

**Search for  $\Upsilon(1S,2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$  and  $e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  at  $\sqrt{s} = 10.52, 10.58,$   
and  $10.867$  GeV**

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(Received 6 May 2018; published 14 June 2018)

The first search for double charged charmoniumlike state production in  $\Upsilon(1S)$  and  $\Upsilon(2S)$  decays and in  $e^+e^-$  annihilation at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV is conducted using data collected with the Belle detector at the KEKB asymmetric energy electron-positron collider. No significant signals are observed in any of the studied modes, and the 90% credibility level upper limits on their product branching fractions in  $\Upsilon(1S)$  and  $\Upsilon(2S)$  decays [ $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$  ( $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ )] and the product of Born cross section and branching fraction for  $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  ( $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$ ) at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV are determined. Here,  $Z_c$  refers to the  $Z_c(3900)$  and  $Z_c(4200)$  observed in the  $\pi J/\psi$  final state, the  $Z_{c1}(4050)$  and  $Z_{c2}(4250)$  in the  $\pi\chi_{c1}(1P)$  final state, and the  $Z_c(4050)$  and  $Z_c(4430)$  in the  $\pi\psi(2S)$  final state.

DOI: [10.1103/PhysRevD.97.112004](https://doi.org/10.1103/PhysRevD.97.112004)

## I. INTRODUCTION

In the past decade, many experiments, at both lepton and hadron colliders, reported evidence for a large number of new particles having exotic properties that are difficult to accommodate within the conventional quark model, especially those that couple to heavy quarkonium and have nonzero charge [1]. Since the discovery of the  $Z_c^+(4430)(\rightarrow \pi^+\psi(2S))$  [2], more charged charmoniumlike resonances have been observed, including  $Z_{c1}^+(4050)$  and  $Z_{c2}^+(4250)(\rightarrow \pi^+\chi_{c1}(1P))$  [3],  $Z_c^+(3900)$  [4,5] and  $Z_c^+(4200)(\rightarrow \pi^+J/\psi)$  [6], and  $Z_c^+(4050)(\rightarrow \pi^+\psi(2S))$  [7]. The preferred assignment of the quantum numbers for  $Z_c^+(4430)$ ,  $Z_c^+(3900)$ , and  $Z_c^+(4200)$  are  $J^P = 1^+$  [6,8,9]. Although these charged charmoniumlike candidates are widely interpreted unconventional  $c\bar{c}$  states, it remains a significant challenge to determine their internal dynamics [10].

Considerable efforts in theory have been devoted to interpreting these states as tetraquarks, molecules, hybrids, or hadrocharmonia [11–13]. To distinguish among these explanations, experimental input is needed, such as that from the study of double  $Z_c^\pm$  production in  $e^+e^-$  annihilation [14,15]. For  $e^+e^- \rightarrow Z_c^+ Z_c^-$ , the dependence on  $s$  (the  $e^+e^-$  c.m. energy squared) of the electromagnetic form factor,  $F_{Z_c^+ Z_c^-}$ , is  $1/s^3$  for a  $Z_c$  state with tetraquark

structure or  $1/s$  for a  $Z_c$  system of two tightly bound diquarks [14,15].

In this paper, we report the search for  $Z_c$ -pair production in  $\Upsilon(1S)$  and  $\Upsilon(2S)$  decays as well as in  $e^+e^-$  annihilation at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV. We measure the production rates of  $\Upsilon(1S)$  and  $\Upsilon(2S)$  decay into a pair of  $Z_c$  states and the Born cross sections of  $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV. Here,  $Z_c^+$  denotes  $Z_c^+(3900)$ ,  $Z_c^+(4200)$ ,  $Z_{c1}^+(4050)$ ,  $Z_{c2}^+(4250)$ ,  $Z_c^+(4050)$ , or  $Z_c^+(4430)$ . In these searches, the decay modes considered are  $Z_c^+(3900)/Z_c^+(4200) \rightarrow \pi^+J/\psi$ ,  $Z_{c1}^+(4050)/Z_{c2}^+(4250) \rightarrow \pi^+\chi_{c1}(1P)$ , and  $Z_c^+(4050)/Z_c^+(4430) \rightarrow \pi^+\psi(2S)$  [16]. Although it is possible to search for combinations of any two of the enumerated  $Z_c$  states, only both  $Z_c$  states decaying into the same final states are currently studied in this study; the information of the other final-state combinations can be estimated from the results reported here.

## II. DATA SAMPLE AND THE BELLE DETECTOR

This analysis utilizes  $(5.74 \pm 0.09)$  fb $^{-1}$  data collected at the  $\Upsilon(1S)$  peak [(102 ± 3) million  $\Upsilon(1S)$  events],  $(24.91 \pm 0.35)$  fb $^{-1}$  data collected at the  $\Upsilon(2S)$  peak [(158 ± 4) million  $\Upsilon(2S)$  events],  $(89.5 \pm 1.3)$  fb $^{-1}$  data collected at  $\sqrt{s} = 10.52$  GeV,  $(711.0 \pm 10.0)$  fb $^{-1}$  data collected at  $\sqrt{s} = 10.58$  GeV [ $\Upsilon(4S)$  peak], and  $(121.4 \pm 1.7)$  fb $^{-1}$  data collected at  $\sqrt{s} = 10.867$  GeV [ $\Upsilon(5S)$  peak]. All the data were collected with the Belle detector [17] operating at the KEKB asymmetric-energy  $e^+e^-$  collider [18]. The Belle detector is a large solid-angle magnetic spectrometer that consists of a silicon vertex detector, a 50-layer central drift chamber (CDC), an array

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TABLE I. The world-average (nominal) values of  $Z_c^\pm$  masses (in MeV/ $c^2$ ) and widths (in MeV) [23].

$Z_c$ states	$Z_c$ labels in Ref. [23]	Mass	Width
$Z_c^+(3900)$	$X^+(3900)$	$3886.6 \pm 2.4$	$28.1 \pm 2.6$
$Z_c^+(4200)$	$X^+(4200)$	$4196_{-32}^{+35}$	$370_{-150}^{+100}$
$Z_{c1}^+(4050)$	$X^+(4050)$	$4051_{-40}^{+24}$	$82_{-28}^{+50}$
$Z_{c2}^+(4250)$	$X^+(4250)$	$4248_{-50}^{+190}$	$177_{-70}^{+320}$
$Z_c^+(4050)$	$X^+(4055)$	$4054 \pm 3.2$	$45 \pm 13$
$Z_c^+(4430)$	$X^+(4430)$	$4478_{-18}^{+15}$	$181 \pm 31$

of aerogel threshold Cherenkov counters (ACC), a barrel-like arrangement of time-of-flight scintillation counters (TOF), and an electromagnetic calorimeter comprised of CsI(Tl) crystals (ECL) located inside a superconducting solenoid coil that provides a 1.5T magnetic field. An iron flux-return yoke instrumented with resistive plate chambers located outside the coil is used to detect  $K_L^0$  mesons and to identify muons. A detailed description of the Belle detector can be found in Refs. [17,19].

We generate large signal Monte Carlo (MC) samples with the EVTGEN event generator [20] to understand the signal event topology and to estimate the signal selection efficiency. The angular distribution for  $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$  and  $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  at  $\sqrt{s} = 10.52, 10.58$ , and 10.867 GeV is simulated assuming an  $e^+e^- \rightarrow A\bar{A}$  decay mode ( $A$  denoting an axial-vector state), i.e.,  $dN/d\cos\theta_{Z_c} \propto 1 - \cos^2\theta_{Z_c}$  [21], where  $\theta_{Z_c}$  is the polar angle of the  $Z_c$  in the  $e^+e^-$  c.m. system. Initial-state radiation (ISR) is taken into account by assuming that the cross sections follow a  $1/s^2$  dependence in  $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  reactions. For the  $Z_c$  pair in one MC event, one  $Z_c$  decays into  $\pi^+ J/\psi$ ,  $\pi^\pm \chi_{c1}(1P)$ , or  $\pi^\pm \psi(2S)$  using the phase space model, while the other is simulated with the inclusive decays using PYTHIA [22]. The masses and widths of the  $Z_c^\pm$  states are set according to the latest world-average values [23], as summarized in Table I.

### III. COMMON EVENT SELECTION CRITERIA

For well-reconstructed charged tracks, the impact parameters perpendicular to and along the beam direction with respect to the nominal interaction point are required to be less than 0.5 and 4 cm, respectively, and the transverse momentum in the laboratory frame is required to be larger than 0.1 GeV/ $c$ . We require the number of well-reconstructed charged tracks to be greater than four to suppress the significant background from quantum electrodynamics processes. For charged tracks, information from different detector subsystems including specific ionization in the CDC, time measurements in the TOF, and the response of the ACC is combined to form the likelihood  $\mathcal{L}_i$  for particle species  $i$ , where  $i = \pi, K$ , or  $p$  [24]. Charged tracks with  $R_K = \mathcal{L}_K / (\mathcal{L}_K + \mathcal{L}_\pi) < 0.4$  are considered to be pions. With this condition, the pion identification

efficiency is 97%, and the kaon to pion misidentification rate is about 4%. A similar likelihood ratio is defined as  $R_e = \mathcal{L}_e / (\mathcal{L}_e + \mathcal{L}_{\text{non-}e})$  [25] for electron identification and  $R_\mu = \mathcal{L}_\mu / (\mathcal{L}_\mu + \mathcal{L}_K + \mathcal{L}_\pi)$  [26] for muon identification. An ECL cluster is treated as a photon candidate if it is isolated from the projected path of charged tracks in the CDC, and its energy is greater than 50 MeV.

For the lepton pair  $\ell^+ \ell^-$  ( $\ell = e$  or  $\mu$ ) used to reconstruct the  $J/\psi$ , both tracks must have  $R_e > 0.95$  in the  $e^+e^-$  mode, while one track must have  $R_\mu > 0.95$  and the other  $R_\mu > 0.05$  in the  $\mu^+\mu^-$  mode. The lepton-pair identification efficiencies for  $e^+e^-$  and  $\mu^+\mu^-$  are 96% and 93%, respectively. We apply a lepton veto to the bachelor pion candidate by requiring  $R_\mu < 0.95$  and  $R_e < 0.95$ . To reduce the effect of bremsstrahlung and final-state radiation, photons detected in the ECL within a 50 mrad cone of the original electron or positron direction are absorbed into the  $e^+/e^-$  four-momentum.

### IV. SEARCH FOR $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$ AND $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$ AT $\sqrt{s} = 10.52, 10.58$ , AND 10.867 GEV WITH $Z_c^+ \rightarrow \pi^+ J/\psi$

We search for the production of double- $Z_c(3900)$  states, double- $Z_c(4200)$  states, and  $Z_c(3900)$ -plus- $Z_c(4200)$  states in  $\Upsilon(1S, 2S)$  decays and  $e^+e^-$  annihilation at  $\sqrt{s} = 10.52, 10.58$ , and 10.867 GeV, where one  $Z_c$  decays into  $\pi^+$  and  $J/\psi$  ( $\rightarrow \ell^+ \ell^-$ ) and the other decays inclusively.

After applying the aforementioned common event selections, the invariant-mass distributions of the  $J/\psi$  candidates from the five data samples are shown in Figs. 1(a)–1(e). Clear  $J/\psi$  signals are observed. The  $J/\psi$  signal region is defined as  $|M_{\ell^+\ell^-} - m_{J/\psi}| < 0.03$  GeV/ $c^2$  ( $> 98\%$  signal events are reserved according to the MC simulation), where  $m_{J/\psi}$  is the nominal mass [23], while the  $J/\psi$  mass sidebands are  $2.97 < M_{\ell^+\ell^-} < 3.03$  GeV/ $c^2$  or  $3.17 < M_{\ell^+\ell^-} < 3.23$  GeV/ $c^2$  (twice the width of the signal region). In order to improve the  $J/\psi$  momentum resolution, a mass-constrained fit is applied to the  $J/\psi$  candidates in the signal region.

For the events with the lepton-pair mass within the  $J/\psi$  signal region, Figs. 2(a)–2(c) show the recoil-mass spectra of the  $Z_c^+ (\rightarrow \pi^+ J/\psi)$  states from  $\Upsilon(1S)$  decays in signal MC samples. The  $Z_c^+$  shapes are described by Breit-Wigner (BW) functions convolved with Gaussian functions. The solid arrows show the required recoil-mass signal region.

To suppress the background level effectively for  $\Upsilon(1S)$  decays, we require the recoil mass of  $\pi^+ J/\psi$ ,  $M_{\text{miss}}(\pi^+ J/\psi)$ , to satisfy  $|M_{\text{miss}}(\pi^+ J/\psi) - m_{Z_c^-(3900)}| < 0.03$  GeV/ $c^2$  for the double- $Z_c(3900)$  mode or  $|M_{\text{miss}}(\pi^+ J/\psi) - m_{Z_c^-(4200)}| < 0.21$  GeV/ $c^2$  for the double- $Z_c(4200)$  and  $Z_c(3900)$ -plus- $Z_c(4200)$  modes, where  $m_{Z_c^-(3900)}$  and  $m_{Z_c^-(4200)}$  are the nominal masses [23] and  $M_{\text{miss}}(\pi^+ J/\psi) = \sqrt{(p_{e^+e^-} - p_{\pi^+ J/\psi})^2}$  ( $p_{e^+e^-}$  and  $p_{\pi^+ J/\psi}$

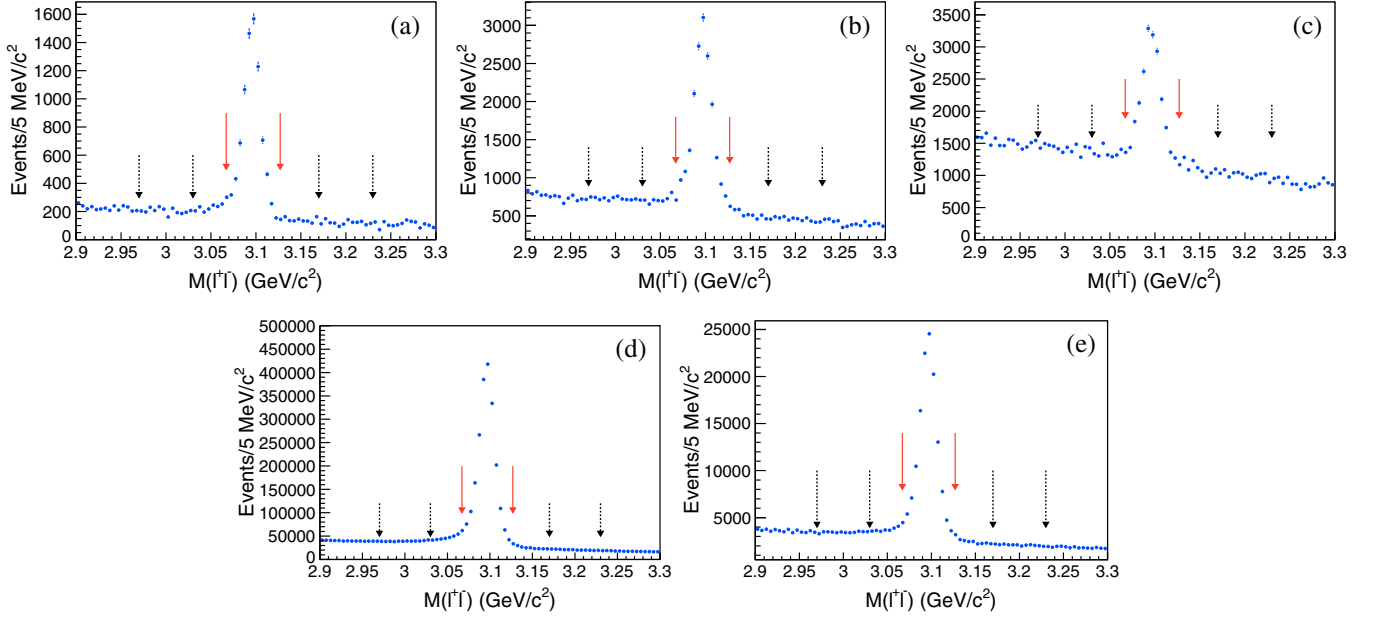


FIG. 1. The invariant-mass distributions of the  $J/\psi$  candidates from (a)  $\Upsilon(1S)$ , (b)  $\Upsilon(2S)$ , (c)  $\sqrt{s} = 10.52$  GeV, (d)  $\sqrt{s} = 10.58$  GeV, and (e)  $\sqrt{s} = 10.867$  GeV data samples. The solid red arrows show the  $J/\psi$  signal region, and the dashed black arrows show the  $J/\psi$  mass sideband regions.

are the four-momenta of the  $e^+e^-$  and  $\pi^+J/\psi$  systems). These requirements maximize  $S/\sqrt{S+B}$ , where  $S$  is the number of fitted signal events in signal MC samples assuming  $\mathcal{B}(\Upsilon(1S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi) = 10^{-5}$  and  $B$  is the number of estimated background events in the  $Z_c^+$  signal region using inclusive MC samples. The same requirements also maximize  $S/\sqrt{S+B}$  for  $\Upsilon(2S)$  decays and  $e^+e^-$  reactions at  $\sqrt{s} = 10.52, 10.58,$  and  $10.867$  GeV.

After the application of the above requirements, Figs. 3(a)–3(c) show the invariant-mass distributions of the  $Z_c^+$  states from  $\Upsilon(1S)$  decays in the MC signal sample, where the solid lines show the fitted results with a BW function convolved with a Gaussian function as  $Z_c^+$  signal shapes. Here, all combinations of  $\pi^+J/\psi$  are retained, as is done for the modes  $\pi^+\chi_{c1}$  and  $\pi^+\psi(2S)$ . Fewer than 1% of events have multiple entries here. Based on the fitted

results, the signal selection efficiencies are obtained and listed in Table II. Note that the efficiencies of the double- $Z_c(3900)$  and double- $Z_c(4200)$  modes are determined from the sum of the reconstructed  $Z_c^+$  and  $Z_c^-$  signals. This applies as well for double- $Z_c$  production in the other data sets with the same  $Z_c^-$ .

Figures 4(a)–4(o) show the  $\pi^+J/\psi$  invariant-mass spectra from the five data samples together with the backgrounds from the normalized  $J/\psi$  mass sideband events. There are no evident signals for  $Z_c^+$  states at the expected positions. No peaking backgrounds are found in the  $J/\psi$  mass sideband events or in the  $\Upsilon(1S, 2S)$  and  $e^+e^- \rightarrow q\bar{q}$  ( $q = u, d, s, c$ ) inclusive MC samples. There is a wide background enhancement at around  $4.5$   $\text{GeV}/c^2$  in the  $\pi^+J/\psi$  invariant-mass distribution for  $e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$  at  $\sqrt{s} = 10.58$  GeV, as shown in Fig. 4(k), arising from  $B$  decays to a  $J/\psi$  meson.

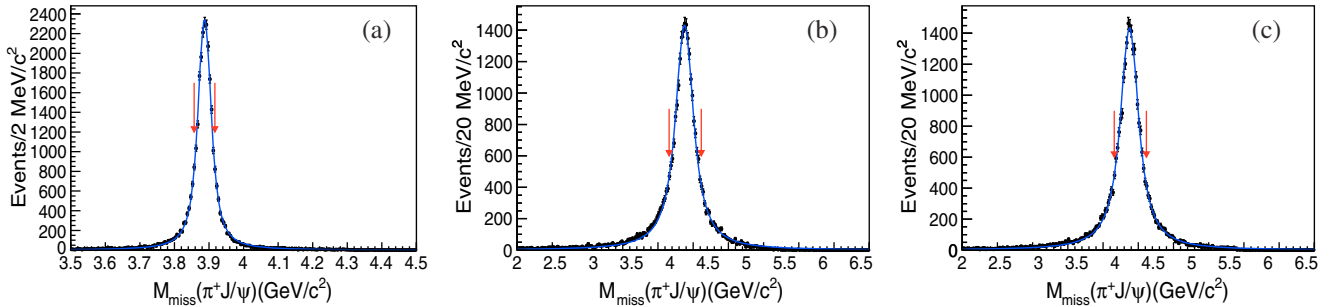


FIG. 2. The recoil-mass spectra of the  $Z_c^+$  ( $\rightarrow \pi^+J/\psi$ ) states from MC simulated (a)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(3900)$ , (b)  $\Upsilon(1S) \rightarrow Z_c^+(4200)Z_c^-(4200)$ , and (c)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$

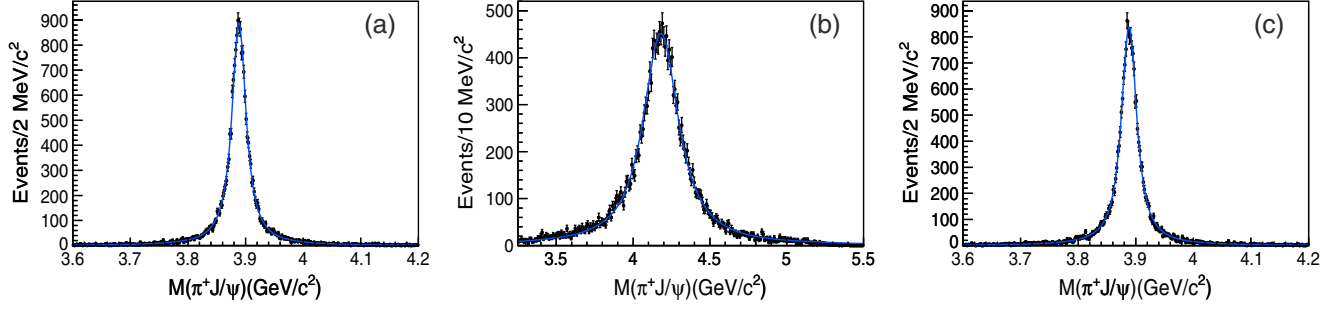


FIG. 3. The invariant-mass distributions of the  $Z_c^+$  ( $\rightarrow \pi^+ J/\psi$ ) states from MC simulated (a)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(3900)$ , (b)  $\Upsilon(1S) \rightarrow Z_c^+(4200)Z_c^-(4200)$ , and (c)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$

Although these background events can be removed by requiring that the  $J/\psi$  momentum in the  $e^+e^-$  c.m. frame be larger than  $2.5 \text{ GeV}/c$  ( $P_{J/\psi}^* > 2.5 \text{ GeV}/c$ ), we retain them since the requirement of  $P_{J/\psi}^* > 2.5 \text{ GeV}/c$  decreases the signal efficiency by 40%, and this distinct feature is well described by a third-order Chebyshev polynomial function.

An unbinned extended maximum likelihood fit to each  $\pi^+ J/\psi$  invariant-mass spectrum is performed to extract the signal and background yields in the five data samples. The  $Z_c^+$  signal shapes used in the fits are BW functions

convolved with a Gaussian function as signal probability density functions [the parameters of the BW function being fixed to the nominal masses and widths of  $Z_c^+$  states and Gaussian functions being fixed to those from the fits to MC signal samples with mass resolutions of 4 and 6  $\text{MeV}/c^2$  for  $Z_c^+(3900)$  and  $Z_c^+(4200)$ , respectively]. For the backgrounds, a second-order Chebyshev polynomial function is adopted, except for  $e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$  at  $\sqrt{s} = 10.58 \text{ GeV}$ , where a third-order polynomial is used. The fitted results are shown in Fig. 4 and summarized in Tables II and III.

TABLE II. Summary of the 90% C.L. upper limits on  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  for  $Z_c^+(3900)Z_c^-(3900)$ ,  $Z_c^+(4200)Z_c^-(4200)$ , and  $Z_c^+(3900)Z_c^-(4200) + c.c.$ , where  $N^{\text{fit}}$  is the signal yield,  $N^{\text{UL}}$  is the 90% C.L. upper limit on the number of signal events,  $\epsilon(\%)$  is the selection efficiency,  $\Sigma(\sigma)$  is the statistical signal significance,  $\sigma_{\text{syst}}(\%)$  is the total systematic uncertainty, and  $\mathcal{B}$  and  $\mathcal{B}^{\text{UL}}$  are the branching fraction and the corresponding 90% C.L. upper limit on the branching fraction in units of  $10^{-6}$ .

Mode	$N^{\text{fit}}$	$N^{\text{UL}}$	$\epsilon(\%)$	$\Sigma(\sigma)$	$\sigma_{\text{syst}}(\%)$	$\mathcal{B}(\Upsilon \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$	$\mathcal{B}^{\text{UL}}(\Upsilon \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$
$\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(3900)$	$0.9 \pm 4.3$	9.2	44.1	0.2	25.2	$0.2 \pm 0.7$	1.8
$\Upsilon(1S) \rightarrow Z_c^+(4200)Z_c^-(4200)$	$50.9 \pm 42.4$	117.1	44.6	1.2	27.8	$9.4 \pm 8.3$	22.3
$\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$	$3.0 \pm 10.1$	22.0	22.6	0.3	20.5	$1.1 \pm 3.7$	8.1
$\Upsilon(2S) \rightarrow Z_c^+(3900)Z_c^-(3900)$	$-1.7 \pm 3.0$	7.4	41.1	...	16.9	$-0.2 \pm 0.4$	1.0
$\Upsilon(2S) \rightarrow Z_c^+(4200)Z_c^-(4200)$	$58.0 \pm 47.9$	125.2	42.2	1.2	31.4	$7.3 \pm 6.4$	16.7
$\Upsilon(2S) \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$	$11.2 \pm 11.5$	29.2	21.3	1.0	16.8	$2.8 \pm 2.9$	7.3

TABLE III. Summary of the 90% C.L. upper limits on  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  for  $Z_c^+(3900)Z_c^-(3900)$ ,  $Z_c^+(4200)Z_c^-(4200)$ , and  $Z_c^+(3900)Z_c^-(4200) + c.c.$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867 \text{ GeV}$ , where  $N^{\text{fit}}$  is the signal yield,  $N^{\text{UL}}$  is the 90% C.L. upper limit on the number of signal events,  $\epsilon(\%)$  is the selection efficiency,  $\Sigma(\sigma)$  is the statistical signal significance,  $\sigma_{\text{syst}}(\%)$  is the total systematic uncertainty,  $\sigma$  is the Born cross section  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-})$ , and  $\sigma^{\text{UL}}$  is the corresponding 90% C.L. upper limit in units of fb.

Mode	$\sqrt{s} \text{ (GeV)}$	$N^{\text{fit}}$	$N^{\text{UL}}$	$\epsilon(\%)$	$\Sigma(\sigma)$	$\sigma_{\text{syst}}(\%)$	$\sigma \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$	$\sigma^{\text{UL}} \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$	10.52	$-4.9 \pm 3.6$	7.2	41.5	...	10.3	$-1.6 \pm 1.2$	2.3
$e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$	10.52	$-27.5 \pm 57.8$	82.8	43.7	...	34.2	$-8.5 \pm 18.1$	26.5
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$	10.52	$-0.5 \pm 15.0$	28.4	21.0	...	22.9	$-0.3 \pm 9.7$	18.3
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$	10.58	$11.8 \pm 13.0$	32.2	41.5	0.9	12.7	$0.5 \pm 0.5$	1.3
$e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$	10.58	$132.1 \pm 173.0$	390.1	43.4	0.8	35.4	$5.1 \pm 6.9$	15.5
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$	10.58	$-7.7 \pm 39.4$	63.4	20.8	...	20.7	$-0.6 \pm 3.2$	5.1
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$	10.867	$-1.4 \pm 4.6$	9.0	41.5	...	17.0	$-0.3 \pm 1.1$	2.2
$e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$	10.867	$-0.2 \pm 41.6$	93.7	43.7	...	33.2	$-0.1 \pm 9.4$	21.9
$e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + c.c.$	10.867	$30.3 \pm 16.7$	53.9	20.5	1.9	16.3	$14.6 \pm 8.4$	26.6

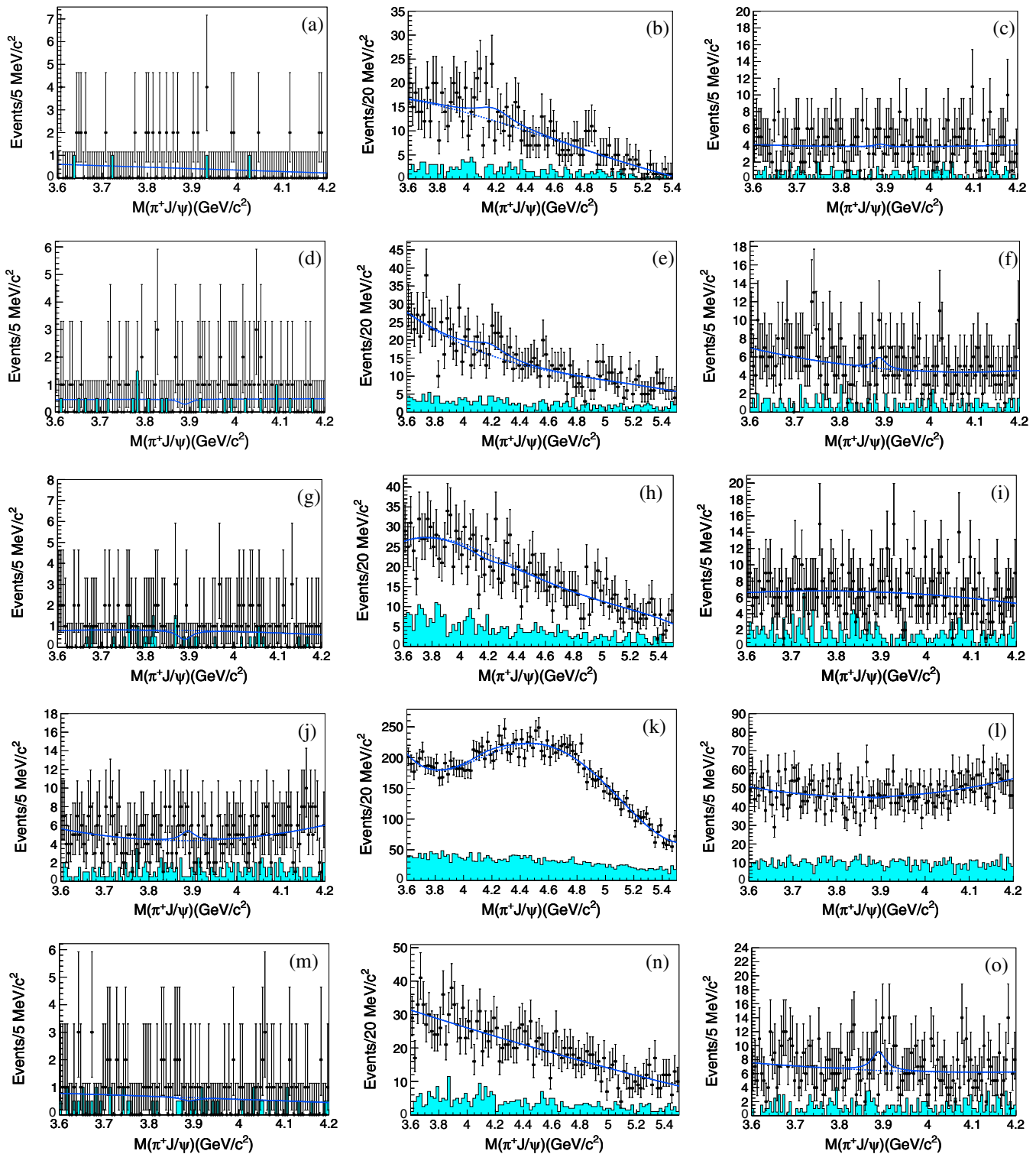


FIG. 4. The invariant-mass spectra of the  $Z_c^+$  ( $\rightarrow \pi^+ J/\psi$ ) states in the processes (a)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(3900)$ ; (b)  $\Upsilon(1S) \rightarrow Z_c^+(4200)Z_c^-(4200)$ ; (c)  $\Upsilon(1S) \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$ ; (d)  $\Upsilon(2S) \rightarrow Z_c^+(3900)Z_c^-(3900)$ ; (e)  $\Upsilon(2S) \rightarrow Z_c^+(4200)Z_c^-(4200)$ ; (f)  $\Upsilon(2S) \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$ ; (g)–(i)  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$ ,  $e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$ , and  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$  at  $\sqrt{s} = 10.52$  GeV; (j)–(l)  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$ ,  $e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$ , and  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$  at  $\sqrt{s} = 10.58$  GeV; and (m)–(o)  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(3900)$ ,  $e^+e^- \rightarrow Z_c^+(4200)Z_c^-(4200)$ , and  $e^+e^- \rightarrow Z_c^+(3900)Z_c^-(4200) + \text{c.c.}$  at  $\sqrt{s} = 10.867$  GeV. The solid curves are the best fits described in the text, the dotted lines are the fitted backgrounds, and the shaded histograms are from the normalized  $J/\psi$  mass sideband events.

The  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  and  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  are calculated using

$$\frac{N_{\text{fit}}}{N_{\Upsilon(1S,2S)} \times \varepsilon \times \mathcal{B}_{\text{decay}}} \quad (1)$$

and

$$\frac{N_{\text{fit}} \times |1 - \prod|^2}{\mathcal{L} \times \mathcal{B}_{\text{decay}} \times \varepsilon \times (1 + \delta)_{\text{ISR}}}, \quad (2)$$

respectively, where  $N_{\text{fit}}$  is the fitted  $Z_c^+$  signal yield,  $N_{\Upsilon(1S,2S)}$  is the total number of  $\Upsilon(1S, 2S)$  events,  $\varepsilon$  is the corresponding selection efficiency,  $\mathcal{L}$  is the integrated luminosity,  $\mathcal{B}_{\text{decay}} = \mathcal{B}(J/\psi \rightarrow \ell^+ \ell^-)$  [23],  $(1 + \delta)_{\text{ISR}}$  is the radiative correction factor, and  $|1 - \prod|^2$  is the vacuum polarization factor. The radiative correction factors  $(1 + \delta)_{\text{ISR}}$  are 0.650, 0.657, and 0.654 for  $\sqrt{s} = 10.52$ , 10.58, and 10.867 GeV, respectively, calculated using formulas given in Ref. [27]; the corresponding values of  $|1 - \prod|^2$  are 0.931, 0.930, and 0.929 [28]. In the calculation of  $(1 + \delta)_{\text{ISR}}$ , we assume that the  $s$  dependence of the cross section is  $\sigma(e^+e^- \rightarrow A\bar{A}) \propto 1/s^2$ . The product branching fractions and the product of the Born cross section and the branching fraction for the studied modes are listed in Tables II and III.

The statistical significances of the  $Z_c^+$  signals are calculated using  $\sqrt{-2 \ln(\mathcal{L}_0/\mathcal{L}_{\text{max}})}$ , where  $\mathcal{L}_0$  and  $\mathcal{L}_{\text{max}}$  are the likelihoods of the fits without and with signal, respectively. The values are summarized in Tables II and III. Since the statistical significance in each case is less than  $3\sigma$ , upper limits on the signal yields ( $N^{\text{UL}}$ ), the product branching fractions [ $\mathcal{B}^{\text{UL}}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$ ], and the product of the Born cross section and the branching fraction [ $\sigma^{\text{UL}}(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$ ], are determined at the 90% credibility level (C.L.) [29] by solving the equation  $\int_0^{\text{UL}} \mathcal{L}(x) dx / \int_0^{\infty} \mathcal{L}(x) dx = 0.9$ , where  $x$  is the assumed signal yield, product branching fraction, or product of the Born cross section and the branching fraction and  $\mathcal{L}(x)$  is the corresponding maximized likelihood for the data. To take into account the systematic uncertainties discussed in Sec. VI, the likelihood is convolved with a Gaussian function of which the width equals the corresponding systematic uncertainty.

The determined 90% C.L. upper limits of  $N^{\text{UL}}$ ,  $\mathcal{B}^{\text{UL}}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$ , and  $\sigma^{\text{UL}}(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  at  $\sqrt{s} = 10.52$ , 10.58, and 10.867 GeV are listed in Tables II and III, together with the signal yields ( $N^{\text{fit}}$ ), the selection efficiencies ( $\varepsilon$ ), the statistical significances ( $\Sigma$ ), the systematic uncertainties ( $\sigma_{\text{sys}}$ ) (discussed below), and the central values of  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$  and  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ J/\psi)$ , with the total errors being the sum in quadrature of the statistical and systematic errors.

## V. SEARCH FOR $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$ AND $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$ AT $\sqrt{s} = 10.52, 10.58$ , AND 10.867 GEV WITH $Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P), \pi^+ \psi(2S)$

In the five data sets, we search for the production of double- $Z_{c1}(4050)$  states, double- $Z_{c2}(4250)$  states, and  $Z_{c1}(4050)$ -plus- $Z_{c2}(4250)$  states, where one of the  $Z_c$  decays into  $\pi^+$  and  $\chi_{c1}(1P)$  and the other decays inclusively, and double- $Z_c(4050)$  states, double- $Z_c(4430)$  states, and  $Z_c(4050)$ -plus- $Z_c(4430)$  states, where one of the  $Z_c$  decays into  $\pi^+$  and  $\psi(2S)$  and the other decays inclusively. The  $\chi_{c1}(1P)$  and  $\psi(2S)$  are reconstructed via their decays into  $\gamma J/\psi$  and  $\pi^+ \pi^- J/\psi$ , respectively, with  $J/\psi \rightarrow \ell^+ \ell^-$ .

After requiring the mass of the lepton pair to be within the  $J/\psi$  signal region ( $|M_{\ell^+ \ell^-} - m_{J/\psi}| < 0.03 \text{ GeV}/c^2$ ), Figs. 5(a)–5(e) and 6(a)–6(e) show the invariant-mass distributions of the  $\chi_{c1}(1P)$  and  $\psi(2S)$  candidates from five data samples. Clear  $\chi_{c1}(1P)$  signals could be seen in  $\Upsilon(1S)$  decay and in  $e^+e^-$  annihilation at  $\sqrt{s} = 10.58$  and 10.867 GeV, and evidence for  $\chi_{c1}(1P)$  is observed in  $\Upsilon(2S)$  decay; clear  $\psi(2S)$  signals are observed in  $\Upsilon(1S, 2S)$  decay and in  $e^+e^-$  annihilation at  $\sqrt{s} = 10.52$ , 10.58, and 10.867 GeV. We define the  $\chi_{c1}(1P)$  and  $\psi(2S)$  signal and sideband regions as follows. The  $\chi_{c1}(1P)$  signal region is  $|M_{\gamma J/\psi} - m_{\chi_{c1}(1P)}| < 0.04 \text{ GeV}/c^2$  ( $> 98\%$  signal events are reserved according to the MC simulation), where  $m_{\chi_{c1}(1P)}$  is the nominal mass [23], while the  $\chi_{c1}(1P)$  mass sidebands are  $3.35 < M_{\gamma J/\psi} < 3.43 \text{ GeV}/c^2$  or  $3.62 < M_{\gamma J/\psi} < 3.70 \text{ GeV}/c^2$  (twice the width of the signal region). The  $\psi(2S)$  signal region is  $|M_{\pi^+ \pi^- J/\psi} - m_{\psi(2S)}| < 0.007 \text{ GeV}/c^2$  ( $> 98\%$  signal events are reserved according to the MC simulation), where  $m_{\psi(2S)}$  is the nominal mass [23], while the  $\psi(2S)$  mass sidebands are  $3.65 < M_{\pi^+ \pi^- J/\psi} < 3.664 \text{ GeV}/c^2$  or  $3.71 < M_{\pi^+ \pi^- J/\psi} < 3.724 \text{ GeV}/c^2$  (twice the width of the signal region).

As before, for  $\Upsilon(1S)$  decays, we optimize the requirements of recoil mass of  $\pi^+ \chi_{c1}(1P)$  and  $\pi^+ \psi(2S)$  by maximizing  $S/\sqrt{S+B}$ . The optimized requirements are  $|M_{\text{miss}}(\pi^+ \chi_{c1}(1P)) - m_{Z_{c1}^-(4050)}| < 0.08 \text{ GeV}/c^2$  for the  $Z_{c1}^+(4050)Z_{c1}^-(4050)$  mode and  $|M_{\text{miss}}(\pi^+ \chi_{c1}(1P)) - m_{Z_{c2}^-(4250)}| < 0.13 \text{ GeV}/c^2$  for the  $Z_{c2}^+(4250)Z_{c2}^-(4250)$  and  $Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$  modes; they are  $|M_{\text{miss}}(\pi^+ \psi(2S)) - m_{Z_c^-(4050)}| < 0.06 \text{ GeV}/c^2$  for the  $Z_c^+(4050)Z_c^-(4050)$  mode and  $|M_{\text{miss}}(\pi^+ \psi(2S)) - m_{Z_c^-(4430)}| < 0.12 \text{ GeV}$  for the  $Z_c^+(4430)Z_c^-(4430)$  and  $Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  modes. For  $\Upsilon(2S)$  decays and  $e^+e^-$  reactions at  $\sqrt{s} = 10.52$ , 10.58, and 10.867 GeV, the optimized requirements are the same as  $\Upsilon(1S)$ .

After the application of these requirements, Figs. 7(a)–7(o) and 8(a)–8(o) show the  $\pi^+ \chi_{c1}(1P)$  and  $\pi^+ \psi(2S)$  invariant-mass spectra from the five data samples, together with the backgrounds from the normalized  $\chi_{c1}(1P)$  or  $\psi(2S)$  mass sideband events. There are no evident signals for  $Z_c^+$  states at the expected position. No peaking backgrounds are found in the  $\chi_{c1}(1P)$  or  $\psi(2S)$

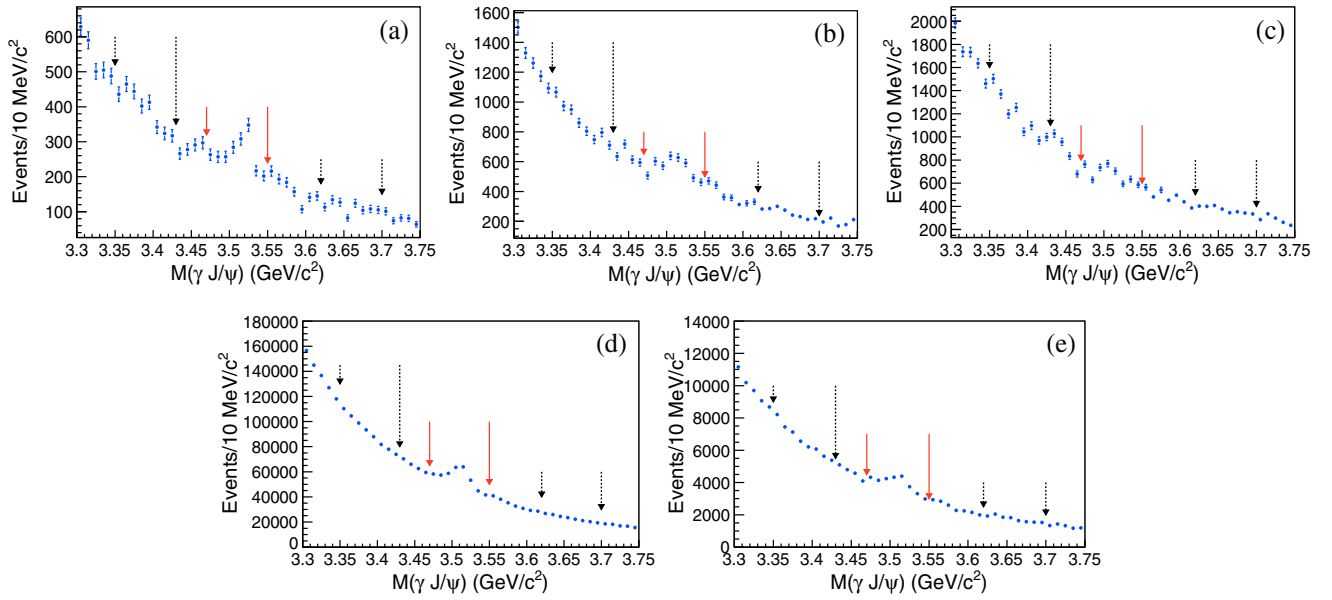


FIG. 5. The invariant-mass distributions of the  $\chi_{c1}(1P)$  candidates from (a)  $\Upsilon(1S)$ , (b)  $\Upsilon(2S)$ , (c)  $\sqrt{s} = 10.52$  GeV, (d)  $\sqrt{s} = 10.58$  GeV, and (e)  $\sqrt{s} = 10.867$  GeV data samples. The solid red arrows show the  $\chi_{c1}(1P)$  signal region, and the dashed black arrows show the  $\chi_{c1}(1P)$  mass sideband regions.

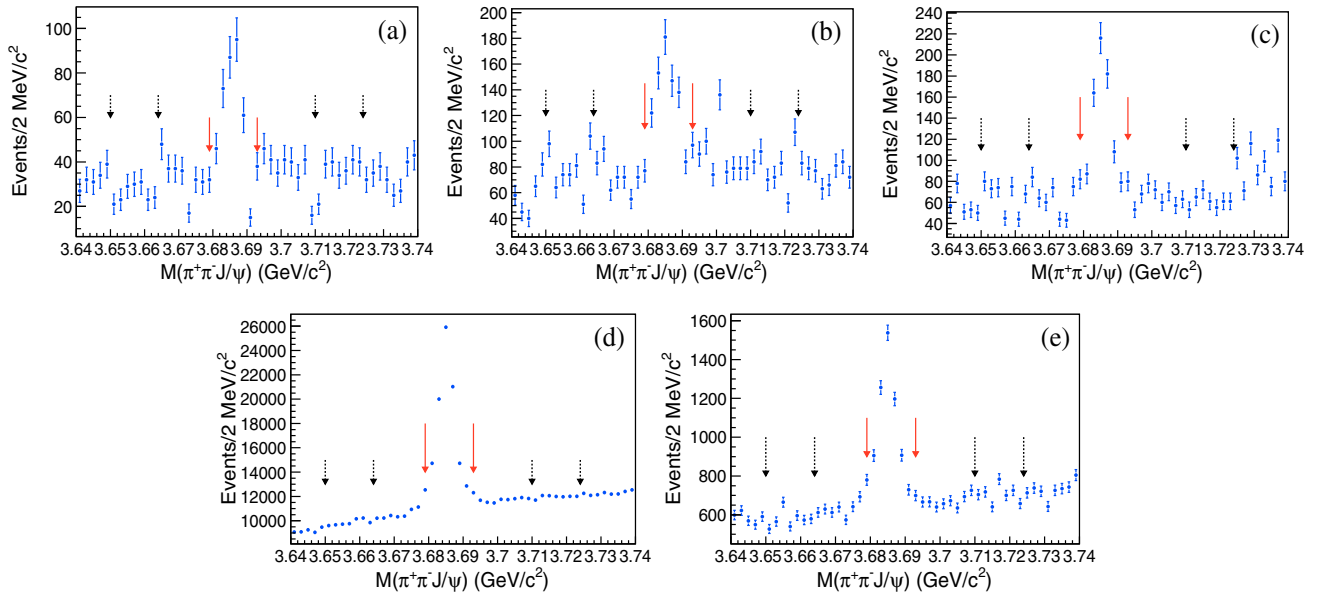


FIG. 6. The invariant-mass distributions of the  $\psi(2S)$  candidates from (a)  $\Upsilon(1S)$ , (b)  $\Upsilon(2S)$ , (c)  $\sqrt{s} = 10.52$  GeV, (d)  $\sqrt{s} = 10.58$  GeV, and (e)  $\sqrt{s} = 10.867$  GeV data samples. The solid red arrows show the  $\psi(2S)$  signal region, and the dashed black arrows show the  $\psi(2S)$  mass sideband regions.

mass sideband events or in the  $\Upsilon(1S, 2S)$  and  $e^+e^- \rightarrow q\bar{q}$  inclusive MC samples.

An unbinned extended maximum likelihood fit to the  $\pi^+\chi_{c1}(1P)$  ( $\pi^+\psi(2S)$ ) invariant-mass spectra is performed to extract the signal and background yields from  $\Upsilon(1S, 2S)$  decays and  $e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  reactions at  $\sqrt{s} = 10.52$ , 10.58, and 10.867 GeV. The  $Z_c^+$  signal probability density

functions used in the fits are the convolutions of a BW function with a Gaussian function [the parameters of BW function being fixed to the nominal masses and widths of  $Z_c^+$  states and Gaussian functions being fixed to those from the fits to signal MC samples with mass resolutions of 15, 16, 10, and 11  $\text{MeV}/c^2$  for  $Z_{c1}^+(4050)$ ,  $Z_{c2}^+(4250)$ ,  $Z_c^+(4050)$ , and  $Z_c^+(4250)$ , respectively]. For

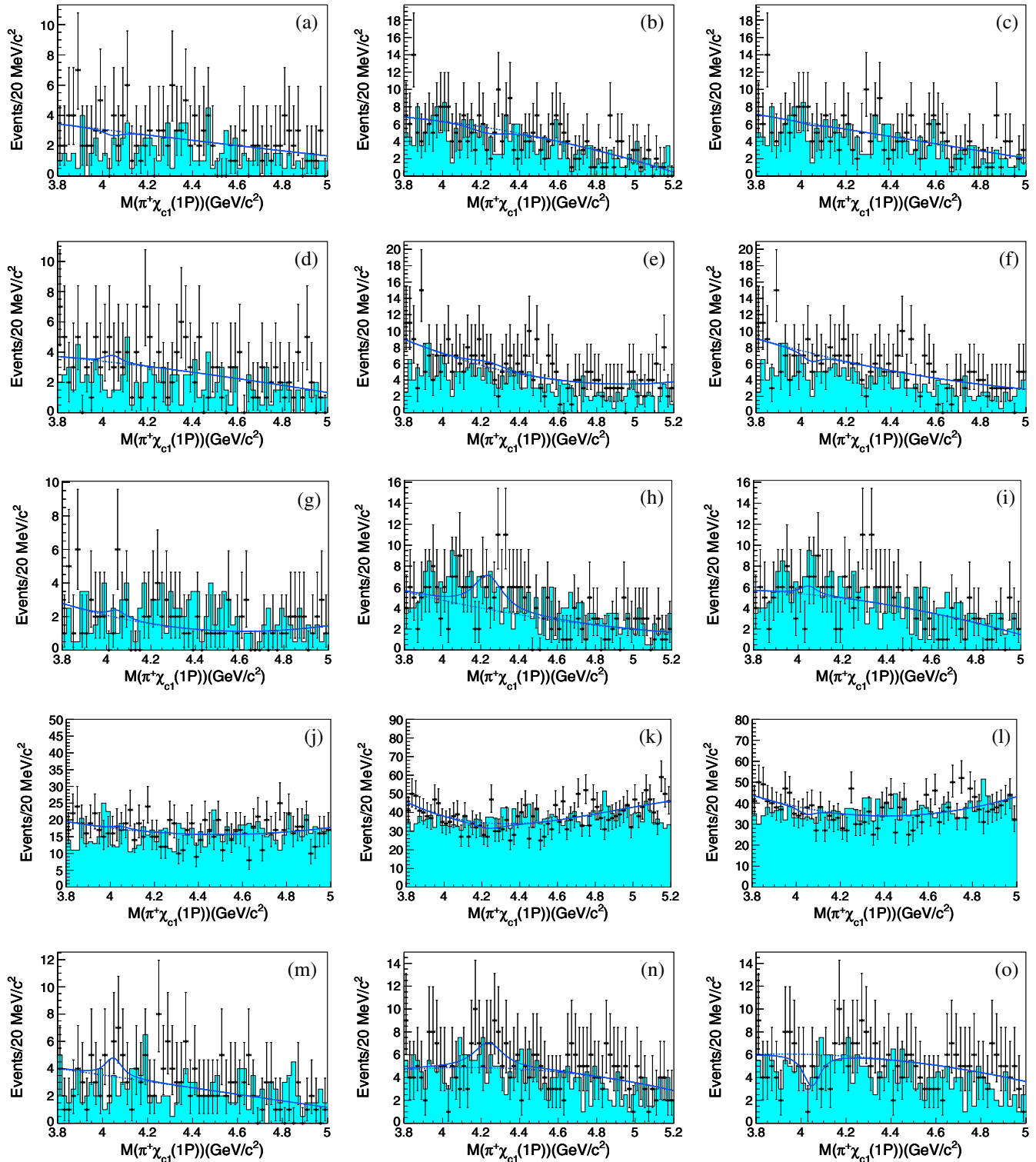


FIG. 7. The invariant-mass spectra of the  $Z_c^+$  [ $\rightarrow \pi^+\chi_{c1}(1P)$ ] states in the processes (a)  $\Upsilon(1S) \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$ ; (b)  $\Upsilon(1S) \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$ ; (c)  $\Upsilon(1S) \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$ ; (d)  $\Upsilon(2S) \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$ ; (e)  $\Upsilon(2S) \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$ ; (f)  $\Upsilon(2S) \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$ ; (g)–(i)  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$ ,  $e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$ , and  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$  at  $\sqrt{s} = 10.52$  GeV; (j)–(l)  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$ ,  $e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$ , and  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$  at  $\sqrt{s} = 10.58$  GeV; and (m)–(o)  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$ ,  $e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$ , and  $e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250)$  at  $\sqrt{s} = 10.867$  GeV. The solid curves are the best fits described in the text, the dotted lines are the fitted backgrounds, and the shaded histograms are from the normalized  $\chi_{c1}(1P)$  mass sideband events.

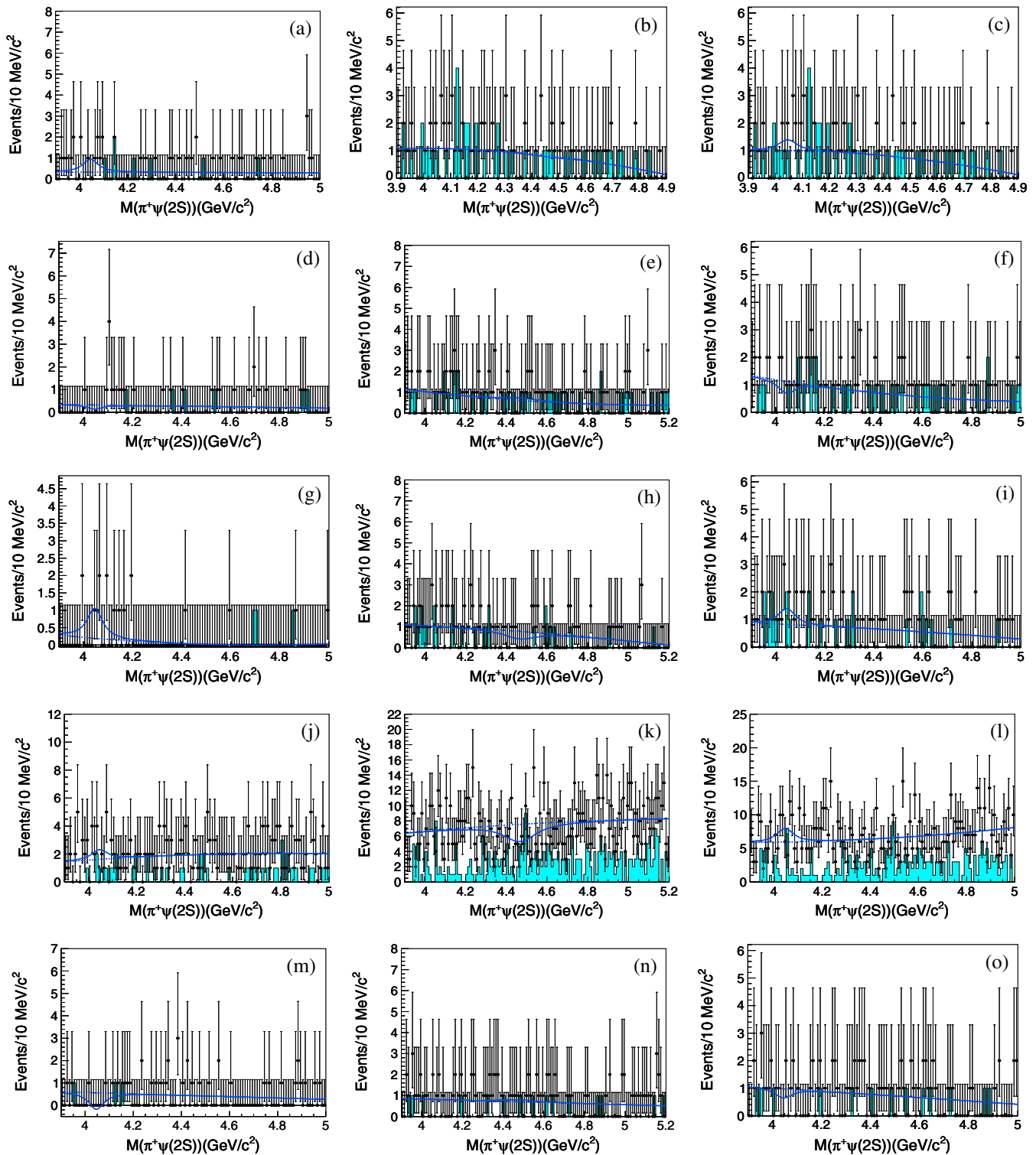


FIG. 8. The invariant-mass spectra of the  $Z_c^+$  [ $\rightarrow \pi^+\psi(2S)$ ] states in the processes (a)  $\Upsilon(1S) \rightarrow Z_c^+(4050)Z_c^-(4050)$ ; (b)  $\Upsilon(1S) \rightarrow Z_c^+(4430)Z_c^-(4430)$ ; (c)  $\Upsilon(1S) \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$ ; (d)  $\Upsilon(2S) \rightarrow Z_c^+(4050)Z_c^-(4050)$ ; (e)  $\Upsilon(2S) \rightarrow Z_c^+(4430)Z_c^-(4430)$ ; (f)  $\Upsilon(2S) \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$ ; (g)–(i)  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$ ,  $e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$ , and  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  at  $\sqrt{s} = 10.52$  GeV; (j)–(l)  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$ ,  $e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$ , and  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  at  $\sqrt{s} = 10.58$  GeV; and (m)–(o)  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$ ,  $e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$ , and  $e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  at  $\sqrt{s} = 10.867$  GeV. The solid curves are the best fits described in the text, the dotted lines are the fitted backgrounds, and the shaded histograms are from the normalized  $\psi(2S)$  mass sideband events.

TABLE IV. Summary of the 90% C.L. upper limits on  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$  for  $Z_{c1}^+(4050)Z_{c1}^-(4050)$ ,  $Z_{c2}^+(4250)Z_{c2}^-(4250)$ , and  $Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$  to  $\pi^+ \chi_{c1}(1P) + \text{anything}$  and  $Z_c^+(4050)Z_c^-(4050)$ ,  $Z_c^+(4430)Z_c^-(4430)$ , and  $Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  to  $\pi^+ \psi(2S) + \text{anything}$ , where  $N^{\text{fit}}$  is the signal yield,  $N^{\text{UL}}$  is the 90% C.L. upper limit on the number of signal events,  $\epsilon(\%)$  is the selection efficiency,  $\Sigma(\sigma)$  is the statistical signal significance,  $\sigma_{\text{syst}}(\%)$  is the total systematic uncertainty, and  $\mathcal{B}$  and  $\mathcal{B}^{\text{UL}}$  are the branching fraction and the corresponding 90% C.L. upper limit on the branching fraction in units of  $10^{-6}$ .

Mode	$N^{\text{fit}}$	$N^{\text{UL}}$	$\epsilon(\%)$	$\Sigma(\sigma)$	$\sigma_{\text{syst}}(\%)$	$\mathcal{B}(\Upsilon \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$	$\mathcal{B}^{\text{UL}}(\Upsilon \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$
$\Upsilon(1S) \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$	$-2.1 \pm 7.2$	13.1	21.6	...	41.3	$-2.4 \pm 8.1$	15.8
$\Upsilon(1S) \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$	$-8.3 \pm 14.1$	21.7	20.9	...	34.5	$-9.7 \pm 16.8$	26.6
$\Upsilon(1S) \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$	$-1.3 \pm 10.2$	18.1	10.1	...	25.9	$-3.1 \pm 24.3$	44.2
$\Upsilon(2S) \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$	$2.9 \pm 7.7$	16.4	20.5	0.2	38.7	$2.2 \pm 5.9$	13.5
$\Upsilon(2S) \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$	$8.1 \pm 15.8$	32.6	19.2	0.5	34.0	$6.6 \pm 13.2$	26.7
$\Upsilon(2S) \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$	$-6.3 \pm 10.5$	16.1	9.4	...	18.0	$-10.5 \pm 17.6$	27.2
$\Upsilon(1S) \rightarrow Z_c^+(4050)Z_c^-(4050)$	$6.7 \pm 5.3$	14.4	16.0	1.4	16.4	$10.0 \pm 8.1$	23.3
$\Upsilon(1S) \rightarrow Z_c^+(4430)Z_c^-(4430)$	$-1.5 \pm 7.6$	13.6	16.2	...	22.0	$-2.2 \pm 11.3$	20.3
$\Upsilon(1S) \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$	$4.4 \pm 5.7$	14.3	7.7	0.8	28.8	$13.6 \pm 18.1$	45.5
$\Upsilon(2S) \rightarrow Z_c^+(4050)Z_c^-(4050)$	$-1.9 \pm 6.4$	10.8	15.1	...	16.1	$-1.9 \pm 6.6$	11.1
$\Upsilon(2S) \rightarrow Z_c^+(4430)Z_c^-(4430)$	$3.4 \pm 9.6$	20.0	15.3	0.3	17.8	$3.4 \pm 9.7$	20.3
$\Upsilon(2S) \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$	$-4.9 \pm 6.2$	10.2	7.5	...	26.1	$-10.1 \pm 13.1$	21.1

the backgrounds, a second-order Chebyshev polynomial function is adopted. The fitted results are shown in Figs. 7 and 8 and summarized in Tables IV and V.

Marginal signals are seen in each fit. The largest statistical signal significance is  $2.6\sigma$  from  $e^+e^- \rightarrow Z_{c2}(4250) + Z_{c2}(4250)$  at  $\sqrt{s} = 10.52$  GeV. We use an

ensemble of simulated experiments to estimate the probability that background fluctuations alone would produce signals as significant as those seen in the data. We generate  $\pi^+ \chi_{c1}(1P)$  mass spectra based on a fitted second-order Chebyshev polynomial function alone with the same observed events in data and search for the most

TABLE V. Summary of the 90% C.L. upper limits on  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$  for  $Z_{c1}^+(4050)Z_{c1}^-(4050)$ ,  $Z_{c2}^+(4250)Z_{c2}^-(4250)$ , and  $Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$  to  $\pi^+ \chi_{c1}(1P) + \text{anything}$  and  $Z_c^+(4050)Z_c^-(4050)$ ,  $Z_c^+(4430)Z_c^-(4430)$ , and  $Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$  to  $\pi^+ \psi(2S) + \text{anything}$  at  $\sqrt{s} = 10.52, 10.58$ , and  $10.867$  GeV where  $N^{\text{fit}}$  is the signal yield,  $N^{\text{UL}}$  is the 90% C.L. upper limit on the number of signal events,  $\epsilon(\%)$  is the selection efficiency,  $\Sigma(\sigma)$  is the statistical signal significance,  $\sigma_{\text{syst}}(\%)$  is the total systematic uncertainty,  $\sigma$  is the Born cross section  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-})$ , and  $\sigma^{\text{UL}}$  is the corresponding 90% C.L. upper limit in units of fb.

Mode	$\sqrt{s}$ (GeV)	$N^{\text{fit}}$	$N^{\text{UL}}$	$\epsilon(\%)$	$\Sigma(\sigma)$	$\sigma_{\text{syst}}(\%)$	$\sigma \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$	$\sigma^{\text{UL}} \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$	10.52	$1.2 \pm 6.5$	13.2	20.9	0.2	28.3	$2.3 \pm 12.4$	25.0
$e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$	10.52	$40.9 \pm 16.8$	65.1	19.4	2.6	32.9	$83.9 \pm 44.1$	143.9
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$	10.52	$5.2 \pm 10.4$	21.5	9.5	0.5	33.0	$21.7 \pm 44.1$	93.2
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$	10.58	$4.1 \pm 18.9$	36.3	20.5	0.2	21.9	$1.0 \pm 4.6$	8.8
$e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$	10.58	$-35.2 \pm 48.3$	25.7	19.2	...	45.8	$-9.0 \pm 13.1$	7.1
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$	10.58	$-18.0 \pm 24.8$	34.5	9.8	...	45.0	$-9.1 \pm 13.2$	18.2
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c1}^-(4050)$	10.867	$8.6 \pm 8.5$	23.0	19.4	1.0	26.0	$12.9 \pm 13.2$	35.7
$e^+e^- \rightarrow Z_{c2}^+(4250)Z_{c2}^-(4250)$	10.867	$27.7 \pm 16.1$	49.5	18.5	1.7	27.0	$43.6 \pm 28.0$	82.0
$e^+e^- \rightarrow Z_{c1}^+(4050)Z_{c2}^-(4250) + \text{c.c.}$	10.867	$-17.5 \pm 8.6$	9.4	9.1	...	28.5	$-55.7 \pm 31.6$	30.8
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$	10.52	$9.4 \pm 15.5$	18.1	15.0	1.1	23.4	$24.5 \pm 40.8$	47.7
$e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$	10.52	$-9.7 \pm 8.4$	10.5	15.0	...	16.9	$-25.3 \pm 22.3$	29.7
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$	10.52	$6.5 \pm 7.2$	18.7	7.5	0.9	17.3	$33.9 \pm 38.0$	97.9
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$	10.58	$7.7 \pm 9.3$	23.5	15.0	0.7	16.5	$2.5 \pm 3.0$	7.6
$e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$	10.58	$-60.5 \pm 27.8$	22.9	14.6	...	12.7	$-20.1 \pm 9.6$	8.3
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$	10.58	$22.8 \pm 17.2$	48.5	7.3	1.3	19.5	$15.1 \pm 11.8$	32.2
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4050)$	10.867	$-8.0 \pm 3.4$	5.2	14.2	...	20.8	$-16.1 \pm 7.6$	10.8
$e^+e^- \rightarrow Z_c^+(4430)Z_c^-(4430)$	10.867	$2.7 \pm 8.2$	16.7	14.0	0.3	22.1	$5.5 \pm 16.7$	35.2
$e^+e^- \rightarrow Z_c^+(4050)Z_c^-(4430) + \text{c.c.}$	10.867	$-3.7 \pm 5.7$	9.1	7.0	...	21.1	$-15.1 \pm 23.4$	39.1

TABLE VI. Relative systematic errors (%) on the measurements of the branching fractions for  $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(0)-}$  and the Born cross sections for  $e^+ e^- \rightarrow Z_c^+ Z_c^{(0)-}$  at  $\sqrt{s} = 10.52, 10.58,$  and  $10.867$  GeV, where  $Z_c^+ \rightarrow \pi^+ + c\bar{c}$  ( $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ ).

Mode	Tracking	Particle identification	Photon	MC statistics	Branching fraction	Resonance parameter	Fit	Recoil Mass	$N_{\Upsilon(1S)/\Upsilon(2S)}/$ Luminosity	SUM
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(3900) (\rightarrow \pi^+ J/\psi) \times$ $Z_c^-(3900)$	1.4	3.9	...	1.0	1.1	10.1/5.9	20.5/13.8	9.6/6.1	2.0/2.3	25.2/16.9
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4200) (\rightarrow \pi^+ J/\psi) \times$ $Z_c^-(4200)$	1.4	3.9	...	1.0	1.1	18.1/22.1	9.4/10.2	18.1/19.2	2.0/2.3	27.8/31.4
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(3900) (\rightarrow \pi^+ J/\psi) \times$ $Z_c^-(4200) + c.c.$	1.4	3.9	...	1.0	1.1	5.4/6.8	3.6/9.6	18.8/10.9	2.0/2.3	20.5/16.8
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4050) \times$ $(\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4050)$	1.4	3.9	2.0	1.0	3.8	15.3/6.0	35.2/36.8	14.0/8.3	2.0/2.3	41.3/38.7
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4250) \times$ $(\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4250)$	1.4	3.9	2.0	1.0	3.8	20.9/14.2	17.8/14.5	20.0/26.6	2.0/2.3	34.5/34.0
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4050) \times$ $(\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4250) + c.c.$	1.4	3.9	2.0	1.0	3.8	6.4/3.1	7.7/10.1	23.1/13.2	2.0/2.3	25.9/18.0
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4050) \times$ $(\rightarrow \pi^+ \psi(2S)) Z_c^-(4050)$	1.8	4.6	...	1.0	1.5	1.9/1.3	9.6/14.8	12.0/2.5	2.0/2.3	16.5/16.1
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4430) \times$ $(\rightarrow \pi^+ \psi(2S)) Z_c^-(4430)$	1.8	4.6	...	1.0	1.5	7.2/3.8	19.5/13.8	4.7/9.1	2.0/2.3	22.0/17.9
$\Upsilon(1S)/\Upsilon(2S) \rightarrow Z_c^+(4050) \times$ $(\rightarrow \pi^+ \psi(2S)) Z_c^-(4430) + c.c.$	1.8	4.6	...	1.0	1.5	7.4/8.5	16.0/8.8	22.2/22.3	2.0/2.3	28.8/26.1
$e^+ e^- \rightarrow Z_c^+(3900) (\rightarrow \pi^+ J/\psi) \times$ $Z_c^-(3900)$ at 10.52/10.58/10.867 GeV	1.4	3.9	...	1.0	1.1	2.8/5.3/7.2	5.6/9.8/12.1	6.8/3.9/8.4	1.4	10.3/12.7/17.0
$e^+ e^- \rightarrow Z_c^+(4200) (\rightarrow \pi^+ J/\psi) Z_c^-(4200)$ at 10.52/10.58/10.867 GeV	1.4	3.9	...	1.0	1.1	17.8/13.1/11.9	16.8/16.1/27.5	23.4/28.3/13.4	1.4	34.2/35.4/33.2
$e^+ e^- \rightarrow Z_c^+(3900) (\rightarrow \pi^+ J/\psi) Z_c^-(4200) +$ c.c. at 10.52/10.58/10.867 GeV	1.4	3.9	...	1.0	1.1	5.8/7.6/5.0	1.0/1.4/4.7	21.6/18.6/14.1	1.4	22.9/20.7/16.3
$e^+ e^- \rightarrow Z_c^+(4050) (\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4050) \times$ $Z_c^+(4250) (\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4250) \times$ at 10.52/10.58/10.867 GeV	1.4	3.9	2.0	1.0	3.8	5.4/11.7/11.4	21.0/15.7/20.3	17.1/7.5/9.7	1.4	28.3/21.9/26.0
$e^+ e^- \rightarrow Z_c^+(4250) (\rightarrow \pi^+ \chi_{c1}(1P)) Z_c^-(4250) \times$ at 10.52/10.58/10.867 GeV	1.4	3.9	2.0	1.0	3.8	15.7/23.1/15.4	11.8/22.4/12.5	25.6/32.0/17.3	1.4	32.9/45.8/27.0
$e^+ e^- \rightarrow Z_c^+(4050) (\rightarrow \pi^+ \chi_{c1}(1P)) \times$ $Z_c^-(4250) + c.c.$ at 10.52/10.58/10.867 GeV	1.4	3.9	2.0	1.0	3.8	9.1/7.9/10.4	21.9/32.7/12.2	22.1/29.1/22.8	1.4	33.0/45.0/28.5
$e^+ e^- \rightarrow Z_c^+(4050) (\rightarrow \pi^+ \psi(2S)) \times$ $Z_c^-(4050)$ at 10.52/10.58/10.867 GeV	1.8	4.6	...	1.0	1.5	6.8/3.2/12.3	12.1/10.9/12.3	18.1/10.7/10.2	1.4	23.4/16.5/20.9
$e^+ e^- \rightarrow Z_c^+(4430) (\rightarrow \pi^+ \psi(2S)) Z_c^-(4430)$ at 10.52/10.58/10.867 GeV	1.8	4.6	...	1.0	1.5	7.2/5.3/6.1	11.7/6.5/13.3	8.4/7.9/15.6	1.4	17.0/12.8/22.1
$e^+ e^- \rightarrow Z_c^+(4050) (\rightarrow \pi^+ \psi \times$ $(2S)) Z_c^-(4430) + c.c.$ at 10.52/10.58/10.867 GeV	1.8	4.6	...	1.0	1.5	4.3/8.1/5.1	7.5/9.6/9.7	14.0/13.9/17.2	1.4	17.3/19.5/21.1

significant fluctuation in each mass spectrum in the studied mass range. From these pure-background spectra, we obtain the distribution for  $-2\ln(\mathcal{L}_0/\mathcal{L}_{\max})$  and compare it with the signal in the data. In a total of 10,000 simulations, we find 1373 trials with a  $-2\ln(\mathcal{L}_0/\mathcal{L}_{\max})$  value greater than or equal to the value obtained in the data. The resulting  $p$  value is 0.1373, corresponding to a significance of  $1.5\sigma$ .

Using the same method as described in the previous section, we determine the 90% C.L. upper limits on  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$  and  $\sigma(e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV. The values are listed in Tables IV and V, together with the signal yields ( $N^{\text{fit}}$ ), the 90% C.L. upper limits on the numbers of signal events ( $N^{\text{UL}}$ ), the selection efficiencies ( $\epsilon$ ), the signal significances ( $\Sigma$ ), the systematic uncertainties ( $\sigma_{\text{sys}}$ ) discussed below, and the central values of  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$  and  $\sigma(e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)/\pi^+ \psi(2S))$ , with the total errors being the sum in quadrature of the statistical and systematic errors.

## VI. SYSTEMATIC ERRORS

The following sources of systematic errors are taken into account in the branching fraction and Born cross section measurements.

The systematic uncertainty due to charged-track reconstruction is determined from a study of partially reconstructed  $D^{*+} \rightarrow D^0(\rightarrow K_S^0 \pi^+ \pi^-) \pi^+$  decays and is 0.35% per track. Based on the measurements of the particle identification efficiencies of lepton pairs from  $\gamma\gamma \rightarrow \ell^+ \ell^-$  events and pions from a low-background sample of  $D^*$  events, the MC simulation yields uncertainties of 3.6% for each lepton pair and 1.3% for each pion. The photon reconstruction contributes 2.0% per photon, as determined from radiative Bhabha events.

The MC statistical errors are estimated using the yields of selected and generated events; these are 1.0% or less. Errors on the branching fractions of the intermediate states are taken from Ref. [23]. The uncertainties of the branching fractions for  $J/\psi \rightarrow \ell^+ \ell^-$ ,  $\chi_{c1}(1P) \rightarrow \gamma J/\psi$ , and  $\psi(2S) \rightarrow \pi^+ \pi^- J/\psi$  are 1.1%, 3.6% and 0.9%, respectively. The trigger efficiency, evaluated from simulation, is approximately 100% with a negligible uncertainty. We generate MC signal samples by assuming the  $Z_c^+ \rightarrow \pi^+ \chi_{c1}(1P)$  decays are  $P$ -wave and find the differences between  $P$ -wave and  $S$ -wave in the selection efficiencies are less than 1% and are thus neglected.

The difference in the 90% C.L. upper limit on the signal yield when the mass and width of each  $Z_c^\pm$  state are varied by  $1\sigma$  is used as an estimate of the systematic uncertainty

associated with the mass and width uncertainties. By changing the order of the background polynomial and the range of the fit, the decay-mode-dependent relative difference in the 90% C.L. upper limit on signal yields is obtained and taken as the systematic error due to the uncertainty of the fit.

To estimate the systematic uncertainties due to the recoil-mass requirements for  $Z_c^+ \rightarrow \pi^+ + c\bar{c}$  [ $cc = J/\psi, \chi_{c1}(1P), \psi(2S)$ ], we optimize the recoil-mass regions again by changing the  $Z_c^+$  widths by  $1\sigma$  and take the differences in the 90% C.L. upper limits on the signal yields as systematic uncertainties.

Changing the  $s$  dependence of the cross sections of  $e^+ e^- \rightarrow Z_c^+ + Z_c^{(\prime)-}$  from  $1/s^2$  to  $1/s^6$ , the radiative correction factors  $(1 + \delta)_{\text{ISR}}$  become 0.651, 0.659, and 0.655 for  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV, respectively. The differences are less than 1% and are thus neglected.

The uncertainties on the total numbers of  $\Upsilon(1S)$  and  $\Upsilon(2S)$  events are 2.0% and 2.3%, respectively, which are mainly due to imperfect simulations of the charged-track multiplicity distributions from inclusive hadronic MC events. Finally, the total luminosity is determined to 1.4% precision using wide-angle Bhabha scattering events.

All the systematic uncertainties are summarized in Table VI for the measurements of  $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$  and  $e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  at  $\sqrt{s} = 10.52, 10.58, 10.867$  GeV, where  $Z_c^+ \rightarrow \pi^+ + c\bar{c}$  ( $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ ). Assuming all the sources are independent and adding them in quadrature, the total systematic uncertainties are obtained for each mode.

## VII. CONCLUSION

In summary, using the large data samples of  $102 \times 10^6$   $\Upsilon(1S)$  events;  $158 \times 10^6$   $\Upsilon(2S)$  events; and  $89.5 \text{ fb}^{-1}$ ,  $702.6 \text{ fb}^{-1}$ , and  $121.1 \text{ fb}^{-1}$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV, respectively, collected by Belle, we search for  $\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}$  and  $e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV with  $Z_c^+ \rightarrow \pi^+ + c\bar{c}$  [ $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ ]. No clear signals are observed in the studied modes. We determine the 90% C.L. upper limits on  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$  and  $\sigma(e^+ e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$  at  $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867$  GeV. The results are displayed graphically in Figs. 9 and 10. Due to G-parity conservation, our studied processes are electromagnetic, i.e., can only proceed through a virtual photon, which then transforms into a light-quark pair ( $u\bar{u}$ ) or ( $d\bar{d}$ ). In this case, the dynamical suppression is much larger due to the production of two ( $c\bar{c}$ ) pairs. The expected production cross section should be much lower than the double charmonium processes like  $e^+ e^- \rightarrow J/\psi \eta_c$  [30]. The reported upper limits are not in contradiction with the naive expectation.

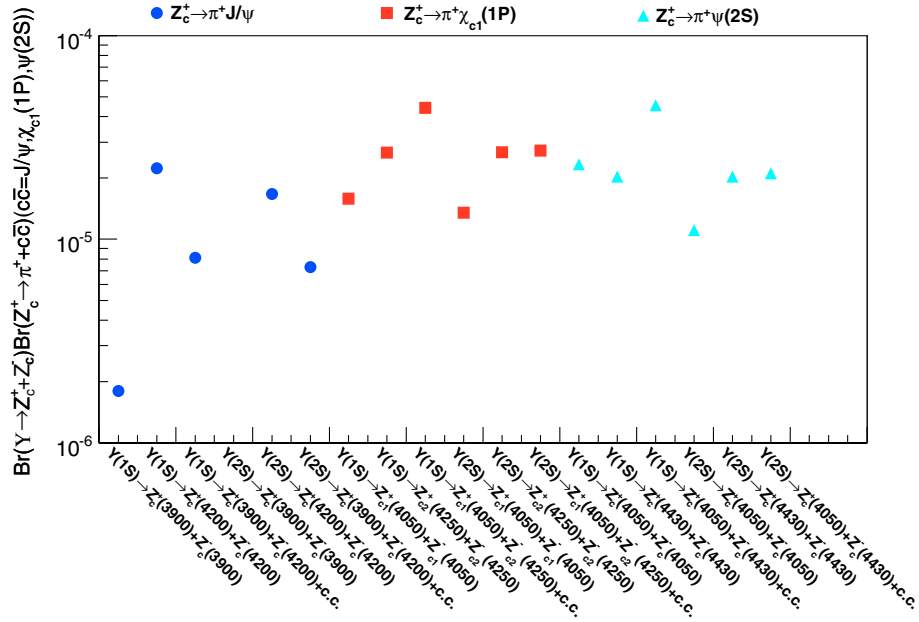


FIG. 9. The 90% C.L. upper limits on  $\mathcal{B}(\Upsilon(1S, 2S) \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$  [ $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ ]. The blue circles represent the results for modes  $Z_c^+ \rightarrow \pi^+ + J/\psi$ , red boxes represent the results for  $Z_c^+ \rightarrow \pi^+ + \chi_{c1}(1P)$ , and cyan triangles represent results for  $Z_c^+ \rightarrow \pi^+ + \psi(2S)$ .

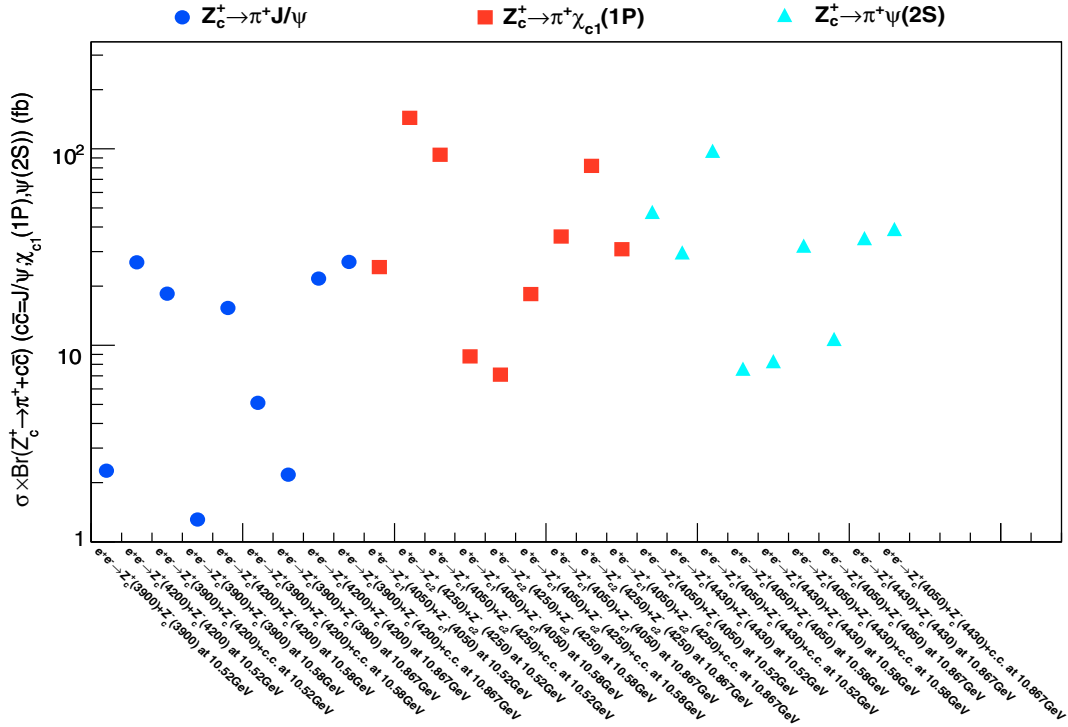


FIG. 10. The 90% C.L. upper limits on  $\sigma(e^+e^- \rightarrow Z_c^+ Z_c^{(\prime)-}) \times \mathcal{B}(Z_c^+ \rightarrow \pi^+ + c\bar{c})$  [ $c\bar{c} = J/\psi, \chi_{c1}(1P), \psi(2S)$ ] at  $\sqrt{s} = 10.52, 10.58$ , and  $10.867$  GeV. The blue circles represent the results for modes  $Z_c^+ \rightarrow \pi^+ + J/\psi$ , red boxes represent results for  $Z_c^+ \rightarrow \pi^+ + \chi_{c1}(1P)$ , and cyan triangles represent results for  $Z_c^+ \rightarrow \pi^+ + \psi(2S)$ .

## ACKNOWLEDGMENTS

We thank the KEKB group for the excellent operation of the accelerator; the KEK cryogenics group for the efficient

operation of the solenoid; and the High Energy Accelerator Research Organisation (KEK) computer group, the National Institute of Informatics, and the Pacific Northwest National Laboratory Environmental Molecular

Sciences Laboratory computing group for valuable computing and Science Information NETWORK 5 network support. We acknowledge support from the Ministry of Education, Culture, Sports, Science, and Technology of Japan, the Japan Society for the Promotion of Science, and the Tau-Lepton Physics Research Center of Nagoya University; the Australian Research Council; Austrian Science Fund under Grant No. P 26794-N20; the National Natural Science Foundation of China under Contracts No. 11435013, No. 11475187, No. 11521505, No. 11575017, No. 11675166, No. 11705209, and No. 11761141009; Key Research Program of Frontier Sciences, Chinese Academy of Sciences, Grant No. QYZDJ-SSW-SLH011; the CAS Center for Excellence in Particle Physics; Fudan University Grants No. JIH5913023, No. IDH5913011/003, No. JIH5913024, and No. IDH5913011/002; the Ministry of Education, Youth and Sports of the Czech Republic under Contract No. LTT17020; the Carl Zeiss Foundation, the Deutsche Forschungsgemeinschaft, the Excellence Cluster Universe, and the VolkswagenStiftung; the Department of Science

and Technology of India; the Istituto Nazionale di Fisica Nucleare of Italy; National Research Foundation of Korea Grants No. 2014R1A2A2A01005286, No. 2015R1A2A2A01003280, No. 2015H1A2A1033649, No. 2016R1D1A1B01010135, No. 2016K1A3A7A09005 603, and No. 2016R1D1A1B02012900; Radiation Science Research Institute, Foreign Large-size Research Facility Application Supporting project, and the Global Science Experimental Data Hub Center of the Korea Institute of Science and Technology Information; the Polish Ministry of Science and Higher Education and the National Science Center; the Ministry of Education and Science of the Russian Federation and the Russian Foundation for Basic Research; the Slovenian Research Agency; Ikerbasque, Basque Foundation for Science, Basque Government (Grant No. IT956-16) and Ministry of Economy and Competitiveness (Juan de la Cierva), Spain; the Swiss National Science Foundation; the Ministry of Education and the Ministry of Science and Technology of Taiwan; and the United States Department of Energy and the National Science Foundation.

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