

Strong gravitational radiation from a simple dark matter model

Camilo Garcia Cely, DESY



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Beyond General Relativity, Beyond Cosmological Standard Model

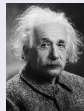
1st July, 2019

In collaboration with Iason Baldes

Based on JHEP 1905 (2019) 190

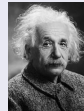
Gravitational Waves (GWs)

- Predicted by Poincaré (1905).



This talk

Gravitational Waves (GWs)

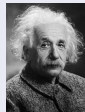


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$$\square h_{\mu\nu} = -16\pi G T_{\mu\nu}$$

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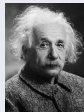
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- First-order phase transitions in the Early Universe produce GWs. Witten (1984).

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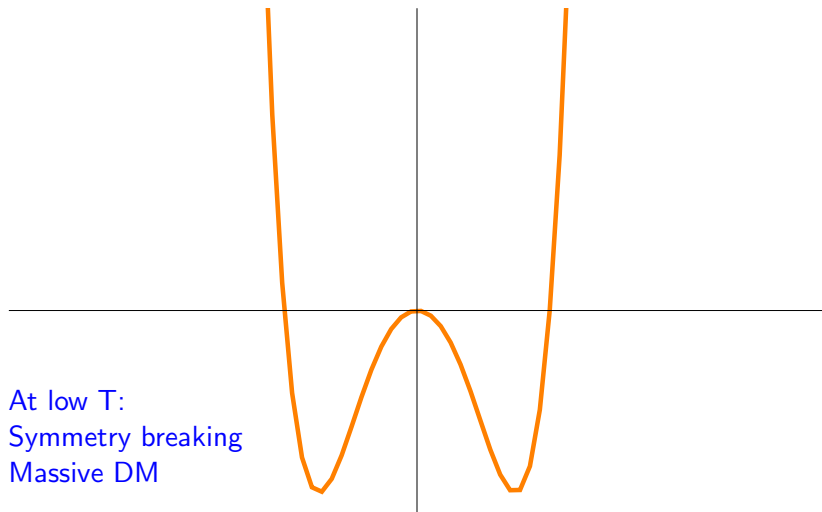
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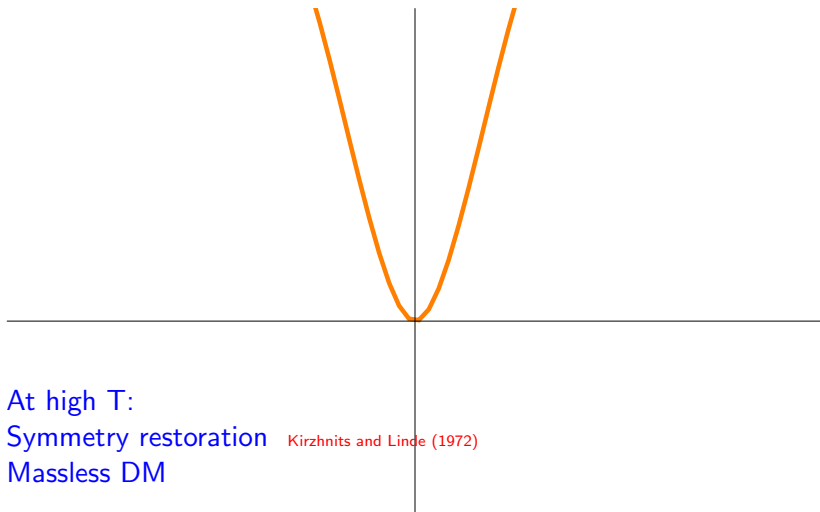
This talk

- Discuss a **simple scenario** where dark matter in the Early Universe undergoes a first-order phase transition.

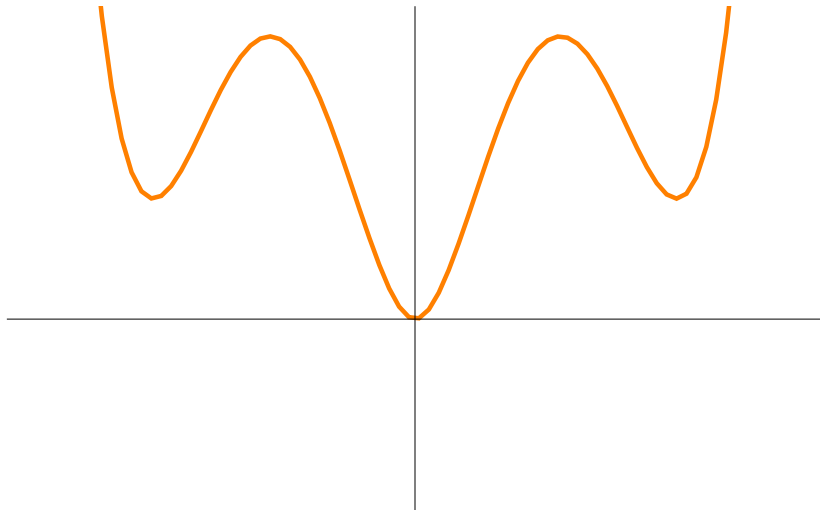
First-order phase transition



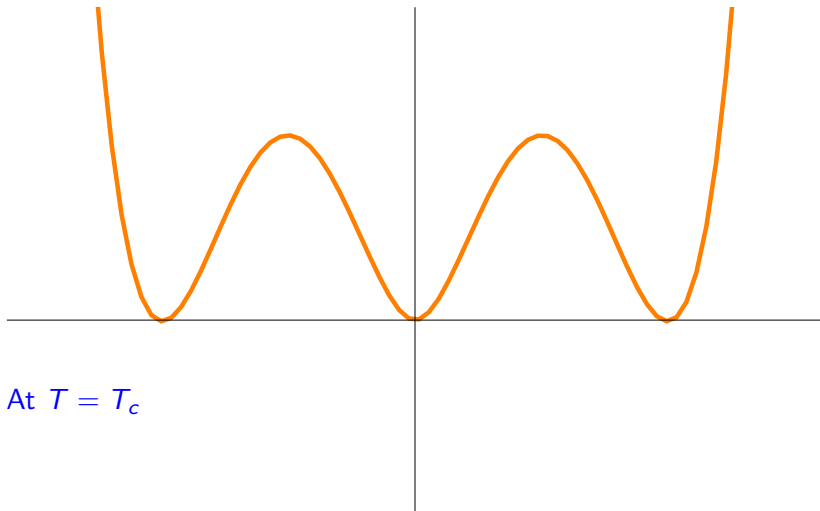
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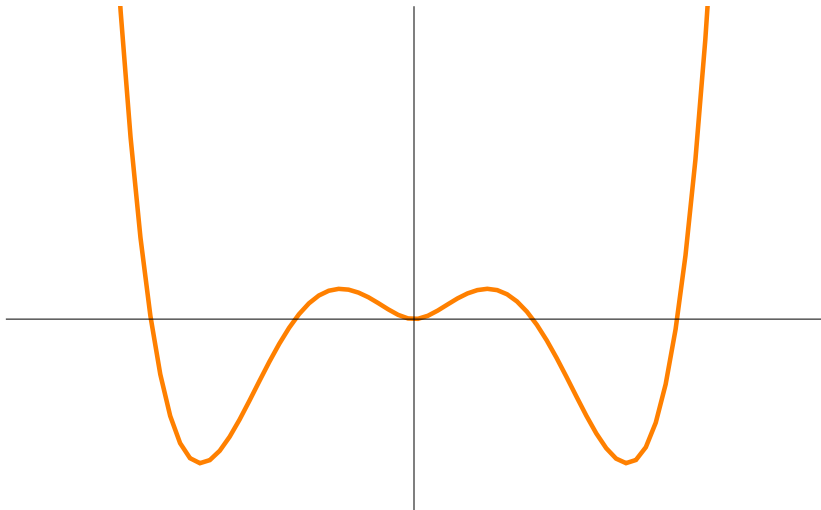
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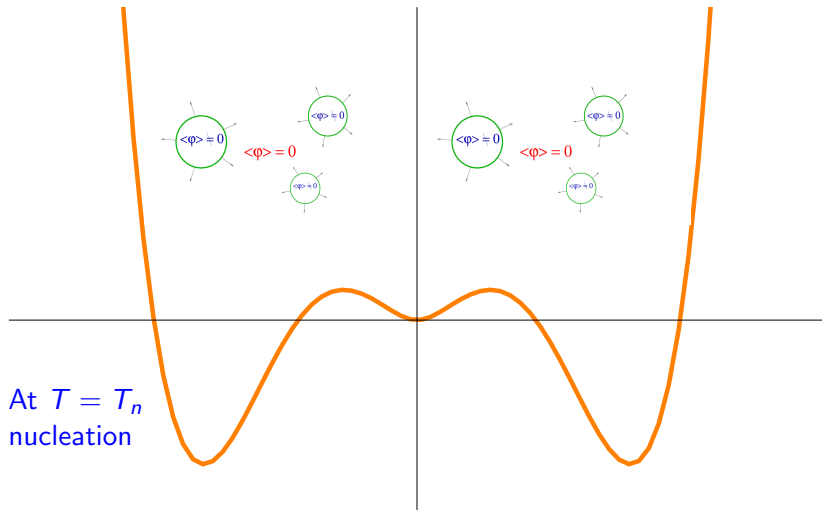
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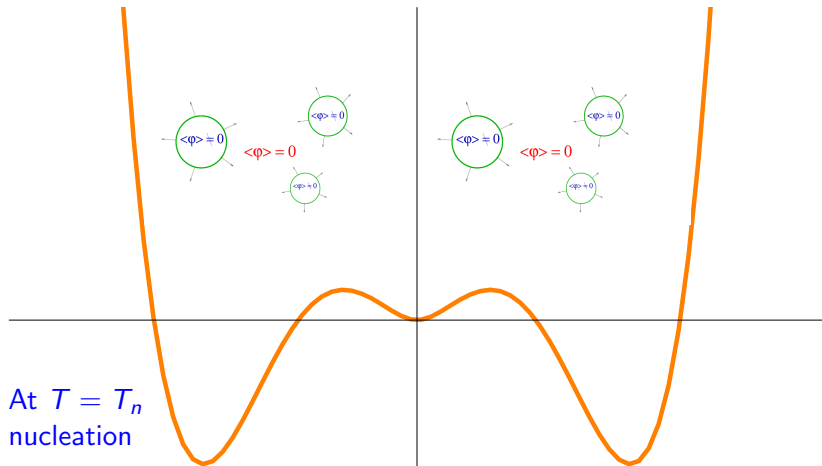
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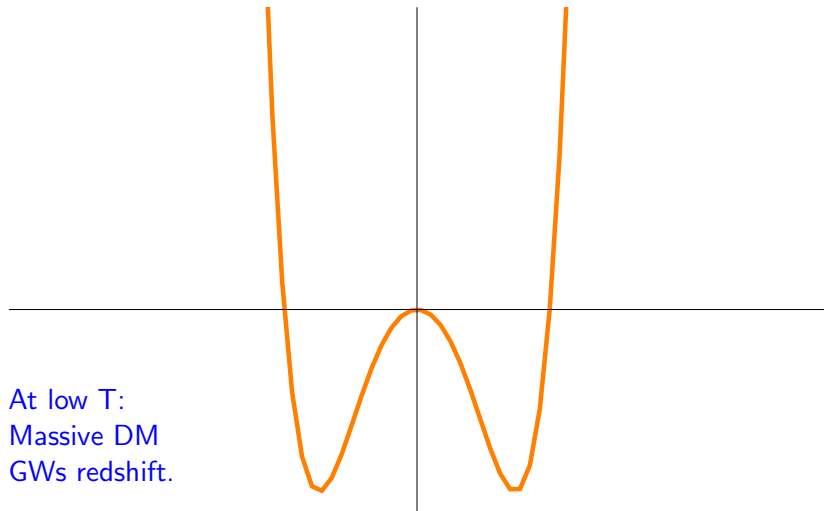


First-order phase transition



This produces produces gravitational waves E. Witten (1984)

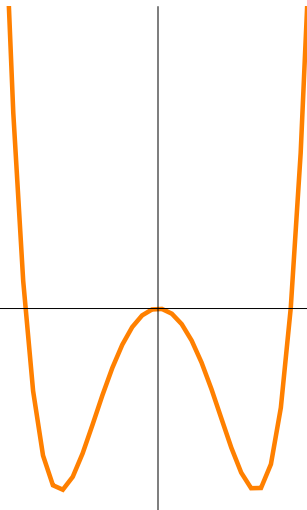
First-order phase transition



First-order phase transition

$$m_{\text{DM}} \sim 1 \text{ TeV}$$
$$\rightarrow f \sim 10^{-2} \text{ Hz}$$

At low T:
Massive DM
GWs redshift.



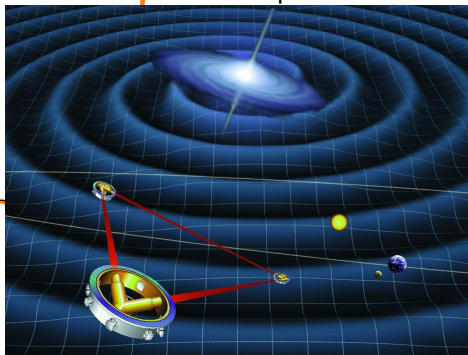
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Laser Interferometer Space Antenna



Caprini et al (2015)

- A few studies along these lines:

P. Schwaller, PRL 115 (2015), Baldes JCAP (2017), Chao et al, JHEP (2017)
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- Unless the dark matter properties are **closely** related to the phase transition with the latter being determined by a **handful** of parameters.
- **Objective:** Find a dark matter model that overcomes these difficulties.

Mimic the EW sector: DM as gauge bosons

Field	$SU(3)$	$SU(2)$	$U(1)_Y$	$SU(2)_D$
H	1	2	$\frac{1}{2}$	1
H_D	1	1	0	2

Local $SU(2)_D$ \rightarrow Global $SO(3)$
 Gauge Fields A'_μ \rightarrow Massive Fields A_μ
 Dark doublet H_D \rightarrow Higgs-like h_D

Hambye (JHEP 2009)

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High temperatures

→

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Global $SO(3)$ Massive Fields A_μ Higgs-like h_D

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Hambye (JHEP 2009) Phase transition in the Early Universe!!!!!!!!!!!!

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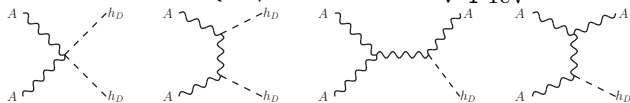
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- DM mass

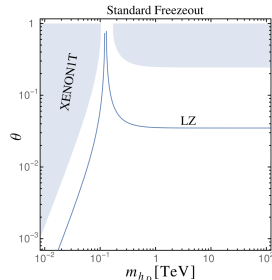
Four parameters:

- Mass of the extra scalar
- DM mass
- DM coupling which is fixed
by the relic density (via freeze-out):

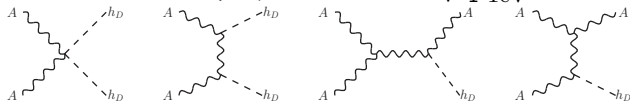
$$\left\{ \begin{array}{l} g_D \approx 0.9 \times \sqrt{\frac{m_A}{1 \text{ TeV}}} \\ v_\eta \approx 2.2 \text{ TeV} \times \sqrt{\frac{m_A}{1 \text{ TeV}}} \end{array} \right.$$


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- Mixing angle constrained by direct detection: $\theta \lesssim 0.1$.

GW spectrum

Phase transition parameters

$$T_n = 0.48 \text{ TeV}$$

$$\eta_n = 3.8 \text{ TeV}$$

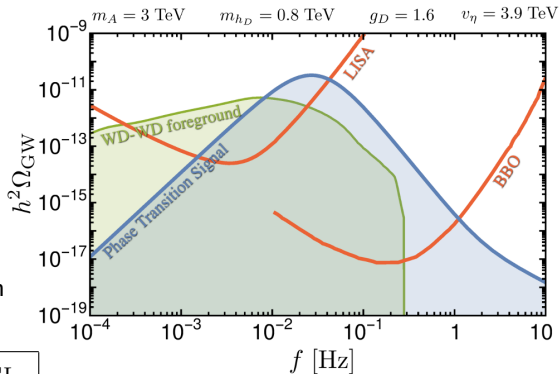
$$\alpha = 0.29, \sim (\text{latent heat})$$

$$\beta/H = 290 \sim (\text{fq. scale})$$

Simulations give Ω_{GW} from them

Caprini et al (2015)

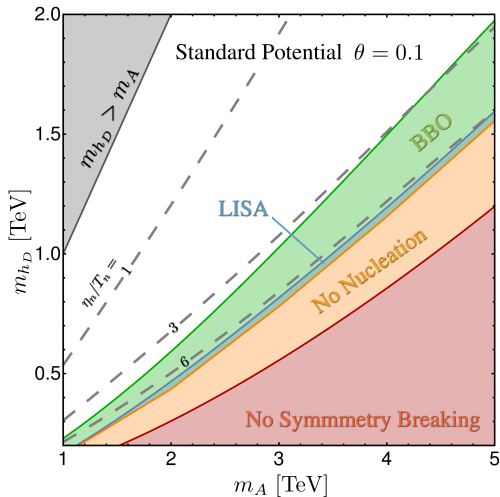
	SNR	SNR _{FGL}
LISA	15	1.8
BBO	3.7×10^5	2.3×10^3



Baldes, CGC 2018

Parameter space for $\text{SNR} > 5$.

$$\text{SNR} = \sqrt{t_{\text{obs}} \int \left[\frac{h^2 \Omega_{\text{GW}}(f)}{h^2 \Omega_{\text{sens}}(f)} \right]^2 df}$$



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Radiative effects break the $SU(2)_D$ symmetry Coleman-Weinberg (1973)

λ_2 runs to negative values.

Baldes, CGC 2018

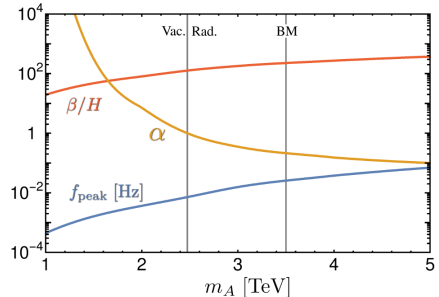
- Only one free parameter after taking the relic density into account.
- Scale-invariant potential
→ strong signal.
- There is a large amount of supercooling

See Marek Lewicki's talk

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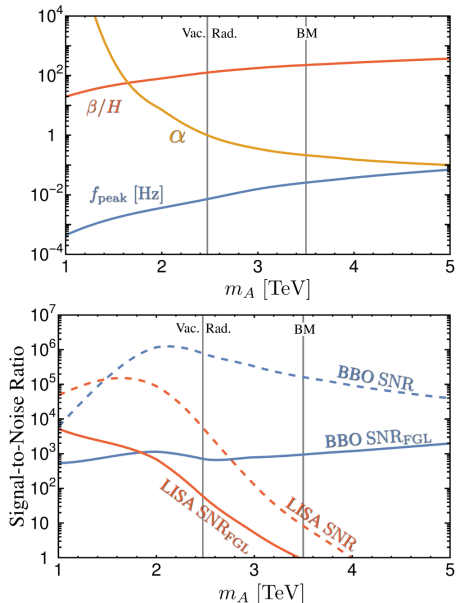
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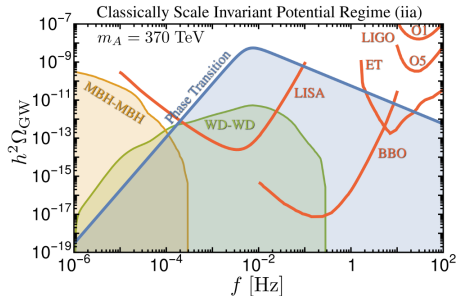


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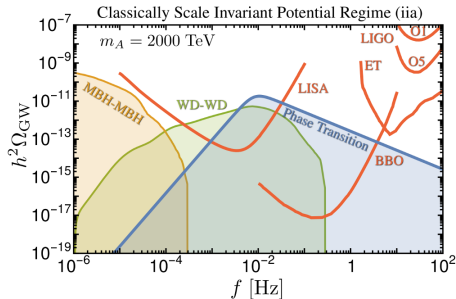
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Dark Sector Parameters	m_A	370	TeV
	m_{h_D}	59	TeV
	v_η	780	TeV
	g_D	0.95	-
	θ	10^{-9}	-
Phase Transition	T_n	2.6	GeV
	$T_{\text{infl.}}$	43	TeV
	T_{RH}	13	TeV
	η_n	$\simeq v_\eta$	-
	α	10^{16}	-
SNR	β/H	6.7	-
	LISA	10^4	-
	LISA(FGL)	270	-
	BBO	10^8	-
	BBO(FGL)	10^7	-



Dark Sector Parameters	m_A	2000	TeV
	m_{h_D}	330	TeV
	v_η	4100	TeV
	g_D	0.98	-
	θ	10^{-11}	-
Phase Transition	T_n	32	GeV
	$T_{\text{infl.}}$	230	TeV
	T_{RH}	1.0	TeV
	η_n	$\simeq v_\eta$	-
	α	10^{15}	-
SNR	β/H	7.1	-
	LISA	44	-
	LISA(FGL)	1.0	-
	BBO	10^5	-
	BBO(FGL)	10^5	-

Conclusions

- We have explored the possibility of DM from a hidden $SU(2)_D$ gauge group. This implies a phase transition that will result in detectable gravitational waves.
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Thanks for your attention