hhh}^{(0), \text{SM}}}$, *i.e.* the prediction for the trilinear self-coupling of the SM-like Higgs boson in a given BSM model, at a given order in perturbation theory, relative to the tree-level prediction for the trilinear Higgs self-coupling in the SM. κ_λ is known to be very sensitive to BSM effects and can easily deviate by several hundred percent from 1 if higher-order corrections are taken into account [1, 2]. The current experimental constraints on κ_λ via gluon- and vector-boson-fusion induced double-Higgs production (including also information from single-Higgs processes) are rather limited, $-0.4 < \kappa_\lambda^{\text{exp.}} < 6.3$ [3], but better constraints are expected at the high-luminosity LHC (HL-LHC), *i.e.* $0.1 < \kappa_\lambda < 2.3$ [4] for the case of the SM. In the context of the Two Higgs Doublet Model (THDM) it was shown in Ref. [5] that current limits on κ_λ can be used to constrain the BSM parameter space that would otherwise be allowed by all other state-of-the-art experimental and theoretical constraints. It is to be expected that many more models will be probed by this new type of experimental constraint.

A quite large number of one- and two-loop studies for κ_λ in supersymmetric (SUSY) [6, 7] as well in non-supersymmetric [1, 2, 8] extensions of the SM exist. However, only a tiny fraction of these results can be conveniently evaluated using public tools, such as H-COUP [9], BSMP [10] or NMSSMCALC [7]. On the other hand, there are many other BSM models for which higher-order predictions for κ_λ are missing. Furthermore, the difficulty in a consistent and reliable computation of κ_λ increases rapidly with the complexity of the considered model. This strongly motivates the automation of the calculation of κ_λ including higher-order corrections.

In these proceedings we summarise the new tool anyH3 which has been introduced in Ref. [11]. The program is capable of computing all spin 0, 1, $1/2$ self-energies as well as all scalar one-, two-, and three-point functions at full one-loop level in arbitrary renormalisable theories. Therefore, it contains all ingredients to obtain one-loop predictions for the renormalised trilinear Higgs self-coupling in a large class of BSM models and renormalisation schemes.

2. Generic λ_{hhh} calculation

In the following we briefly review the ingredients for a *generic* calculation of λ_{hhh} . For a more detailed discussion see Ref. [11]. We compute the renormalised three-point function $\hat{\Gamma}_{hhh}(p_1, p_2, p_3)$ diagrammatically. Intermediate steps such as *e.g.* the tensor reduction at the level of a given model can be very time-consuming but at the same time are quite repetitive. For this reason, we performed the calculation of all possible Feynman diagrams using a *generic* renormalisable lagrangian: We assumed the most general Lorentz, coupling and mass structures. The resulting generic expressions depend on (scalar) loop functions, masses and generic couplings, which are mapped onto the concrete model under consideration at run-time. The vertex counter-term is constructed in an automated way from the tree-level coupling and can take into account contributions from on-shell (OS) mass and electroweak vacuum expectation value (VEV) counterterms. If the tree-level prediction additionally depends on *e.g.* mixing angles, which by default are renormalised $\overline{\text{MS}}$, it can be useful to define custom (user-provided) renormalisation conditions.

3. The anyH3 (anyBSM) library

The installation steps are explained in the anyBSM online documentation:

<https://anybsm.gitlab.io>

which also contains further information and concrete examples. All information about the model is stored using the UFO format [12]. The SM-like Higgs boson (and all other SM particles) can either be specified by the user or be automatically determined based on their PDG ID and numerical mass values. For the renormalisation, the user has to specify which fields and parameters are renormalised OS or $\overline{\text{MS}}$, or to define their own custom counterterms.

It should be stressed that our approach can in principle also be applied to the calculation of other (pseudo) observables. Therefore, all initial steps described above and in Ref. [11] are actually organised in a more general library, called anyBSM, which consists of several sub-modules. At the end of the module-chain several classes can be defined, of which anyH3 is the first, making use of all the tools provided by anyBSM to construct a UV-finite and consistent observable.

Code examples: Importing the main module, loading (for instance) the THDM-II UFO model (which is shipped with the code) and computing λ_{hhh} requires only three commands:

```
from anyBSM import anyBSM
THDM = anyBSM('THDMII')
THDM.lambdahhh()
```

By default, all external legs carry no momentum. This can be changed with appropriate arguments and is demonstrated in Section 5.

Command-line interface: Calling *e.g.*, the command “anyBSM THDMII” one can obtain the same result as in the examples above. To list explanations of all possible options, one can add the “-h” flag to the command.

Mathematica interface: The Mathematica interface is controlled in an analogous way to the Python library

```
<<anyBSM`
LoadModel["THDMII"]
lambda = lambdahhh[]
lambda["total"] - lambda["treelevel"]//.UVparts
```

where lambda consists of analytical rather than numerical results. The fourth line will evaluate to zero after some trivial simplifications and is a nice demonstration of UV-finiteness, one of the cross-checks discussed in the next section.

4. Code validation and mass-splitting effects

The code was validated by performing non-trivial cross-checks for all (currently 14) UFO models shipped with the code. Among the strongest checks are UV-finiteness (see *e.g.* the example above) and direct comparison with results available in the literature. Another consistency check is whether the correct decoupling behaviour occurs. This is demonstrated numerically for a subset of models

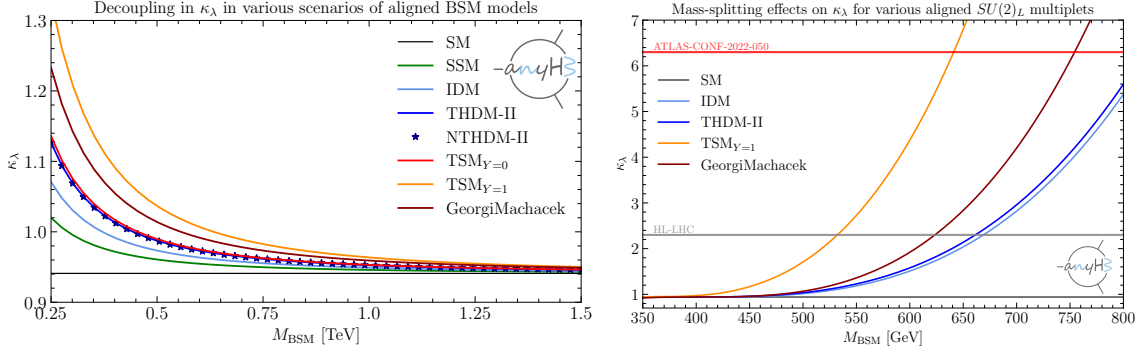


Figure 1: Left: Decoupling property of an excerpt of (UFO) models which are shipped with anyH3. Input parameters as in Ref. [11]. Right: demonstration of mass splitting effects discussed in the text.

in Fig. 1 (left). However, away from the decoupling limit, phenomenologically interesting effects can occur if *e.g.* mass splittings are present. In the simplest such cases, the BSM particles receive their mass entirely via their interaction with the SM-like Higgs boson $M_{\text{BSM}}^2 \propto v^2 \lambda_{hh} \Phi_{\text{BSM}} \Phi_{\text{BSM}}$ implying large quartic couplings in the case of large BSM masses. This well-known effect [5] is demonstrated for a variety of $SU(2)_L$ extensions in Fig. 1 (right), see Ref. [11] for more details. It was checked (using a sub-module of the anyBSM library) that the models stay perturbative in the shown ranges.

5. Momentum dependence in the THDM type I

In Fig. 2 we show the momentum dependence of κ_λ in the THDM-I for two benchmark points featuring $\kappa_\lambda < 6.3$ (left) and $\kappa_\lambda > 6.3$ (right). Since the integration of the total double-Higgs production cross-section peaks around $\sqrt{p^2} \approx 400$ GeV one can expect that the conclusion for the two points (allowed and disallowed, respectively) is not altered by the momentum dependent effects. This is a test which now can be performed with anyH3 on a point-by-point basis.

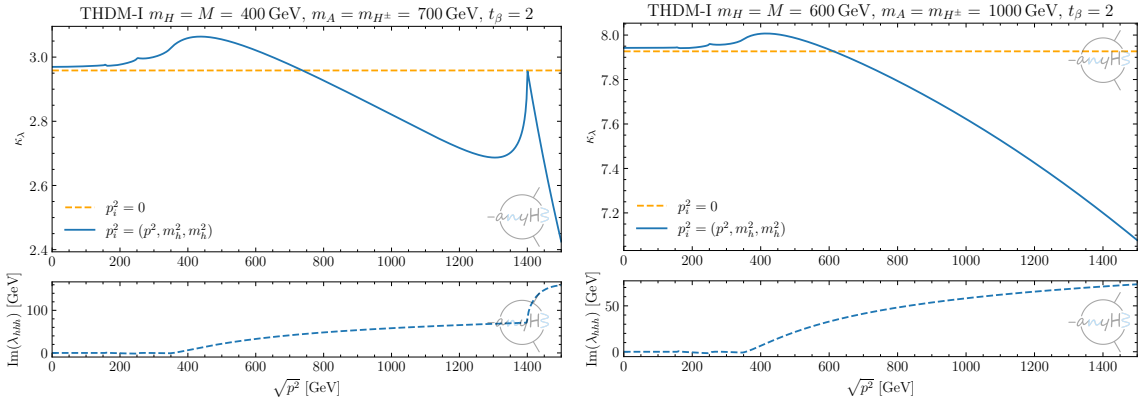


Figure 2: Dependence of κ_λ in the THDM-I on the external momentum $\sqrt{p^2}$ for a parameter point with small (left) and large κ_λ (right). It is assumed that one external leg carries the momentum p while the other two legs are on-shell.

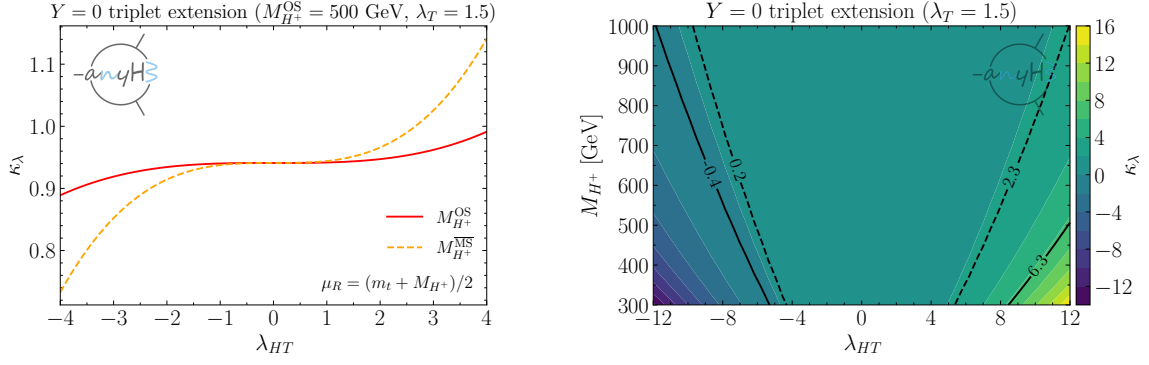


Figure 3: One-loop prediction for κ_λ in the real triplet extended SM.

6. Renormalisation scheme dependence in the TSM

The flexibility in the choice of different renormalisation schemes also allows one to estimate the size of missing higher-order corrections. This is demonstrated in Fig. 3 (left) for the example of the real-triplet extended SM (TSM). The triplet mass, entering κ_λ at one-loop order, is either renormalised OS (red solid) or $\overline{\text{MS}}$ (orange dashed). Thus, the difference of the two κ_λ predictions, here as a function of the doublet-triplet coupling, gives an estimate of the two-loop corrections. In addition, we plot κ_λ in dependence of the triplet mass M_{H^+} in Fig. 3 (right). The solid (dashed) contours show the exclusion limit by the current LHC constraint (HL-LHC projections).

7. Summary

We presented the libraries anyH3 and anyBSM, which can be used to compute the trilinear Higgs self-couplings in arbitrary renormalisable QFTs at the full one-loop order. The renormalisation can be performed in an automated way in both $\overline{\text{MS}}$ and OS as well as custom schemes. We extensively tested the program using different UFO models by checking UV-finiteness, reproducing literature results and the correct decoupling behaviour — using different renormalisation schemes. We have demonstrated that large corrections to the coupling modifier κ_λ are possible in many BSM scenarios featuring mass-splittings among the BSM states. We also obtain predictions for non-zero external momenta, which can be incorporated into the computation of double-Higgs production cross sections. Finally, we showed that the library can be used to obtain an estimate of the size of unknown higher-order corrections in a fast and convenient way.

Acknowledgements H.B. acknowledges support from the Alexander von Humboldt foundation. J.B., M.G., and G.W. acknowledge support by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) under Germany’s Excellence Strategy — EXC 2121 “Quantum Universe” — 390833306. J.B. is supported by the DFG Emmy Noether Grant No. BR 6995/1-1. This work has been partially funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) — 491245950.

References

- [1] S. Kanemura, Y. Okada, E. Senaha and C.P. Yuan, *Phys. Rev.* **D70** (2004) 115002 [[hep-ph/0408364](#)].

- [2] J. Braathen and S. Kanemura, *Phys. Lett. B* **796**, 38–46 (2019) [1903.05417]; J. Braathen and S. Kanemura, *Eur. Phys. J. C* **80**, no.3, 227 (2020) [1911.11507]; J. Braathen, S. Kanemura and M. Shimoda, *JHEP* **03**, 297 (2021) [2011.07580].
- [3] ATLAS collaboration, *Constraining the Higgs boson self-coupling from single- and double-Higgs production with the ATLAS detector using pp collisions at $\sqrt{s} = 13$ TeV*, 2211.01216.
- [4] PHYSICS OF THE HL-LHC WORKING GROUP collaboration, 1902.00134.
- [5] H. Bahl, J. Braathen and G. Weiglein, *Phys. Rev. Lett.* **129** (2022) 231802 [2202.03453].
- [6] V.D. Barger, M.S. Berger, A.L. Stange and R.J.N. Phillips, *Phys. Rev.* **D45** (1992) 4128; W. Hollik and S. Penaranda, *Eur. Phys. J. C* **23** (2002) 163 [hep-ph/0108245]; A. Dobado, M.J. Herrero, W. Hollik and S. Penaranda, *Phys. Rev.* **D66** (2002) 095016 [hep-ph/0208014]; K.E. Williams and G. Weiglein, *Phys. Lett. B* **660** (2008) 217 [0710.5320]; K.E. Williams, H. Rzehak and G. Weiglein, *Eur. Phys. J. C* **71** (2011) 1669 [1103.1335]; M. Brucherseifer, R. Gavin and M. Spira, *Phys. Rev.* **D90** (2014) 117701 [1309.3140].
- [7] D.T. Nhung, M. Mühlleitner, J. Streicher and K. Walz, *JHEP* **11** (2013) 181 [1306.3926]; M. Mühlleitner, D.T. Nhung and H. Ziesche, *JHEP* **12** (2015) 034 [1506.03321]; C. Borschensky, T.N. Dao, M. Gabelmann, M. Mühlleitner and H. Rzehak, *Eur. Phys. J. C* **83** (2023) no.2, 118 2210.02104.
- [8] M. Aoki, S. Kanemura, M. Kikuchi and K. Yagyu, *Phys. Rev.* **D87** (2013) 015012 [1211.6029]; S. Kanemura, M. Kikuchi and K. Yagyu, *Nucl. Phys.* **B896** (2015) 80 [1502.07716]; S. Kanemura, M. Kikuchi and K. Yagyu, *Nucl. Phys.* **B907** (2016) 286 [1511.06211]; A. Arhrib, R. Benbrik, J. El Falaki and A. Jueid, *JHEP* **12** (2015) 007 [1507.03630]; S. Kanemura, M. Kikuchi and K. Sakurai, *Phys. Rev.* **D94** (2016) 115011 [1605.08520]; S. Kanemura, M. Kikuchi and K. Yagyu, *Nucl. Phys.* **B917** (2017) 154 [1608.01582]; S.-P. He and S.-h. Zhu, *Phys. Lett.* **B764** (2017) 31 [1607.04497]; S. Kanemura, M. Kikuchi, K. Sakurai and K. Yagyu, *Phys. Rev.* **D96** (2017) 035014 [1705.05399]; E. Senaha, *Phys. Rev.* **D100** (2019) 055034 [1811.00336]; C.-W. Chiang, A.-L. Kuo and K. Yagyu, *Phys. Rev. D* **98** (2018) 013008 [1804.02633]; J.E. Falaki, *Phys. Lett. B* **840** (2023) 137879 [2301.13773]; H. Bahl, W.H. Chiu, C. Gao, L.-T. Wang and Y.-M. Zhong, *Eur. Phys. J. C* **82** (2022) 944 [2207.04059].
- [9] S. Kanemura, M. Kikuchi, K. Sakurai and K. Yagyu, *Comput. Phys. Commun.* **233** (2018) 134 [1710.04603]; S. Kanemura, M. Kikuchi, K. Mawatari, K. Sakurai and K. Yagyu, *Comput. Phys. Commun.* **257** (2020) 107512 [1910.12769].
- [10] P. Basler and M. Mühlleitner, *Comput. Phys. Commun.* **237** (2019) 62 [1803.02846]; P. Basler, M. Mühlleitner and J. Müller, *Comput. Phys. Commun.* **269** (2021) 108124 [2007.01725].
- [11] H. Bahl, J. Braathen, M. Gabelmann and G. Weiglein, [2305.03015].
- [12] C. Degrande, C. Duhr, B. Fuks, D. Grellscheid, O. Mattelaer and T. Reiter, *Comput. Phys. Commun.* **183** (2012) 1201 [1108.2040]; L. Darmé et al., *Eur. Phys. J. C* **83** (2023) 631 [2304.09883].