

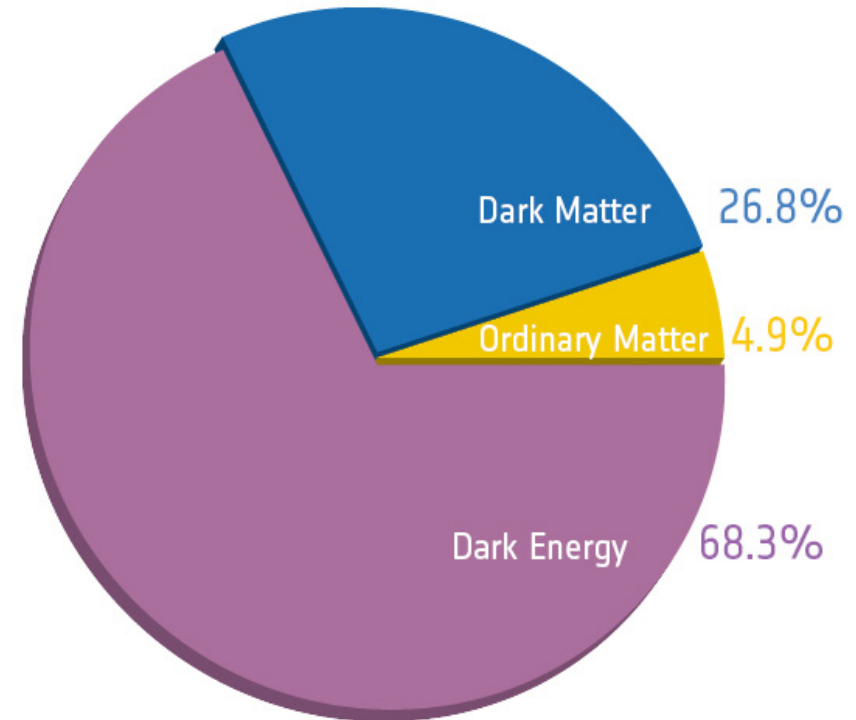
Dark-Matter Axion Mass.

Andreas Ringwald (DESY)

Workshop on Ultralight Dark Matter and Axions
Ann Arbor, MI, USA
5-7 March 2018

Strong Case for Particles Beyond the Standard Model

- > Standard Model (SM) of particle physics describes properties of known matter and forces to a great precision
- > SM describes only about 15% of matter content of Universe

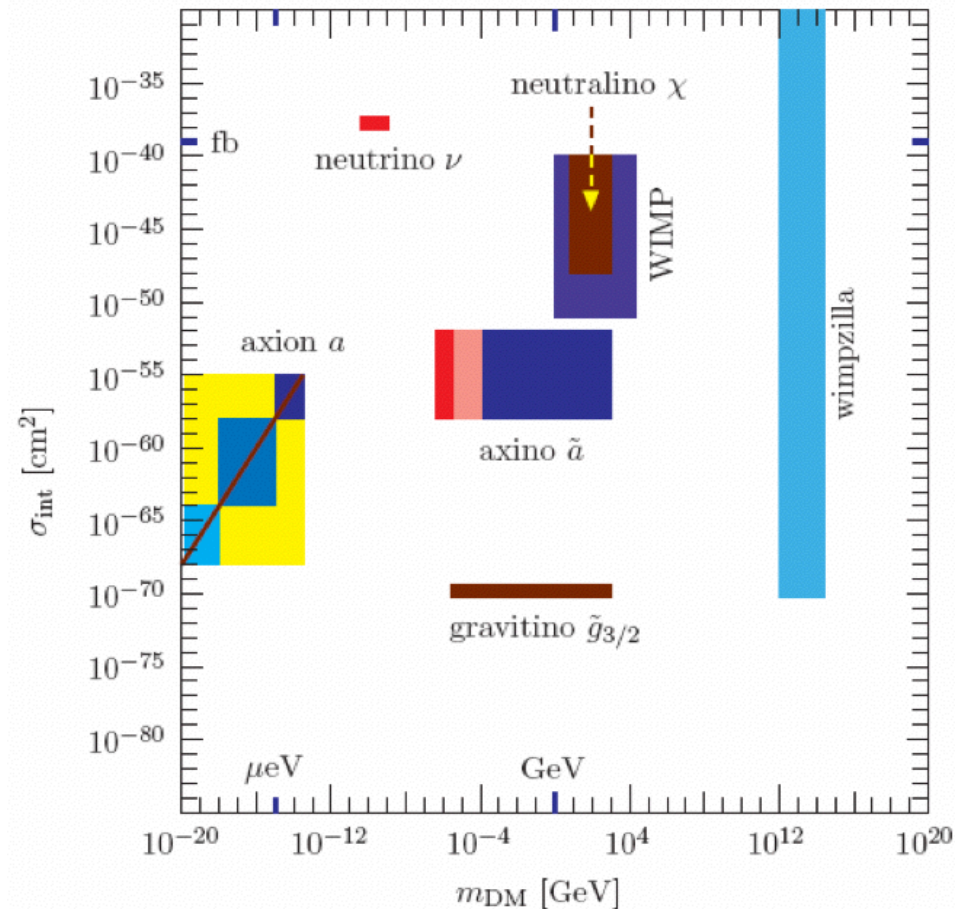


[PLANCK]



Strong Case for Particles Beyond the Standard Model

- Standard Model (SM) of particle physics describes properties of known matter and forces to a great precision
- SM describes only about 15% of matter content of Universe
- Theorists have proposed plenitude of DM candidates, spanning a large parameter space in mass and interaction strength
- Best motivated candidates those which occur in SM extensions solving also other problems, such as
 - Hierarchy problem: [Neutralino](#) in MSSM
 - Strong CP problem: [Axion](#) in PQSM



[Kim, Carosi 10]



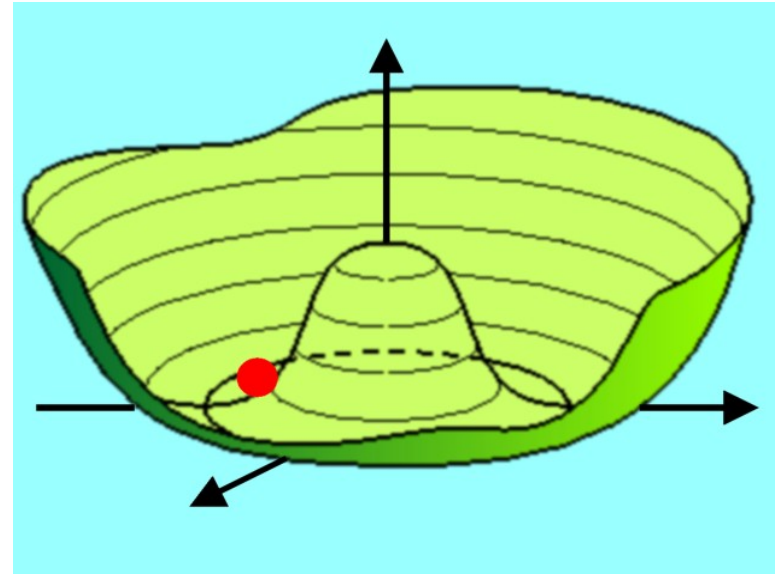
PQSM: Peccei-Quinn Extensions of the Standard Model

> A singlet complex scalar field σ featuring a global $U(1)_{\text{PQ}}$ symmetry is added to SM

> Symmetry is broken by vev $\langle \sigma \rangle = v_{\text{PQ}}/\sqrt{2}$

$$\sigma(x) = \frac{1}{2} (v_{\text{PQ}} + \rho(x)) e^{iA(x)/v_{\text{PQ}}}$$

- Excitation of modulus: $m_\rho \sim v_{\text{PQ}}$
- Excitation of angle: NGB $m_A = 0$



[Raffelt]

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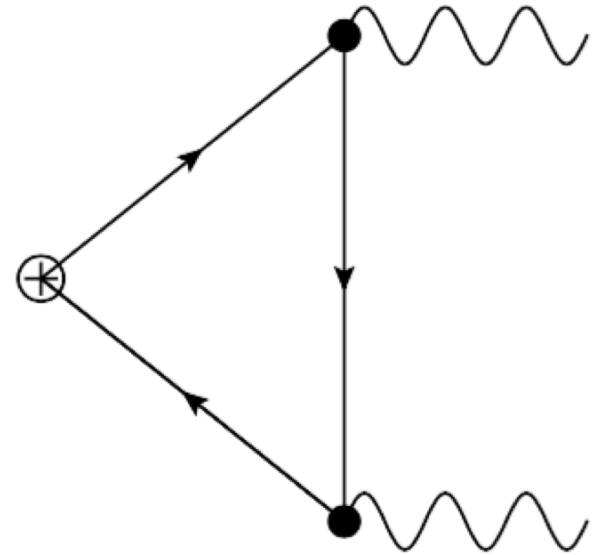
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➤ Quarks (SM or extra) carry PQ charges such that $U(1)_{\text{PQ}}$ is anomalously broken due to gluonic triangle anomaly:

$$\partial_\mu J_{U(1)_{\text{PQ}}}^\mu \supset -\frac{\alpha_s}{8\pi} N G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$



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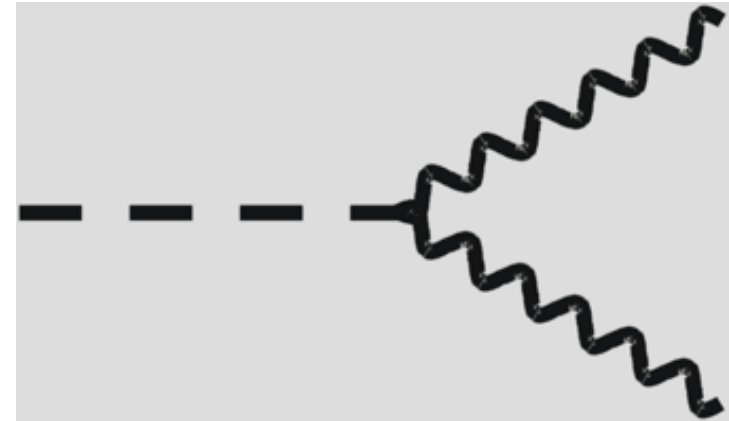
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- > Low energy effective field theory at energies below v ($\ll v_{\text{PQ}}$), but above Λ_{QCD} , [Peccei,Quinn 77; Weinberg 78; Wilczek 78]

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} \frac{A(x)}{f_A} G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}; \quad f_A = v_{\text{PQ}}/N$$

[Kim 79; Shifman, Vainshtein, Zakharov 80; Zhitnitsky 80; Dine, Fischler, Srednicki 81; ...]



Axionic Solution of Strong CP Problem

- > Can eliminate QCD theta term,

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} [\bar{\theta} + \theta(x)] G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$

$$\theta(x) \equiv A(x)/f_A$$

by shift $\theta(x) \rightarrow \theta(x) - \bar{\theta}$

- > Effective potential at energies below Λ_{QCD} ,

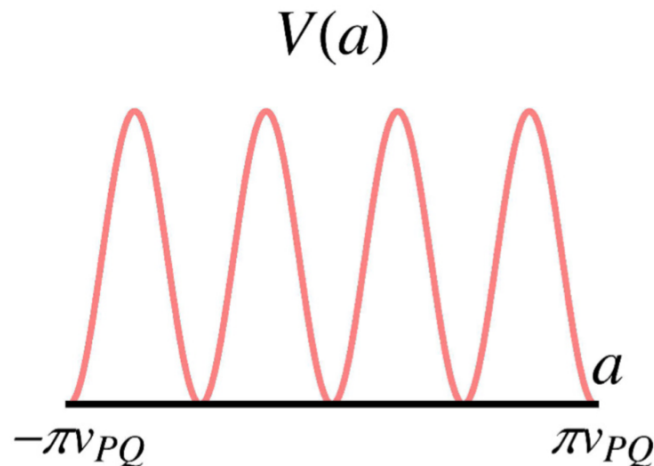
$$V(\theta) = \Sigma (m_u + m_d) \left(1 - \frac{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \theta}}{m_u + m_d} \right)$$

$$\Sigma \equiv -\langle \bar{u}u \rangle = -\langle \bar{d}d \rangle$$

[Di Vecchia, Veneziano '80;
Leutwyler, Smilga 92]

has absolute minimum at $\theta = 0$ and thus predicts vanishing vev, $\langle \theta(x) \rangle = 0$

[Peccei, Quinn 77]



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[Peccei, Quinn 77]

- > Particle excitation: pseudo NG boson “axion”

[Weinberg 78; Wilczek 78]

- > Mass: $m_A \simeq \frac{\sqrt{\Sigma}}{f_A} \sqrt{\frac{m_u m_d}{m_u + m_d}} \simeq \frac{m_\pi f_\pi}{f_A} \frac{\sqrt{m_u m_d}}{m_u + m_d} \simeq 6 \text{ meV} \left(\frac{10^9 \text{ GeV}}{f_A} \right)$

- Precisely (NLO CPT; Lattice QCD):

$$m_A = 57.0(7) \left(\frac{10^{11} \text{ GeV}}{f_A} \right) \mu\text{eV}$$

[Grilli di Cortona et al. '16 ;
Borsanyi et al. '16]

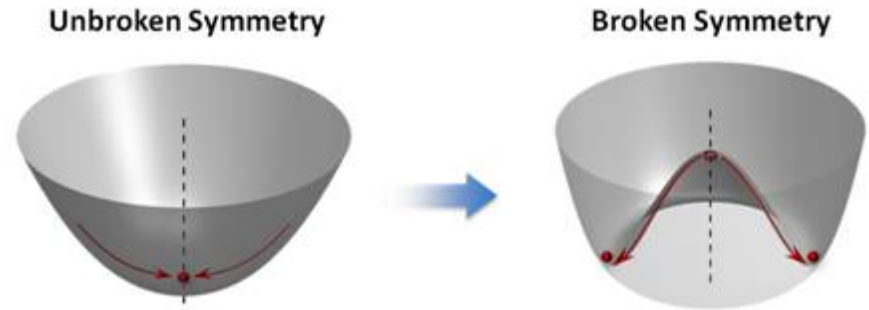


Axion Cold Dark Matter

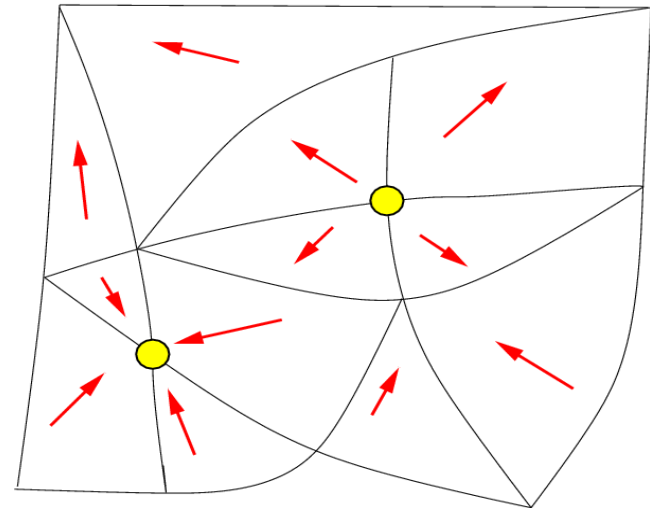
- > Axion field is born after PQ symmetry breaking

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

- > In causally connected regions at phase transition, axion takes random initial values



[Peking University]



[Uhlmann et al. '10]

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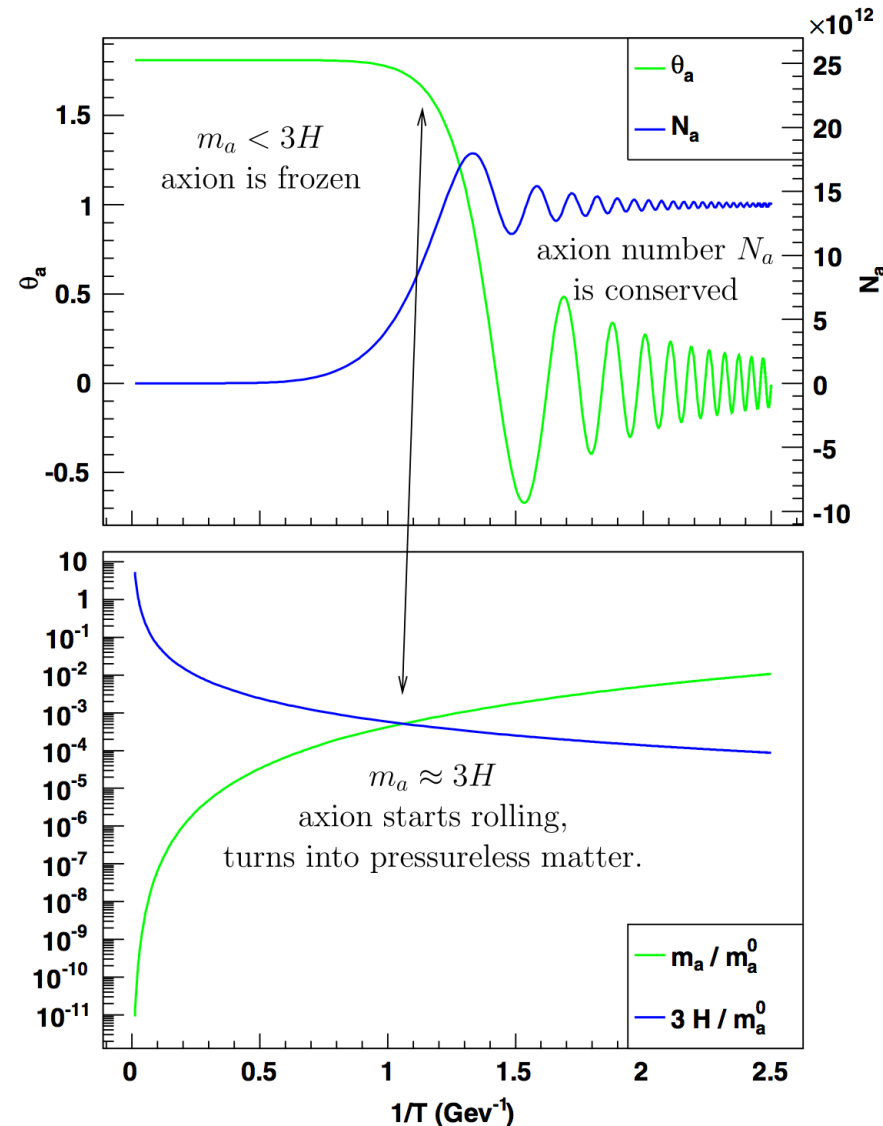
- > Axion field is born after PQ symmetry breaking

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

- > In causally connected regions at phase transition, axion takes random initial values and sticks
- > Later when $H(T) \sim m_A(T)$, axion field starts to oscillate around minimum of potential; behaves like cold dark matter:

$$w_A = p_A / \rho_A \simeq 0$$

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]



[Wantz,Shellard '09]



Axion Cold Dark Matter

➤ Precise determination of evolution needs QCD input:

1. Equation of state at temperatures around 1 GeV: determines

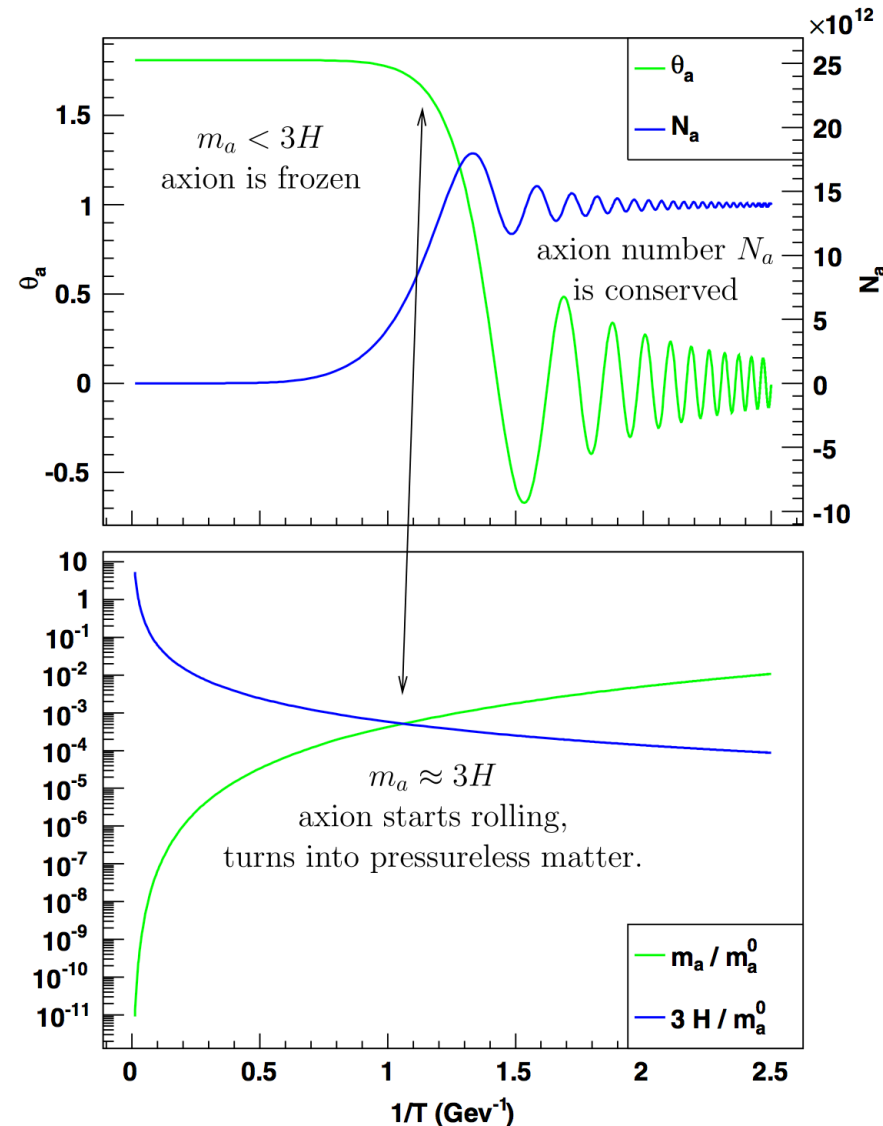
$$H(T)$$

2. Topological susceptibility:

$$\chi(T) \equiv \int d^4x \langle q(x)q(0) \rangle_T$$

determines

$$m_A^2(T) = \chi(T)/f_A^2$$



[Wantz, Shellard '09]



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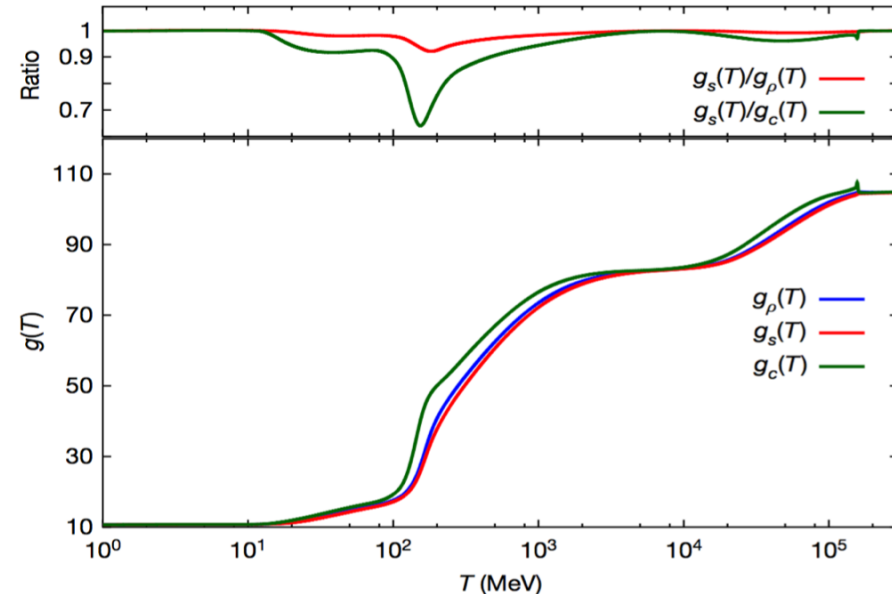
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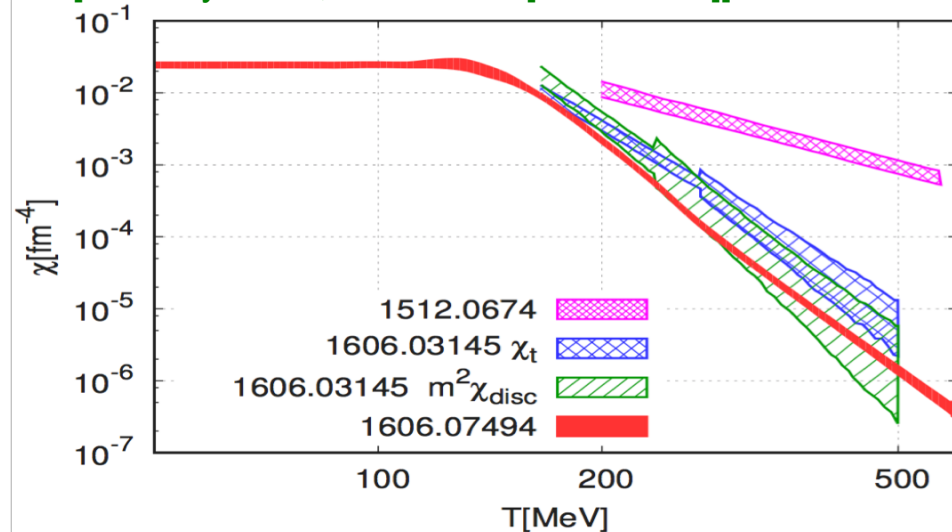
$$m_A^2(T) = \chi(T)/f_A^2$$

➤ Both obtained recently with high precision via lattice QCD

- Temperature slope of $\chi(T)$ remarkably close to prediction from dilute instanton gas; latter underestimates normalization by factor of order ten



[Borsanyi et al., Nature '16 [1606.0794]]



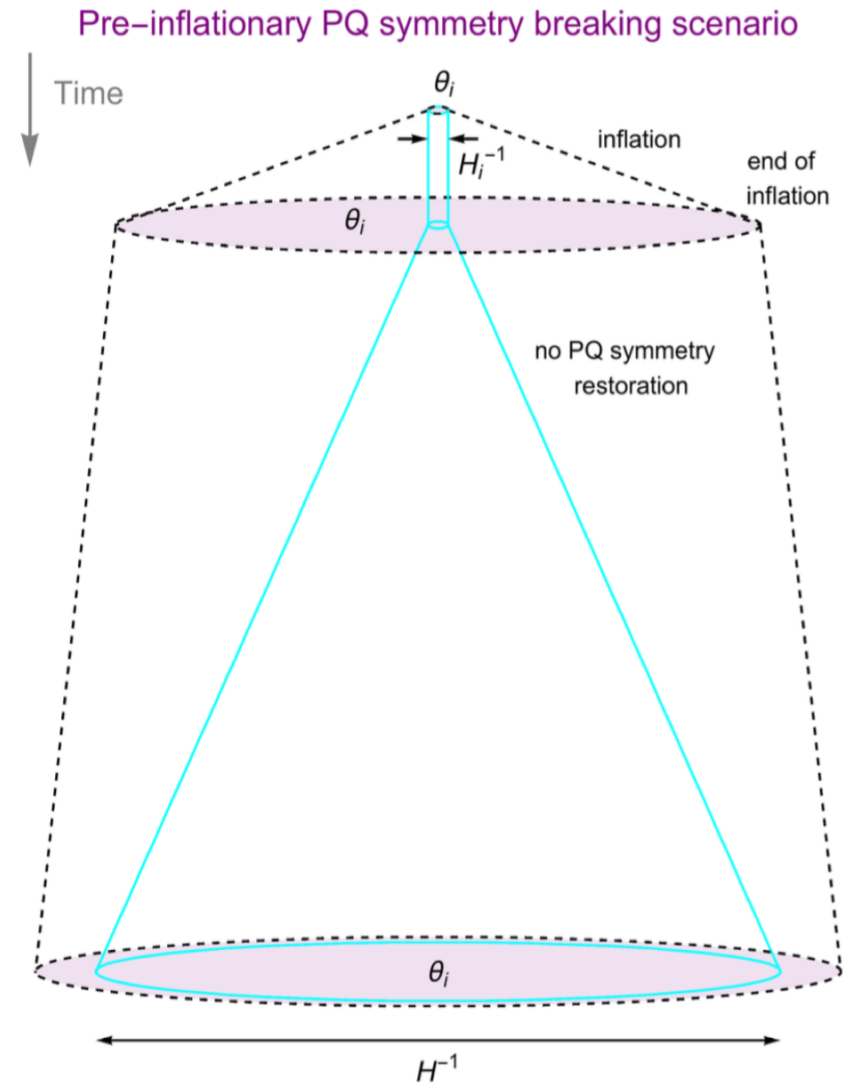
Axion Cold Dark Matter

- If PQ symmetry broken before or during inflation ($f_A > H_I/(2\pi)$) and not restored afterwards (pre-inflationary PQ breaking scenario):

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,....]

- Axion CDM density depends on single initial value in patch which becomes observable universe and f_A :

$$\begin{aligned}\Omega_A^{\text{vr}} h^2 &\approx 0.12 \left(\frac{f_A}{9 \times 10^{11} \text{ GeV}} \right)^{1.165} \theta_i^2 \\ &\approx 0.12 \left(\frac{6 \mu\text{eV}}{m_A} \right)^{1.165} \theta_i^2 ,\end{aligned}$$



[Saikawa]

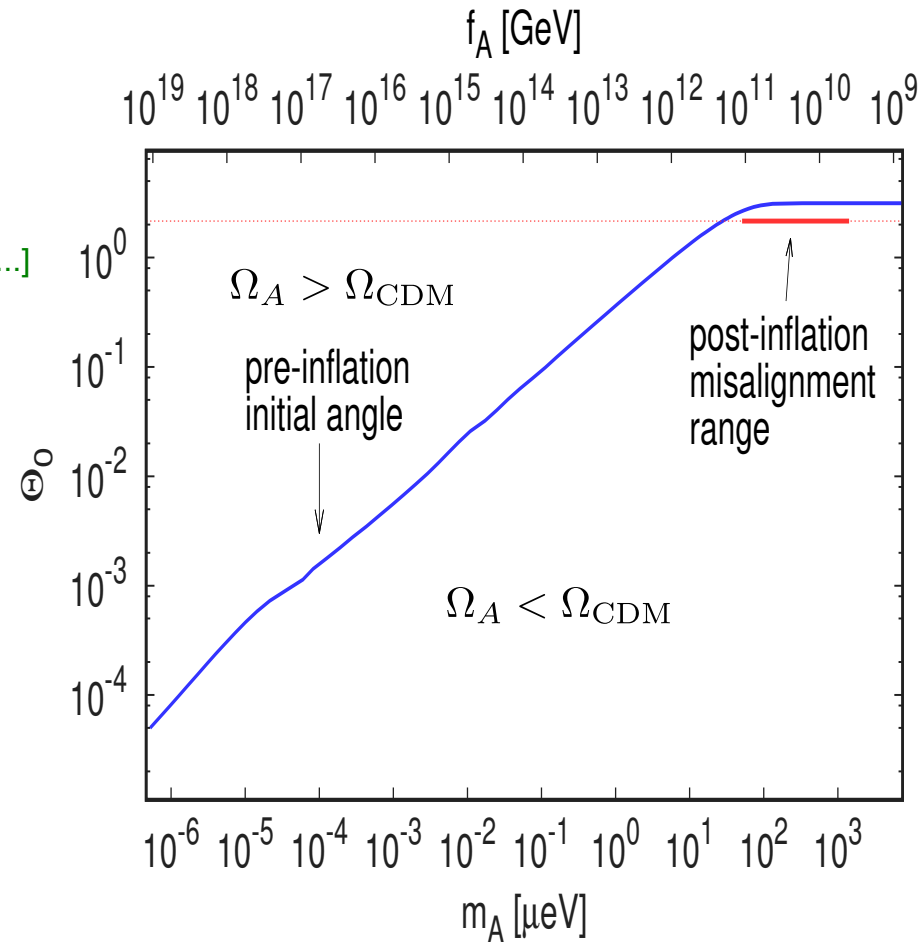
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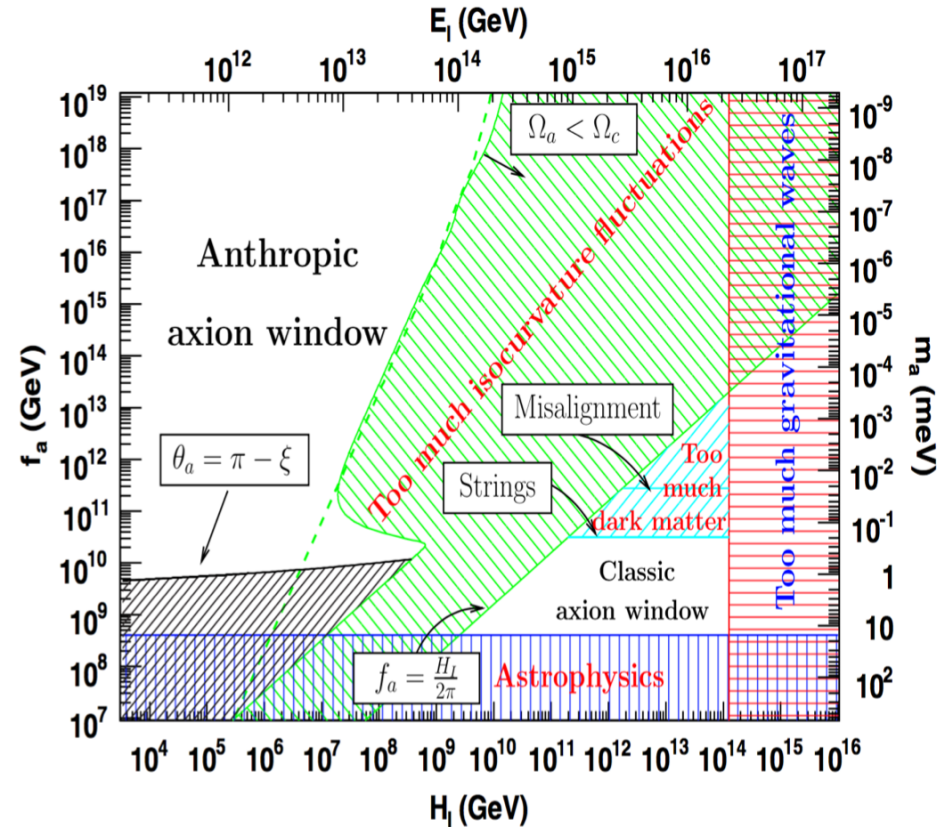
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- Upper bound on scale of inflation from isocurvature fluctuations produced by the axion during inflation and not erased afterwards



[Wilczek,Turner '91; Beltran et al. 06;
Hertzberg,Tegmark,Wilczek 08; Visinelli,Gondolo 09;
Hamann et al. 09; **Wantz,Shellard 09**]

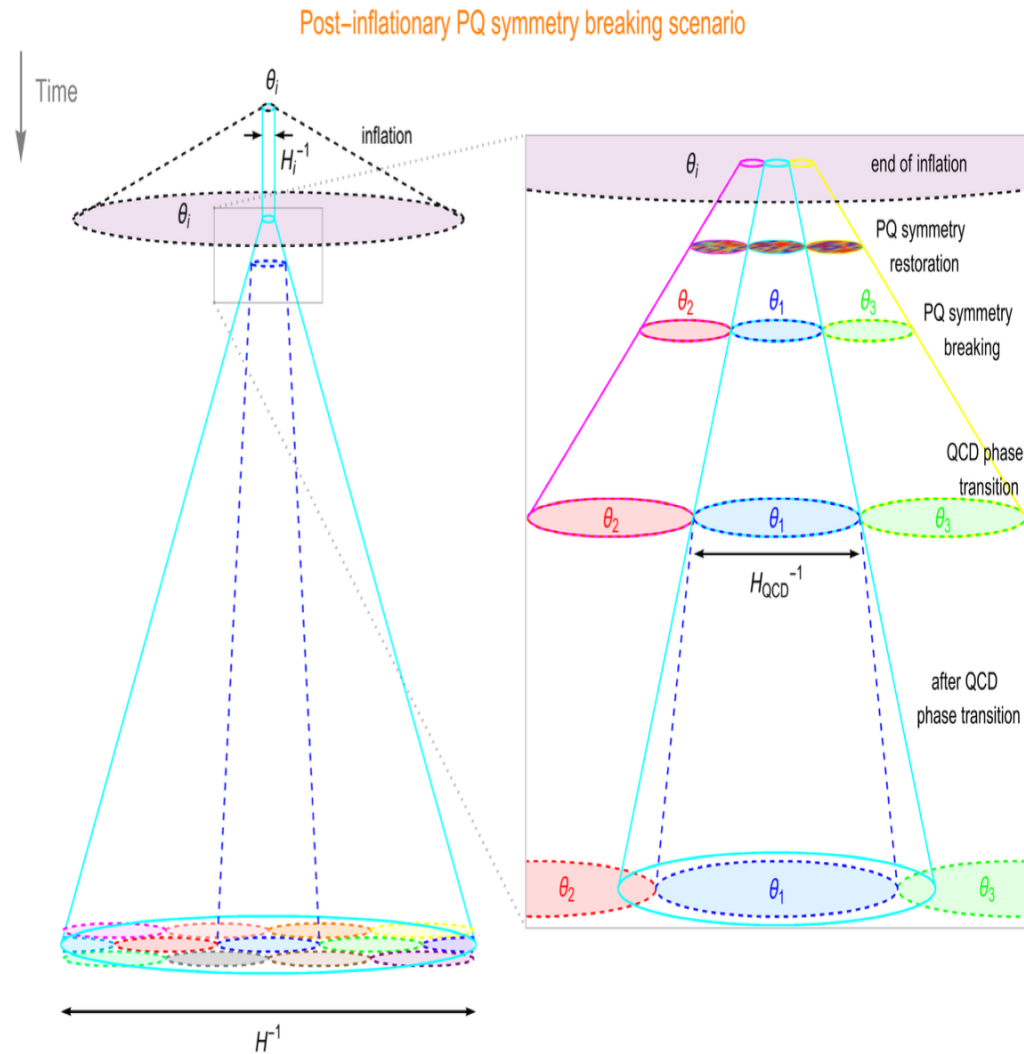


Axion Cold Dark Matter

> If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario):

- Random initial axion field values in regions of post-inflationary causal contact; naive average:

$$\Omega_A^{\text{vr}} h^2 \approx 0.12 \left(\frac{30 \mu\text{eV}}{m_A} \right)^{1.165}$$

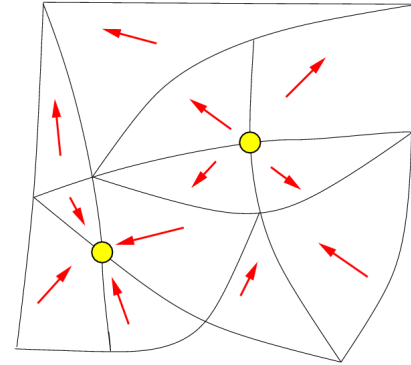


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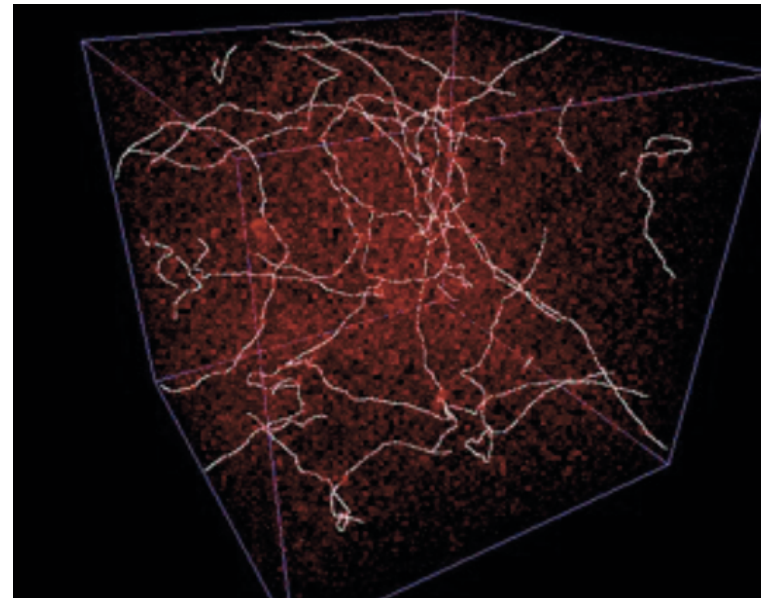


Axion Cold Dark Matter

- If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario):
 - Random initial axion field values in regions of post-inflationary causal contact
 - Network of axionic cosmic strings created by Kibble mechanism



[Uhlmann et al. '10]



[Hiramatsu et al.]

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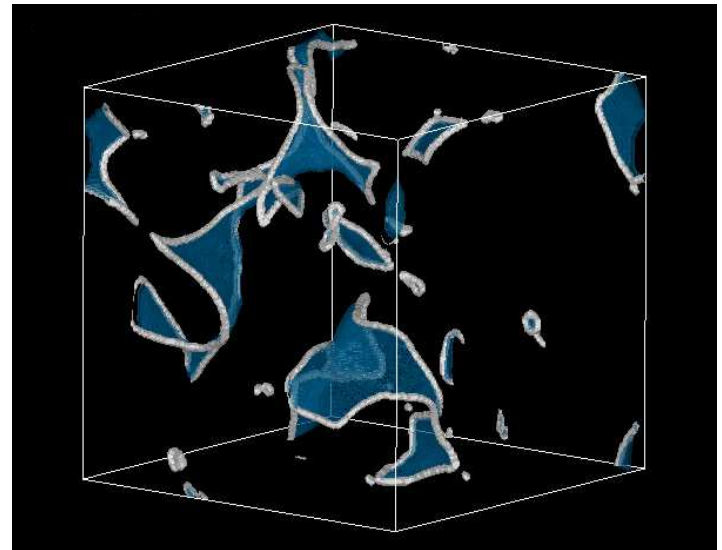
- Around QCD phase transition, axion potential develops,

$$V(A, T) = \chi(T) \left[1 - \cos \left(N \frac{A}{v_{\text{PQ}}} \right) \right]$$

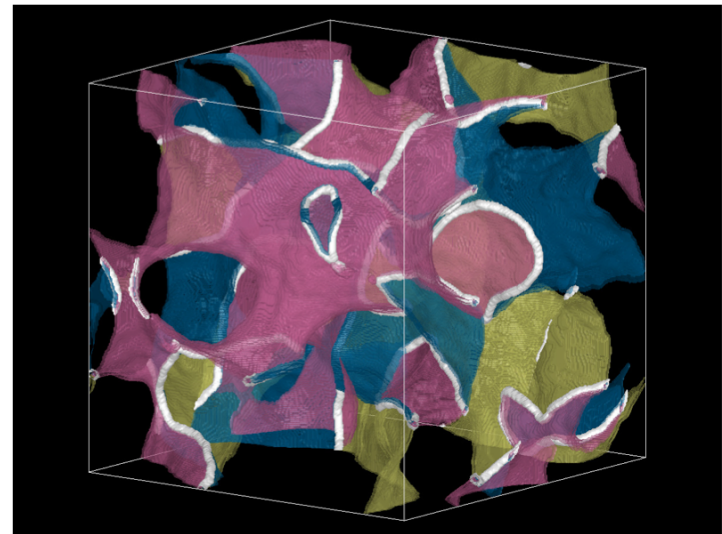
N domain walls end at string

- $N = 1$: String-wall system decays
- $N > 1$: Domain wall problem

$N = 1$



$N = 3$



[Hiramatsu et al.]

Axion Cold Dark Matter

> If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario):

- For $N = 1$, exploiting results from field theoretic lattice simulations, updated to latest determination of topological susceptibility, find CDM explained for

$$m_A \approx (30\text{--}200) \mu\text{eV}$$

[Hiramatsu et al. 11,12,13;
Kawasaki,Saikawa,Segikuchi 15;
Borsanyi et al. 16;
Ballesteros et al. 16]

Large uncertainty due to extrapolation of string tension to physical value

- New method allows to simulate directly at physical string tension, resulting in

$$m_A = (26.2 \pm 3.4) \mu\text{eV}$$

[Klaer,Moore `17]



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- For $N > 1$, domain wall problem can be avoided if PQ symmetry explicitly broken, e.g. by Planck suppressed operators, $\mathcal{L} \supset g M_{\text{P}}^4 (\sigma/M_{\text{P}})^{\mathcal{N}} + \text{h.c.}$, for $\mathcal{N} = 9, 10$,

$$0.2 \text{ meV} \lesssim m_A \lesssim 50 \text{ meV}$$

[AR,Saikawa `16;
Giannotti,Irastorza,Redondo,AR,Saikawa`17]

May employ discrete symmetry to forbid lower dimensional operators e.g. [Dias et al. `14]

A DFSZ axion ($N = 6$) in this mass range explains excessive stellar energy losses

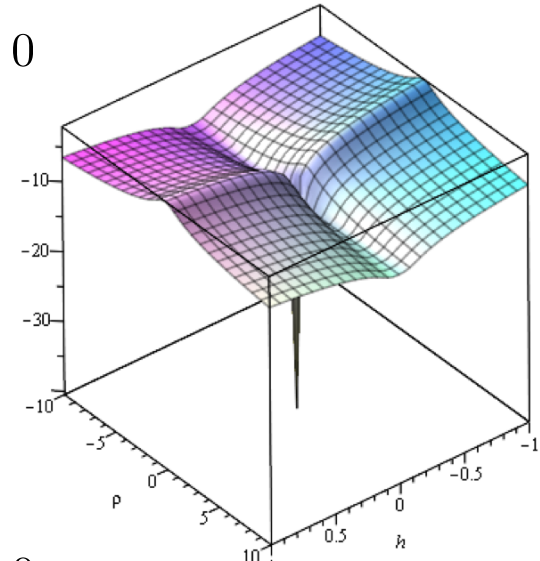


Unifying Inflation and Dark Matter with PQ Field

- > Saxion, $|\sigma| = \rho/\sqrt{2}$, or mixture with Higgs modulus, may play role of inflaton, if it has non-minimal coupling to gravity: [Fairbairn,Hogan,Marsh `14; Ballesteros,Redondo,AR,Tamarit `16]

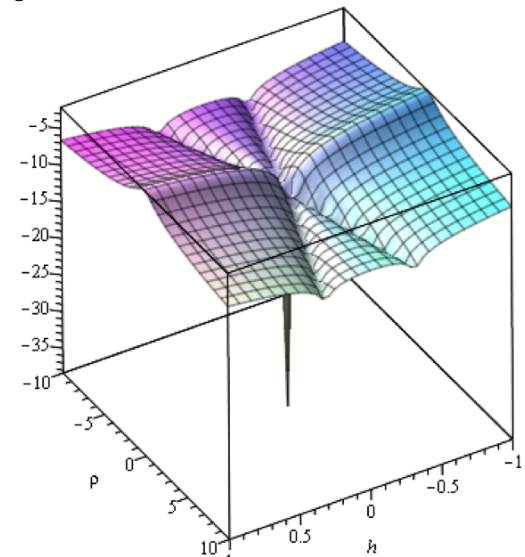
$$S \supset - \int d^4x \sqrt{-g} \xi_\sigma \sigma^* \sigma R$$

$$\lambda_{H\sigma} > 0$$



$$\lambda_{H\sigma} < 0$$

[Ballesteros,Redondo, AR,Tamarit, 1610.01639]



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- > CMB observables fit by

$$\xi \simeq 2 \times 10^5 \sqrt{\lambda} \gtrsim 10^{-3}$$

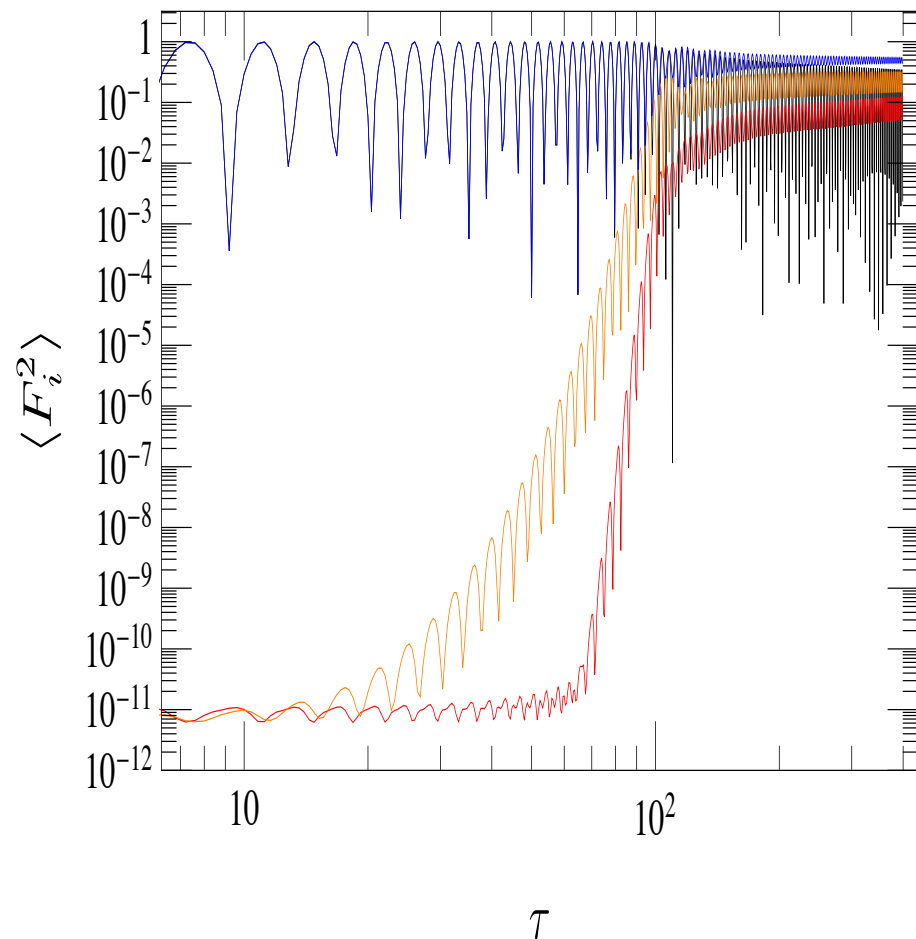
- > PQ symmetry restored after inflation already in preheating stage

- > Dark-matter axion mass:

$$30 \mu\text{eV} \lesssim m_A \lesssim 50 \text{ meV}$$

- > Large reheating temperature: 10^{10} GeV

- > Axion dark radiation: $\Delta N_\nu^{\text{eff}} \simeq 0.03$

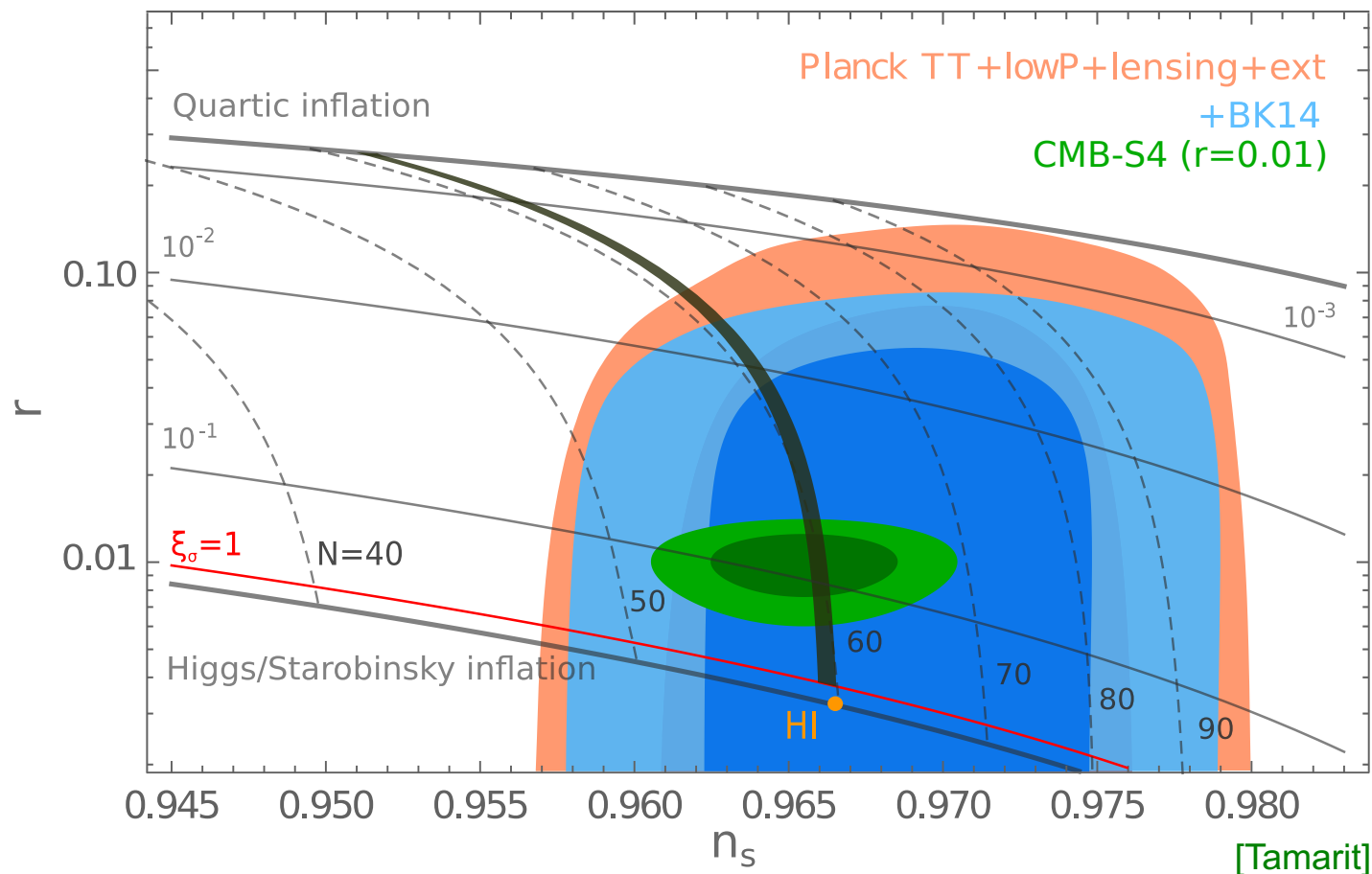


[Ballesteros, Redondo, AR, Tamarit '16]



Unifying Inflation and Dark Matter with PQ Field

- > Sharp prediction of r vs n_s for fixed pivot scale, e.g. $k_0 = 0.002 \text{ Mpc}^{-1}$



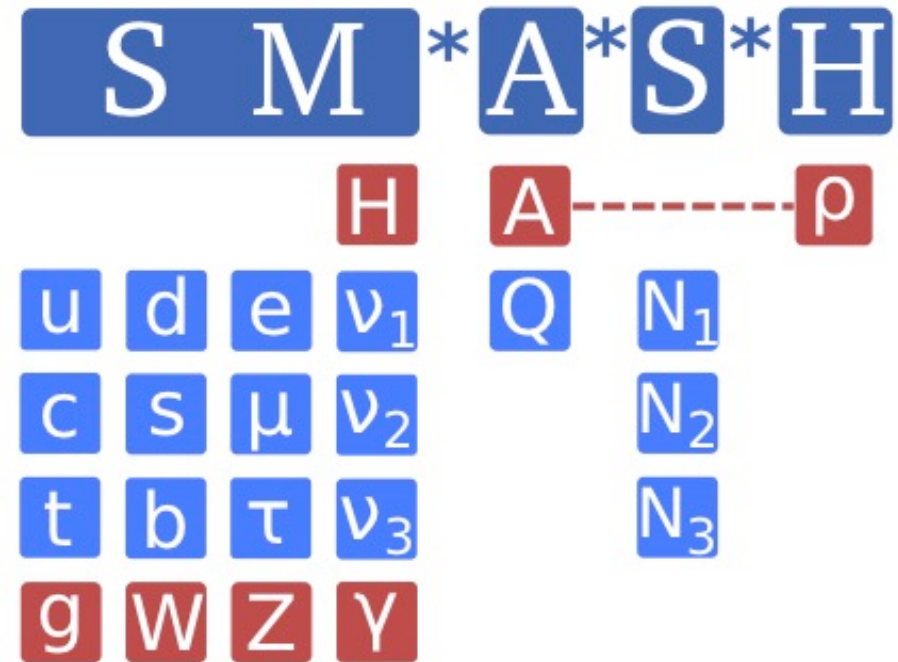
- > Can be probed by next generation CMB experiments (e.g. CMB-S4)

Unifying Inflation, Dark Matter, and Seesaw with PQ Field

> Augmenting axion models with three SM singlet neutrinos, getting their Majorana masses also through the vev $v_\sigma = N f_A$

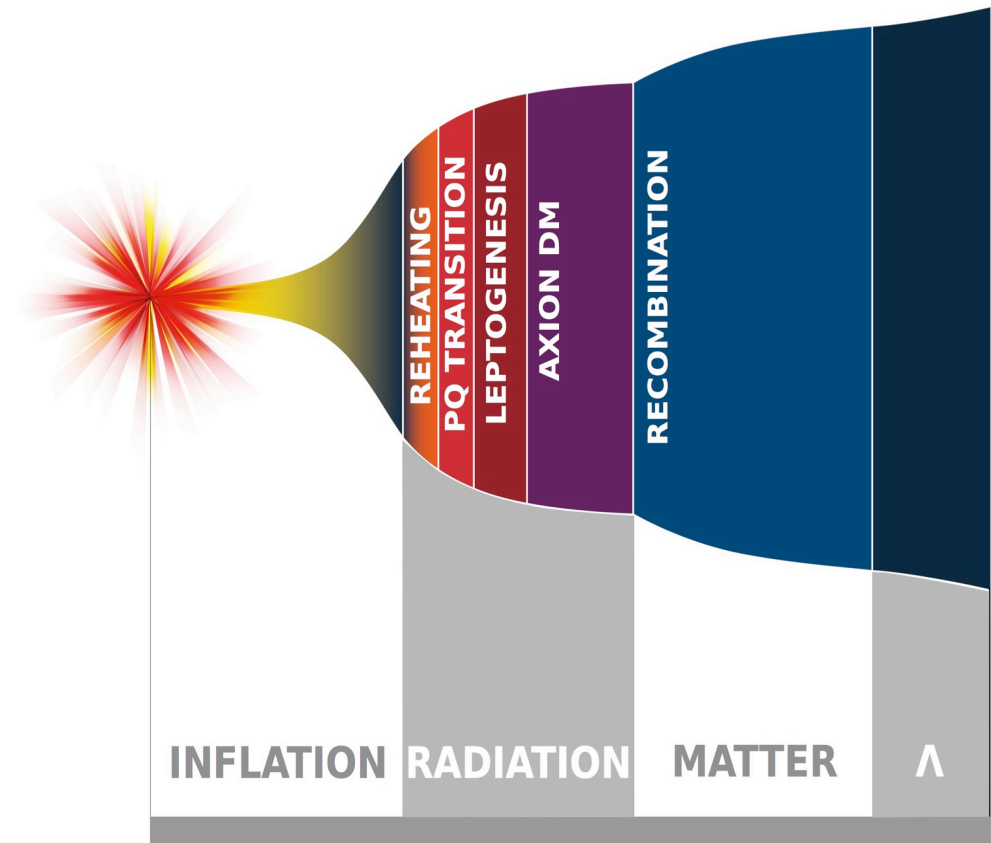
- no strong CP problem
- dark matter
- inflation
- neutrino masses and mixing
- baryogenesis via leptogenesis

[Dias et al. '14; Ballesteros et al. '16]



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- [Dias et al. '14; Ballesteros et al. '16]
- Complete and consistent history of the universe from inflation to now



[desy.de]

Unifying Inflation, Dark Matter, and Seesaw with PQ Field

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[Dias et al. '14; Ballesteros et al. '16]

- > Complete and consistent history of the universe from inflation to now

- > $SO(10)$ GUT SMASH?

[Ernst, AR, Tamarit '18 and in prep.]

- > Minimal $SO(10) \times U(1)_{PQ}$ models:

	16_F	126_H	10_H	210_H	45_H	S	10_F	N
Model 1	1	-2	-2	4	-	-	-	3
Model 2.1	1	-2	-2	0	4	-	-	3
Model 2.2	1	-2	-2	0	4	-	-2	1
Model 3.1	1	-2	-2	0	-	4	-	3
Model 3.2	1	-2	-2	0	-	4	-2	1

- > 16_F automatically features:

- Neutrino masses and mixing
- Baryogenesis via leptogenesis

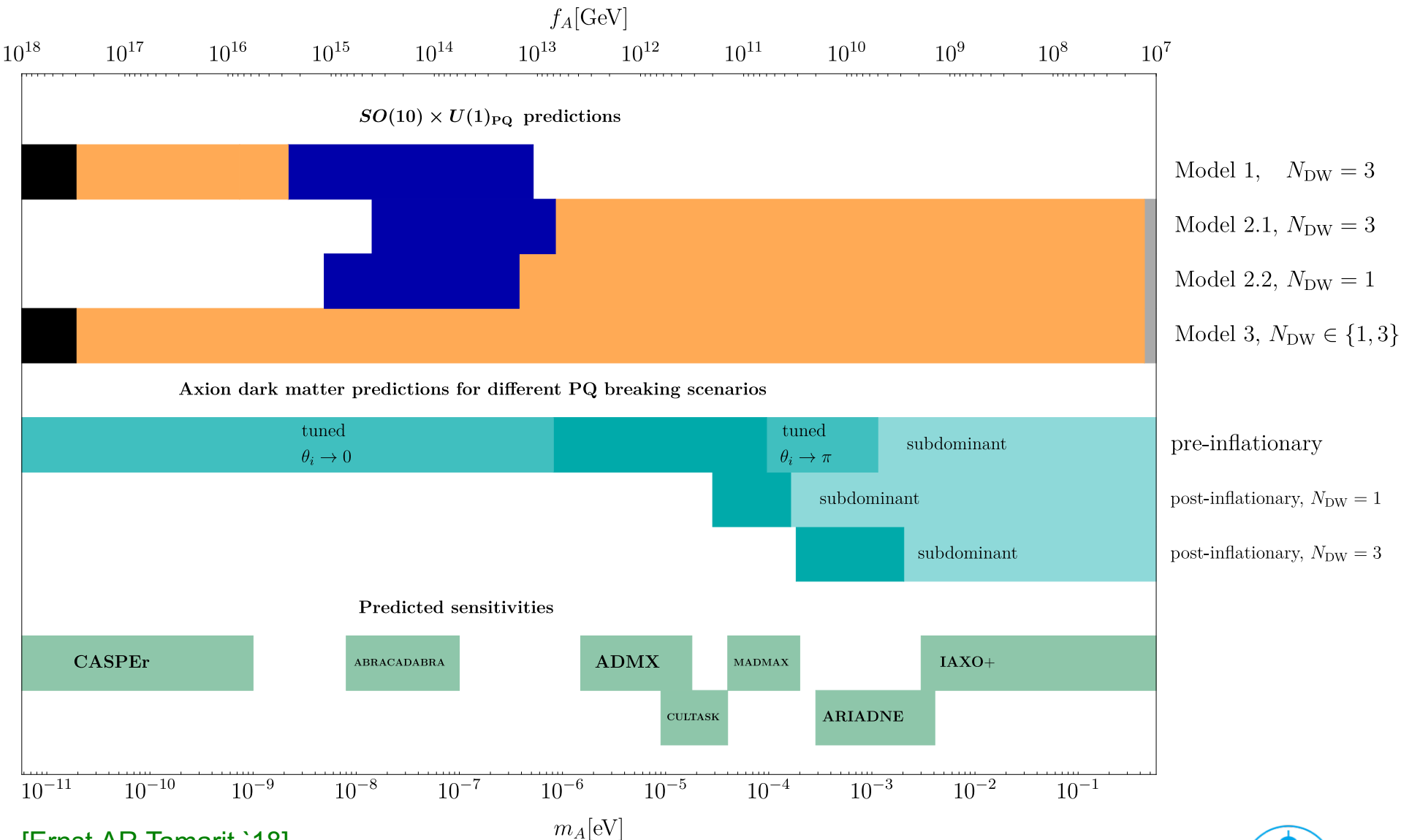
- > PQ extension adds

[Bajc et al. 06;
Altarelli, Meloni 13;
Babu, Khan 15]

- Predictivity of fermion masses/mixing
- Solution of strong CP problem
- Axion dark matter



Unifying Inflation, Dark Matter, and Seesaw with PQ Field

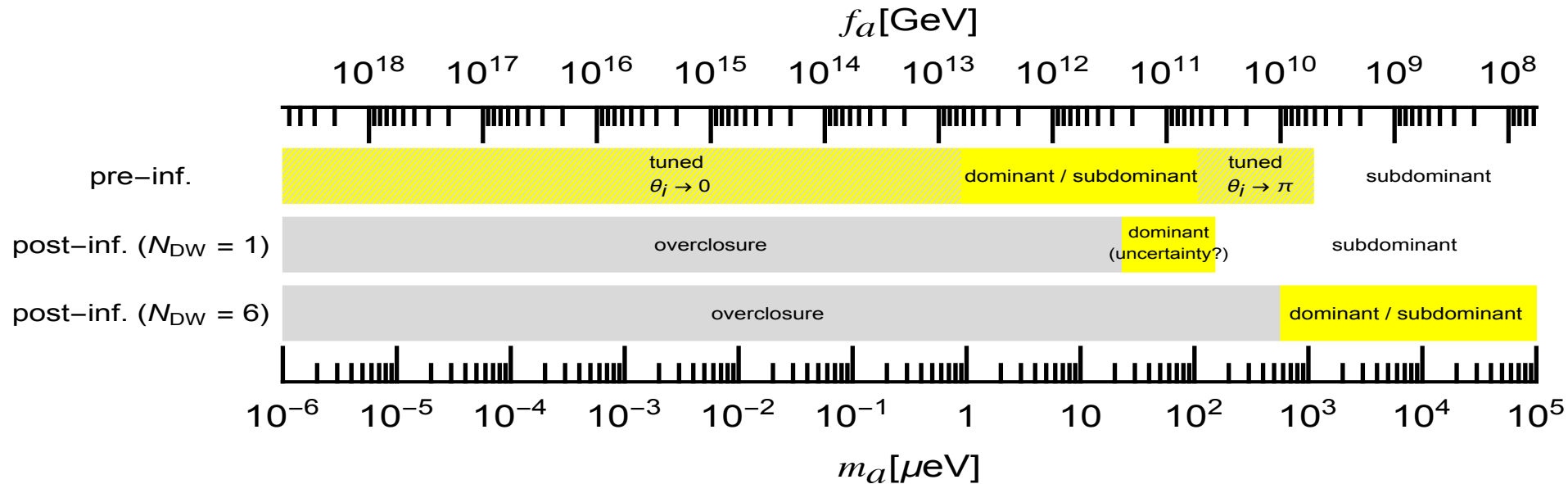


[Ernst, AR, Tamarit '18]



Summary

> Dark-matter axion mass spans a huge range:



Summary

➤ Strong motivation for current und upcoming axion DM experiments:

