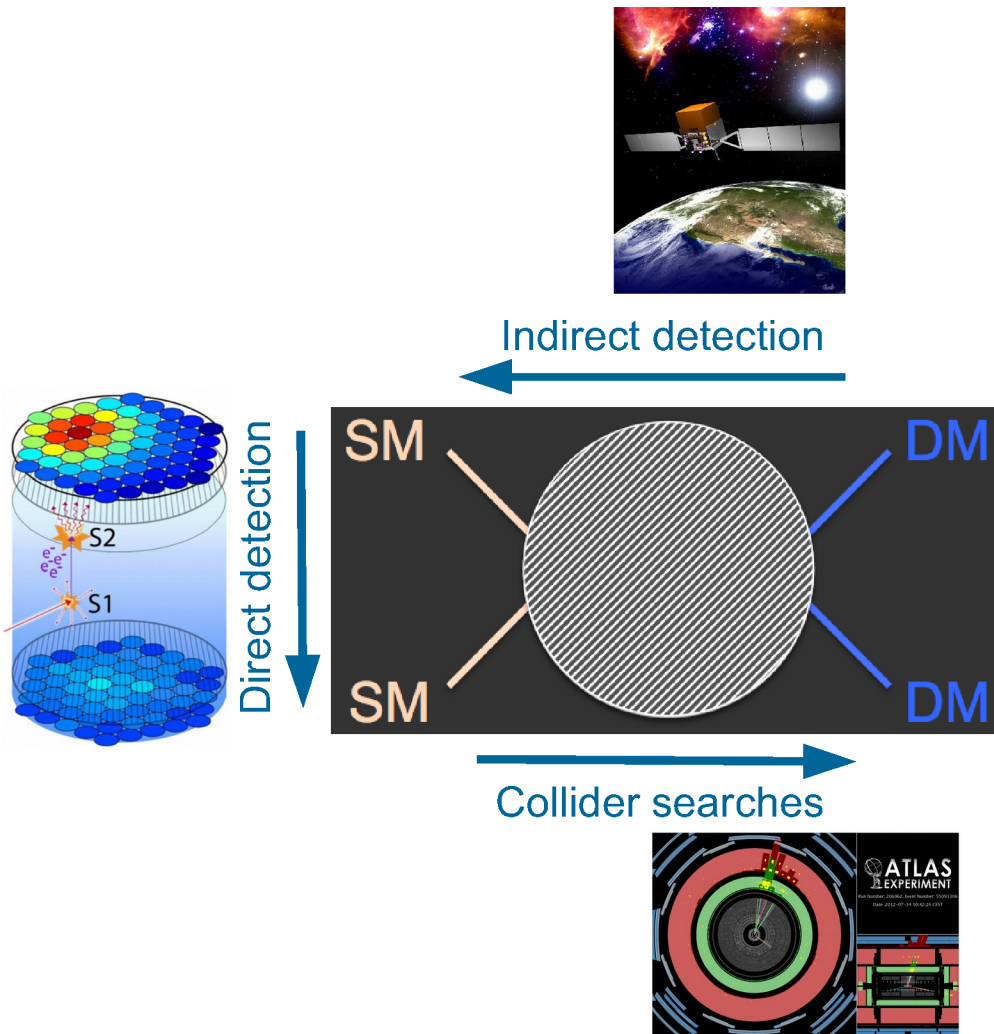


On self-interacting dark matter: Current status and perspectives

Kai Schmidt-Hoberg

How to constrain the properties of dark matter?

- > Our 'usual way' to search for dark matter



How to constrain the properties of dark matter?

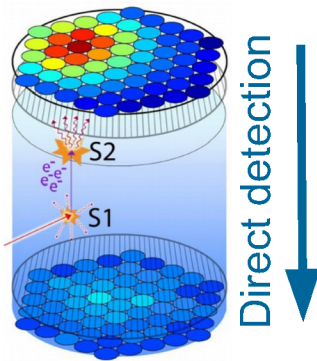
> Our 'usual way' to search for dark matter

cf. Talks by
H M Lee, C Spethmann, D Huang

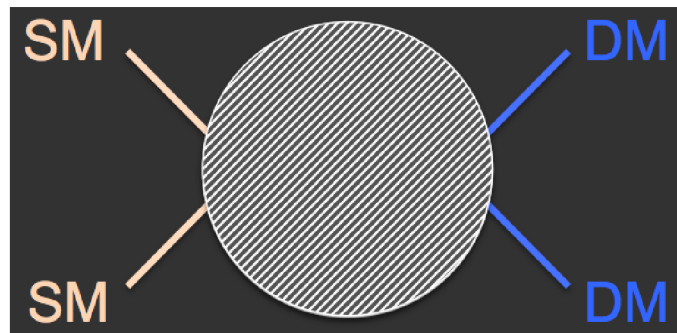
> A fourth way...



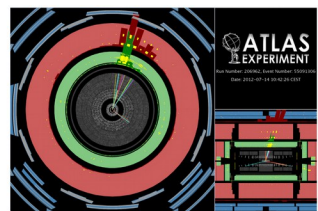
Indirect detection



Direct detection



Collider searches



This Talk:

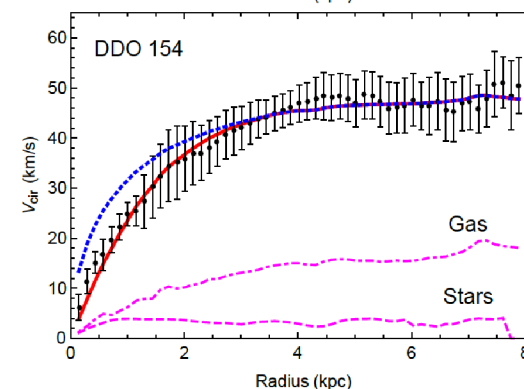
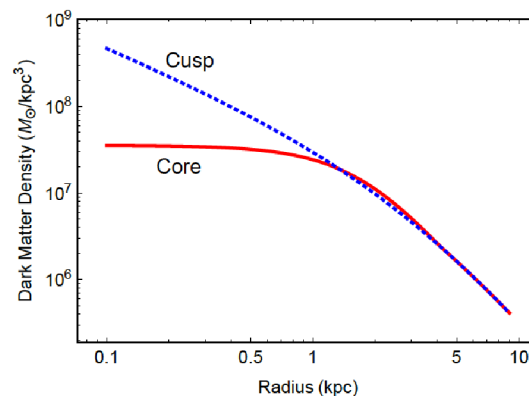
DM self-interactions

Motivation: Cosmology

cf. talk by Pran Nath

- The collisionless cold dark matter paradigm fits perfectly at large scales
- There are however various discrepancies between N-body simulations of collisionless cold DM and astrophysical observations on galactic scales:

- Cusp-vs-core problem



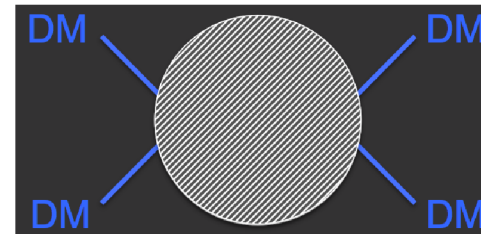
Tulin, Yu: 1705.02358

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 - Missing-satellite problem

DM self-interactions may solve some (or all) of these problems



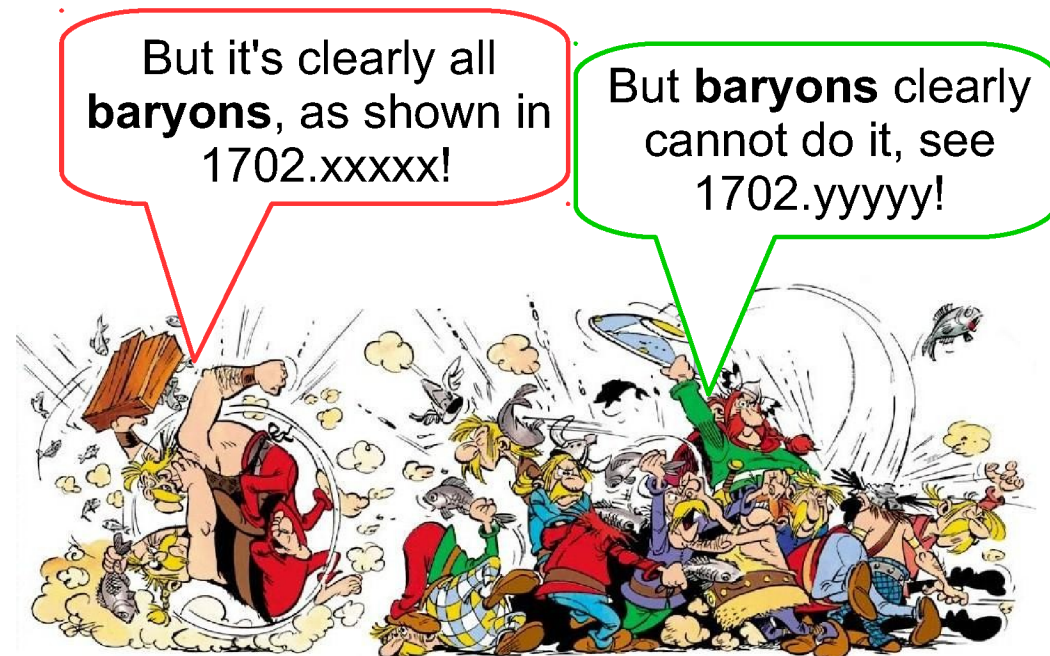
Spergel & Steinhard: astro-ph/9909386
Arsen, Bringmann, Pfrommer, 1205.5809

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Arsen, Bringmann, Pfrommer, 1205.5809

Motivation: Particle physics

- Dark sector often assumed to be simple, mainly because we don't know much...
- Large self-interactions are natural in models with a more complex dark sector (e.g. with a new gauge group)

- Strongly interacting DM

Carlson, Machacek, Hall (1992)
Kusenko, Steinhardt: astro-ph/0106008

- New light mediator in the dark sector

Feng, Kaplinghat, Yu: arXiv:0905.3039
Buckley & Fox: arXiv:0911.3898
Loeb & Weiner: arXiv:1011.6374

- Bonus: We can potentially study the dark sector even if DM has highly suppressed couplings to Standard Model particles.

How large a cross section?

- To be observable on astrophysical scales, self-interaction cross sections have to be large, typically

$$\sigma / m_\chi \sim 1 \text{ cm}^2/\text{g} \sim 2 \text{ barns}/\text{GeV}$$

- The nucleon nucleon scattering cross section ~ 20 barns at low energies
- The typical cross section of a WIMP is 20 orders of magnitude smaller!
- **Potential impact:** Evidence for DM self-interactions on astrophysical scales would rule out most popular models for DM, such as supersymmetric WIMPs, gravitinos, axions...

Constraints on self-interactions

- Various astrophysical observations give constraints on SIDM:

- Bullet cluster Randall et al 0704.0261



- Subhalo evaporation rate

Gnedin, Ostriker: astro-ph/0010436

- Halo ellipticity

Miralda-Escude (2002)

- Core density in clusters and dwarfs

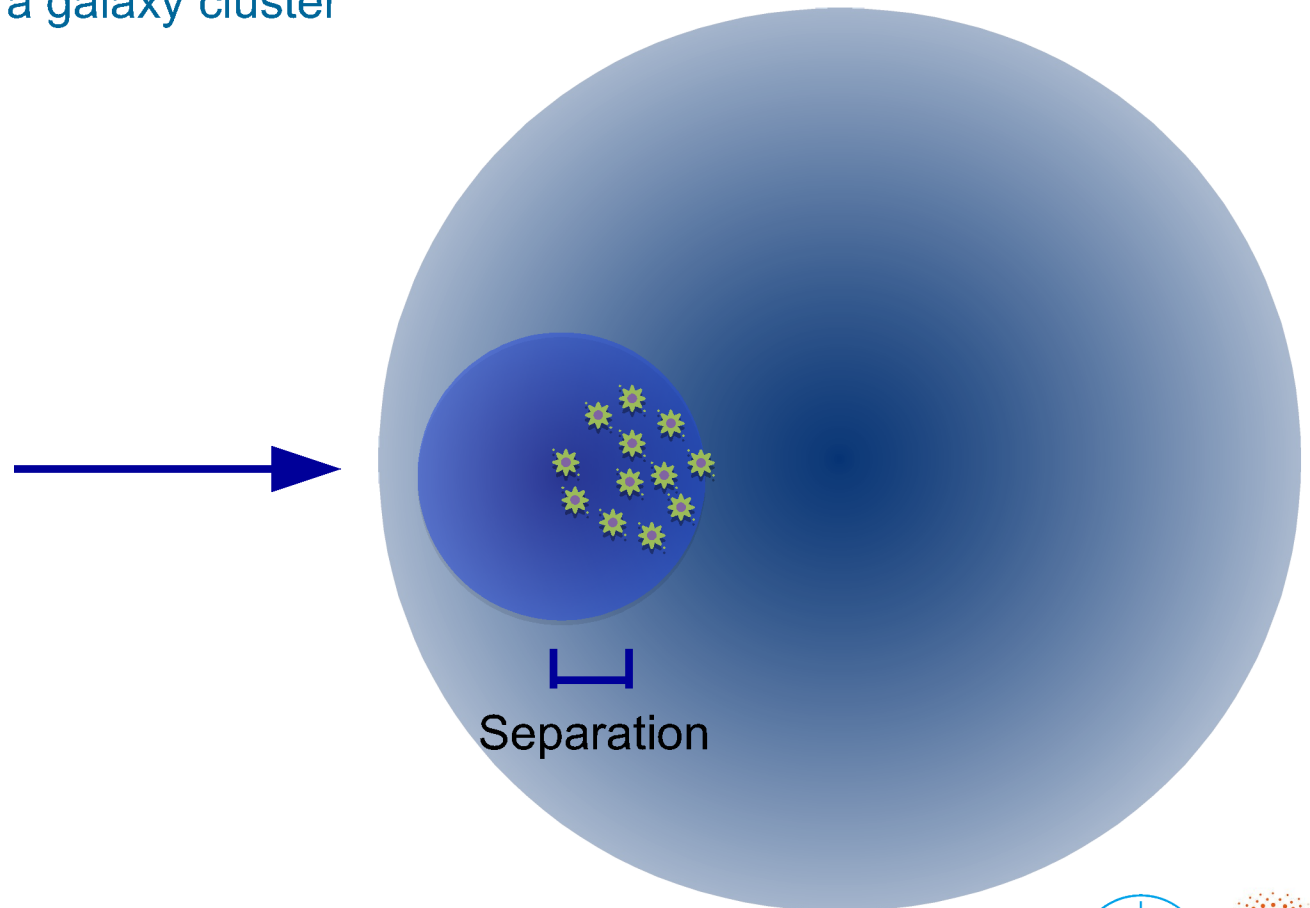
Yoshida et al.: astro-ph/0006134

Dave et al.: astro-ph/0006218

- SIDM probed at different velocities in different systems
→ a handle on the velocity dependence of the self scattering cross section!

Smoking gun?

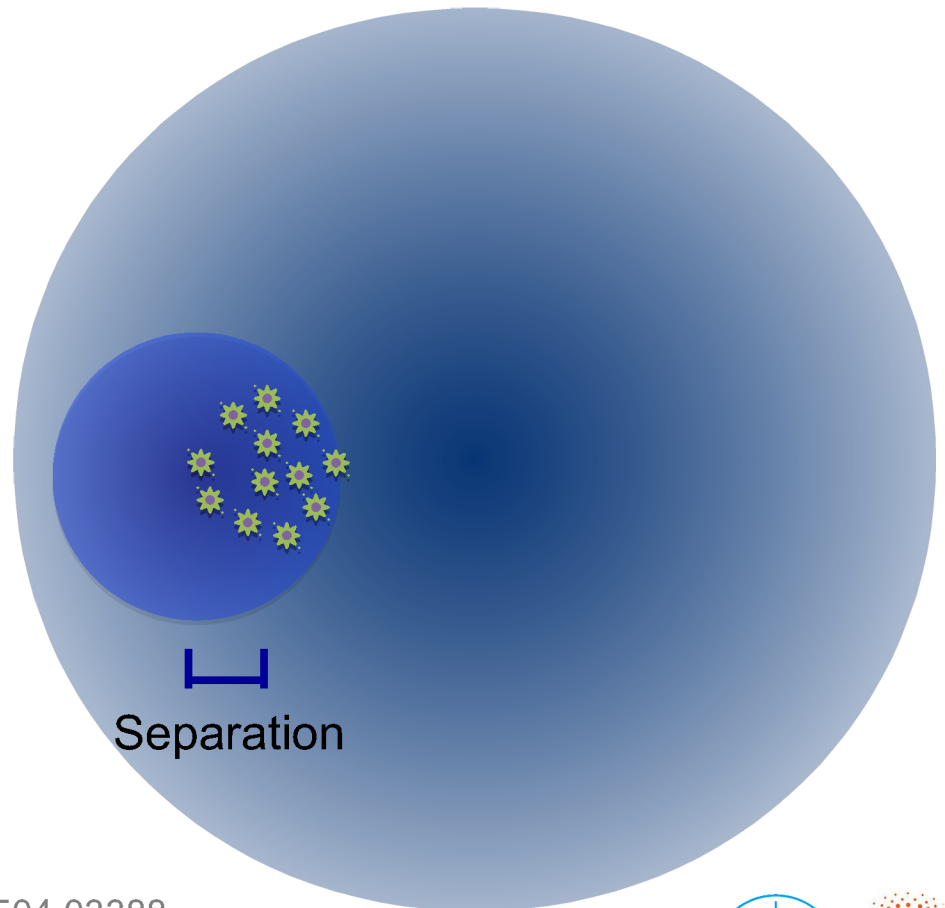
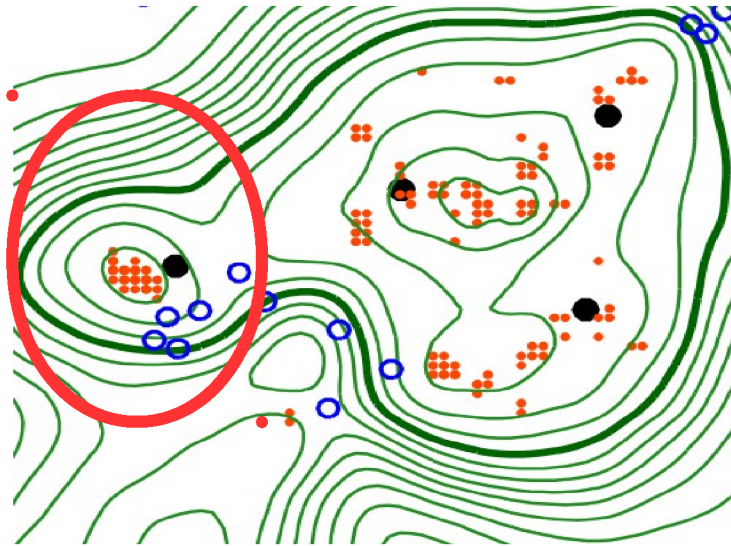
- Smoking gun signal? Separation between dark matter halo and stars of a galaxy falling into a galaxy cluster



Smoking gun?

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Observed offset: $1.62 \pm 0.48 \text{ kpc}$



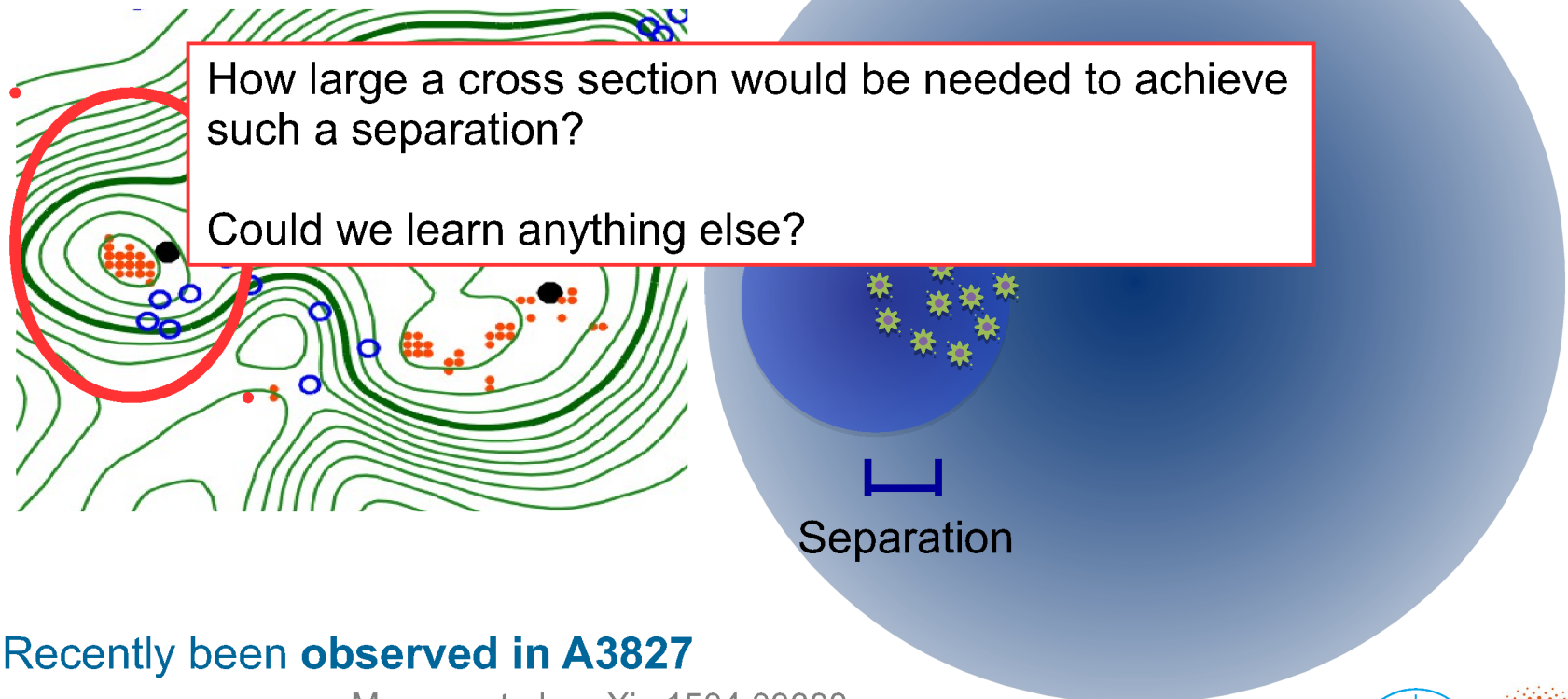
- Recently been **observed in A3827**

Massey et al., arXiv:1504.03388

Smoking gun?

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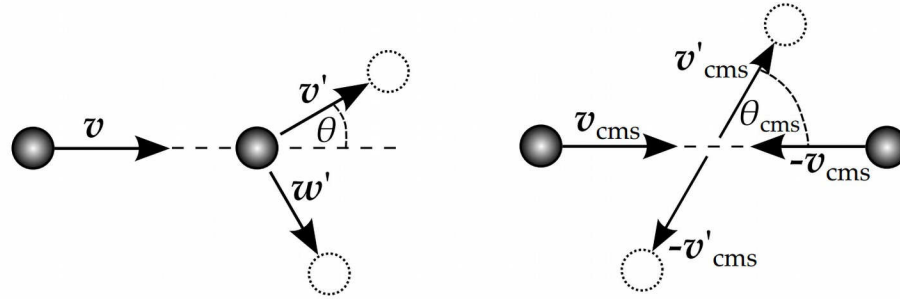
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- Recently been **observed in A3827**

Massey et al., arXiv:1504.03388

Frequent vs. rare scatters



The momentum transfer in a collision of two DM particles is completely fixed by the scattering angle. The effective momentum transfer is given by

$$\sigma_T = 2\pi \int_{-1}^1 \frac{d\sigma}{d\Omega} (1 - |\cos \theta|) d \cos \theta$$

This is the quantity typically studied

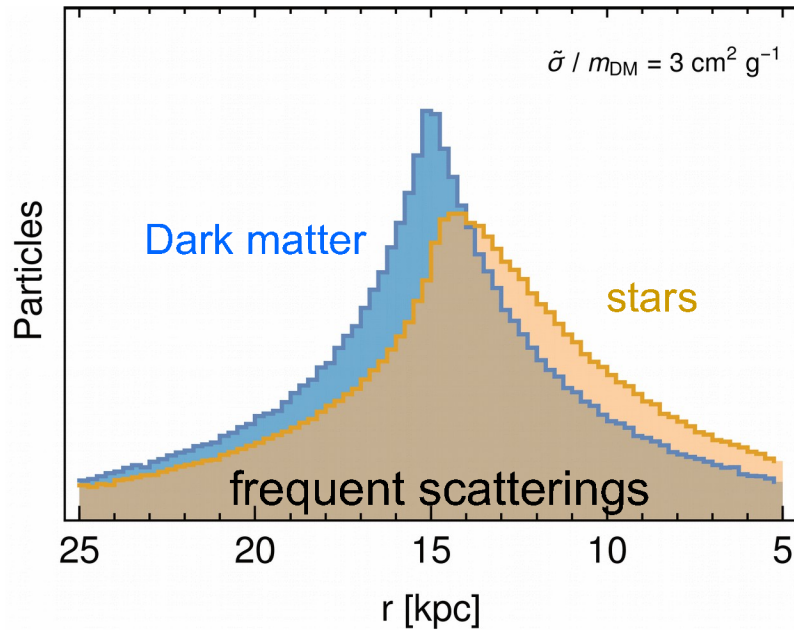
However, this is not all that matters...

Kahlhoefer et al, 1308.3419

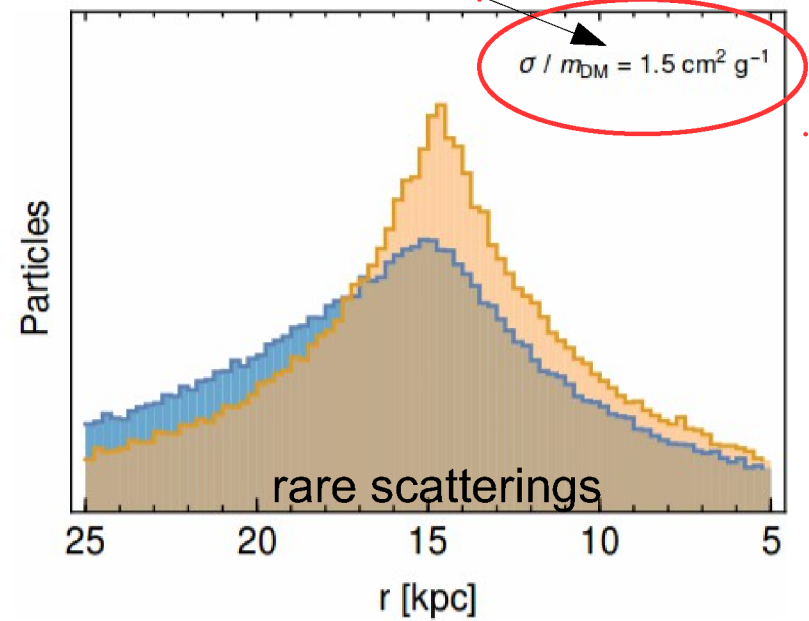
Can be obtained with **rare scatters and large momentum transfer** (e.g. isotropic scattering) or **frequent scatters with small momentum transfer** (e.g. long range interactions)

Infalling galaxy in A3827

Kahlhoefer et al, 1504.06576

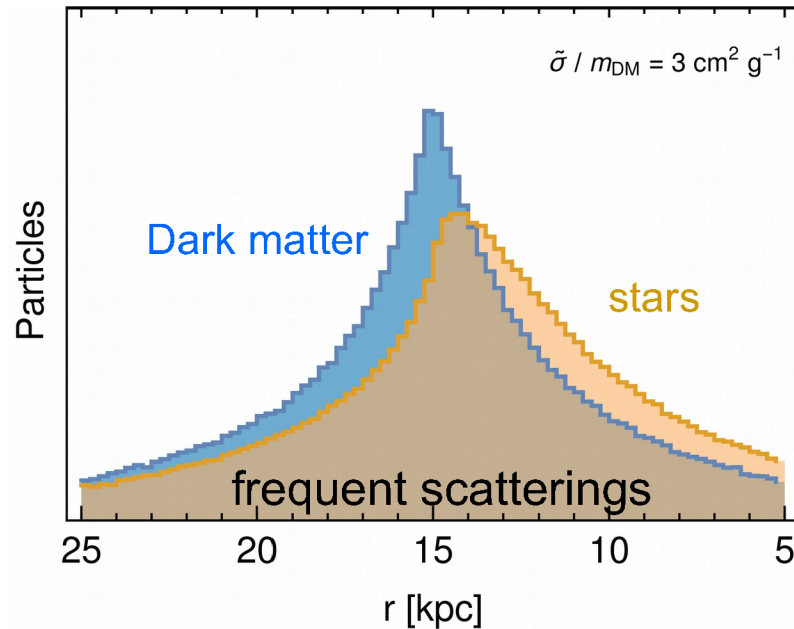


A3827: In (some) tension with upper bounds

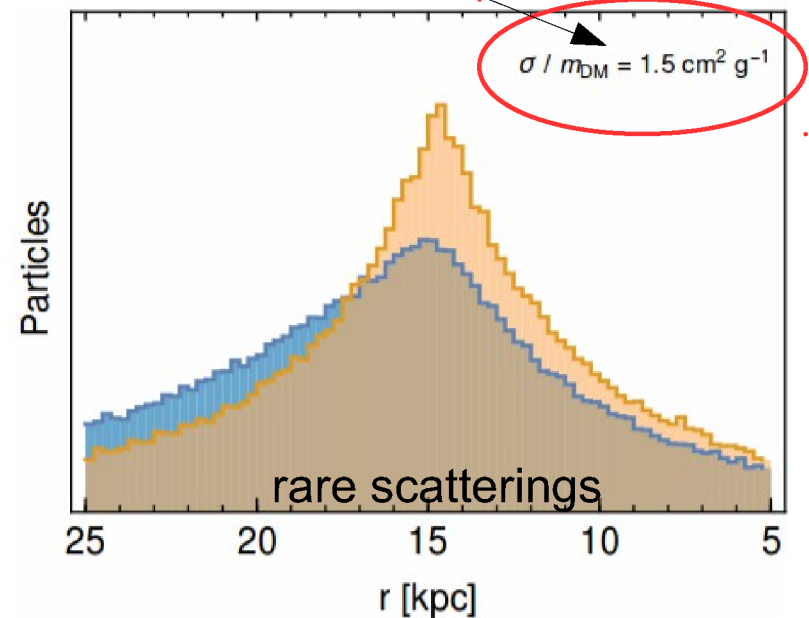


Distinguishing different types of SIDM

Kahlhoefer et al, 1504.06576



A3827: In (some) tension with upper bounds



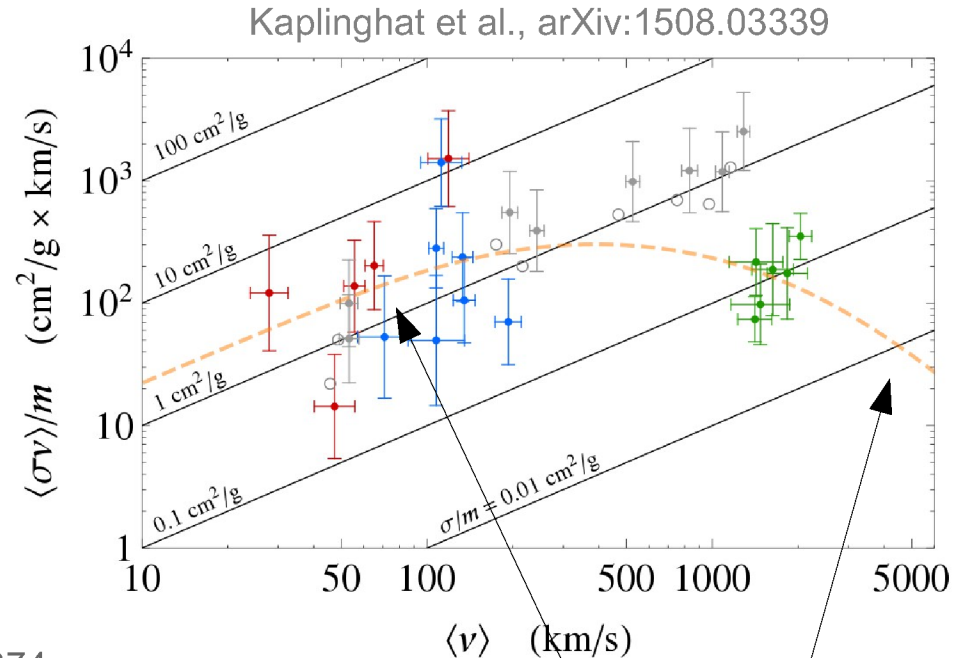
- **Effective drag force:** the DM subhalo retains its shape, while the distribution of stars are both shifted and deformed.
- **Contact interactions:** the DM subhalo is deformed due to the scattered DM particles leaving the subhalo in the backward direction.
- Potentially distinguishable!

Velocity dependent self-interactions

- Idea: Relate core size of different systems to SIDM cross section
- DM self-interactions seem to depend on the typical relative velocity of DM particles.
- Simplest realisation
→ light mediator!

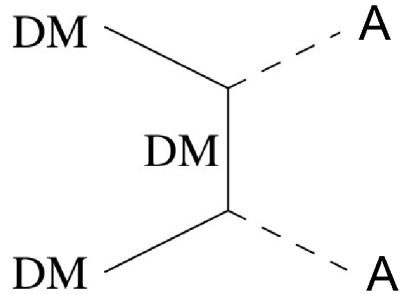
Loeb & Weiner: arXiv:1011.6374

- Consider a mediator with mass $m_{\text{med}} \sim m_{\text{DM}} v_{\text{DM}}$:
 - Scattering for small momentum transfer ($q < m_{\text{med}}$) proportional to $1/m_{\text{med}}^4$
 - Scattering for large momentum transfer ($q > m_{\text{med}}$) proportional to $1/q^4$



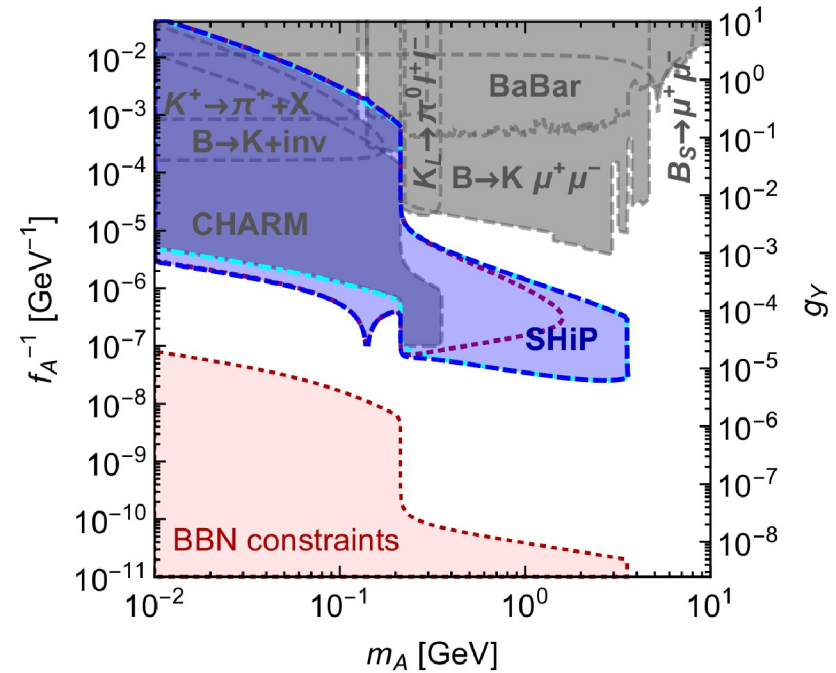
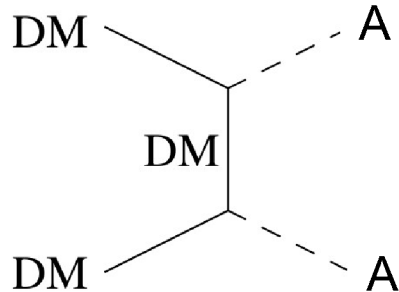
A new light mediator

- The relic abundance is typically set by annihilations into pairs of mediators (so-called dark sector freeze-out):



A new light mediator

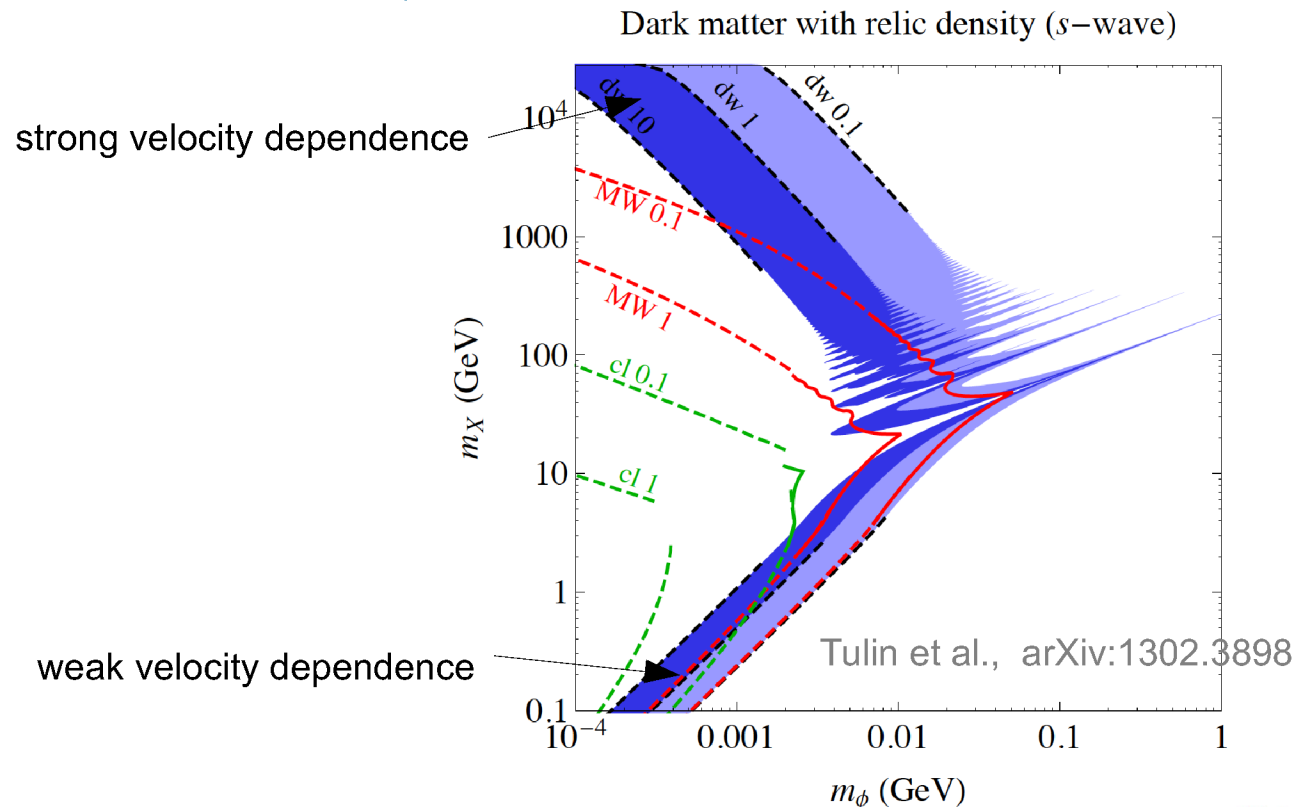
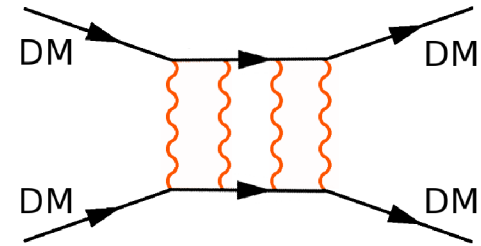
- The relic abundance is typically set by annihilations into pairs of mediators (so-called dark sector freeze-out):



- To avoid overclosing the Universe, the mediator should ultimately decay, so its couplings to SM states cannot be arbitrarily small

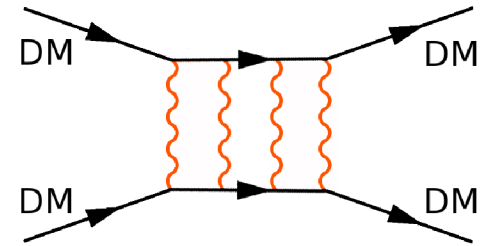
Enhancement of DM self-interactions

- DM self-interactions are enhanced also by non-perturbative effects due to multiple mediator exchange.
- Scalar and vector mediators particularly interesting

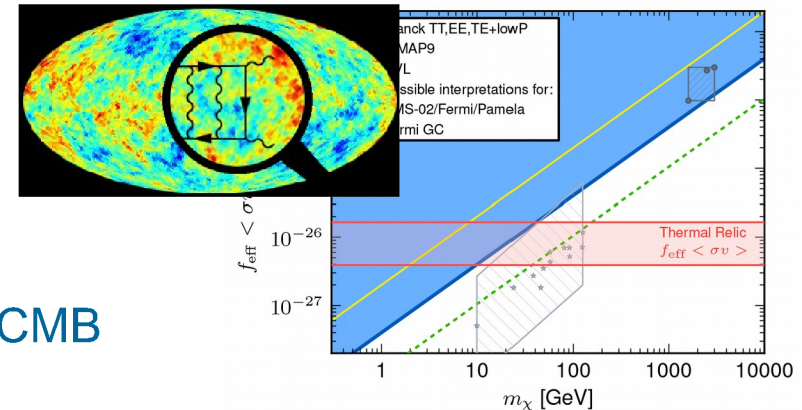


Enhancement of DM self-interactions

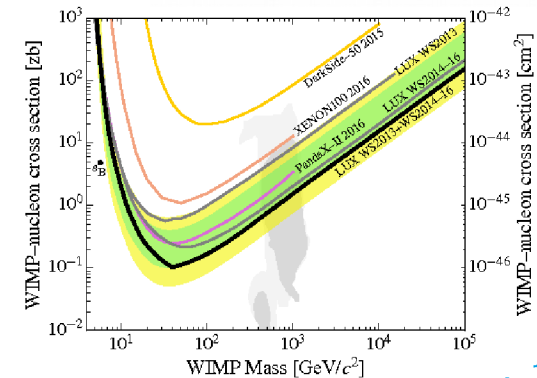
- DM self-interactions are enhanced also by non-perturbative effects due to multiple mediator exchange.
- Scalar and vector mediators particularly interesting
- In this case also Sommerfeld enhancement of annihilations



→ very strong reionisation bounds from the CMB for s-wave annihilation



- DM-nucleon scattering cross section also strongly enhanced for light mediators



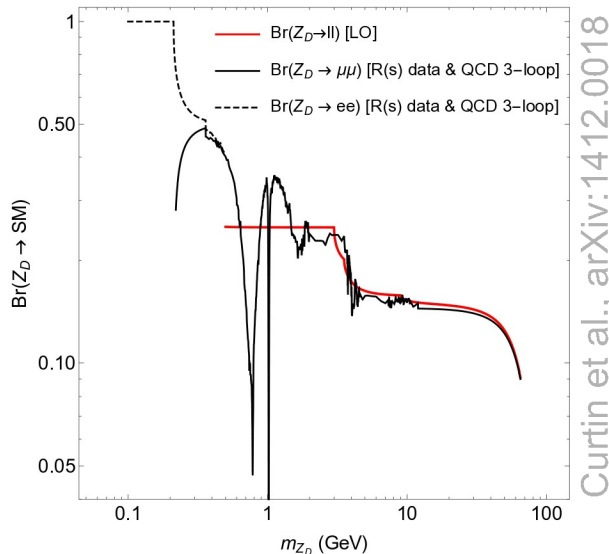
Vector mediators

- Example: A new gauge boson from a spontaneously broken U(1)' gauge group that mixes with the neutral gauge bosons of the Standard Model.

$$\mathcal{L} \supset -g_\chi^V \phi^\mu \bar{\chi} \gamma_\mu \chi - \frac{1}{2} \sin \epsilon B_{\mu\nu} \phi^{\mu\nu} - \delta m^2 \phi^\mu Z_\mu$$

Kinetic mixing:
Mediator obtains photon-like couplings

Mass mixing:
Mediator obtains Z-like couplings



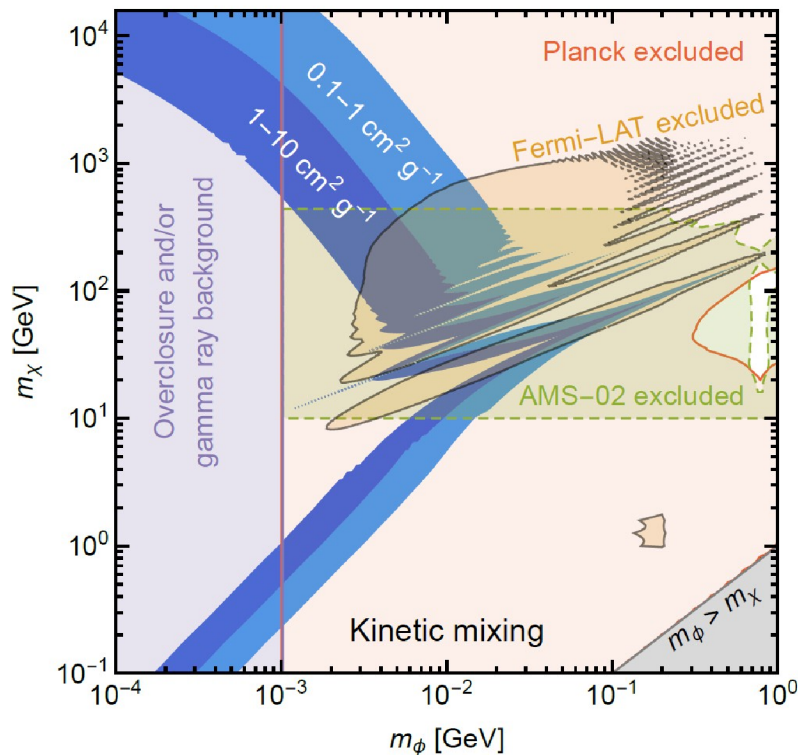
➤ Main difference:

- A gauge boson with kinetic mixing is effectively stable below the electron threshold.
- Mass mixing induces sizeable decay rates into neutrinos

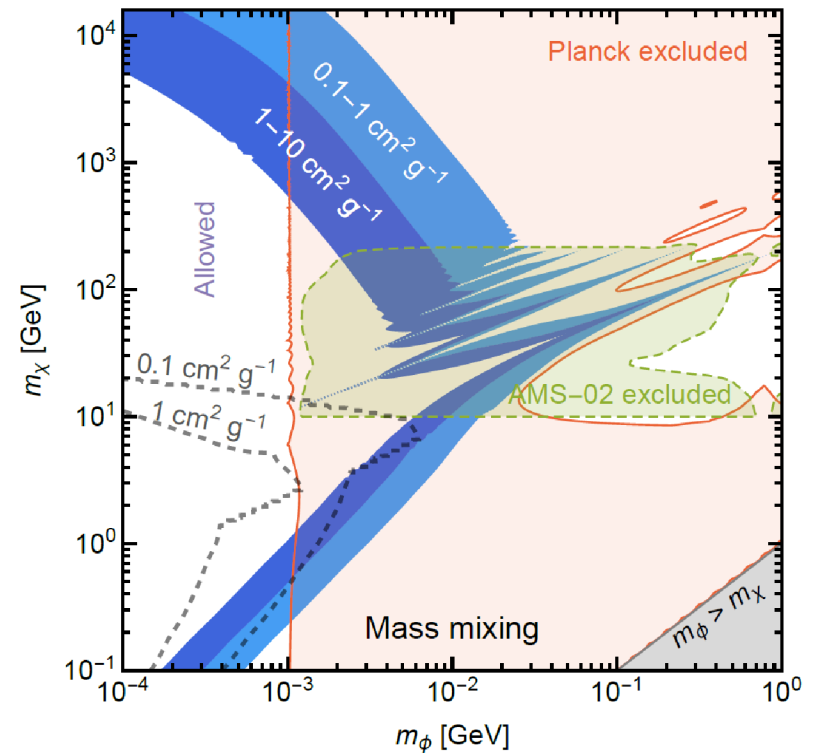
Constraints on vector mediators

- For vector mediators, DM annihilation proceeds via s-wave:
 - Large Sommerfeld enhancement for small velocities
 - g_x fixed by relic density – essentially independent of coupling to SM

Bringmann et al., arXiv:1612.00845

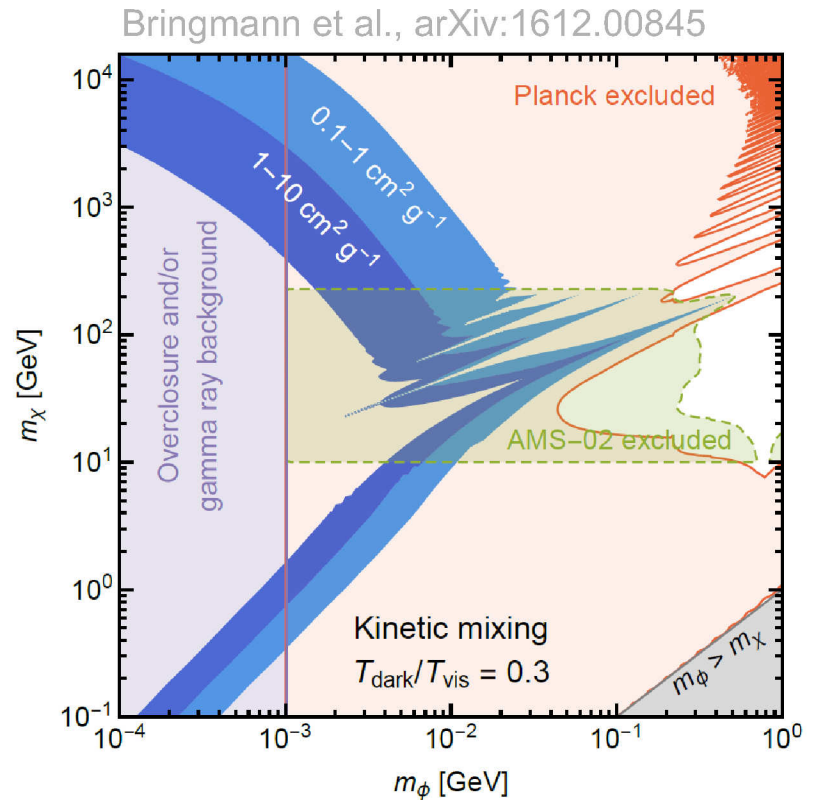


Bringmann et al., arXiv:1612.00845



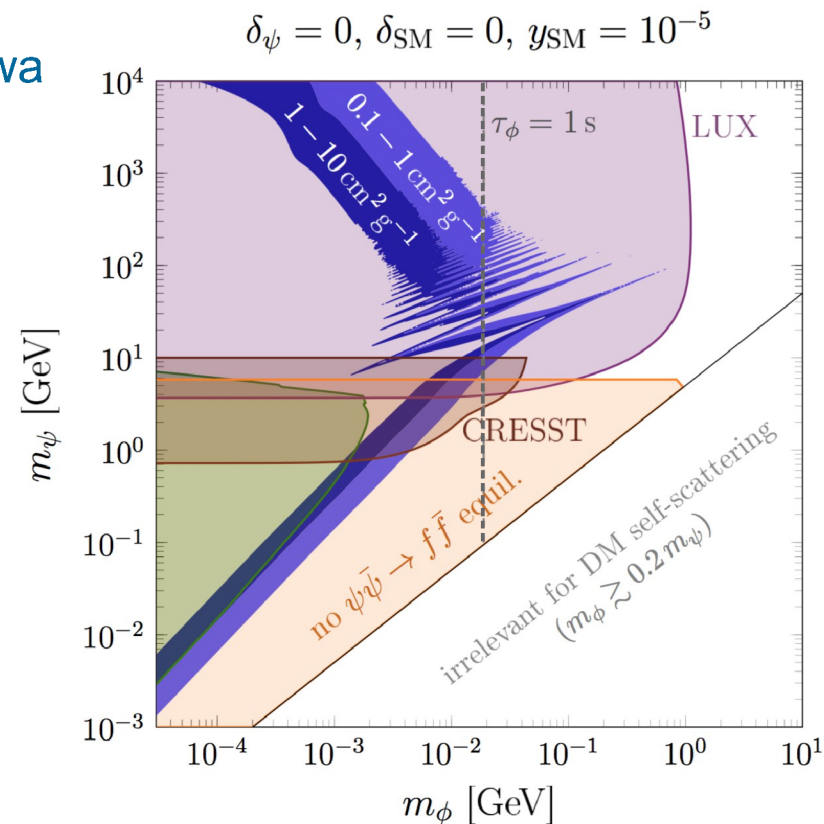
Constraints on vector mediators

- Only assumption: The two sectors have the same temperature during freeze out.
- But even for different temperatures in the two sectors there are very strong constraints.



Constraints on scalar mediators

- For fermionic DM and scalar mediators annihilation proceeds via p-wave
- No constraints from indirect detection or the CMB.
- Direct detection constraints are very strong for scalar mediators.
- Lifetime rather long due to Yukawa suppression
- Naive BBN bound: $\tau < 1$ s
- Impossible to satisfy all requirements and have large self-interaction cross sections.

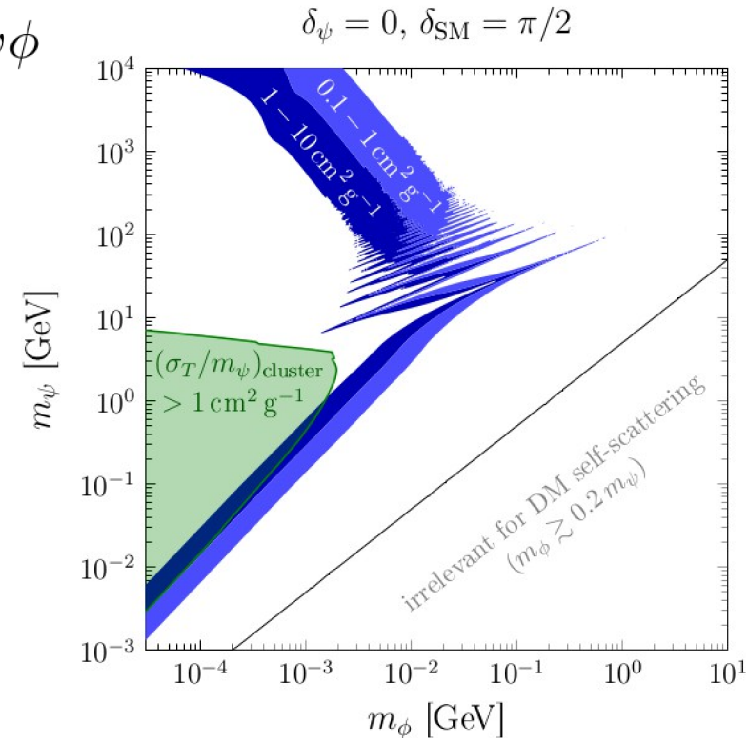


1704.02149

A mixed mediator (CP violation)

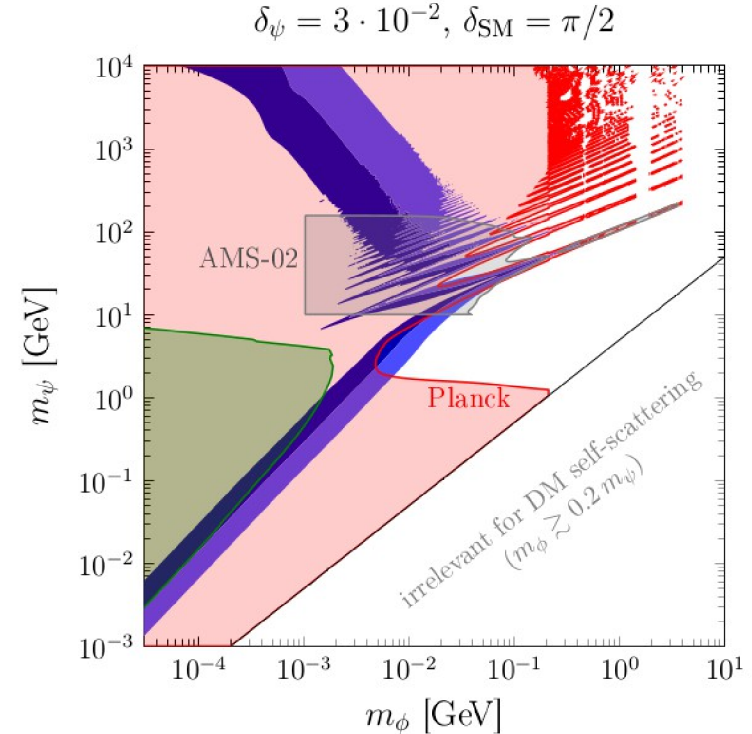
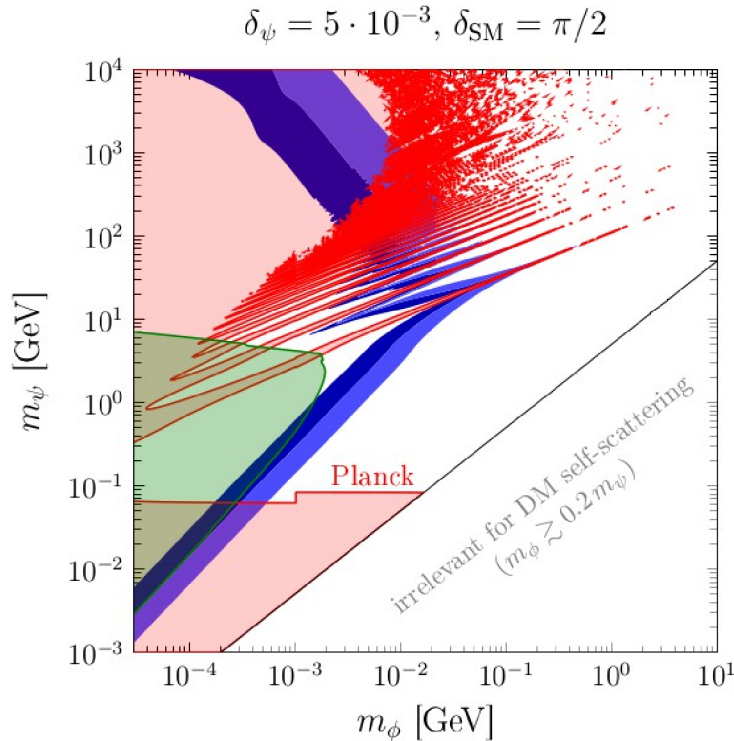
$$\mathcal{L}_{\text{DM}} \supset y_\psi \cos \delta_\psi \bar{\psi}\psi\phi + y_\psi \sin \delta_\psi i\bar{\psi}\gamma^5\psi\phi$$

- For $\delta_\psi \sim 0$ (like a scalar) DM self-interactions can be large.
- For $\delta_{\text{SM}} \sim \pi/2$ (like a pseudoscalar) direct detection constraints are strongly suppressed.
- Large allowed parameter space!
- Constraints on the CP-violating phase δ_{SM} (e.g. from electron EDMs) can be satisfied even for very light mediators as long as y_{SM} is sufficiently small ($y_{\text{SM}} \ll 10^{-2}$).



The return of CMB constraints

- Central problem: The fact that annihilation can only proceed via p-wave was a consequence of CP conservation.
- As soon as δ_ψ is not exactly zero, s-wave annihilation is again possible and will receive large Sommerfeld enhancement.



Future directions for light mediators

- There are a number of ways to evade the various constraints
 - Inert decays of the mediator, for example into (sterile) neutrinos
 - Thermalization via a different mechanism (possibly leading to different temperatures during freeze-out)
 - No thermalization (DM production via the freeze-in mechanism)
Bernal et al., arXiv:1510.08063
 - Suppressed couplings to quarks (to evade direct detection constraints)
- Nevertheless, constraints from BBN, direct detection and the CMB are very generic and will generally be relevant to any model of DM interacting via a new light mediator.
- Exciting phenomenology and interesting model-building challenges!

Summary

- Self interacting dark matter could solve some problems of the collisionless cold dark matter paradigm and can arise naturally in more complex dark sectors
- Orthogonal handle on properties of DM: We can potentially study the dark sector even if DM has highly suppressed couplings to Standard Model particles.
- Can potentially distinguish effective drag forces (from frequent self-interactions) and rare self-interactions
- Also could infer the velocity dependence of the cross section.
- The simplest possibilities (scalar or vector mediator coupling to fermionic dark matter with no additional new states) are in strong tension with direct and indirect detection experiments.
- One simple way out is spontaneous CP violation in the dark sector
- Huge possible impact, ruling out WIMPs, axions, gravitinos,...

Thank you!