

Numerical Implementation of Harmonic Polylogarithms to Weight $w = 8$

J. Ablinger^a, J. Blümlein^b, M. Round^{a,b}, and C. Schneider^a,

^a *Research Institute for Symbolic Computation (RISC),
Johannes Kepler University, Altenbergerstraße 69, A-4040 Linz, Austria*

^b *Deutsches Elektronen-Synchrotron, DESY,
Platanenallee 6, D-15738 Zeuthen, Germany*

Abstract

We present the FORTRAN-code HPOLY.f for the numerical calculation of harmonic polylogarithms up to $w = 8$ at an absolute accuracy of $\sim 10^{-15}$ or better. Using algebraic and argument relations the numerical representation can be limited to the range $x \in [0, \sqrt{2} - 1]$. We provide replacement files to map all harmonic polylogarithms to a basis and the usual range of arguments $x \in]-\infty, +\infty[$ to the above interval analytically. We also briefly comment on a numerical implementation of real valued cyclotomic harmonic polylogarithms.

Program Summary

Title of program: HPOLY.f

Version: 1.00 September 2018

Catalogue number:

Program obtainable from: <https://www-zeuthen.desy.de/~blumlein/>

e-mail: Johannes.Bluemlein@desy.de

Licensing provisions: CC by NC 3.0

Programming Language: Fortran

Supplementary Material: none

Journal Reference of previous version: none

Does the new version supersede the previous version?: no

Reasons for the new version: none

Summary of revisions: no revisions

Solution method: Bernoulli-improvement of series expansions, Ref. [1].

Additional comments including Restrictions and Unusual features: none

Computers: all

Operating system: all

Processor used: Intel(R) Core(TM) i7-6600U CPU @ 2.60GHz. All run times given refer to this processor.

Memory required for the executable: 57 kB

No. of bits per word: 64

No. of lines in distributed program: 1291

Other programs called: none.

External files needed: (a) Input data: B51.m, B52.m, B53.m, B54.m, B55.m, B61.m, B62.m, B63.m, B64.m, B65.m, B66.m, B71.m, B72.m, B73.m, B74.m, B75.m, B76.m, B77.m, B81.m, B82.m, B83.m, B84.m, B85.m, B86.m, B87.m, B88.m, TTT2.m, TTT3.m, TTT4.m, TTT5.m, TTT6.m, TTT7.m, TTT8.m

(b) Bases: LIST1.m, LIST2.m, LIST3.m, LIST4.m, LIST5.m, LIST6.m, LIST7.m, LIST8.m, LIST.m

(c) **Mathematica** replacement files: hrel, LISTMTX.m, LISTOMX.m, LISTOOX.m,

Keywords: harmonic polylogarithms, to weight $w = 8$, numerical implementation, Nielsen integrals, classical polylogarithms.

Nature of physical problem: Harmonic polylogarithms form key building blocks in the analytic calculation of many multi-loop and multi-leg calculations in high energy physics and other disciplines of modern physics. Their numerical representation is provided.

Typical running time: $\sim 10^{-2}$ sec for 1318 calls of the basic harmonic polylogarithms.

References: [1] G. 't Hooft and M.J.G. Veltman, *Scalar One Loop Integrals*, Nucl. Phys. B **153** (1979) 365.

1 Introduction

Many analytic higher order and multi-leg calculations result in representations based on harmonic polylogarithms (HPLs) [2–4]. This function space has been introduced in Ref. [5]. It is formed by the Volterra-iterative integrals over the alphabet

$$f_{-1}(x) = \frac{1}{1+x}, \quad f_0(x) = \frac{1}{x}, \quad f_1(x) = \frac{1}{1-x}, \quad (1)$$

with

$$H_{b,\vec{a}}(x) = \int_0^x dy f_b(y) H_{\vec{a}}(y), \quad H_{\emptyset} = 1, \quad (2)$$

$$H_{\underbrace{0,\dots,0}_n}(x) = \frac{1}{n!} H_0^n(x). \quad (3)$$

By definition all HPLs are zero for $x = 0$. The harmonic polylogarithms are generalizations of the classical polylogarithms [6, 7] and the Nielsen integrals [8–10]. They obey algebraic [11] and argument relations [5]. Furthermore, they are dual structures, by the Mellin transform

$$\mathbf{M}[f(x)](N) = \int_0^1 dx x^{N-1} f(x), \quad (4)$$

to the nested harmonic sums [12, 13]

$$S_{b,\vec{a}}(N) = \sum_{k=1}^N \frac{(\text{sign}(b))^k}{k^{|b|}} S_{\vec{a}}(k), \quad S_{\emptyset} = 1, \quad b, a_i \in \mathbb{Z} \setminus \{0\}, N \in \mathbb{N} \setminus \{0\}. \quad (5)$$

In the past numerical representations of a series of classical polylogarithms and Nielsen integrals were given in [9] and have been in wide use for decades. The generalization to harmonic polylogarithms has been performed in [14] up to weight $\mathbf{w} = 4$ and in a private version to $\mathbf{w} = 5$. An implementation using the package **Ginac** [15] has been given in [16]. The **Mathematica** packages **HPL** [17] and **HarmonicSums** [18–22] do also allow the numerical evaluation of harmonic polylogarithms. There is also an implementation in **Maple** [23]. In various applications one needs the numerical representation of higher weight harmonic polylogarithms, e.g. to weight $\mathbf{w} = 6$ in Refs. [24–26], which were given in [27]. Numerical implementations of the analytic continuation of harmonic sums to complex arguments have been given in [28, 29].

In the foreseeable future fast numerical implementations of harmonic polylogarithms are also needed for higher weights. In the present paper we present a **FORTRAN** implementation for harmonic polylogarithms up to weight $\mathbf{w} = 8$. If one compares the computational times between **FORTRAN** implementations and the ones in **Ginac** [16] or **Mathematica** [17, 18] the latter ones are somewhat slower. In many phenomenological and experimental applications users prefer to work with numerical programmes only. The other existing programmes have of course also advantages against pure numeric implementations, since they can be tuned to arbitrary precision and allow analytic functional mappings.

The number of harmonic polylogarithms grows rapidly with the weight. Therefore it is necessary to give numerical representations only after the algebraic relations have been used. Furthermore, it is sufficient to represent the harmonic polylogarithms in the core interval $x \in [0, \sqrt{2} - 1]$ since the representations for the remaining real domain can be given by functional relations.

In Section 2 we describe the argument mappings and the numerical implementation for the harmonic polylogarithms. Technical aspects of the code and its use are discussed in Section 3. In Section 4 we comment on the corresponding numerical implementation for cyclotomic harmonic polylogarithms¹ and Section 5 contains the conclusions. Numerical examples are given in Appendix A. Appendix B contains the list of auxiliary files attached to this paper to perform preparatory calculations in **Mathematica**. Others are subsidiary files used in the initialization of the **FORTRAN**-code.

¹Beyond these functions various more special functions contribute in analytic calculations in Quantum Field Theories. For recent surveys see [30–32].

2 Mapping Harmonic Polylogarithms and Numerical Implementation

In order to keep the numerical implementation as compact as possible the user is asked to perform mappings on his expressions beforehand into the main interval $x \in [0, \sqrt{2} - 1]$. This can be done in **Mathematica** using subsidiary files attached to this paper, which have been created using the **Mathematica** package **HarmonicSums** [18–22]. For this purpose we work with bases of harmonic polylogarithms for the different weights. These are listed in the files **LIST1.m**, ..., **LIST8.m** attached to this paper. The file **LIST.m** contains the combined basis. The number of basis elements per weight is given in Table 1.

weight	# of basis elements
1	3
2	3
3	8
4	18
5	48
6	116
7	312
8	810

Table 1: Number of basis elements as a function of the weight.

The number of basis elements can be calculated using a formula by Witt, cf. [11], for a 3-letter alphabet

$$N_{\text{bas}}(\mathbf{w}) = \frac{1}{\mathbf{w}} \sum_{d|\mathbf{w}} \mu\left(\frac{\mathbf{w}}{d}\right) 3^d, \quad (6)$$

where d are the divisors of $\mathbf{w}$.² Here μ denotes the Möbius function for $N \in \mathbb{N} \setminus \{0\}$

$$\mu(N) = \begin{cases} 1 & N = 1 \\ (-1)^k & \text{if } N \text{ is a factor of } k \text{ different primes} \\ 0 & \text{else.} \end{cases} \quad (7)$$

We have e.g.

$$\text{LIST1.m} = \{H[1, x], H[0, x], H[-1, x]\} \quad (8)$$

$$\text{LIST2.m} = \{H[0, 1, x], H[-1, 1, x], H[-1, 0, x]\} \quad (9)$$

$$\begin{aligned} \text{LIST3.m} = & \{H[0, 1, 1, x], H[0, 0, 1, x], H[-1, 1, 1, x], H[-1, 1, 0, x], H[-1, 0, 1, x], H[-1, 0, 0, x], \\ & H[-1, -1, 1, x], H[-1, -1, 0, x]\}. \end{aligned} \quad (10)$$

²We remark that all harmonic polylogarithms up to $\mathbf{w} = 3$ can be expressed in terms of classical polylogarithms, allowing for rational function arguments [6], see also [29], and therefore have a numerical representation using these functions.

The relations of all other HPLs are stored in the replacement list

$$\text{hrel.m} \quad [12.8 \text{ Mbyte}] \quad (11)$$

attached to this paper.

Usually one has expressions with main argument x of the harmonic polylogarithms in $x \in] - \infty, +\infty[$. We provide the following algebraic transformation tables of the harmonic polylogarithms, which can be used to map between different kinematic regions. To do this the above basis representation is required:

$$\text{LISTOMX.m} \quad x \mapsto -x \quad [2.6 \text{ Mbyte}] \quad (12)$$

$$\text{LISTOOX.m} \quad x \mapsto \frac{1}{x} \quad [13.1 \text{ Mbyte}] \quad (13)$$

$$\text{LISTMTX.m} \quad x \mapsto \frac{1-x}{1+x} \quad [13.7 \text{ Mbyte}]. \quad (14)$$

These files can be loaded in any **Mathematica** program to carry out the respective transformations. One finally ends up with a representation in the interval

$$x \in [0, \sqrt{2} - 1], \quad (15)$$

for which the numeric implementation of **HPOLY** is provided. In course of these mappings also special constants contribute, cf. [33]. They are given by

$$\left\{ \ln(2), \zeta_2, \zeta_3, \text{Li}_4\left(\frac{1}{2}\right), \zeta_5, \text{Li}_5\left(\frac{1}{2}\right), \text{Li}_6\left(\frac{1}{2}\right), s_6, \zeta_7, \text{Li}_7\left(\frac{1}{2}\right), s_{7a}, s_{7b}, \text{Li}_8\left(\frac{1}{2}\right), s_{8a}, s_{8b}, s_{8c}, s_{8d} \right\}, \quad (16)$$

with

$$\zeta_k = S_k(\infty), \quad k \geq 2 \quad (17)$$

$$\text{Li}_k(x) = \sum_{l=1}^{\infty} \frac{x^l}{l^k} \quad (18)$$

$$s_6 = S_{-5,-1}(\infty) \quad (19)$$

$$s_{7a} = S_{-5,1,1}(\infty) \quad (20)$$

$$s_{7b} = S_{5,-1,-1}(\infty) \quad (21)$$

$$s_{8a} = S_{5,3}(\infty) \quad (22)$$

$$s_{8b} = S_{-7,-1}(\infty) \quad (23)$$

$$s_{8c} = S_{-5,-1,-1,-1}(\infty) \quad (24)$$

$$s_{8d} = S_{5,-1,1,1}(\infty). \quad (25)$$

Now the Bernoulli-improvement [1], see also [14], of the respective series expansion is performed. To this end we substitute

$$\text{term1}(x) \mapsto \text{term1}(x)|_{x \rightarrow 1-e^{-v}} \quad (26)$$

$$\text{term2}(x) \mapsto \text{term2}(x)|_{x \rightarrow e^u - 1}, \quad (27)$$

with

$$u = \ln(1+x) \quad (28)$$

$$v = -\ln(1-x) \quad (29)$$

and expand into series in u and v up to a certain order N_{max} , which we choose $N_{max} = 40$. This representation converges much faster than the corresponding Taylor series. The representations obtained are power series in u and v and logarithms of u and v to a certain power. At weight w there are contributions of up to $O(\ln^{w-1}(u))$ and of $O(\ln^{w-2}(v))$. In Appendix A we give one example for illustration.

One stores now the corresponding constants appearing in the expansion for each weight w in a data file. These constants are read into a very compact **FORTRAN** programme. Herewith we obtain a representation of all basis **HPLs** up to $w = 8$ for $x \in \mathbb{R}$, after arguments from all other regions are mapped using the commands described above. This allows to work with real representations only. Complex valued **HPLs** are already separated into their real and imaginary part by the above transformations.

For the purpose of possible later extensions to even higher weight we list a series of commands of **HarmonicSums**, with which we created the present tables and which can be used also for further extensions. For the arguments $x \in]-\infty, 0]$ one first applies the procedure

$$\text{TransformH}[H_{\bar{a}}(f(x)), x] \quad (30)$$

of **HarmonicSums**.

Example:

$$\begin{aligned} \text{TransformH}[H[-1, 0, 0, -x], x] &= -H[1, x]H[0, 0, -x] + H[0, -x]H[0, 1, x] - H[0, 0, 1, x] \\ &= -\frac{1}{2}H[1, x](H[0, x] + i\pi)^2 + (H[0, x] + i\pi)H[0, 1, x] \\ &\quad - H[0, 0, 1, x]. \end{aligned} \quad (31)$$

Here one observes (3) and $H[0, -x] = H[0, x] + i\pi$.

The same command maps $x \in [1, \infty]$ to $x \in [0, 1]$.

Example:

$$\text{TransformH}\left[H\left[-1, 0, 0, \frac{1}{x}\right]\right] = H[-1, 0, 0, x] - \frac{1}{6}H[0, x]^3. \quad (32)$$

Now only two regions have to be considered:

$$x \in [0, \sqrt{2} - 1] \quad \text{and} \quad x \in [\sqrt{2} - 1, 1]. \quad (33)$$

The expression in the second region is obtained from the first one by the mapping

$$x \mapsto \frac{1-x}{1+x}. \quad (34)$$

Example:

$$\begin{aligned} \text{TransformH}\left[H\left[-1, 0, 0, \frac{1-x}{1+x}\right]\right] &= \frac{3}{4}\zeta_3 - \frac{1}{6}H[-1, x]^3 - H[-1, x]H[-1, 1, x] \\ &\quad + H[-1, -1, 1, x] - H[-1, 1, 1, x]. \end{aligned} \quad (35)$$

In the case one is leaving the basis one uses

$$\text{ReduceToHBasis}[expression] \quad (36)$$

to get back to the basis again. The corresponding substitution rules are tabulated in the file `hrel.m` to weight $w = 8$. The above replacements are illustrated by a `Mathematica` notebook `HTransform.nb` provided in the ancillary files to this paper.

In the case we have HPLs over a two-letter alphabet, i.e. $\{0, -1\}$ or $\{0, 1\}$, only one branch cut is present, i.e. $[-1, -\infty[$ or $[1, \infty[$. In all other cases one has to consider a cut separated representation of the respective HPL. This cut separation is performed by using the command

$$\text{HSeparateCuts}[H_{\vec{a}}(x)]. \quad (37)$$

Example:

$$\begin{aligned} \text{HSeparateCuts}[H[-1, 1, 1, x]] &= \{\text{term1}, \text{term2}\} \\ &= \left\{ -\frac{1}{2} \ln^2(2) H[-1, x] + H[-1, 1, 1, x], \frac{1}{2} \ln^2(2) H[-1, x] \right\}. \end{aligned} \quad (38)$$

This command is needed to obtain the representations (26,27). The sum of the two terms yields the original expression again.

3 Description of the code

The code `HPOLY` provides the calculation of all HPLs, using the basis representation and argument mappings described in Section 2 in the region $x \in [0, \sqrt{2} - 1]$. It shall be compiled using

`gfortran hpoly.f.`

in `LINUX`. There is a user-routine `UHPOLY` through which the corresponding calculations can be made. Through the initialization routine `UHPOLYIN` the maximal weight of the calculation is chosen setting the variable `IW`. If the calculation of an HPL is intended at a higher weight than $w = 8$, the code cannot perform the calculation. Choosing `IW = 8` the initialization time is longest with ~ 0.27 sec for reading in the constant tables. Through `UHPOLY` the user can arrange special output formats. Only the routines `UHPOLYIN` and `UHPOLY` are free to be modified by the user.

A built in logic checks whether the required HPL is contained in the basis and whether the variable is in the correct range. The routine `HPOLYIN` performs necessary initializations.

The main routine is

```

      PROGRAM HPOLY
      *
      *-----
      *
      *    numerical evaluation of HPLs up to weight w = 8
      *
      *    J. Bluemlein, 17.9.2018
      *
      *-----
      *
      WRITE(6,*)
      &'**** HPOLY: calculation of HPLS up to w=8 ****'
      CALL UHPOLYIN

```

```

*
CALL HPOLYIN
WRITE(6,*) '**** HPOLY: initialization finished ****'
*
CALL UHPOLY
*
WRITE(6,*) '**** HPOLY: finished ****'
STOP
END
SUBROUTINE UHPOLYIN
*
*-----
*
IMPLICIT NONE
INTEGER IW
COMMON/WEIGHT/ IW
*
IW = 8
*
RETURN
END
SUBROUTINE UHPOLY
*
*-----
*
IMPLICIT NONE
*
REAL*8 H1,H2,H3,H4,H5,H6,H7,H8
REAL*8 T(8),X
INTEGER K
*
EXTERNAL H1,H2,H3,H4,H5,H6,H7,H8
*
CALL CPU_TIME(START)
X = 0.3D0
CALL HPOLYC(X)
*
T(1) = H1(-1,X)
T(2) = H2(0,1,X)
T(3) = H3(-1,0,0,X)
T(4) = H4(-1,-1,1,0,X)
T(5) = H5(-1,-1,1,0,1,X)
T(6) = H6(-1,0,-1,1,1,1,X)
T(7) = H7(-1,-1,1,1,0,1,0,X)
T(8) = H8(-1,0,-1,0,-1,0,1,1,X)
*
DO 1 K = 1,8
WRITE(6,*) 'K, X , T=', K,X,T(K)

```

```

1      CONTINUE
      CALL CPU_TIME(END)
      WRITE(6,*) 'CPU TIME=', END-START, '    SEC'
*
      RETURN
      END
*-----

```

Running the above example produces the following output values:

HPL	value	accuracy
H[-1,x]	0.26236426446749106e-0	7.96e-18
H[0,1,x]	0.32612951007547608e-0	1.05e-17
H[-1,0,0,x]	0.81699704232693138e-0	1.20e-16
H[-1,-1,1,0,x]	-0.10536957058865759e-0	1.20e-16
H[-1,-1,1,0,1,x]	0.27247014022231675e-3	3.97e-17
H[-1,0,-1,1,1,1,x]	0.45411840144185533e-5	8.46e-19
H[-1,-1,1,1,0,1,0,x]	-0.80698691040978487e-4	3.86e-18
H[-1,0,-1,0,-1,0,1,1,x]	0.57153046109648109e-6	3.47e-18

Table 2: The result of a test calculation of the code HPOLY.f with $x = 0.3$.

The following output is obtained by running the code default:

```

*****
*   HPOLY.f: J. Bluemlein 17.09.2018   *****
*****
**** HPOLY: calculation of HPLS up to w=8 ****
**** HPLs up to weight w =           8
READ-IN TIME=  0.263834999999999999          SEC
**** HPOLY: initialization finished ****
K,X,T=         1  0.299999999999999999          0.26236426446749106
K,X,T=         2  0.299999999999999999          0.32612951007547608
K,X,T=         3  0.299999999999999999          0.81699704232693138
K,X,T=         4  0.299999999999999999          -1.0536957058865759E-002
K,X,T=         5  0.299999999999999999          2.7247014022231675E-004
K,X,T=         6  0.299999999999999999          4.5411840144185533E-006
K,X,T=         7  0.299999999999999999          -8.0698691040978487E-005
K,X,T=         8  0.299999999999999999          5.7153046109648109E-007
CPU TIME=      2.2499999999997522E-004          SEC
*****
**** HPOLY: calculation      finished ****
*****

```

The absolute accuracy has been determined comparing with the corresponding result obtained using **Ginac** [16]. We have compared the numerical values produced by **HPOLY** and using **Ginac** forming the difference in a **Maple** programme, which can handle floating point number inputs straightforwardly.

4 Cyclotomic Harmonic Polylogarithms

In the following we would like to make a few brief remarks on the numerical evaluation of cyclotomic harmonic polylogarithms [20].

These functions emerge in massive calculations from 3-loops onward. Usually their main argument x is located in the interval $[-1, 1]$. This is going to be the range we are considering in the following. The new letters are of cyclotomy 3,4 and 6 :

$$\left\{ \frac{1}{1+x+x^2}, \frac{x}{1+x+x^2}, \frac{1}{1+x^2}, \frac{x}{1+x^2}, \frac{1}{1-x+x^2}, \frac{x}{1-x+x^2} \right\}. \quad (39)$$

Here cyclotomy labels the degree of the highest cyclotomic polynomial, cf. [20]. While here the transformations $x \mapsto 1/x$ and $x \mapsto -x$ are class preserving, the transformation $x \mapsto (1-x)/(1+x)$ is not. There is also no reason, why the above Bernoulli-improvement should accelerate the convergence of the expansions, given the structure of the above letters.

Yet one would like to have a series representation in the interval $[-1, 1]$, see e.g. Ref. [25, 26]. One possibility is given by Taylor series in x around $x = 0, 1$ and $x = -1$, which can be calculated analytically using the command

$$\text{HarmonicSumsSeries}[f(x), \{x, 0, n\}] \quad (40)$$

of **HarmonicSums**, where **n** denotes the highest power in the expansion. The series are intended to hold at a convergence radius of $\rho = 1/2$ at double precision, i.e. $53 \cdot \ln_{10}(2) \approx 15.955$ digits [34].

x	HPL	run time Fortran [sec]	run time Ginac [sec]	accuracy
-1.0	-0.19456824628601374558e+0	6.0e-6	2.4e-4	5.5e-18
-0.8	-0.10723997730030421271e+0	6.0e-6	2.9e-4	4.5e-18
-0.4	-0.16724871547938951751e-1	8.0e-6	3.3e-4	5.2e-18
0.4	0.25064429005154578702e-1	8.0e-6	2.7e-4	2.3e-18
0.8	0.19221120777067257053e+0	8.0e-6	6.2e-3	6.0e-18
1.0	0.34039020776540250510e+0	2.9e-5	6.0e-3	1.5e-17

Table 3: Example of the numerical calculation of an cyclotomic HPL $H[\{6, 1\}, 0, -1, x]$ comparing the run times in FORTRAN and Ginac.

The expansion around $x = -1$ may be performed by mapping

$$\left\{ \frac{1}{1+x+x^2} \leftrightarrow \frac{1}{1-x+x^2}, \frac{x}{1+x+x^2} \leftrightarrow -\frac{x}{1-x+x^2} \right\} \quad (41)$$

more effectively. In Table 3 we illustrate the run times needed to calculate the cyclotomic harmonic polylogarithm

$$\mathcal{H}[\{6, 1\}, 0, -1, x] = \int_0^x dy_1 \frac{y_1}{1 - y_1 + y_1^2} \int_0^{y_1} \frac{dy_2}{y_2} \int_0^{y_2} \frac{dy_2}{1 + y_2}. \quad (42)$$

We note that some of the cyclotomic **HPLs** can be complex for $x \in [-1, 0[$. The accuracy of the respective representation can be tested comparing the numerical result with the complex-valued representations provided by **Ginac**, cf. Ref. [16]. To have these explicit series representations is of importance since they usually perform faster than the algorithm based on a Hölder convolution in [16]. However, the latter one has the advantage that it can be extended to arbitrary precision.

In Table 3 we summarize the performance for a series of values using both implementations by considering the example of $\mathcal{H}(\{6, 1\}, 0, -1, x)$. The function is graphically illustrated in Figure 1.

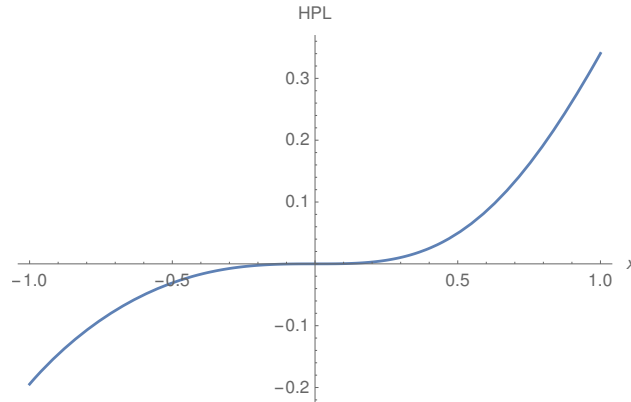


Figure 1: The function $\mathcal{H}[\{6, 1\}, 0, -1, x]$ in the region $x \in [-1, 1]$.

In the case of the planar contributions to the 3-loop massive Form factors [25, 26] 206 real-valued cyclotomic polylogarithms up to $w=6$ contribute. Their numerical representation is given in Ref. [26].³

5 Conclusions

We presented a compact **FORTRAN** code which calculates the **HPLs** numerically up to weight $w = 8$. Computer-readable replacement files are provided to allow the analytic representation of a corresponding polynomial of **HPLs** over the chosen algebraic basis, mapping the emerging arguments $x \in \mathbb{R}$ into the interval $x \in [0, \sqrt{2} - 1]$ in a **Mathematica** session during a preparatory step. For the implementation of the **HPLs** we used the Bernoulli-improvement [1]. The numerical accuracy of the representation is $\sim 10^{-15}$ or better. We also compared run times with corresponding implementations in **Ginac**. Because we limit the accuracy to about double precision, the **FORTRAN** implementation is faster. Numerical representations using **Mathematica** are usually slower, cf. Table 4. This is of advantage for larger scale phenomenological and experimental applications. The present code extends earlier work of the authors of Ref. [14] to far higher weight.

³In Refs. [35] the numerical representations are based on **Ginac**.

We briefly commented on the numerical implementation of cyclotomic harmonic polylogarithms, which are important in massive higher loop calculations, as e.g. for the massive Form factors from 3-loop order [25, 26, 35]. Here, precise Taylor series around $x = 0, 1, -1$ providing sufficient convergence in a circle of radius $\rho = 1/2$ within $x \in [-1, 1]$ serve the purpose. Detailed representations are given in Ref. [26].

The use of the FORTRAN-code `HPOLY.f` and of the attached subsidiary files requires the citation of the present paper. Discovered bugs shall be reported to `Johannes.Bluemlein@desy.de`.

A Numerical examples

A.1 Harmonic Polylogarithms

In the following we show the explicit representation for one example, $H[-1, 1, 0, x]$ and give some numerical illustration. This harmonic polylogarithm is represented by

$$\begin{aligned}
H[-1, 1, 0, x] &= 1.51561421398405852765343970478u - 0.173286795139986327354308030365u^2 \\
&- 9.62704417444368485301711279803e-3u^3 + 4.81352208722184242650855639902e-5u^5 \\
&- 5.45750803539891431576933832088e-7u^7 + 7.95886588495675004383028505128e-9u^9 \\
&- 1.31551502230690083369095620682e-10u^{11} + 2.34789839498327069523925944587e-12u^{13} \\
&- 4.41717498332597960030540850887e-14u^{15} + 8.63804169263087598091731324426e-16u^{17} \\
&- 1.74017681474319135406058068735e-17u^{19} + 3.58929549165142423790521910252e-19u^{21} \\
&- 7.5465410635862794583439948632e-21u^{23} + 1.6120829755323789347574740235e-22u^{25} \\
&- 3.49013136556204594004342294876e-24u^{27} + 7.64298689734656581222853583268e-26u^{29} \\
&- 1.69034929890524597704956962836e-27u^{31} + 3.77081955969303929350628544829e-29u^{33} \\
&- 8.47603948488133928506876824739e-31u^{35} + 1.91812087138166865720209408063e-32u^{37} \\
&- 4.36689003100059767808970725046e-34u^{39} \\
&- 1.51561421398405852765343970478v + 2.45753828564099545590515613688e-1v^2 \\
&- 7.28443369411397782763491496943e-3v^3 - 4.665222471648288365706911026793e-3v^4 \\
&+ 7.18180060085976020176571697142e-4v^5 + 1.0998660373605051065035495605e-5v^6 \\
&- 1.15456136926644459648518129105e-5v^7 + 2.2668351288044216255466542194e-7v^8 \\
&+ 2.10153210137148470728310501335e-7v^9 - 1.14687914794188966052447108046e-8v^{10} \\
&- 3.89908427576215576063242315571e-9v^{11} + 3.66138107063939833413172284743e-10v^{12} \\
&+ 7.06461980894443084509781017833e-11v^{13} - 1.01701961998442644609478487375e-11v^{14} \\
&- 1.20382742279859366919809526002e-12v^{15} + 2.62615442112180864875738033e-13v^{16} \\
&+ 1.80691839540417471780902160508e-14v^{17} - 6.45036155764417783408860256621e-15v^{18} \\
&- 1.94331333196745041013938386874e-16v^{19} + 1.52063344250754104057653980337e-16v^{20} \\
&- 5.82510897939787787036716572212e-19v^{21} - 3.44907968543220341667785192179e-18v^{22} \\
&+ 1.26829440977979122960207894152e-19v^{23} + 7.51151317182163069584552487431e-20v^{24} \\
&- 5.40148573489695792030638540445e-21v^{25} - 1.55998517649156934189241423249e-21v^{26} \\
&+ 1.76849481787857125449286908312e-22v^{27} + 3.04523398314277895864778604675e-23v^{28}
\end{aligned}$$

$$\begin{aligned}
& -5.13172809239911634765975392106e-24v^{29} - 5.41954452181938287426807053132e-25v^{30} \\
& + 1.38155545866828806840477032923e-25v^{31} + 8.13686152516356159595967623328e-27v^{32} \\
& - 3.51809188882467289354219400405e-27v^{33} - 7.47063186973607435240189362655e-29v^{34} \\
& + 8.54615019188533962991092101997e-29v^{35} - 1.08160378090075874429236755202e-30v^{36} \\
& - 1.98536129319620756482783417757e-30v^{37} + 9.18943710177493828604598228192e-32v^{38} \\
& + 4.40086958481705100275797309929e-32v^{39} - 3.65462415959228934623749630635e-33v^{40} \\
& + \left[0.693147180559945309417232121458u + 0.50000000000000000000000000000000u^2 \right. \\
& + 0.33333333333333333333333333333333u^3 + 0.25000000000000000000000000000000u^4 \\
& + 0.216666666666666666666666666666667u^5 + 0.20833333333333333333333333333333u^6 \\
& + 0.214682539682539682539682539683u^7 + 0.232291666666666666666666666666667u^8 \\
& + 0.260653659611992945326278659612u^9 + 0.300834986772486772486772486772u^{10} \\
& + 0.355101661455828122494789161456u^{11} + 0.426919505070546737213403880071u^{12} \\
& + 0.521158552417232972788528344084u^{13} + 0.644462450647346480679814013147u^{14} \\
& + 0.805794410709134584795960457336u^{15} + 1.01720115026784328619646079964u^{16} \\
& + 1.29486269565703958818843661670u^{17} + 1.66052621302802363001326846124u^{18} \\
& + 2.14346104513832728152640246783u^{19} + 2.78312455814376537110798639484u^{20} \\
& + 3.63279999838074563025362227705u^{21} + 4.76456594745939637957111115559u^{22} \\
& + 6.27609256313695770804541041481u^{23} + 8.29994698224797632624920886217u^{24} \\
& + 11.0163489628560348344708790596u^{25} + 14.6706757083827785850262439117u^{26} \\
& + 19.5975102694203842148444445452u^{27} + 26.2537143880730842374878981161u^{28} \\
& + 35.2639584799959710344981360061u^{29} + 47.4834621531241956309082413623u^{30} \\
& + 64.0845324452896992786471204045u^{31} + 86.6760348342256170319013958612u^{32} \\
& + 117.468474369604443607935643308u^{33} + 159.502292173007654388803768318u^{34} \\
& + 216.963842161082342229133897163u^{35} + 295.623066963675442979916734791u^{36} \\
& + 403.440206590022104677573187301u^{37} + 551.407438745662532802587646249u^{38} \\
& + 754.717250342131567887848557378u^{39} + 1034.38549262485015775284310232u^{40} \left. \right]
\end{aligned}$$

$$\times H[0, u]$$

$$\begin{aligned}
& + \left[6.93147180559945309417232121458e-1v - 1.93147180559945309417232121458e-1v^2 \right. \\
& + 2.64805138932786427505654547915e-2v^3 - 9.09445606607418535334798246358e-4v^4 \\
& - 2.32690966735029895126514843946e-4v^5 + 2.04869918599655702150619197958e-5v^6 \\
& + 3.27244532693809559445428120669e-6v^7 - 4.76697494718168113031752060736e-7v^8 \\
& - 4.79445670857266595292321218748e-8v^9 + 1.09995268113063922275364870359e-8v^{10} \\
& + 6.30079730899918561829533455058e-10v^{11} - 2.4918346059372030398198902277e-10v^{12} \\
& - 5.34007883918865414818687371719e-12v^{13} + 5.5156739926803903227757686507e-12v^{14} \\
& - 6.28857786197662796320841934221e-14v^{15} - 1.18725105791642065839718279224e-13v^{16} \\
& + 5.26578132632148224808742741623e-15v^{17} + 2.46845929257087823714218639074e-15v^{18} \\
& - 1.99041996903111499913398470405e-16v^{19} - 4.90370649831330005463993519578e-17v^{20} \\
& + 6.12701864575634669387293564239e-18v^{21} + 9.12721047200658523982514560198e-19v^{22}
\end{aligned}$$

$$\begin{aligned}
& -1.70288180301515490338143682483e-19v^{23} - 1.5276386825065569828307657806e-20v^{24} \\
& + 4.42944449090124929492159879193e-21v^{25} + 2.05134350399689853005153453942e-22v^{26} \\
& - 1.09492124504286623636232861158e-22v^{27} - 1.10399957724647642108199227876e-24v^{28} \\
& + 2.58812712884495900131556798953e-24v^{29} - 6.10104748471770267083048214269e-26v^{30} \\
& - 5.85427217667691779189384255622e-26v^{31} + 3.39895981945488747450890358933e-27v^{32} \\
& + 1.26193521067501965399016009955e-27v^{33} - 1.22550544038337779976290204537e-28v^{34} \\
& - 2.56460584430131563394858221159e-29v^{35} + 3.76155735706687563080727351239e-30v^{36} \\
& + 4.80146080059348793629186205266e-31v^{37} - 1.05400577554603083138457524555e-31v^{38} \\
& - 7.83689686147728007727555462251e-33v^{39} + 2.77092887621278377430721664171e-33v^{40} \Big] \\
& \times H[0, v] \tag{43}
\end{aligned}$$

for $x \in [0, \sqrt{2} - 1]$.

x	HPL	run time Fortran [sec]	run time Ginac [sec]	run time Mathematica [sec]
-20.0	5.369919763979762 -18.46370249603318 i	2.5e-5	0.02923	—
-0.9	-1.652038279906588 +3.344002738868969 i	2.0e-5	0.01917	3.45
-0.2	-0.067890106575246 +0.068215824899983 i	8.6e-5	0.00133	0.94
0.2	-0.058464914759637	5.4e-5	0.00189	0.84
0.9	-0.550223509450311	8.2e-5	0.01310	1.05
50.0	-18.95831087429180	3.5e-5	0.01260	2.08

Table 4: Some numerical values for $H[-1, 1, 0, x]$ and the run times in the FORTRAN implementation, by Ginac, and Mathematica.

The following mappings hold :

$$\begin{aligned}
H[-1, 1, 0, z] &= \frac{1}{2} \ln(2) \zeta_2 - \zeta_3 + \zeta_2 H[-1, x] + \frac{1}{6} H[-1, x]^3 - H[-1, x] H[0, -1, x] \\
&\quad - H[-1, x] H[0, 1, x] + H[-1, -1, 1, x] + 2H[0, -1, -1, x] + H[0, -1, 1, x] \\
&\quad + H[0, 1, -1, x] \text{ for } z \in [\sqrt{2} - 1, 1], \quad z = \frac{1-x}{1+x}. \tag{44}
\end{aligned}$$

$$\begin{aligned}
H[-1, 1, 0, -x] &= H[-1, x] H[0, -x] H[1, x] - H[0, -x] H[-1, 1, x] - H[1, x] H[0, -1, x] \\
&\quad + H[0, -1, 1, x] \tag{45}
\end{aligned}$$

$$H \left[-1, 1, 0, \frac{1}{x} \right] = 3 \ln(2) \zeta_2 + \frac{3}{4} \zeta_3 - 2 \zeta_2 H[-1, x] + 2 \zeta_2 H[0, x] - \frac{1}{2} H[-1, x] H[0, x]^2 + \frac{1}{6} H[0, x]^3$$

$$\begin{aligned}
& -H[0, x]H[-1, 1, x] + H[0, x]H[0, -1, x] + H[-1, x]H[0, 1, x] \\
& + H[0, x]H[0, 1, x] - H[0, 0, -1, x] - 2H[0, 0, 1, x] - H[0, 1, -1, x].
\end{aligned} \tag{46}$$

Here $H[0, -x] = H[0, x] + i\pi$ and one can see that the function is real for $x \geq 0$ and complex for $x \leq 0$.

In Table 4 we present some numerical results and compare the run times of **HPOLY**, **Ginac** [16], and **Mathematica**. In the latter case the harmonic polylogarithm is represented by a double-integral because of its last letter being 0 and has been computed using the **Mathematica**-command **NIntegrate**. In Figure 2 the behaviour of $H[-1, 1, 0, x]$ in the range $x \in [-3, 3]$ is illustrated for the real and imaginary part of this function.

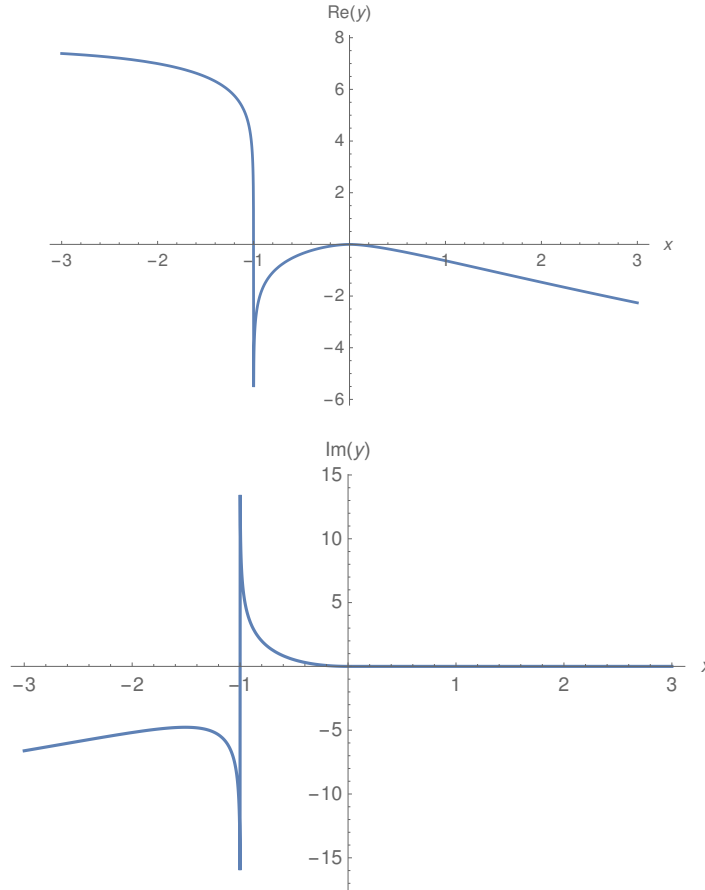


Figure 2: Real and imaginary part of the function $H[-1, 1, 0, x]$.

We have also measured the run times for all HPLs in **HPOLY** and **Ginac** at $x = 0.3$, yielding $9 \cdot 10^{-3}$ sec and 2.7 sec, respectively.

A.2 Cyclotomic Harmonic Polylogarithms

We discuss the representation of the cyclotomic harmonic polylogarithm $H[\{6, 1\}, 0, -1, x]$, i.e. a real representation w.r.t. the indices of this iterated integral, and give some numerical illustrations. The treatment is very similar to the case of the usual harmonic polylogarithms, since only more letters (regular on $x \in [-1, 1]$) are added to the alphabet. In the following we display

the first 40 expansion coefficients, which show that the series converges only slowly, requiring the expansion coefficients $|\xi| \leq 1/2$ to reach double precision, with $\xi \in x, x+, x-$. One obtains

$$\begin{aligned} & H[\{6, 1\}, 0, -1, x] \\ &= 0.333333333333333333333333333333x^3 + 0.187500000000000000000000000000x^4 \\ &\quad - 0.027777777777777777777777777777x^5 - 0.158564814814814814814814814815x^6 \\ &\quad - 0.110357142857142857142857142857x^7 + 0.018888888888888888888888888889x^8 \\ &\quad + 0.104891030486268581506676744772x^9 + 0.0777283163265306122448979591837x^{10} \\ &\quad - 0.0140354938271604938271604938272x^{11} - 0.0784727996136726295456454186613x^{12} \\ &\quad - 0.0599245154196452897751599050300x^{13} + 0.0111221750331598816447301295786x^{14} \\ &\quad + 0.0627097540455293813047170800529x^{15} + 0.0487396137126484903707680930458x^{16} \\ &\quad - 0.00919812040212473112906013338914x^{17} - 0.0522282286799164597292395420194x^{18} \\ &\quad - 0.0410673611468378472909181683198x^{19} + 0.00783709173477453784228900692598x^{20} \\ &\quad + 0.0447519895041858538174032906623x^{21} + 0.0354795429496550987872585893358x^{22} \\ &\quad - 0.00682496786953676661431421089295x^{23} - 0.0391495967330932143454954592724x^{24} \\ &\quad - 0.0312290280570659818475651680142x^{25} + 0.00604340419447986488429425811813x^{26} \\ &\quad + 0.0347946003882639271548437004232x^{27} + 0.0278873718387793265650809033939x^{28} \\ &\quad - 0.00542193299602834607361795351938x^{29} - 0.0313119326191574937536749057194x^{30} \\ &\quad - 0.0251913826144454848304917188798x^{31} + 0.00491606270049386474230914922487x^{32} \\ &\quad + 0.0284632560663618316507637169020x^{33} + 0.0229704965080628462565091141932x^{34} \\ &\quad - 0.00449635129343592282670836238196x^{35} - 0.0260898396790974516252031566278x^{36} \\ &\quad - 0.0211093407201794247209341019341x^{37} + 0.00414255307533152470774948058168x^{38} \\ &\quad + 0.0240818739410040077679013044162x^{39} + 0.0195270886515233773838902417075x^{40} \end{aligned}$$

for $|x| \leq 1/2$,

$$\begin{aligned} H[\{6, 1\}, 0, -1, x] = & 0.340390207765402505097420897861 - 0.822467033424113218236207583323x_- \\ & + 0.346573590279972654708616060729x_-^2 + 0.306346874568028624314941214684x_-^3 \\ & + 0.0380088949293707529895132089211x_-^4 - 0.156620128473319660687641363698x_-^5 \\ & - 0.156597541296457541255631321152x_-^6 - 0.0225554460774175569791333271786x_-^7 \\ & + 0.0976519307918251785767956228722x_-^8 + 0.104325369067180060242342866096x_-^9 \\ & + 0.0157646326300601883740158725548x_-^{10} - 0.0710280097948583028355514206622x_-^{11} \\ & - 0.0782470790060635646220138334389x_-^{12} - 0.0121277817768820409689235610102x_-^{13} \\ & + 0.0558072840567786665216825950946x_-^{14} + 0.0625974915019450287029214667407x_-^{15} \\ & + 0.00985375412765262691406233891207x_-^{16} - 0.0459589676261932861599464151009x_-^{17} \\ & - 0.0521645876903227753783336731631x_-^{18} - 0.00829790297360334455491929168585x_-^{19} \\ & + 0.0390651204796707863269908727563x_-^{20} + 0.0447125028842087830347950513173x_-^{21} \\ & + 0.00716637038578137582472531686648x_-^{22} - 0.0339696701394558426216208414143x_-^{23} \\ & - 0.0391234400919570434982680106204x_-^{24} - 0.00630640596933985512905044780475x_-^{25} \\ & + 0.0300500928025446660489690952202x_-^{26} + 0.0347763911870470964246495341404x_-^{27} \end{aligned}$$

$$\begin{aligned}
& +0.00563071961290383166332573519540x_-^{28} - 0.0269414625137767299362368745328x_-^{29} \\
& -0.0312987520688579112914323618433x_-^{30} - 0.00508581126349987627019732720481x_-^{31} \\
& +0.0244157004030054693980770411665x_-^{32} + 0.0284534109716415871863511989614x_-^{33} \\
& +0.00463706321081657316300068964177x_-^{34} - 0.0223229260827576788486381082623x_-^{35} \\
& -0.0260822933906759433489536861150x_-^{36} - 0.00426108511264431308690488692722x_-^{37} \\
& +0.0205605898130653384249909302377x_-^{38} + 0.0240759631298542819670313546736x_-^{39} \\
& +0.00394150372919578889368783808396x_-^{40}
\end{aligned} \tag{48}$$

for $x \in [1/2, 1]$ and

$$\begin{aligned} & -0.194568246286013745578625203082 \\ & + \left(-0.548311355616075478824138388882x_+ + 0.250000000000000000000000000000x_+^2 \right. \\ & + 0.107219780253638016165645006172x_+^3 + 0.271740944494877713834930138883e-1x_+^4 \\ & - 0.257505086150775649670496049413e-2x_+^5 - 0.943456337336485738803035905363e-2x_+^6 \\ & - 0.816905052458110731788689690942e-2x_+^7 - 0.520795277987738305198622658940e-2x_+^8 \\ & - 0.277060313984829741036690222526e-2x_+^9 - 0.126851938027123900831848220794e-2x_+^{10} \\ & - 0.503667726779323063780400849847e-3x_+^{11} - 0.180235498282985631923185530178e-3x_+^{12} \\ & - 0.733668756985816187679438098227e-4x_+^{13} - 0.517342767257388203755181868192e-4x_+^{14} \\ & - 0.529768547494945595363266938968e-4x_+^{15} - 0.541523823048175872246118351091e-4x_+^{16} \\ & - 0.504956623487350579432257136818e-4x_+^{17} - 0.435057422798886284192642625177e-4x_+^{18} \\ & - 0.355976912750893316702486185276e-4x_+^{19} - 0.283747696506692566232934066468e-4x_+^{20} \\ & - 0.224862637086975580231554304003e-4x_+^{21} - 0.179654244130490719256568063724e-4x_+^{22} \\ & - 0.145743279416543605609014732902e-4x_+^{23} - 0.120215246792584256329867181528e-4x_+^{24} \\ & - 0.100603516738056680791340289733e-5x_+^{25} - 0.851341006823932654661521938303e-5x_+^{26} \\ & - 0.726349827322799151408472485058e-5x_+^{27} - 0.623549500679759632941394636300e-5x_+^{28} \\ & - 0.538028643232232014650606829135e-5x_+^{29} - 0.466381278035087853538774282359e-5x_+^{30} \\ & - 0.406070457715978200292192742285e-5x_+^{31} - 0.355097711186231539907222170265e-5x_+^{32} \\ & - 0.311840966987030702179107122442e-5x_+^{33} - 0.274971225781815571600317219459e-5x_+^{34} \\ & - 0.243400658296993141391406093405e-5x_+^{35} - 0.216242026558948610498083647252e-5x_+^{36} \\ & - 0.192773190450817788358838757683e-5x_+^{37} - 0.172405883754813902219634478337e-5x_+^{38} \\ & - 0.154659185525022548769109014074e-5x_+^{39} - 0.139137829546522993754067930066e-5x_+^{40} \\ & + \left[-0.166666666666666666666666666667x_+^2 - 0.555555555555555555555555555555e-1x_+^3 \right. \\ & + 0.166666666666666666666666666667e-1x_+^5 + 0.166666666666666666666666666667e-1x_+^6 \\ & + 0.119047619047619047619047619048e-1x_+^7 + 0.724206349206349206349206349206e-2x_+^8 \\ & + 0.401234567901234567901234567901e-2x_+^9 + 0.214285714285714285714285714286e-2x_+^{10} \\ & + 0.119047619047619047619047619048e-2x_+^{11} + 0.748556998556998556998556998557e-3x_+^{12} \\ & + 0.549450549450549450549450549451e-3x_+^{13} + 0.448955806098663241520384377527e-3x_+^{14} \\ & + 0.382395382395382395382395382395e-3x_+^{15} + 0.326756576756576756576756576757e-3x_+^{16} \\ & + 0.276765717942188530423824541472e-3x_+^{17} + 0.232655869910771871556185281675e-3x_+^{18} \\ & + 0.195199537304800462695199537305e-3x_+^{19} + 0.164375742936114453142316919407e-3x_+^{20} \end{aligned}$$

$$\begin{aligned}
& +0.139449719281652054761298458777e-3x_+^{21} + 0.119375300560214717265097225694e-3x_+^{22} \\
& +0.131134477762556911836761076230e-3x_+^{23} + 0.897896648502382786943293067940e-4x_+^{24} \\
& +0.787312103101576785787312103102e-4x_+^{25} + 0.694430190195227419411231291907e-4x_+^{26} \\
& +0.615646514102514860305899426862e-4x_+^{27} + 0.548299963598119473742033978943e-4x_+^{28} \\
& +0.490370738953877725049414746362e-4x_+^{29} + 0.440279180491117866083633973683e-4x_+^{30} \\
& +0.396759437506466368945146534281e-4x_+^{31} + 0.358780629639696323771643547090e-4x_+^{32} \\
& +0.325495384282835477024478780285e-4x_+^{33} + 0.296203528090464701585299046187e-4x_+^{34} \\
& +0.270324458622603140106269524853e-4x_+^{35} + 0.247375088378105046141745858401e-4x_+^{36} \\
& +0.226951904770740284572302948414e-4x_+^{37} + 0.208716366275858749420441086658e-4x_+^{38} \\
& +0.192383083866313435144166242061e-4x_+^{39} + 0.177710310170514453671986061289e-4x_+^{40} \Big] \\
& \times H[-1, x] \Big) \tag{49}
\end{aligned}$$

for $x \in [-1, -1/2]$. Here we used the notation $x_{\pm} = (1 \pm x)$.

B List of the attached files

The following files are attached to this paper. The analytic files are given in **Mathematica** format and can be easily converted into **Maple** or **FORM** [36] format.

File name	contents
LISTn.m	Bases per weight n , n = 1...8
LIST.m	Total basis
hrel.m	Replacement rules to basis for all HPLs up to weight w =8.
LISTOMX.m	Replacement rules: $x \mapsto -x$ for the basis elements
LISTOOX.m	Replacement rules: $x \mapsto 1/x$ for the basis elements
LISTMTX.m	Replacement rules: $x \mapsto (1-x)/(1+x)$ for the basis elements
HTransform.nb	an example notebook for variable transformation
HPOLY.f	FORTTRAN code (1298 lines)
TTTn.m	Data files containing the numerical constants needed at the different weights, n = 2...8. the different weights
Bkl.m	Data files to identify the basis elements: k =1 .. k =1, k = 5,6,7,8

Table 5: The list of attached files

The data files **TTTn.m** contain 360, 1600, 5040, 17280, 51040, 162240, 486000 constants for **w** = 2...8. Initially the external files for reading or writing are all set to **STATUS="NEW"**, as e.g.

```

OPEN(UNIT=14, FILE="TTT2.m", FORM="FORMATTED", STATUS="NEW",
&      ACTION="READ")

```

In later use the standard input files have to be set to `STATUS="OLD"`. For convenience we provide two versions of `HPOLY.f`. `HPOLYnew.f` contains the file status: `"NEW"` and `HPOLY.f` the file status `"OLD"`.

Acknowledgment. We would like to thank R. Germundson, P. Marquard, N. Rana, and K. Schönwald for discussions and N. Rana for an introduction to and help with using `Ginac`. This work was supported in part by the Austrian Science Fund (FWF) grant SFB F50 (F5009-N15), by the bilateral project DNTS-Austria 01/3/2017 (WTZ BG03/2017), funded by the Bulgarian National Science Fund and OeAD (Austria), by ‘Innovatives OÖ 2020’ of the Upper Austrian Government, by the EU TMR network SAGEX Marie Skłodowska-Curie grant agreement No. 764850 and COST action CA16201: Unraveling new physics at the LHC through the precision frontier. M. Round thanks DESY and the KMP Berlin for support.

References

- [1] G. 't Hooft and M.J.G. Veltman, Scalar One Loop Integrals, Nucl. Phys. B **153** (1979) 365–401.
- [2] T. Gehrmann and E. Remiddi, Two loop master integrals for $\gamma^* \rightarrow 3$ jets: The non-planar topologies, Nucl. Phys. B **601** (2001) 287–317 [hep-ph/0101124].
- [3] J.A.M. Vermaseren, A. Vogt and S. Moch, The third-order QCD corrections to deep-inelastic scattering by photon exchange, Nucl. Phys. B **724** (2005) 3–182 [hep-ph/0504242].
- [4] J. Ablinger, A. Behring, J. Blümlein, A. De Freitas, A. von Manteuffel and C. Schneider, The 3-loop pure singlet heavy flavor contributions to the structure function $F_2(x, Q^2)$ and the anomalous dimension, Nucl. Phys. B **890** (2014) 48–151 [arXiv:1409.1135 [hep-ph]].
- [5] E. Remiddi and J.A.M. Vermaseren, Harmonic polylogarithms, Int. J. Mod. Phys. A **15** (2000) 725–754 [hep-ph/9905237].
- [6] L. Lewin, Dilogarithms and associated functions, (Macdonald, London, 1958);
L. Lewin, Polylogarithms and associated functions, (North Holland, New York, 1981).
- [7] A. Devoto and D.W. Duke, Table of integrals and formulae for Feynman diagram calculations, Riv. Nuovo Cim. **7N6** (1984) 1–39
- [8] N. Nielsen, Der Eulersche Dilogarithmus und seine Verallgemeinerungen, Nova Acta Leopold. **XC** (1909) Nr. 3, 125–211.
- [9] K.S. Kölbig, J.A. Mignoco and E. Remiddi, On Nielsen’s generalized polylogarithms and their numerical calculation, BIT **10** (1970) 38–74.
- [10] K.S. Kölbig, Nielsen’s generalized polylogarithms, SIAM J. Math. Anal. **17** (1986) 1232–1258.
- [11] J. Blümlein, Algebraic relations between harmonic sums and associated quantities, Comput. Phys. Commun. **159** (2004) 19–54 [hep-ph/0311046].

- [12] J.A.M. Vermaseren, Harmonic sums, Mellin transforms and integrals, *Int. J. Mod. Phys. A* **14** (1999) 2037–2076 [hep-ph/9806280].
- [13] J. Blümlein and S. Kurth, Harmonic sums and Mellin transforms up to two loop order, *Phys. Rev. D* **60** (1999) 014018 [hep-ph/9810241].
- [14] T. Gehrmann and E. Remiddi, Numerical evaluation of harmonic polylogarithms, *Comput. Phys. Commun.* **141** (2001) 296–312 [hep-ph/0107173].
- [15] C.W. Bauer, A. Frink and R. Kreckel, Introduction to the GiNaC framework for symbolic computation within the C++ programming language, *J. Symb. Comput.* **33** (2000) 1–12 [cs/0004015 [cs-sc]].
- [16] J. Vollinga and S. Weinzierl, Numerical evaluation of multiple polylogarithms, *Comput. Phys. Commun.* **167** (2005) 177–194 [hep-ph/0410259].
- [17] D. Maitre, HPL, a mathematica implementation of the harmonic polylogarithms, *Comput. Phys. Commun.* **174** (2006) 222–240 [hep-ph/0507152]; Extension of HPL to complex arguments, *Comput. Phys. Commun.* **183** (2012) 846 [hep-ph/0703052].
- [18] J. Ablinger, The package `HarmonicSums`: computer algebra and analytic aspects of nested sums, *PoS (LL2014)* 019; A Computer Algebra Toolbox for Harmonic Sums Related to Particle Physics, Diploma Thesis, J. Kepler University Linz, 2009, arXiv:1011.1176 [math-ph].
- [19] J. Ablinger, Computer Algebra Algorithms for Special Functions in Particle Physics, Ph.D. Thesis, J. Kepler University Linz, 2012, arXiv:1305.0687 [math-ph].
- [20] J. Ablinger, J. Blümlein and C. Schneider, Harmonic sums and polylogarithms generated by cyclotomic polynomials, *J. Math. Phys.* **52** (2011) 102301 [arXiv:1105.6063 [math-ph]].
- [21] J. Ablinger, J. Blümlein and C. Schneider, Analytic and algorithmic aspects of generalized harmonic sums and polylogarithms, *J. Math. Phys.* **54** (2013) 082301 [arXiv:1302.0378 [math-ph]].
- [22] J. Ablinger, J. Blümlein, C.G. Raab and C. Schneider, Iterated binomial sums and their associated iterated integrals, *J. Math. Phys.* **55** (2014) 112301 [arXiv:1407.1822 [hep-th]].
- [23] H. Frellesvig, Generalized polylogarithms in `Maple`, arXiv:1806.02883 [hep-th].
- [24] S. Moch, B. Ruijl, T. Ueda, J.A.M. Vermaseren and A. Vogt, Four-loop non-singlet splitting functions in the planar limit and beyond, *JHEP* **1710** (2017) 041 [arXiv:1707.08315 [hep-ph]].
- [25] J. Ablinger, J. Blümlein, P. Marquard, N. Rana and C. Schneider, Heavy quark form factors at three loops in the planar limit, *Phys. Lett. B* **782** (2018) 528–532 [arXiv:1804.07313 [hep-ph]].
- [26] J. Ablinger, J. Blümlein, P. Marquard, N. Rana and C. Schneider, Automated Solution of First Order Factorizable Systems of Differential Equations in One Variable, *Nucl. Phys. B*, in print, [arXiv:1810.12261 [hep-ph]].
- [27] J. Ablinger, J. Blümlein, M. Round and C. Schneider, Special functions, transcendentals and their numerics, *PoS (RADCOR 2017)* 010 [arXiv:1712.08541 [hep-th]].

- [28] J. Blümlein, Analytic continuation of Mellin transforms up to two loop order, *Comput. Phys. Commun.* **133** (2000) 76–104 [hep-ph/0003100];
 J. Blümlein and S.O. Moch, Analytic continuation of the harmonic sums for the 3-loop anomalous dimensions, *Phys. Lett. B* **614** (2005) 53–61 [hep-ph/0503188];
 Structural Relations of harmonic sums and Mellin transforms at weight $w = 6$, *Clay Math. Proc.* **12** (2010) 167–188 [arXiv:0901.0837 [math-ph]];
 A.V. Kotikov and V.N. Velizhanin, Analytic continuation of the Mellin moments of deep inelastic structure functions, hep-ph/0501274.
- [29] J. Blümlein, Structural relations of harmonic sums and Mellin transforms up to weight $w = 5$, *Comput. Phys. Commun.* **180** (2009) 2218–2249 [arXiv:0901.3106 [hep-ph]];
- [30] J. Ablinger and J. Blümlein, Harmonic sums, polylogarithms, special numbers, and their generalizations, in : **Computer Algebra in Quantum Field Theory: Integration, Summation and Special Functions** (Texts & Monographs in Symbolic Computation) (Springer, Wien, 2013) 1–29, Eds. C. Schneider and J. Blümlein, arXiv:1304.7071 [math-ph];
 J. Ablinger, J. Blümlein and C. Schneider, Generalized harmonic, cyclotomic, and binomial sums, their polylogarithms and special numbers, *J. Phys. Conf. Ser.* **523** (2014) 012060 [arXiv:1310.5645 [math-ph]].
- [31] C. Duhr, Mathematical aspects of scattering amplitudes, in: **TASI Lectures Journeys Through the Precision Frontier: Amplitudes for Colliders**, Eds. L. Dixon and F. Petriello, p. 419–476, (World Scientific, Singapore, 2015), [arXiv:1411.7538 [hep-ph]].
- [32] J. Blümlein and C. Schneider, Analytic computing methods for precision calculations in quantum field theory, *Int. J. Mod. Phys. A* **33** (2018) no.17, 1830015 [arXiv:1809.02889 [hep-ph]]; Computer algebra tools for Feynman integrals and related multi-sums, *PoS (LL2018)* 052 [arXiv:1809.06168 [cs.SC]].
- [33] J. Blümlein, D.J. Broadhurst and J.A.M. Vermaseren, The multiple zeta value data mine, *Comput. Phys. Commun.* **181** (2010) 582–625 [arXiv:0907.2557 [math-ph]].
- [34] https://de.wikipedia.org/wiki/Doppelte_Genauigkeit
- [35] R.N. Lee, A.V. Smirnov, V.A. Smirnov and M. Steinhauser, Three-loop massive form factors: complete light-fermion and large- N_c corrections for vector, axial-vector, scalar and pseudo-scalar currents, *JHEP* **1805** (2018) 187 [arXiv:1804.07310 [hep-ph]]; Three-loop massive form factors: complete light-fermion corrections for the vector current, *JHEP* **1803** (2018) 136 [arXiv:1801.08151 [hep-ph]];
 J. Henn, A.V. Smirnov, V.A. Smirnov and M. Steinhauser, Massive three-loop form factor in the planar limit, *JHEP* **1701** (2017) 074 [arXiv:1611.07535 [hep-ph]].
- [36] J.A.M. Vermaseren, New features of FORM, math-ph/0010025.