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# Evidence for X(3872) in PbPb collisions and studies of its prompt production at $\sqrt{s_{\text{NN}}} = 5.02$ TeV

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## Abstract

The first evidence for X(3872) production in relativistic heavy ion collisions is reported. The X(3872) production is studied in lead-lead (PbPb) collisions at a center-of-mass energy of  $\sqrt{s_{\text{NN}}} = 5.02$  TeV per nucleon pair, using the decay chain  $X(3872) \rightarrow J/\psi \pi^+ \pi^- \rightarrow \mu^+ \mu^- \pi^+ \pi^-$ . The data were recorded with the CMS detector in 2018 and correspond to an integrated luminosity of  $1.7 \text{ nb}^{-1}$ . The measurement is performed in the rapidity and transverse momentum ranges  $|y| < 1.6$  and  $15 < p_T < 50 \text{ GeV}/c$ . The significance of the inclusive X(3872) signal is 4.2 standard deviations. The prompt X(3872) to  $\psi(2S)$  yield ratio is found to be  $\rho^{\text{PbPb}} = 1.08 \pm 0.49 (\text{stat}) \pm 0.52 (\text{syst})$ , to be compared with typical values of 0.1 for pp collisions. This result provides a unique experimental input to theoretical models of the X(3872) production mechanism, and of the nature of this exotic state.

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The  $X(3872)$ , also known as  $\chi_{c1}(3872)$ , is an exotic particle that was first observed by the Belle Collaboration [1], and then confirmed and studied by other experiments at electron-positron [2, 3] and hadron colliders [4–8]. The quantum numbers of the  $X(3872)$  have been determined to be  $J^{PC} = 1^{++}$  by the CDF [9] and LHCb [10] Collaborations. However, the nature of this particle is still not fully understood and interpretations in terms of conventional charmonium (a bound state of charm-anticharm quarks),  $D^*(2010)^0 \bar{D}^0$  molecules [11], tetraquark states [12], and their admixture [13] have been proposed. The production and survival of the  $X(3872)$  in a quark-gluon plasma (QGP), a deconfined state of quarks and gluons [14, 15], is expected to depend upon the  $X(3872)$ 's internal structure [16, 17]. Thus, the recent large data set of lead-lead (PbPb) collisions at a center-of-mass energy of  $\sqrt{s_{NN}} = 5.02$  TeV per nucleon pair, delivered by the Large Hadron Collider (LHC) at CERN at the end of 2018, opened new opportunities to probe the nature of this exotic state [18–20].

It is expected that in relativistic heavy ion collisions, the formation of the QGP could enhance or suppress the production of the  $X(3872)$  particle. Coalescence mechanisms could enhance the  $X(3872)$  production yield [16, 19]. These mechanisms can be modeled via the overlap of the density matrix of the constituents in an emission source with the Wigner function of the produced particle [21]. Therefore, the enhancement of the  $X(3872)$  production in the QGP would depend on the spatial configuration of the exotic state. Moreover, a longer distance between the quarks and antiquarks that constitute the state could also lead to a higher  $X(3872)$  dissociation rate, similar to that from the mechanism of quarkonium suppression in heavy ion collisions [22]. Therefore, the study of the  $X(3872)$  state in the QGP may be used as a tool to distinguish a compact tetraquark configuration with a radius  $\sim 0.3$  fm from a molecular state with a radius greater than 1.5 fm [23]. Such a measurement would be complementary to the recent evidence for the radiative decay  $X(3872) \rightarrow \psi(2S)\gamma$  in proton-proton (pp) collisions reported by LHCb Collaboration [24], which does not support a pure  $D^*(2010)^0 \bar{D}^0$  molecular interpretation. In addition, measurements of prompt  $X(3872)$  production could provide an interesting test of the statistical hadronization model, which assumes that the produced matter is in thermodynamic equilibrium [25].

In this Letter, the first evidence for  $X(3872)$  production in PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV is reported. The PbPb sample corresponds to an integrated luminosity of  $1.7 \text{ nb}^{-1}$ . The  $X(3872)$  candidates are reconstructed through the decay chain  $X(3872) \rightarrow J/\psi \pi^+ \pi^- \rightarrow \mu^+ \mu^- \pi^+ \pi^-$ . At the LHC energies, the inclusive  $X(3872)$  yields in pp and PbPb collisions contain a significant nonprompt contribution coming from b hadron decays [8]. The nonprompt  $X(3872)$  component is related to the medium-modified beauty hadron production in heavy ion collisions, which is out of the scope of this paper. Here, we focus on the prompt component, from charm quark fragmentation, for which the ratio  $\rho^i$  ( $i$  is pp or PbPb) between the corrected yields of  $X(3872)$  and  $\psi(2S)$  mesons (where the  $\psi(2S)$  is reconstructed with the same final-state particles in order to reduce systematic uncertainties) is presented:

$$\rho^i = \frac{N_i^{X(3872) \rightarrow J/\psi \pi \pi}}{N_i^{\psi(2S) \rightarrow J/\psi \pi \pi}}. \quad (1)$$

The ratios in pp and PbPb collisions are connected to the nuclear modification factors  $R_{AA}^{X(3872)}$  and  $R_{AA}^{\psi(2S)}$  (the meson yield ratio in nucleus-nucleus and pp interactions normalized by the number of inelastic nucleon-nucleon collisions) via the following relation:

$$\rho^{\text{PbPb}} = \rho^{\text{pp}} \frac{R_{AA}^{X(3872)}}{R_{AA}^{\psi(2S)}}. \quad (2)$$

The central feature of the CMS detector is a superconducting solenoid, which provides a magnetic field of 3.8 T. Within the solenoid volume there are a silicon pixel and strip tracker that measures the trajectories of the charged particles in the pseudorapidity range  $|\eta| < 2.5$ , a lead-tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter. For charged particles of transverse momentum  $1 < p_T < 10 \text{ GeV}/c$  and  $|\eta| < 1.4$ , the typical track resolutions are 1.5% in  $p_T$  and 25–90 (45–150)  $\mu\text{m}$  in the transverse (longitudinal) impact parameter [26]. Muons are measured in the range  $|\eta| < 2.4$  with a muon spectrometer which is located outside the solenoid and uses three technologies: drift tubes, cathode strip chambers, and resistive-plate chambers. The muon reconstruction algorithm starts by finding tracks in the muon spectrometer, which are then matched with tracks reconstructed in the silicon tracker to form muon candidates. Matching muons to tracks measured in the tracker results in a relative  $p_T$  resolution for muons with  $20 < p_T < 100 \text{ GeV}/c$  of 1.3–2.0% in the barrel ( $|\eta| < 1.2$ ) and better than 6% in the endcaps ( $1.6 < |\eta| < 2.4$ ) [27]. The hadron forward (HF) calorimeter uses steel as an absorber and quartz fibers as sensitive material. The two halves of the HF calorimeters are located 11.2 m away from the interaction point, one on each end, together providing coverage in the range  $3.0 < |\eta| < 5.2$ . In this analysis, the HF information is used for performing an offline event selection and determining centrality (the degree of overlap between the two colliding nuclei). The results refer to the collisions with 0–90% centrality, i.e., the top 90% events based on the total transverse energy deposition in both HF detectors [28], which corresponds to the 90% of collisions having the largest overlap of the two nuclei. A detailed description of the CMS experiment and coordinate system can be found in Ref. [29].

Events of interest were selected in real-time using the CMS two-tiered trigger system: the first level (L1), composed of custom hardware processors [30], and the high-level trigger (HLT), consisting of a farm of processors running a version of the full event reconstruction software optimized for fast processing [31]. The selection required the presence of two muon candidates, with at least one muon reconstructed in the outer muon spectrometer, and one muon reconstructed using information from both the outer muon spectrometer and the inner tracker. The dimuon candidate invariant mass is required to be  $1 < m_{\mu\mu} < 5 \text{ GeV}/c^2$ . For the offline analysis, events have to pass a set of selection criteria designed to reject events from background processes (beam-gas collisions, beam scraping events and electromagnetic interactions) as described in Ref. [32]. Events are required to have at least one reconstructed primary interaction vertex formed by two or more tracks, with a distance from the center of the nominal interaction point of less than 15 cm along the beam axis. The shapes of the clusters in the pixel detector have to be compatible with those expected from particles produced by a PbPb collision [33]. In order to select hadronic collisions, the PbPb events are also required to have at least two towers (i.e., a geometrically defined group of calorimeter cells) in each of the HF detectors with total energy deposits of more than 4 GeV per tower. This analysis is restricted to events within centrality 0–90%, for which the hadronic event selection is fully efficient.

Dedicated PbPb  $X(3872)$  and  $\psi(2S)$  Monte Carlo (MC) simulated samples were generated in order to estimate the acceptance and selection efficiencies, to study the background components, and to evaluate systematic uncertainties. The PYTHIA 8 v212 [34] Tune CP5 [35] was used to generate the  $X(3872)$  and  $\psi(2S)$  signals at  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ . Since the  $X(3872)$  cannot be generated by PYTHIA, the  $\chi_{c1}(1P)$  particle is used, with a modified mass of  $3.8716 \text{ GeV}/c^2$  [36]. The  $\chi_{c1}(1P)$  has the quantum numbers  $J^{PC} = 1^{++}$ , identical to those of the  $X(3872)$ . The EVTGEN 1.3.0 package [37] was used to force the  $X(3872)$  and  $\psi(2S)$  to decay into  $J/\psi \pi^+ \pi^-$  (assuming the  $\rho$  resonance dominates the pion pair spectrum [7, 8]), followed by the  $J/\psi$  meson decaying into two muons. Final-state radiation was generated using PHOTOS 2.0 [38]. The  $\chi_{c1}(1P) \rightarrow J/\psi \pi^+ \pi^-$  decay is generated with EVTGEN. The samples with prompt (fragment-

ing in charm quarks) and nonprompt (originating from b hadron decays)  $\psi(2S)/X(3872)$  production are generated separately. Each PYTHIA event is embedded in a PbPb collision event generated with HYDJET 1.8 [39], which is tuned to reproduce global event properties such as the charged-hadron  $p_T$  spectrum and particle multiplicity.

The  $X(3872)$  signal is extracted in the following steps. Each muon candidate must be matched to a triggered muon and have  $p_T^\mu > 3.5 \text{ GeV}/c$  in the interval  $|\eta^\mu| < 1.2$ ,  $p_T^\mu > (5.47 - 1.89|\eta^\mu|) \text{ GeV}/c$  in the interval  $1.2 < |\eta^\mu| < 2.1$ , or  $p_T^\mu > 1.5 \text{ GeV}/c$  in the forward region  $2.1 < |\eta^\mu| < 2.4$ . Two muons of opposite sign, with an invariant mass within  $\pm 150 \text{ MeV}/c^2$  of the world-average  $J/\psi$  meson mass [36] are selected to reconstruct a  $J/\psi$  candidate. The opposite-sign muon pairs are fitted with a common vertex constraint and are kept if the  $\chi^2$  probability of the fit is greater than 1%, thus reducing the background from charm and beauty hadron semileptonic decays. The  $X(3872)$  and  $\psi(2S)$  candidates are built by combining the  $J/\psi$  candidates with two additional tracks, which have  $p_T > 0.9 \text{ GeV}/c$ ,  $|\eta| < 2.4$  and pass a high purity selection, assumed to be produced by two pions. Then, a kinematic fit to the  $J/\psi \pi^+ \pi^-$  system is performed, requiring that the four tracks originate from a common vertex and forcing the mass of the dimuon pair to be equal to the nominal  $J/\psi$  mass [36]. The selection is further optimized, using a boosted decision tree (BDT) algorithm [40], in order to maximize the statistical significance of  $X(3872)$ . The  $X(3872)$  decay vertex probability, the radial distance between the pion and the  $J/\psi$  candidate momentum vectors ( $\sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ , where  $\phi$  is the azimuthal angle), and the  $p_T$  of each pion are used in the BDT algorithm to distinguish signal and combinatorial background formed by random combinations of tracks. The  $X(3872)$  signal samples are taken from simulation, while the background samples for the multivariate training are taken from data sidebands of the  $X(3872)$  meson invariant mass. The BDT selection determined in this way is applied to entries in the whole invariant mass range from  $3.62$  to  $4 \text{ GeV}/c^2$  in data.

The raw inclusive yields of  $X(3872)$  and  $\psi(2S)$  are extracted by an extended unbinned maximum-likelihood fit. A double-Gaussian function with a common mean but independent widths is used to model the signal component for each of the  $X(3872)$  and  $\psi(2S)$  peaks. This was preferred to a single-Gaussian or a Breit-Wigner function since it described better (i.e., superior  $\chi^2$  of the fit) the signal shape in MC simulations. For describing the combinatorial background, mostly produced by the random combination of a  $J/\psi$  candidate with tracks that are not coming from  $X(3872)$  or  $\psi(2S)$  decay, a 4th-order polynomial is used, which gives the best fit in terms of  $\chi^2$  per degrees of freedom and stability during all studies. For the signal, only the magnitude of the two peaks is left free in the fit, the rest (the mean and widths of the two Gaussian functions, as well as their relative contribution to the signal yield in either  $X(3872)$  or  $\psi(2S)$  peaks) are set to the values derived from simulation. The five parameters of the combinatorial background are all allowed to float. The invariant mass range considered for the fit is  $3.62$  to  $4 \text{ GeV}/c^2$ . The invariant mass fits for both the inclusive and nonprompt samples, with BDT selection optimized for  $X(3872)$ , are shown in Fig. 1. The significance of the inclusive  $X(3872)$  signal against background-only hypothesis is 4.2 standard deviations. This significance is calculated as the square root of the logarithm of the profile likelihood ratio where the signal is zero. It considers all the systematic uncertainties related to modeling the signal component of the  $X(3872)$  peak, as detailed below.

The contribution from b hadron decays is subtracted from the inclusive result using the “pseudo-proper” decay length  $l_{xy}$ , defined as the distance in the transverse plane,  $L_{xy}$ , between the vertex formed by the 4-tracks and the closest reconstructed primary vertex in the  $z$  direction (along the beam axis), corrected by the transverse Lorentz boost of the candidate:  $l_{xy} = L_{xy} m_{J/\psi \pi \pi} c / |\vec{p}_T|$ . The prompt-component fraction ( $f_{\text{prompt}}$ ) is estimated using a cutoff-

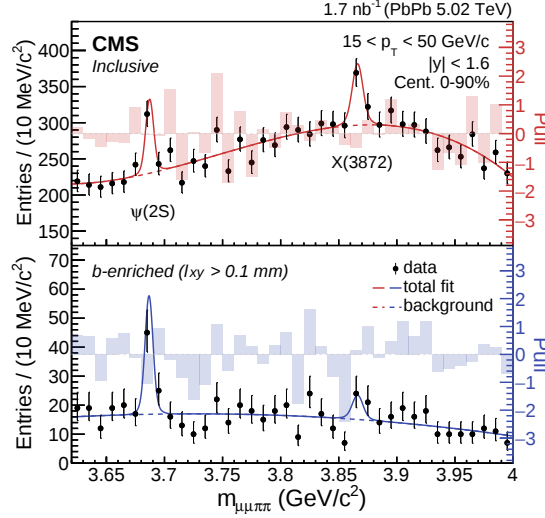


Figure 1: Invariant mass distribution of  $m_{\mu\mu\pi\pi}$  in PbPb collisions, for the inclusive (upper) and b-enriched (bottom) samples. The vertical lines represent statistical uncertainties in the data. The results of the unbinned maximum-likelihood fits for the signal+background and background alone are also shown by the solid and dashed lines, respectively. The pull distribution is represented by the shaded bars.

based method in the following way. Since the  $l_{xy}$  of the prompt component was found in MC studies to be smaller than 0.1 mm, the prompt fraction,  $f_{\text{prompt}}$ , can be derived from: i) the raw inclusive yield,  $N_{\text{incl}}$ , obtained from the fit to the invariant mass distributions of all candidates, shown in Fig. 1, and ii)  $N_{\text{b-enr}}$ , the “b-enriched” yield, obtained from a fit to the invariant mass distribution only containing candidates that passed the selection  $l_{xy} > 0.1$  mm, also shown in Fig. 1. In addition,  $N_{\text{b-enr}}$  has to be corrected to account for nonprompt candidates, with  $l_{xy} < 0.1$  mm, that have been missed:  $f_{\text{nonprompt}}^{\text{b-enr}} = N_{\text{nonprompt}}(l_{xy} > 0.1 \text{ mm}) / N_{\text{nonprompt}}$ . The correction was obtained from simulation. The raw prompt fraction is then calculated as:

$$f_{\text{prompt}} = 1 - \frac{N_{\text{b-enr}} / f_{\text{nonprompt}}^{\text{b-enr}}}{N_{\text{incl}}}, \quad (3)$$

separately for X(3872) and  $\psi(2S)$  states. The corrected yield of the prompt component can then be derived as:

$$N^{i \rightarrow J/\psi \pi \pi} = N_{\text{raw}}^i f_{\text{prompt}}^i / (\alpha \epsilon_{\text{reco}} \epsilon_{\text{sel}})^i, \quad (4)$$

where  $i$  is X(3872) or  $\psi(2S)$ ,  $\alpha$  is the acceptance,  $\epsilon_{\text{reco}}$  is the candidate reconstruction efficiency and  $\epsilon_{\text{sel}}$  is the candidate selection efficiency. Since the two states are reconstructed in the same decay channel and are relatively close in mass, their corresponding  $\alpha \epsilon_{\text{reco}}$  values are similar. The choice of the BDT optimization criteria results in  $\epsilon_{\text{sel}}$  being higher for the X(3872) than for the  $\psi(2S)$ .

The measurement of  $\rho^{\text{PbPb}}$  is affected by several sources of systematic uncertainty, arising from the candidate selection, invariant mass fit, and efficiency corrections. To estimate the systematic uncertainty associated with the BDT selection, the BDT cutoff values are varied within the range that allows a robust invariant mass fit procedure, and the maximum difference of the final  $R$  value from the nominal result (40%) is quoted as the systematic uncertainty. The relatively large  $R$  uncertainty associated with BDT cutoff value is the convolution of two causes: the BDT variables distribution differences in data and MC samples for the X(3872) meson, and the statistical limitation of the signal in data.

The uncertainty in the invariant mass fit (8.0%) is calculated by adding in quadrature the maximum deviations from the nominal result to that found using two alternative fitting functions for both signal and background. For the signal, one variation consists of using a triple-Gaussian function, while for the other the signal width of the nominal fit is allowed to float to account for the resolution difference between data and MC. Other choices for the signal shape (e.g., one-Gaussian function) were not considered because of their poor-quality fits. For the background, the fit function is changed once to a third-order polynomial (as an exponential function or lower-order polynomials could not describe the data), and the fit range is also changed from  $3.62\text{--}4\text{ GeV}/c^2$  to  $3.62\text{--}3.9\text{ GeV}/c^2$  to exclude the right-hand shoulder.

The efficiency corrections obtained from simulation are sensitive to how well the  $p_T$  spectrum of the  $X(3872)$  and  $\psi(2S)$  candidates is modeled. The uncertainty related to the simulated  $p_T$  shape is evaluated by comparing the reconstruction and selection efficiencies calculated using the default PYTHIA MC sample, with another MC sample in which the  $p_T$  distributions of  $X(3872)$  and  $\psi(2S)$  are tuned to reproduce the extracted  $X(3872)$  and  $\psi(2S)$   $p_T$  and  $y$  spectra obtained in data, by performing mass fits in bins of  $X(3872)$  and  $\psi(2S)$   $p_T$  and  $y$ . The  $p_T$  and  $y$  spectra of the alternative MC samples are allowed to vary within the statistical uncertainties in data. The mean of the differences between efficiencies from the alternative MC samples and the default PYTHIA MC due to the variation of  $p_T$  and  $y$  spectra, which is 13%, is quoted as the systematic uncertainty.

The uncertainties in the trigger efficiency, in the muon reconstruction and identification are evaluated using single muons from  $J/\psi$  meson decays in both simulated and collision data, with the tag-and-probe method [27, 41]. This combined uncertainty is found to be negligible, below 1%. Scale factors, calculated as the ratio of data to simulated efficiencies as a function of  $p_T^\mu$  and  $\eta^\mu$ , are applied to each dimuon pair on a muon-by-muon basis. The uncertainties of the scaling factors from tag-and-probe studies are quoted as systematic uncertainties.

To estimate the uncertainty in the prompt fraction arising from potential differences between the resolution in data and simulation, a template fit of the  $l_{xy}$  distribution in data is performed using prompt and nonprompt  $l_{xy}$  templates from simulation. Data are binned in  $l_{xy}$ , and an invariant mass fit is performed to extract the inclusive yield in each  $l_{xy}$  bin. This background-subtracted  $l_{xy}$  distribution is then fitted using a two component fit, which includes the prompt and nonprompt  $l_{xy}$  templates from simulation. The widths of the simulated DCA distributions are varied by a floating scale factor, and the best simulated smearing scale factor to match data is determined by minimizing the  $\chi^2$  of the two-component fit. The difference between the ratio of the prompt fractions of  $X(3872)$  to  $\psi(2S)$  using the template fit method and the nominal result (8.1%) is quoted as a systematic uncertainty.

When calculating the uncertainties in the ratio of the acceptance-corrected yields of prompt  $X(3872)$  production over prompt  $\psi(2S)$  production, the uncertainties of  $X(3872)$  and  $\psi(2S)$  yields are assumed to be independent except for the systematic uncertainties from muon reconstruction, efficiencies, and prompt fractions.

The ratio  $\rho^{\text{PbPb}}$  is shown in Fig. 2, together with  $\rho^{\text{PP}}$  measured as a function of  $p_T$ . The pp data were measured at  $\sqrt{s} = 7$  and 8 TeV, in the  $|y| < 1.2$  and  $|y| < 0.75$  intervals, respectively [7, 8, 10]. The 7 TeV result was derived using the CMS Collaboration published ratio of the inclusive yields [7] and prompt fractions [7, 42]. From Fig. 2 it is clear that the prompt  $\rho^{\text{PP}}$  does not depend significantly on collision energy or rapidity. In pp collisions at  $\sqrt{s} = 8$  TeV, in the kinematic range of  $16 < p_T < 22\text{ GeV}/c$ , the  $\rho^{\text{PP}}$  measured by the ATLAS Collaboration is  $0.106 \pm 0.008$  (stat)  $\pm 0.004$  (syst) [8]. This is to be compared to the  $\rho^{\text{PbPb}}$  measured in this Letter,  $\rho^{\text{PbPb}} = 1.08 \pm 0.49$  (stat)  $\pm 0.52$  (syst).

In the interval  $15 < p_T < 20 \text{ GeV}/c$ , the yield of the prompt  $\psi(2S)$  in PbPb collisions was reported to be significantly suppressed with respect to pp collisions,  $R_{AA}^{\psi(2S)} = 0.142 \pm 0.061 \text{ (stat)} \pm 0.020 \text{ (syst)}$  [43]. This leads, using Eq. (2), to an  $R_{AA}^{X(3872)}$  central value larger than 1 (i.e., enhancement of the prompt X(3872) yield in PbPb compared to pp collisions). However, the uncertainties are such that  $R_{AA}^{X(3872)}$  is compatible with 1 within  $\sim 1$  standard deviation, and with  $R_{AA}^{\psi(2S)}$  within  $\sim 2$  standard deviations. Thus, it is possible that in PbPb collisions, the prompt X(3872) yield has either no suppression with respect to pp collisions, or as much suppression as the  $\psi(2S)$  state. The much larger data sample expected in Run 3 at the LHC will answer whether  $R_{AA}^{X(3872)}$  is different from  $R_{AA}^{\psi(2S)}$  and significantly above unity. It may answer whether the  $\psi(2S)$  meson production (a bound state of a c and  $\bar{c}$  quarks, with  $r \sim 0.9 \text{ fm}$ ), is affected differently by the medium produced in PbPb collisions than the X(3872) state (that could be made of c,  $\bar{c}$ , u, and  $\bar{u}$  quarks, with a radius of  $r \sim 0.3 \text{ fm}$  or  $r > 1.5 \text{ fm}$ ), the difference in both size and quark content playing a role into their production mechanisms. It will also answer whether the X(3872) prompt state production is different in PbPb collisions compared to pp collisions. The question whether X(3872) is a tetraquark or a molecule cannot yet be answered, because of the statistical limitation of the data, and the disagreement among models. For example, while the AMPT transport model [16] predicts  $R_{AA}^{\text{molecule}} \gg R_{AA}^{\text{tetraquark}}$  with  $R_{AA}^{\text{molecule}} > 1$ , the TAMU transport model [17] predicts  $R_{AA}^{\text{molecule}} \sim R_{AA}^{\text{tetraquark}}/2$  (albeit, considering only the X(3872) from regeneration processes).

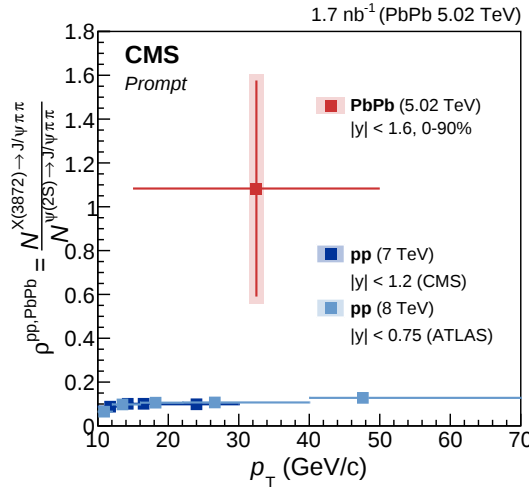


Figure 2: The yield ratio  $\rho^{\text{PbPb}}$  of prompt X(3872) over  $\psi(2S)$  production in PbPb collisions at  $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ . The vertical bars (boxes) correspond to statistical (systematic) uncertainties. The yield ratios  $\rho^{\text{pp}}$  of prompt X(3872) over  $\psi(2S)$  production in pp collisions at  $\sqrt{s} = 8 \text{ TeV}$ , measured by ATLAS [8], and at  $\sqrt{s} = 7 \text{ TeV}$ , measured by CMS [7] are also shown.

In summary, the first evidence for X(3872) production in heavy ion collisions is presented using lead-lead collisions at a center-of-mass energy of  $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$  per nucleon pair, recorded with the CMS detector. The X(3872) state is reconstructed using the decay chain  $X(3872) \rightarrow J/\psi \pi^+ \pi^- \rightarrow \mu^+ \mu^- \pi^+ \pi^-$ . The measurement is performed for transverse momentum values of the X(3872) of  $15 < p_T < 50 \text{ GeV}/c$  and rapidity  $|y| < 1.6$ . The significance of the inclusive X(3872) signal is 4.2 standard deviations. The ratio  $\rho^{\text{PbPb}}$  between the prompt X(3872) and  $\psi(2S)$  yields times their branching fractions into  $J/\psi \pi^+ \pi^-$  is found to be  $1.08 \pm 0.49 \text{ (stat)} \pm 0.52 \text{ (syst)}$ , to be compared with typical values of 0.1 for pp collisions. This result provides a unique experimental input to theoretical models of the X(3872) production mechanism, and of

the nature of this exotic state.

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