

Systematic search for very-high-energy gamma-ray emission from bow shocks of runaway stars

H.E.S.S. Collaboration: H. Abdalla¹, A. Abramowski², F. Aharonian^{3,4,5}, F. Ait Benkhali³, A. G. Akhperjanian^{6,5,*}, T. Andersson¹⁰, E. O. Angüner²¹, M. Arakawa⁴², M. Arrieta¹⁵, P. Aubert²⁴, M. Backes⁸, A. Balzer⁹, M. Barnard¹, Y. Becherini¹⁰, J. Becker Tjus¹¹, D. Berge¹², S. Bernhard¹³, K. Bernlöhr³, R. Blackwell¹⁴, M. Böttcher¹, C. Boisson¹⁵, J. Bolmont¹⁶, P. Bordas³, J. Bregeon¹⁷, F. Brun²⁶, P. Brun¹⁸, M. Bryan⁹, M. Büchele³⁵, T. Bulik¹⁹, M. Capasso²⁸, J. Carr²⁰, S. Casanova^{21,3}, M. Cerruti¹⁶, N. Chakraborty³, R. Chalme-Calvet¹⁶, R. C. G. Chaves^{17,22}, A. Chen²³, J. Chevalier²⁴, M. Chréten¹⁶, M. Coffaro²⁸, S. Colafrancesco²³, G. Cologne²⁵, B. Condon²⁶, J. Conrad^{27,*}, Y. Cui²⁸, I. D. Davids^{1,8}, J. Decock¹⁸, B. Degrange²⁹, C. Deil³, J. Devin¹⁷, P. deWilt¹⁴, L. Dirson², A. Djannati-Atai³⁰, W. Domainko³, A. Donath³, L. O'C. Drury⁴, K. Dutton³², J. Dyks³³, T. Edwards³, K. Egberts³⁴, P. Eger³, J.-P. Ernenwein²⁰, S. Eschbach³⁵, C. Farnier^{27,10}, S. Fegan²⁹, M. V. Fernandes², A. Fiasson²⁴, G. Fontaine²⁹, A. Förster³, S. Funk³⁵, M. Füßling³⁶, S. Gabici³⁰, M. Gajdus⁷, Y. A. Gallant¹⁷, T. Garrigoux¹, G. Giavotto³⁶, B. Giebels²⁹, J. F. Glicenstein¹⁸, D. Gottschall²⁸, A. Goyal³⁷, M.-H. Grondin²⁶, J. Hahn³, M. Haupt^{36,**}, J. Hawkes¹⁴, G. Heinzelmann², G. Henri³¹, G. Hermann³, O. Hervet^{15,44}, J. A. Hinton³, W. Hofmann³, C. Hoischen³⁴, M. Holler²⁹, D. Horns², A. Ivashenko¹, H. Iwasaki⁴², A. Jacholkowska¹⁶, M. Jamroz³⁷, M. Janiak³³, D. Jankowsky³⁵, F. Jankowsky²⁵, M. Jingo²³, T. Jogler³⁵, L. Jouvin³⁰, I. Jung-Richardt³⁵, M. A. Kastendieck², K. Katarzyński³⁸, M. Katsuragawa⁴³, U. Katz³⁵, D. Kerszberg¹⁶, D. Khangulyan⁴², B. Khelifi³⁰, M. Kieffer¹⁶, J. King³, S. Klepser³⁶, D. Klockhov²⁸, W. Kluźniak³³, D. Kolitzus¹³, Nu. Komin²³, K. Kosack¹⁸, S. Krakau¹¹, M. Kraus³⁵, P. P. Krüger¹, H. Laffon²⁶, G. Lamanna²⁴, J. Lau¹⁴, J.-P. Lees²⁴, J. Lefaucheur¹⁵, V. Lefranc¹⁸, A. Lemièrre³⁰, M. Lemoine-Goumard²⁶, J.-P. Lenain¹⁶, E. Leser³⁴, T. Lohse⁷, M. Lorentz¹⁸, R. Liu³, R. López-Coto³, I. Lypova³⁶, V. Marandon³, A. Marcowith¹⁷, C. Mariaud²⁹, R. Marx³, G. Maurin²⁴, N. Maxted¹⁴, M. Mayer⁷, P. J. Meintjes³⁹, M. Meyer²⁷, A. M. W. Mitchell³, R. Moderski³³, M. Mohamed²⁵, L. Mohrmann³⁵, K. Morá²⁷, E. Moulin¹⁸, T. Murach⁷, S. Nakashima⁴³, M. de Naurois²⁹, F. Niederwanger¹³, J. Niemiec²¹, L. Oakes⁷, P. O'Brien³², H. Odaka⁴³, S. Öttl¹³, S. Ohm³⁶, M. Ostrowski³⁷, I. Oya³⁶, M. Padovani¹⁷, M. Panter³, R. D. Parsons³, N. W. Pekeur¹, G. Pelletier³¹, C. Perennes¹⁶, P.-O. Petrucci³¹, B. Peyaud¹⁸, Q. Piel²⁴, S. Pita³⁰, H. Poon³, D. Prokhorov¹⁰, H. Prokoph¹⁰, G. Pühlhofer²⁸, M. Punch^{30,10}, A. Quirrenbach²⁵, S. Raab³⁵, A. Reimer¹³, O. Reimer¹³, M. Renaud¹⁷, R. de los Reyes³, S. Richter¹, F. Rieger^{3,40}, C. Romoli⁴, G. Rowell¹⁴, B. Rudak³³, C. B. Rulten¹⁵, V. Sahakian^{6,5}, S. Saito⁴², D. Salek⁴¹, D. A. Sanchez²⁴, A. Santangelo²⁸, M. Sasaki²⁸, R. Schlickeiser¹¹, F. Schüssler¹⁸, A. Schulz^{36,**}, U. Schwanke⁷, S. Schwemmer²⁵, M. Seglar-Arroyo¹⁸, M. Settimo¹⁶, A. S. Seyffert¹, N. Shafi²³, I. Shilon³⁵, R. Simoni⁹, H. Sol¹⁵, F. Spanier¹, G. Spengler²⁷, F. Spies², L. Stawarz³⁷, R. Steenkamp⁸, C. Stegmann^{34,36}, K. Stycz³⁶, I. Sushch¹, T. Takahashi⁴³, J.-P. Tavernet¹⁶, T. Tavernier³⁰, A. M. Taylor⁴, R. Terrier³⁰, L. Tibaldo³, D. Tiziani³⁵, M. Tluczykont², C. Trichard²⁰, N. Tsuji⁴², R. Tuffs³, Y. Uchiyama⁴², D. J. van der Walt¹, C. van Eldik³⁵, C. van Rensburg¹, B. van Soelen³⁹, G. Vasileiadis¹⁷, J. Veh³⁵, C. Venter¹, A. Viana³, P. Vincent¹⁶, J. Vink⁹, F. Voisin¹⁴, H. J. Völk³, T. Vuillaume²⁴, Z. Wadiasingh¹, S. J. Wagner²⁵, P. Wagner⁷, R. M. Wagner²⁷, R. White³, A. Wierzcholska²¹, P. Willmann³⁵, A. Wörnlein³⁵, D. Wouters¹⁸, R. Yang³, V. Zabalza³², D. Zaborov²⁹, M. Zacharias²⁵, R. Zanin³, A. A. Zdziarski³³, A. Zech¹⁵, F. Zefi²⁹, A. Ziegler³⁵, and N. Żywucka³⁷

(Affiliations can be found after the references)

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ABSTRACT

Context. Runaway stars form bow shocks by ploughing through the interstellar medium at supersonic speeds and are promising sources of non-thermal emission of photons. One of these objects has been found to emit non-thermal radiation in the radio band. This triggered the development of theoretical models predicting non-thermal photons from radio up to very-high-energy (VHE, $E \geq 0.1$ TeV) gamma rays. Subsequently, one bow shock was also detected in X-ray observations. However, the data did not allow discrimination between a hot thermal and a non-thermal origin. Further observations of different candidates at X-ray energies showed no evidence for emission at the position of the bow shocks either. A systematic search in the *Fermi*-LAT energy regime resulted in flux upper limits for 27 candidates listed in the E-BOSS catalogue.

Aims. Here we perform the first systematic search for VHE gamma-ray emission from bow shocks of runaway stars.

Methods. Using all available archival H.E.S.S. data we search for very-high-energy gamma-ray emission at the positions of bow shock candidates listed in the second E-BOSS catalogue release. Out of the 73 bow shock candidates in this catalogue, 32 have been observed with H.E.S.S.

Results. None of the observed 32 bow shock candidates in this population study show significant emission in the H.E.S.S. energy range. Therefore, flux upper limits are calculated in five energy bins and the fraction of the kinetic wind power that is converted into VHE gamma rays is constrained.

Conclusions. Emission from stellar bow shocks is not detected in the energy range between 0.14 and 18 TeV. The resulting upper limits constrain the level of VHE gamma-ray emission from these objects down to 0.1–1% of the kinetic wind energy.

Key words. radiation mechanisms: non-thermal – gamma rays: ISM – stars: early-type – gamma rays: stars

* Wallenberg Academy Fellow.

** Corresponding authors: H.E.S.S. Collaboration, e-mail: contact.hess@hess-experiment.eu

† Deceased.

1. Introduction

Stars with velocities larger than $\sim 30 \text{ km s}^{-1}$ (corrected for solar motion and Galactic rotation) are historically called runaway stars due to their fast movement away from OB associations. Two scenarios for the formation process of runaway stars have been proposed: the dynamical ejection and the binary supernova scenario. [Poveda et al. \(1967\)](#) used simulations to verify that during the collapse of a small cluster, dynamical interactions of the stars can lead to high spatial velocities. [Zwicky \(1957\)](#) suggested that the runaway stars are formed during the supernova explosion in a binary system, where the second star keeps its high spatial velocity due to sudden mass loss during the supernova event. [Hoogerwerf et al. \(2000\)](#) showed that both proposed mechanisms take place in nature by retracing star trajectories. Examples for the supernova scenario and the dynamical ejection scenario are ζ Ophiuchi and AE Aurigae, respectively.

Since these massive OB stars have very fast stellar winds with velocities up to a few thousand kilometer per second, comparable to the shock speed of young supernova remnants, they are promising candidates for the acceleration of particles (electrons/protons) to high energies producing non-thermal emission. Stars moving through the interstellar medium (ISM) at supersonic speeds sweep matter up in their direction of motion and form bow shocks. The swept-up dust in these large-scale bow shocks is heated and ionized by the stellar radiation, which leads to infrared emission. The thermal emission of these com-shaped features was first discovered by [van Buren & McCray \(1988\)](#) using data from the Infrared Astronomical Satellite (IRAS). The first survey of stellar bow shocks was performed by [van Buren et al. \(1995\)](#), followed by the Extensive stellar BOW Shock Survey catalogue (E-BOSS; [Peri et al. 2012](#)).

[Benaglia et al. \(2010\)](#) were the first to report on the detection of non-thermal radio emission from a stellar bow shock, namely BD+43°3654. They introduced an emission model predicting non-thermal photons detectable at radio, X-ray and gamma-ray energies. In this model, charged particles are accelerated up to relativistic energies via Fermi acceleration in the shock wave originating from the supersonic motion of the star. These relativistic particles interact with the ambient matter, photon or magnetic fields and produce non-thermal emission. The bow shock system is composed of two shocks, a slow forward shock with the ISM and a fast reverse shock with the stellar wind in which the relativistic particles are accelerated more efficiently. A more detailed description of the model and further developments can be found in [Benaglia et al. \(2010\)](#), [del Valle & Romero \(2012\)](#), [López-Santiago et al. \(2012\)](#) and [Pereira et al. \(2016\)](#).

Based on this model, several observations of promising bow shock candidates followed, aiming to detect non-thermal emission. The follow-up search by [Terada et al. \(2012\)](#) for a non-thermal X-ray counterpart of BD+43°3654 using data from *Suzaku* revealed no emission in this regime. However, the resulting upper limits imply that the emission model from [Benaglia et al. \(2010\)](#) overestimated either the efficiency of the shock-heating process, leading to electron energies that do not exceed 10 TeV, or the grade of turbulence of the magnetic field in the shock acceleration region. Further X-ray observations of ζ Ophiuchi and BD+43°3654 ([Toalá et al. 2016](#)) resulted in upper limits for non-thermal emission and lead to the conclusion that the intensity of the emission is below the sensitivity of current X-ray satellites. X-ray observations with *XMM-Newton* of AE Aurigae (HIP 24575) revealed for the first time significant emission, but its nature (very hot thermal or non-thermal) could not be firmly determined ([López-Santiago et al. 2012](#)). In the

case of ζ Ophiuchi, [Toalá et al. \(2016\)](#) detected diffuse emission in the vicinity of this candidate, which they attribute to a plasma with a temperature of $2 \times 10^6 \text{ K}$, in agreement with predictions of high plasma temperatures caused by instabilities mixing material between the shocked wind and the photo-ionized gas at the wake of the shock ([Mackey et al. 2015](#)). Recently, ([Toalá et al. 2017](#)) showed that the X-ray emission close to AE Aurigae is point-like and unrelated to the bow shock. They furthermore searched for non-thermal diffuse X-ray emission around 6 well-determined runaway stars and found no evidence for it.

[Del Valle et al. \(2013\)](#) suggested the high-energy (HE, 100 MeV to $\sim 100 \text{ GeV}$) gamma-ray source 2FGL J2030.7+4417 ([Nolan et al. 2012](#)) to be associated with the bow shock of HD 195592. However, the source 2FGL J2030.7+4417 has been identified as a gamma-ray pulsar ([Pletsch et al. 2012](#)) and shows no significant off-pulse emission ([Abdo et al. 2013](#)), a strong indication that the detected photons predominantly originate in the pulsar and not in the bow shock.

A possibility of stellar bow shocks being variable gamma-ray sources was introduced by [del Valle & Romero \(2014\)](#). The predicted variability in the gamma-ray flux originates from inhomogeneities of the ambient medium, leading to changes in the physical properties and thus the luminosity. The expected time-scale of the variations is $\sim 1 \text{ yr}$ and depends on the size and density gradient of the molecular cloud and the velocity of the star.

In the HE gamma-ray regime [Schulz et al. \(2014\)](#) performed the first systematic search for non-thermal emission from bow shocks around runaway stars using 57 months of *Fermi*-LAT data testing the predictions of [del Valle & Romero \(2012\)](#), [Benaglia et al. \(2010\)](#) and [López-Santiago et al. \(2012\)](#). This population study resulted in upper limits for 27 bow shocks including ζ Ophiuchi for which the upper limit on its emission was found to be a factor ~ 5 below the predicted emission from [del Valle & Romero \(2012\)](#).

In this work, we search for very-high-energy (VHE, $E \geq 0.1 \text{ TeV}$) emission from stellar bow shocks using the latest, most comprehensive survey of bow shocks of runaway stars ([Peri et al. 2015](#), second E-BOSS catalogue release) which uses recent infrared data releases, mainly from the Wide-field Infrared Survey Explorer (WISE). The second E-BOSS catalogue release includes bow shocks from literature and serendipitously found ones to complete the sample. It comprises 73 bow shock candidates: 28 candidates from the first E-BOSS catalogue ([Peri et al. 2012](#)), 16 new ones and 29 from recent publications.

We describe the H.E.S.S. observations, data analysis and results of 32 bow shock candidates in Sect. 2. A discussion on the implications of these non-detections is presented in Sect. 3.

2. Observations, data analysis and results

H.E.S.S. is an array of imaging atmospheric Cherenkov telescopes located in the Khomas Highland in Namibia at an altitude of 1800 m above sea level ($23^\circ 16' 18'' \text{ S}$, $16^\circ 30' 00'' \text{ E}$; [Hinton & the HESS Collaboration 2004](#)). The initial configuration of four 12 m telescopes (H.E.S.S. phase I) was extended with a central 28 m telescope in July 2012. This work only uses data from the initial configuration, which provides an energy threshold of $\sim 100 \text{ GeV}$ with an angular resolution better than 0.1° and an energy resolution below 15%. The standard quality selection was used to discard observations during bad weather or instrumental conditions ([Aharonian et al. 2006](#)).

The second E-BOSS catalogue release ([Peri et al. 2015](#)) is the basis for the population study presented in this work. Almost

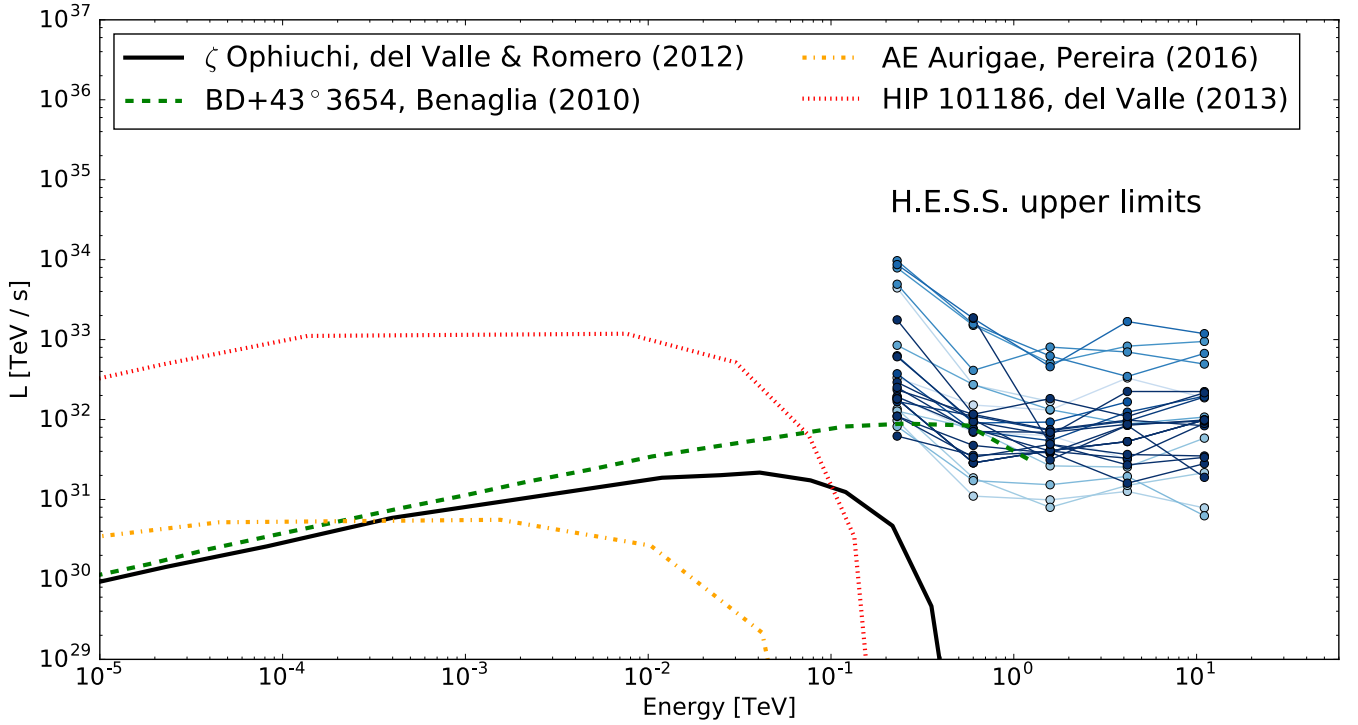


Fig. 1. VHE gamma-ray luminosity upper limits for the 28 bow shock candidates with known distance, compared to model predictions for four different bow shocks (none of these four is in the H.E.S.S. survey sample, see text for details).

50% (32 out of 73) of the candidates in the second E-BOSS catalogue release are covered by H.E.S.S. observations. 27 of the observed bow shocks are located within the Galactic plane, profiting from the nine-year-long H.E.S.S. Galactic Plane Survey (HGPS; H.E.S.S. Collaboration 2018b) of the inner Milky Way. The study presented here complements the population studies on pulsar wind nebulae (H.E.S.S. Collaboration 2018c) and supernova remnants (H.E.S.S. Collaboration 2018a).

The coordinates listed in the second E-BOSS catalogue release are the stars' coordinates for all candidates except the seven serendipitous discoveries “SER1–7” for which the star could not be firmly identified. In these cases, the apex of the bow shock was estimated visually using publicly available WISE data. The four corresponding candidates in the H.E.S.S. sample are marked with *** in Table 1.

The H.E.S.S. analyses are performed for the positions given in the second E-BOSS catalogue release with seven exceptions: The three bow shocks in M 17 have an angular separation of less than 0.1° which is not resolvable for H.E.S.S. due to its point spread function ($\text{PSF} \sim 0.1^\circ$; Aharonian et al. 2006). For these three objects one analysis was performed for the position of M 17-S2 which is in the centre of the three. The two exceptions M 17-S1 and S3 are marked with a * in Table 1, since the coordinates of M 17-S2 are used for the analysis.

To ensure that the defined source region of the analysis covers the bow shock, its size and distance to the star have to be evaluated. If the length of the bow shock listed in the second E-BOSS catalogue release is larger than 0.1° we estimate the bow shock position visually using publicly available WISE data and perform the analysis for this position. This criterion leads to updated coordinates for HIP 32067, HIP 88652, HIP 92865, Star 1 and G2; they are marked with ** in Table 1. For HIP 32067 with a length of $13'$ ($> 2 * 0.1^\circ$) the source region was enlarged from the standard 0.1° to 0.11° .

The data were analysed using the *ImPACT* analysis method described in Parsons & Hinton (2014). The “standard cuts” of this analysis were adopted, including a minimum charge of 60 photoelectrons per shower image and a signal extraction region of 0.1° . A cross-check analysis performed with the *model* analysis method as presented by de Naurois & Rolland (2009) yielded compatible results.

The differential upper limits are presented in Table 1, including the duration of the H.E.S.S. observations (live-time) and the parameters for each star. To avoid potential systematic biases, upper limits are only calculated if more than 10 events are recorded in the OFF regions that are used to estimate the background.

The analyses of all bow shock candidates were performed in a systematic way by using the same analysis cuts and configurations. None of the analysed bow shock candidates showed statistically significant VHE gamma-ray emission at the position of the bow shock; thus, upper limits on the flux are calculated using the method presented by Rolke et al. (2005).

In some cases, the candidates are close to known VHE gamma-ray sources, which leads to significances up to 3σ . However, dedicated analyses of the sky maps and the squared-angular distance distributions of the reconstructed direction of the events with respect to the candidates' source positions clearly showed that the emission is not originating from the bow shock. In these cases upper limits were calculated in the same way as for the rest of the population.

The differential gamma-ray flux upper limits at 95% confidence-level in five energy bins (equally spaced in logarithmic energy) assuming a power-law spectrum of gamma-ray emission ($d\Phi/dE = \Phi_0(E/E_0)^{-\Gamma}$) with a photon index $\Gamma = 2.5$ are presented in Table 1 and in Fig. 1. Assuming different indices ($\Gamma = 2.0$ and $\Gamma = 3.0$) leads to marginal changes in the upper limits of order 10% or less.

Table 1. Differential gamma-ray flux upper limits (95% confidence-level) for bow shocks of runaway stars.

ID	Star	l [$^{\circ}$]	b [$^{\circ}$]	d [pc]	v_{wind} [km s $^{-1}$]	\dot{M} 10 $^{-6}$ M_{\odot} /yr	Live-time (acc.-corr.) [h]	Upper limits $E^2 d\Phi/dE$ [10 $^{-12}$ TeV cm $^{-2}$ s $^{-1}$] energy bins [TeV]					
								0.14–0.37	0.37–0.97	0.97–2.57	2.57–6.78	6.78–17.92	0.1–10
EB8	HIP 25923	210.44	−20.98	900	[1000]	0.06	3.9 (2.2)	2.01	0.71	0.65	0.33	—	0.40
EB13	HIP 32067**	206.20	0.85	2117±367	2960	0.13	21.7 (9.4)	0.62	0.28	0.25	0.62	0.35	0.12
EB15	HIP 38430	243.16	0.36	900	[2570]	0.70	1.8 (0.1)	45.73	2.79	1.73	—	—	0.80
EB17	HIP 72510	318.77	2.77	350	[2545]	0.27	12.8 (2.4)	6.64	0.75	0.67	0.86	0.53	0.58
EB18	HIP 75095	322.68	0.91	800	[1065]	0.14	22.5 (13.9)	1.77	0.24	0.10	0.20	0.28	0.15
EB23	HIP 88652**	15.11	3.36	650	[1535]	0.50	9.2 (2.5)	2.53	1.52	0.52	0.50	1.16	0.72
EB24	HIP 92865**	41.75	3.41	350	[1755]	0.04	3.9 (2.3)	5.60	1.17	1.04	1.33	0.43	0.87
EB32	SER1***	264.78	1.54	—	250	0.03	3.0 (2.0)	2.92	1.06	0.32	0.77	—	0.59
EB33	HIP 44368	263.10	3.90	1900 ± 200 ^a	1100	0.80	7.9 (6.1)	1.97	0.63	0.31	0.20	0.25	0.36
EB36	SER2***	282.48	−2.46	—	—	—	15.7 (7.9)	1.10	0.54	0.24	0.18	0.33	0.22
EB37	RCW 49-S1	284.08	0.43	6100	2800	3.23	51.3 (29.2)	1.77	0.34	0.11	0.19	0.21	0.17
EB38	RCW 49-S2	284.30	0.30	6100	2600	0.60	51.7 (31.5)	1.11	0.09	0.18	0.16	0.11	0.07
EB39	RCW 49-S3	284.34	0.20	6100	2800	2.00	52.2 (33.5)	2.18	0.35	0.14	0.08	0.15	0.16
EB40	SER3***	286.46	−0.34	—	250	0.03	62.2 (29.0)	1.88	0.28	0.17	0.23	0.15	0.16
EB41	J1117-6120	291.88	−0.50	7600	2600	0.60	52.9 (32.3)	1.26	0.27	0.07	0.24	0.17	0.11
EB42	SER7***	347.15	2.36	—	—	—	13.2 (7.7)	0.79	0.13	0.21	0.15	0.28	0.08
EB43	G4	352.57	2.11	1700	2550	0.50	4.8 (1.5)	1.82	0.26	0.27	0.48	—	0.22
EB44	G2**	352.81	1.34	1700	2250	0.40	20.9 (8.7)	1.08	0.21	0.16	0.36	0.57	0.13
EB45	G5	351.65	0.51	1700	2000	0.10	28.0 (11.9)	0.49	0.32	0.21	0.28	0.55	0.18
EB46	G6	353.06	1.29	1700	[1000]	0.10	30.1 (11.6)	0.32	0.14	0.11	0.25	0.29	0.07
EB47	G8	353.16	1.05	1700	[1500]	0.04	34.9 (16.6)	0.68	0.34	0.22	0.28	0.24	0.20
EB48	G1	353.42	0.45	1700	2100	0.20	56.4 (31.1)	0.56	0.22	0.12	0.09	0.26	0.12
EB49	G7	354.03	0.85	1700	[1000]	0.10	38.0 (20.5)	0.85	0.18	0.11	0.05	0.08	0.04
EB50	G3	353.30	0.08	1700	2000	0.40	48.3 (29.9)	0.85	0.29	0.09	0.24	0.05	0.19
EB51	HIP 86768	18.70	11.60	737	[550]	0.03	1.3 (0.4)	9.39	1.45	0.98	3.45	3.44	0.94
EB52	Star 1**	16.99	1.77	1800	2200	0.63	20.6 (13.9)	0.65	0.18	0.18	0.22	0.25	0.12
EB53	M 17-S1*	15.08	0.65	1600	1000	0.03	22.7 (6.6)	0.59	0.09	0.13	0.17	0.32	0.06
EB54	M 17-S2	15.08	0.65	1600	[1500]	0.16	22.7 (6.6)	0.59	0.09	0.13	0.17	0.32	0.06
EB55	M 17-S3*	15.08	0.65	1600	2300	0.25	22.7 (6.6)	0.59	0.09	0.13	0.17	0.32	0.06
EB56	BD -14 5040	16.89	−1.12	1800	400	0.03	111.3 (73.2)	0.28	0.09	0.13	0.09	0.09	0.09
EB57	4U 1907+09	43.74	0.47	4000	2900	0.70	94.2 (63.1)	0.92	0.06	0.10	0.06	0.11	0.04
EB58	HIP 98418	71.60	2.90	529.1	2545	0.24	4.1 (3.0)	—	55.94	1.48	0.81	1.00	1.46

Notes. ID, Star, distance d , wind velocity v_{wind} and mass-loss rate \dot{M} as listed in [Peri et al. \(2015\)](#). The positions with the coordinates l and b denote the ones used for the analysis, which is not in all cases equal to the one in [Peri et al. 2012](#), and [2015](#)) (see text for details). Wind velocity: All values from [Peri et al. \(2012\)](#); brackets indicate values adopted from stars with the same spectral type. The live-time, how long each object was observed with H.E.S.S., is also listed. The acceptance-corrected live-time (the observation time corrected for the non-uniform acceptance across the field of view of the camera) is given in parentheses. (*) The bow shocks M 17-S1, M 17-S2 and M 17-S3 are closer than 0.1 degree and therefore not resolvable by H.E.S.S., the upper limits are calculated for the position of M 17-S2 but valid for all three bow shock candidates. (**) The analysis was done for the bow shock coordinates, see text for more details. (***) The coordinates listed in second E-BOSS catalogue release are the apex coordinates of the bow shock, not the star's. ^(a) Distance uncertainty wrong in second E-BOSS catalogue release (1900 ± 0.1 pc), original paper ([Sadakane et al. 1985](#)) 1.9 ± 0.2 kpc.

3. Discussion and conclusions

There are currently no model predictions published for the bow shocks analysed here. Therefore, the published models for four different bow shocks are shown together with the upper limits from this work in Fig. 1. All four model predictions are based on the model by [del Valle & Romero \(2012\)](#) and were motivated by dedicated searches for non-thermal emission. The comparison of the VHE upper limits with the model prediction for BD+43°3654 (the only confirmed non-thermal emitter) suggests that several candidates of the H.E.S.S. bow shock sample do not emit VHE gamma rays at the level predicted by [Benaglia et al. \(2010\)](#).

Power considerations

For the bow shocks with known stellar parameters and distance, the kinetic power of the wind can be compared to the upper limits of the radiative power at very-high energies. The kinetic power in the wind is given by:

$$P_{\text{wind}} = \frac{1}{2} \dot{M} v_{\text{wind}}^2, \quad (1)$$

with the mass-loss rate \dot{M} and wind velocity v_{wind} listed in Table 1. The integrated upper limit of the VHE radiative power P_{UL} is calculated using the VHE flux upper limits derived in this work (see Sect. 3):

$$P_{\text{UL}} = 4\pi d^2 \int_{E_{\text{min}}}^{E_{\text{max}}} dE (E d\Phi/dE), \quad (2)$$

with the distance d listed in Table 1. The unknown uncertainties of the distances are treated as a systematic caveat here and are not included in the calculation. For this power calculation, the upper limits in the 0.1–10 TeV bin, shown in the last column in Table 1, are used.

Figure 2 shows the ratio of the powers ($P_{\text{UL}}/P_{\text{wind}}$) as a function of the wind power. We constrain the fraction of wind power that is converted into the production of VHE gamma rays in all cases. In five cases we show that less than 0.1% of the wind power is potentially converted into the production of VHE gamma rays, while the majority of the limits constrain the ratio of the powers to <0.1–1%. One should note that not all of the wind's kinetic power (as given in Eq. (1)) is available for particle acceleration: the wind is emitted isotropically, while the bow shock covers only a limited solid angle.

Our systematic population study reveals no evidence for VHE gamma-ray emission from the bow shocks of runaway stars observed in the H.E.S.S. dataset. Together with the HE gamma-ray upper limits by [Schulz et al. \(2014\)](#) and several X-ray upper limits, this challenges the level of predicted non-thermal emission from bow shocks of runaway stars published so far (see Fig. 1 and references therein).

One reason for the non-detection could be that particle acceleration is in general less efficient in bow shocks than in known gamma-ray sources. [Terada et al. \(2012\)](#) concluded that the magnetic fields in the bow shocks of runaway stars might be less turbulent compared to those of pulsar wind nebulae or supernova remnants, where gamma-ray emission is detected in many cases. A lower maximum energy of the accelerated particles or lower photon densities could also explain the non-detections.

For five bow shocks, we calculate that less than 0.1% of the kinetic power of the wind is converted into VHE gamma rays originating from relativistically accelerated particles. This

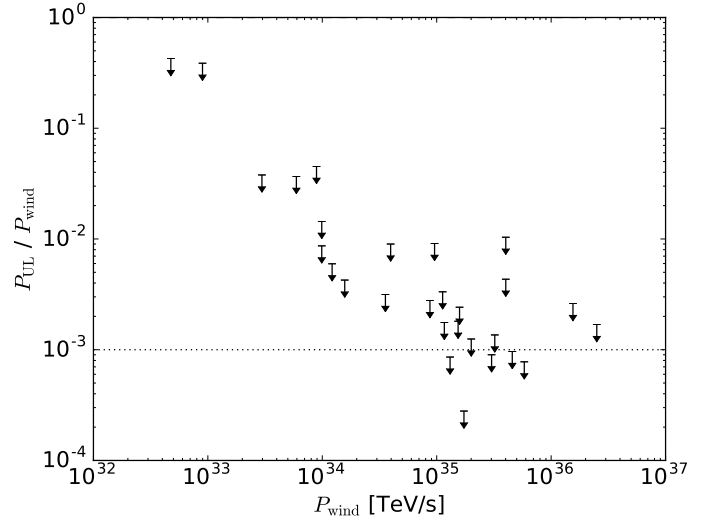


Fig. 2. Ratio of power in VHE gamma rays and power in the wind as a function of wind power for the 28 bow shock candidates where the runaway star is identified. The dotted line depicts 0.1% of the kinetic wind energy.

is roughly the order of magnitude expected from geometrical considerations. For other astrophysical systems, like e.g. novae ([Cheung et al. 2016](#)), the fraction of the total energy in electrons compared to the kinetic energy of the ejected mass is $\sim 0.1\%$.

In general, the search for non-thermal emission from bow shocks of runaway stars proves to be a challenge: so far, only one detection of non-thermal radio emission has been reported ([Benaglia et al. 2010](#)) and upper limits in other radio, X-ray, HE gamma-ray and VHE gamma-ray observations. Our paper presents the first VHE gamma-ray observations of this source class.

Our population study shows that none of the already observed stellar bow shocks listed in the second E-BOSS catalogue release emits VHE gamma rays at a flux level detectable with current imaging atmospheric Cherenkov telescopes. BD+43°3654 could also be unique in this source class as the only bow shock emitting non-thermal radiation.

The future Cherenkov Telescope Array ([Hinton et al. 2013](#)), with approximately 10 times better sensitivity than current instruments and improved angular resolution, might be able to detect VHE gamma-ray emission from stellar bow shocks and understand the physics of these objects.

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- ¹ Centre for Space Research, North-West University, Potchefstroom 2520, South Africa
- ² Universität Hamburg, Institut für Experimentalphysik, Luruper Chaussee 149, 22761 Hamburg, Germany
- ³ Max-Planck-Institut für Kernphysik, PO Box 103980, 69029 Heidelberg, Germany
- ⁴ Dublin Institute for Advanced Studies, 31 Fitzwilliam Place, Dublin 2, Ireland
- ⁵ National Academy of Sciences of the Republic of Armenia, Marshall Baghramian Avenue, 24, 0019 Yerevan, Republic of Armenia
- ⁶ Yerevan Physics Institute, 2 Alikhanian Brothers St., 375036 Yerevan, Armenia
- ⁷ Institut für Physik, Humboldt-Universität zu Berlin, Newtonstr. 15, 12489 Berlin, Germany
- ⁸ University of Namibia, Department of Physics, Private Bag 13301, Windhoek, Namibia
- ⁹ GRAPPA, Anton Pannekoek Institute for Astronomy, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands
- ¹⁰ Department of Physics and Electrical Engineering, Linnaeus University, 351 95 Växjö, Sweden
- ¹¹ Institut für Theoretische Physik, Lehrstuhl IV: Weltraum und Astrophysik, Ruhr-Universität Bochum, 44780 Bochum, Germany
- ¹² GRAPPA, Anton Pannekoek Institute for Astronomy and Institute of High-Energy Physics, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands
- ¹³ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität Innsbruck, 6020 Innsbruck, Austria
- ¹⁴ School of Physical Sciences, University of Adelaide, Adelaide 5005, Australia
- ¹⁵ LUTH, Observatoire de Paris, PSL Research University, CNRS, Université Paris Diderot, 5 Place Jules Janssen, 92190 Meudon, France
- ¹⁶ Sorbonne Universités, UPMC Université Paris 06, Université Paris Diderot, Sorbonne Paris Cité, CNRS, Laboratoire de Physique Nucléaire et de Hautes Energies (LPNHE), 4 place Jussieu, 75252 Paris Cedex 5, France
- ¹⁷ Laboratoire Univers et Particules de Montpellier, Université Montpellier, CNRS/IN2P3, CC 72, Place Eugène Bataillon, 34095 Montpellier Cedex 5, France
- ¹⁸ DSM/Irfu, CEA Saclay, 91191 Gif-Sur-Yvette Cedex, France
- ¹⁹ Astronomical Observatory, The University of Warsaw, Al. Ujazdowskie 4, 00-478 Warsaw, Poland
- ²⁰ Aix-Marseille Université, CNRS/IN2P3, CPPM UMR 7346, 13288 Marseille, France
- ²¹ Instytut Fizyki Jądrowej PAN, ul. Radzikowskiego 152, 31-342 Kraków, Poland
- ²² Funded by EU FP7 Marie Curie, grant agreement No. PIEF-GA-2012-332350
- ²³ School of Physics, University of the Witwatersrand, 1 Jan Smuts Avenue, Braamfontein, 2050 Johannesburg, South Africa
- ²⁴ Laboratoire d'Annecy-le-Vieux de Physique des Particules, Université Savoie Mont-Blanc, CNRS/IN2P3, 74941 Annecy-le-Vieux, France
- ²⁵ Landessternwarte, Universität Heidelberg, Königstuhl, 69117 Heidelberg, Germany
- ²⁶ Université Bordeaux, CNRS/IN2P3, Centre d'Études Nucléaires de Bordeaux Gradignan, 33175 Gradignan, France
- ²⁷ Oskar Klein Centre, Department of Physics, Stockholm University, Albanova University Center, 10691 Stockholm, Sweden
- ²⁸ Institut für Astronomie und Astrophysik, Universität Tübingen, Sand 1, 72076 Tübingen, Germany
- ²⁹ Laboratoire Leprince-Ringuet, Ecole Polytechnique, CNRS/IN2P3, 91128 Palaiseau, France
- ³⁰ APC, AstroParticule et Cosmologie, Université Paris Diderot, CNRS/IN2P3, CEA/Irfu, Observatoire de Paris, Sorbonne Paris Cité, 10, rue Alice Domon et Léonie Duquet, 75205 Paris Cedex 13, France
- ³¹ Univ. Grenoble Alpes, IPAG; CNRS, IPAG, 38000 Grenoble, France
- ³² Department of Physics and Astronomy, The University of Leicester, University Road, Leicester, LE1 7RH, UK
- ³³ Nicolaus Copernicus Astronomical Center, Polish Academy of Sciences, ul. Bartycka 18, 00-716 Warsaw, Poland
- ³⁴ Institut für Physik und Astronomie, Universität Potsdam, Karl-Liebknecht-Strasse 24/25, 14476 Potsdam, Germany
- ³⁵ Friedrich-Alexander-Universität Erlangen-Nürnberg, Erlangen Centre for Astroparticle Physics, Erwin-Rommel-Str. 1, 91058 Erlangen, Germany
- ³⁶ DESY, 15738 Zeuthen, Germany
- ³⁷ Obserwatorium Astronomiczne, Uniwersytet Jagielloński, ul. Orla 171, 30-244 Kraków, Poland
- ³⁸ Centre for Astronomy, Faculty of Physics, Astronomy and Informatics, Nicolaus Copernicus University, Grudziadzka 5, 87-100 Torun, Poland
- ³⁹ Department of Physics, University of the Free State, PO Box 339, 9300 Bloemfontein, South Africa
- ⁴⁰ Heisenberg Fellow (DFG), ITA Universität Heidelberg, 69120 Heidelberg, Germany
- ⁴¹ GRAPPA, Institute of High-Energy Physics, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands
- ⁴² Department of Physics, Rikkyo University, 3-34-1 Nishi-Ikebukuro, Toshima-ku, 171-8501 Tokyo, Japan
- ⁴³ Japan Aerospace Exploration Agency (JAXA), Institute of Space and Astronautical Science (ISAS), 3-1-1 Yoshinodai, Chuo-ku, Sagami-hara, 229-8510 Kanagawa, Japan
- ⁴⁴ Now at Santa Cruz Institute for Particle Physics and Department of Physics, University of California at Santa Cruz, Santa Cruz, CA 95064, USA