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# Frequency doubler and two-color mode of operation at free electron laser FLASH2

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## ABSTRACT

We report on the results of the first operation of a frequency doubler at FLASH2. The scheme uses the feature of the variable gap of the undulator. Undulator is divided in two parts. The second part of the undulator is tuned to the double frequency of the first part. Modulated electron beam enters the second part of the undulator and generates radiation at the 2nd harmonic. Depending on a balance between the gain of undulator sections, frequency doubler allows operation in a two-color mode and operation at shorter wavelengths with respect to standard SASE scheme. The shortest wavelength of 3.1 nm (photon energy 400 eV) has been achieved at FLASH2 with frequency doubler scheme, which is significantly below the design value for the standard SASE option.

**Keywords:** X-ray, Free Electron Laser, undulator, harmonics, frequency doubling

## 1. INTRODUCTION

Free electron laser FLASH is equipped with two undulator beamlines.<sup>1-3</sup> Fixed gap undulator (period 2.73 cm, peak magnetic field 0.48 T, total magnetic length 27 m) is installed in the first beamline, FLASH1. Tuning of the radiation wavelength is performed by means of changing the energy of electrons. The second beam line, FLASH2, is equipped with variable gap undulator. Undulator of FLASH2 is a planar device with period of 3.14 cm. It consists of 12 modules of 2.5 m length (full magnetic length of the undulator is 30 m). Tunability of the radiation at fixed energy of electrons is performed by means of the opening undulator gap. Maximum wavelength is achieved at closed undulator gap (peak magnetic field is 0.96 T and rms value of the undulator parameter  $K = 2$ ). Minimum wavelength is defined by the quality of the electron beam (peak current, emittance and energy spread) allowing to reach saturation within available undulator length. With operating range of the electron beam energies of 0.4 - 1.25 GeV FLASH1 and FLASH2 beamlines cover wavelength range from 4 nm to 52 nm and from 3.5 nm to 90 nm, respectively. Scope of scientific applications will benefit a lot with extension of the operating range of the facility deep into the water window spanning from 4.38 nm and 2.34 nm (K-absorption edges of carbon and oxygen).

One of the techniques to generate shorter wavelength is the second harmonic afterburner.<sup>4-9</sup> Afterburner operating at the 2nd harmonic has been tested successfully at LCLS with 5 final undulator modules retuned to the 2nd harmonic.<sup>9</sup> With long, variable gap undulator of FLASH2 it is possible to implement frequency doubler scheme. Here we demonstrate experimental results from FLASH2: with an appropriate optimization of undulator tuning it becomes possible to operate facility at visibly shorter wavelengths and organize two color mode of operation ( $\omega + 2\omega$ ) with controllable radiation pulse intensities.

## 2. SASE FEL VERSUS FREQUENCY DOUBLER

An overview of the parameter space of SASE FEL can be easily performed with the fitting formulae for the minimum gain length  $L_g$  of the fundamental TEM<sub>00</sub> achieved at the optimum value of the focusing beta function  $\beta_{\text{opt}}$ .<sup>10-12</sup>

$$\begin{aligned} L_g &= 1.67 \left( \frac{I_A}{I} \right)^{1/2} \frac{(\epsilon_n \lambda_w)^{5/6}}{\lambda^{2/3}} \frac{(1 + K^2)^{1/3}}{K A_{JJ}} (1 + \delta), \\ \beta_{\text{opt}} &= 11.2 \left( \frac{I_A}{I} \right)^{1/2} \frac{\epsilon_n^{3/2} \lambda_w^{1/2}}{\lambda K A_{JJ}} (1 + 8\delta)^{-1/3}, \end{aligned} \quad (1)$$

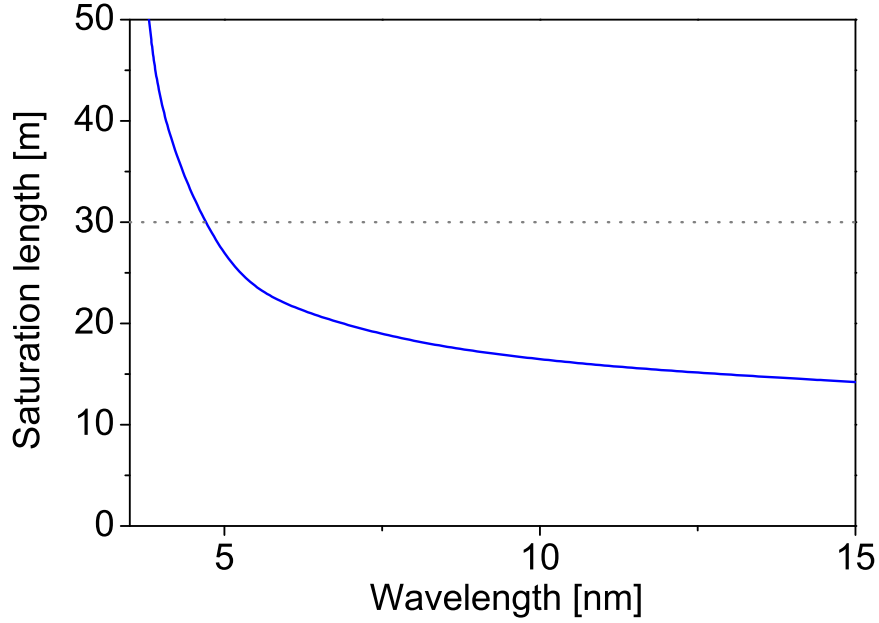


Figure 1. Saturation length of SASE FEL at FLASH2 versus radiation wavelength (blue curve). Gray dashed line at  $z = 30$  m shows magnetic length of FLASH2 undulator. Electron energy is 1080 MeV, beam current is 1500 A, normalized rms emittance  $1 \pi$  mm-mrad, and rms energy spread 0.15 MeV. Calculations have been performed with Eqs. (1) and (3).

with  $\delta$  given by

$$\delta = 131 \frac{I_A}{I} \frac{\epsilon_n^{5/4}}{\lambda^{1/8} \lambda_w^{9/8}} \frac{\sigma_\gamma^2}{(K A_{JJ})^2 (1 + K^2)^{1/8}} .$$

Realization of conditions (1) is referred as optimized FEL amplifier. Typically, for technical reasons, focusing beta function is  $\beta > \beta_{\text{opt}}$  (which is the case of FLASH parameter space), and Eq. (1) is corrected to:

$$L_g(\beta) \simeq L_g(\beta_{\text{opt}}) \left[ 1 + \frac{(\beta - \beta_{\text{opt}})^2 (1 + 8\delta)}{4\beta_{\text{opt}}^2} \right]^{1/6} \quad \text{for } \beta > \beta_{\text{opt}} . \quad (2)$$

Here  $\sigma_\gamma = \sigma_E/mc^2$ ,  $\sigma_E$  is rms energy spread,  $\epsilon_n = \gamma\epsilon$  is normalized rms emittance,  $K$  is rms undulator parameter,  $A_{JJ}$  is averaging factor,  $\lambda$  is radiation wavelength,  $\lambda_w$  is undulator period,  $I$  is the beam current,  $I_A \simeq 17$  kA is Alfven's current.

Saturation length of the optimized SASE FEL is given by:<sup>13</sup>

$$L_{\text{sat}} \simeq 0.6 L_g \ln \left( \frac{N_\lambda L_g}{\lambda_w} \right) . \quad (3)$$

Here  $N_\lambda = I\lambda/c$  is a number of electrons per wavelength. For operating VUV and X-ray SASE FELs one typically has  $L_{\text{sat}} \simeq (10 \pm 1) \times L_g$ .

With fixed energy of the electron beam, the radiation wavelength is reduced by means of opening the undulator gap. Radiation wavelength  $\lambda$  and rms value of the undulator parameter  $K$  are connected by resonance condition:

$$\lambda = \lambda_w \frac{1 + K^2}{2\gamma^2} .$$

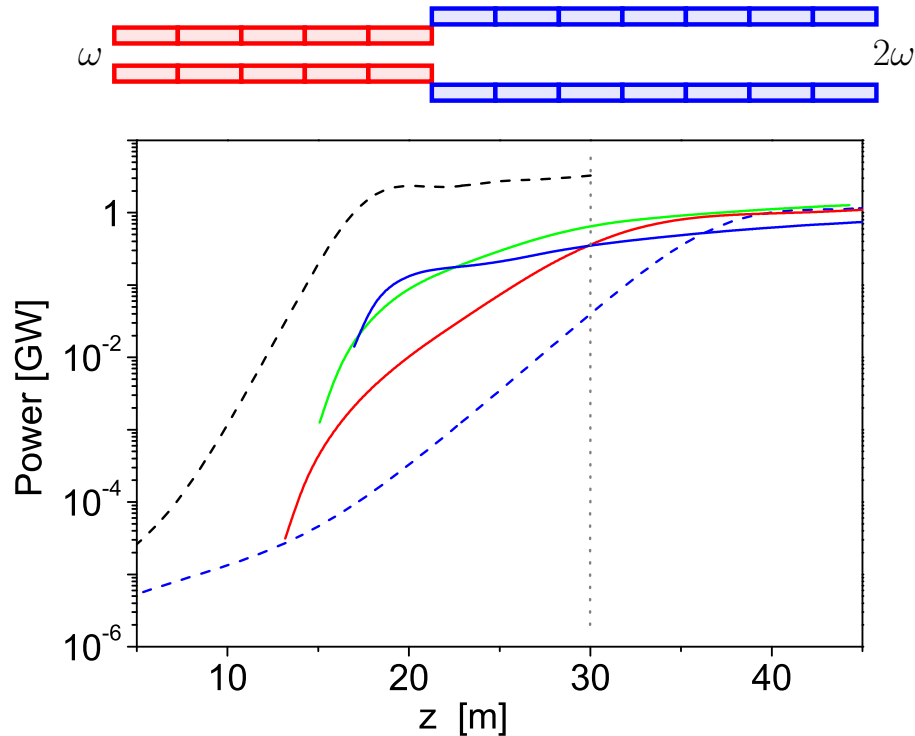


Figure 2. Top: Conceptual scheme of frequency doubler with variable gap undulator. Bottom: Evolution of the radiation power along FLASH2 undulator. Dashed curves correspond to SASE FEL operating at 8 nm (black) and at 4 nm (blue). Solid curves correspond to frequency doubler 8 nm  $\rightarrow$  4 nm for different lengths of the  $\omega$  undulator section (red, green, and blue colors). Gray dashed line at  $z = 30$  m shows magnetic length of FLASH2 undulator. Electron energy is 1080 MeV, beam current is 1500 A, normalized rms emittance  $1 \pi$  mm-mrad, and rms energy spread 0.15 MeV. Simulations have been performed with code FAST.<sup>15</sup>

Undulator field (and undulator parameter) is decreased with the radiation wavelength, and the FEL gain falls down. Saturation length gradually increases as well, and at some point it starts to exceed the undulator length. Then the radiation power quickly falls down with the wavelength (in fact, exponentially) when radiation wavelength approaches to  $\lambda = \lambda_w / (2\gamma^2)$  corresponding zero magnetic field of the undulator.

Let us consider the case of FLASH2 operating at the energy of electron of 108 MeV. Peak current 1500 A, normalized rms emittance  $1 \pi$  mm-mrad, and rms energy spread 0.15 MeV. Dependence of the saturation length on the radiation wavelength is shown in Fig. 1. Saturation length starts to increase rapidly for wavelengths shorter than 5 nm. Figure 2 shows evolution of the radiation power in SASE FEL for two radiation wavelengths, 8 nm and 4 nm. While saturation length at 8 nm is about 18 meters, it becomes to be 40 meters at 4 nm wavelength, and radiation power at the undulator length of 30 meters is two orders of magnitude below the saturation level, which is practically unusable for user applications. Practical limit for the smallest radiation wavelength at FLASH2 achievable with SASE configuration is about 4.5 nm at the energy of 1080 MeV, and 3.5 nm at the maximum energy of 1250 MeV. Thus, conventional SASE FEL allows to touch only the carbon K-edge which is rather unfavorable for the scientific case relying on the whole water window wavelength range.

For longer wavelengths, above 7 nm, saturation length is below 20 meters. This feature gives us an opportunity for application of a frequency doubling scheme. The simplest scheme of the frequency doubler can be organized with variable gap undulator having tunability wavelength range above two (see top plot in Fig. 2). Undulator is divided in two parts. First part of the undulator is tuned to the resonance frequency  $\omega$ , and the second part – to the frequency  $2\omega$ . Amplification process in the first part of the undulator starts from the shot noise in the electron beam, and produces radiation with central frequency  $\omega$ . The length of the first undulator part is such that the amplification process enters nonlinear stage when electron beam density modulation at higher frequency harmonics are induced in the electron beam due to nonlinear harmonic generation mechanism. Then modulated

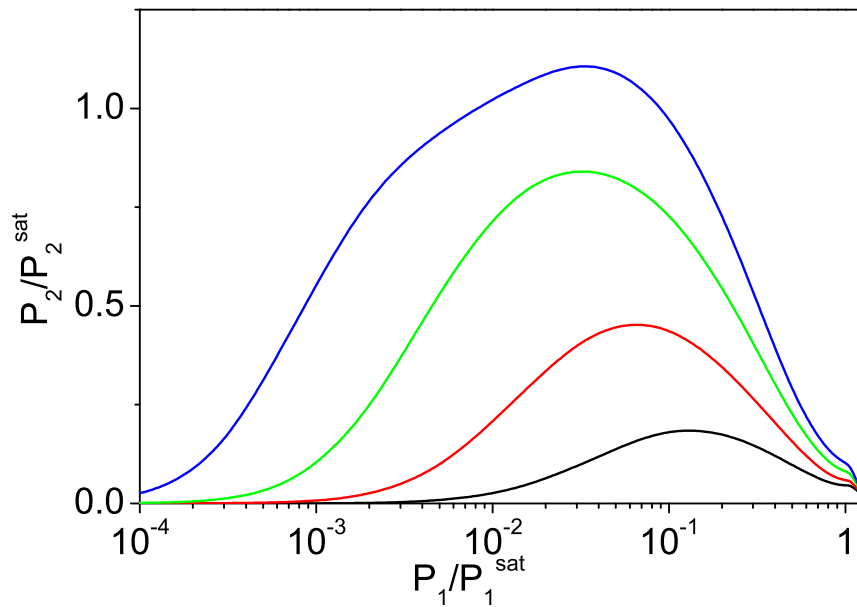


Figure 3. Parameter space of radiation powers from frequency doubler at FLASH2.  $P_1$  is radiation power at  $\omega$  at the exit of the first section, and  $P_2$  is radiation power at  $2\omega$  at the exit of the doubling section. Radiation powers are normalized to the corresponding saturation powers of SASE FEL. Different colors (black, red, green, and blue) denote lengths of the doubling section in terms of the saturation length:  $1/6 z_{\text{sat}}$ ,  $1/3 z_{\text{sat}}$ ,  $1/2 z_{\text{sat}}$ , and  $2/3 z_{\text{sat}}$ . Simulations have been performed with code FAST.<sup>15</sup>

electron beam enters the second part of the undulator tuned to frequency  $2\omega$ , and beam density modulations at frequency  $2\omega$  serve as a seed for amplification process. When these density modulations significantly exceed the shot noise modulations, the length required to reach saturation can be significantly shortened.

The plots in Figure 2 illustrate general features of the operation of the frequency doubling scheme. Black dashed line shows evolution of the radiation power along the  $\omega$ -section. Amplification process stops at some length, and then electron beam enters the  $2\omega$ -section tuned to the doubled frequency. Radiation with frequency  $\omega$  does not interact with the electron beam in the  $2\omega$ -section, and just propagates forward. Radiation power produced in the  $\omega$ -section at the second harmonic frequency is significantly suppressed,<sup>8,14</sup> thus only the beam density modulations at  $2\omega$  frequency can seed the  $2\omega$ -section.

We performed simulations of the amplification process for the full parameter space of the lengths of the  $\omega$  and  $2\omega$  sections, and results are summarized in Fig. 3. Despite simulations have been performed for specific set of electron beam parameters, presentation of the results in normalized form allows to use them in a wider range of parameters. X-coordinate is radiation power from the  $\omega$ -section normalized to the saturation power of SASE FEL operating at frequency  $\omega$ . Y-coordinate is radiation power from the  $2\omega$ -section normalized to the saturation power of SASE FEL operating at frequency  $2\omega$ .

Simple physical considerations explain operation of the frequency doubler scheme. Indeed, for a short length of the  $\omega$ -section, when amplification process stays in the exponential gain regime, the beam density modulations at frequency  $2\omega$  remain on the level of the shot noise. Induced modulations at the frequency  $2\omega$  (and higher harmonics) occur only in the nonlinear stage of amplification. At the onset of nonlinear regime induced modulation at  $2\omega$  is small, and gradually grows when amplification process approaches saturation. At the same time we occur degradation of the electron beam quality due to interaction of the electron beam with radiation. Thus, we deal with two competing effects: growth of the beam modulation, and degradation of the electron beam quality. Lucky factor for the frequency doubling scheme is that the growth of the beam density modulation advances the degradation of the beam quality, since the radiation power should be produced first, and only then it interacts with the beam. When amplification process enters the post-saturation regime, the quality of the beam and the

beam bunching degrade such that the electron beam becomes to be not appropriate for efficient radiation of higher harmonics.

We see that with two-color mode of operation it is possible to produce simultaneously high power at both frequencies with the level of about 20% of the saturation power. When amplification process in the first undulator terminates at the onset of the nonlinear regime, it is possible to generate high radiation power at the second harmonic in the second part of the undulator.

### 3. EXPERIMENTAL RESULTS

Operation of frequency doubler has been successfully demonstrated during several test runs at FLASH2 in 2016. Experiments have been performed at the energy of the electron beam of 1080 MeV and 1230 MeV. Tuning of the frequency doubler scheme with variable gap undulator does not require any special techniques. First, we tune maximum SASE gain in the uniform undulator tuned to frequency  $\omega$ . Analysis of the gain curve and fluctuations of the radiation pulse energy allows to determine optimum length of the  $\omega$ -section. Then we tune remaining sections to the frequency  $2\omega$ , and after adjustment of the phase shifters and electron beam orbit, the frequency doubler starts to generate radiation at the second harmonic.

Figure 4 shows the gain curve of the frequency doubler 8 nm  $\rightarrow$  4 nm. First five undulator sections were tuned to 8 nm, and the last seven sections were tuned to 4 nm. Red and blue curves show radiation pulse energies at 8 nm and 4 nm, respectively. Amplification process in the second part (frequency doubler) exhibits resonance behavior on the frequency detuning (right plot in Fig. 4). Even harmonics of the radiation are strongly suppressed in the planar undulator, thus we get confirmation that the beam bunching at the second harmonic seeds the amplification process of the frequency doubling sections. One can compare experimental and simulation results presented in Fig. 2. Gain curves look pretty much similar. As for the resonance curve, the width of experimental curve is significantly wider than FEL parameter  $\rho$ . This is typical signature of the energy chirp along the electron beam induced by the beam formation system.

Our experience shows that tuning procedure of the frequency doubler is simple and reproducible. In particular, it is possible to tune relative intensities of two colors ( $\omega$ ,  $2\omega$ ) in a wide limits. When tuned to equal intensities, pulse energies were in the range from a few to 10 microjoules. Radiation has been transported to FLASH2 experimental hall for characterization. Figure 5 shows an example for photon beam images and spectra of two color mode of operation for the frequency doubler 9 nm  $\rightarrow$  4.5 nm. With this first experience we can state that this operational mode can be proposed to users.

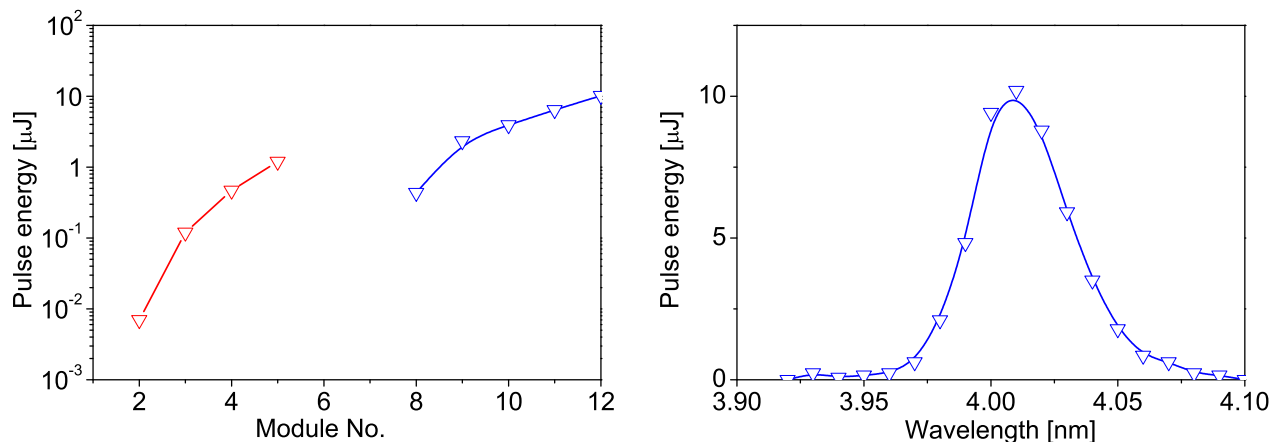


Figure 4. Frequency doubler at FLASH2 (experimental data). Left: gain curve of the frequency doubler at FLASH 2. First part of the undulator (5 modules) is tuned to 8 nm, and second part (7 modules) is tuned to 4 nm. Red and blue color correspond to the radiation wavelength of 8 nm and 4 nm, respectively. Right: radiation pulse energy at the 2nd harmonic versus resonance frequency of the second part of the undulator. Electron energy is 1080 MeV, bunch charge is 300 pC.

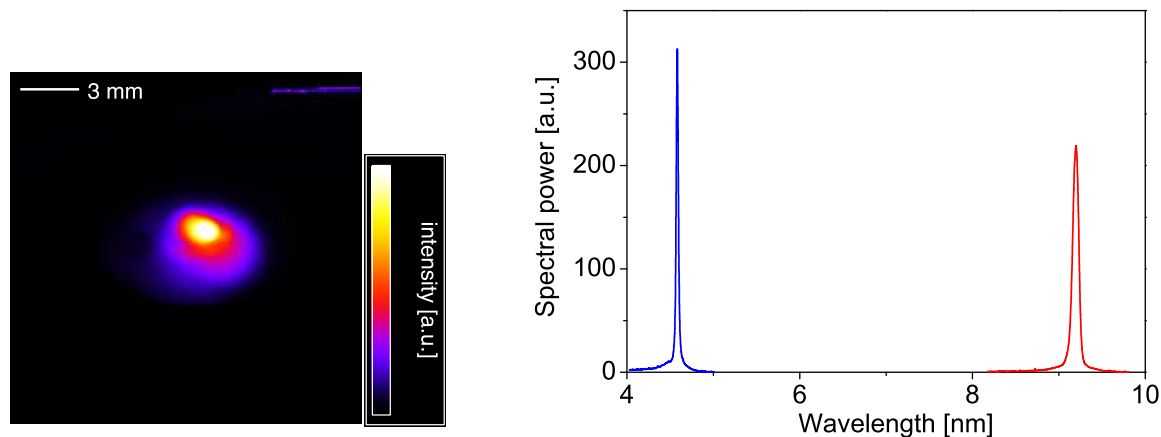


Figure 5. Photon beam image in the experimental hall(left) and radiation spectra (right) of the frequency doubler at FLASH2 (experimental data). Small yellow spot is 4.5 nm (2nd harmonic) radiation, pulse energy is 10 J. Larger blue/pink spot is 9 nm radiation, pulse energy is 10  $\mu$ J. Electron energy is 1080 MeV, bunch charge is 300 pC.

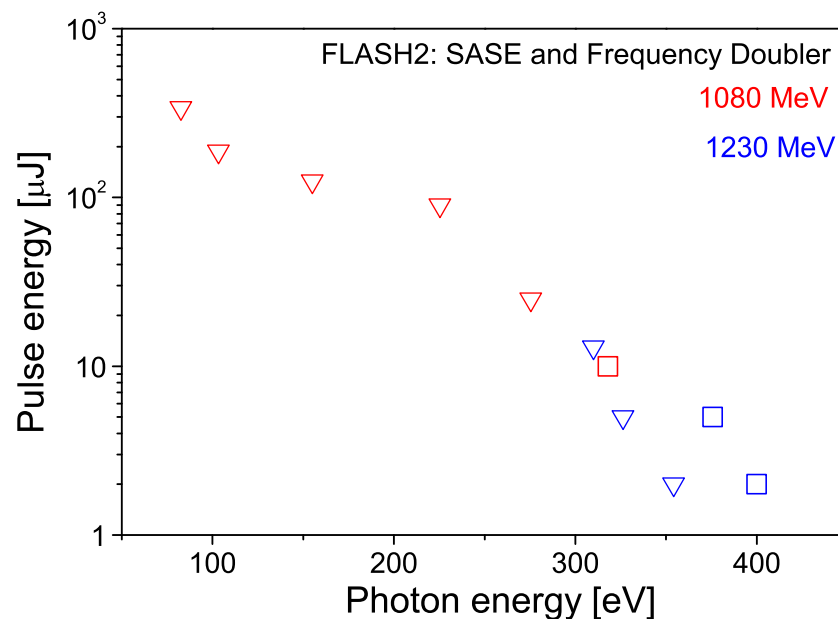


Figure 6. Radiation pulse energy versus photon energy. Experimental data from FLASH2 for electron energy of 1080 MeV and 1230 MeV, bunch charge is 300 pC. Triangles represent SASE, and squares represent frequency doubler.

With fixed electron energy, shortening of the radiation wavelength is achieved by opening the undulator gap. Undulator field is reduced, the FEL gain falls down, and starting from some wavelength the undulator length is not sufficient to reach saturation regime in the uniform undulator. Frequency doubler scheme is capable to generate shorter wavelength radiation than standard SASE FEL. Indeed, the first part of the undulator operates at twice longer wavelength, and saturation is obtained at a half of the full undulator length. Induced beam bunching at the second harmonic is much larger than the shot noise in the electron beam, and it becomes possible to reach saturation on much shorter length of the doubling section. For fair comparison of two options (SASE and frequency doubler) we performed two dedicated runs at FLASH2 operating with electron energies close to the limit. First, standard SASE mode has been optimally tuned at full undulator length, and radiation pulse energy scanned versus radiation wavelength. Then frequency doubler scheme has been tuned optimally with the same electron beam, and the pulse energy scanned. Results of pulse energy measurements are compiled in Fig. 6. We see that visibly shorter wavelengths were reached with the frequency doubler. In particular, photon

energies above Nitrogen K-edge have been demonstrated which significantly exceeds original specifications of FLASH2. The radiation has been successfully transported to FLASH2 experimental hall.

Analysis of experimental results obtained at FLASH2 shows that the frequency doubler scheme works in good agreement with theoretical expectations. Its tuning does not require special efforts, and it starts to work when SASE process tuned along the whole length of the undulator. Two color mode of operation is rather robust for longer wavelengths. Minimum achievable wavelength, strongly depends on the quality of the electron beam which requires dedicated tuning of the beam formation system. However, this is general problem for SASE FEL as well. Presently we reached photon energies just below Nitrogen K-edge (400 eV), and this seems to be close to the limit of the doubling frequency technique at the energy of electron beam of 1230 MeV. Additional installation of two undulator sections ( $2 \times 2.5$  m) may help in more reliable covering of the Nitrogen K-edge range and reaching shorter wavelengths.

#### 4. ACKNOWLEDGEMENTS

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