Observation of $\Xi_c(2930)^0$ and updated measurement of $B^- \rightarrow K^-\Lambda_c^+\bar{\Lambda}_c$ at Belle


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We report the first observation of the $\Xi_c(2930)^0$ charmed-strange baryon with a significance greater than 5σ. The $\Xi_c(2930)^0$ is found in its decay to $K^-\Lambda_c^+$ in $B^- \rightarrow K^-\Lambda_c^+\bar{\Lambda}_c$ decays. The measured mass and width are $[2928.9 \pm 3.0\text{(stat.)}^{+3.8}_{-2.9}\text{(syst.)}]$ MeV/$c^2$ and $[19.5 \pm 8.4\text{(stat.)}^{+8.4}_{-7.6}\text{(syst.)}]$ MeV, respectively, and the product branching fraction is $B(B^- \rightarrow \Xi_c(2930)^0\bar{\Lambda}_c)B(\Xi_c(2930)^0 \rightarrow K^-\Lambda_c^+)$ $= [1.73 \pm 0.45\text{(stat.)} \pm 0.24\text{(syst.)}] \times 10^{-4}$. We also measure $B(B^- \rightarrow K^-\Lambda_c^+\bar{\Lambda}_c) = [4.80 \pm 0.43\text{(stat.)} \pm 0.68\text{(syst.)}] \times 10^{-4}$ with improved precision, and search for the charmonium-like state $Y(4660)$ and its spin partner, $Y_0$, in the $\Lambda_c^+\bar{\Lambda}_c$ invariant mass spectrum. No clear signals of the $Y(4660)$ nor its spin partner are observed and the 90% credibility level
The singly charmed baryon is composed of a charm quark and two light quarks. Charmed baryon spectroscopy provides an excellent ground for studying the dynamics of light quarks in the environment of a heavy quark and offers an excellent laboratory for testing heavy-quark or chiral symmetry of the heavy or light quarks, respectively. Although many new excited states have previously been observed by BaBar, Belle, CLEO and LHCb in the past two decades [1], and many efforts have been made to identify the quantum numbers of these new states and understand their properties, we do not yet have a fully phenomenological model that describes the complicated physics of this sector [2, 3]. Identification and observation of new members in the charmed-baryon family will provide more information to address these open issues.

The Ξ_c(2930) charm-strange baryon has been reported only in the analysis of $B^- \rightarrow K^- \Lambda_c^+ \bar{A}_c$ by BaBar [4], where they claim a signal in the $K^- \Lambda_c^+$ invariant mass distribution with a mass of $[2931 \pm 3(\text{stat.}) \pm 5(\text{syst.})]$ MeV/c² and a width of $[36 \pm 7(\text{stat.}) \pm 11(\text{syst.})]$ MeV. However, despite the results of the fit to their spectrum nor the significance of the signal were given; the Particle Data Group (PDG) lists it as a “one star” state [1]. Despite the weak experimental evidence for the Ξ_c(2930) state, it has been taken into account in many theoretical models, including the chiral quark model [5], the light-cone Quantum Chromodynamics (QCD) sum rule [6], the $3P_0$ mode [7], the constituent quark model [8, 9], and the heavy-hadron chiral perturbation theory [10].

Belle has previously studied $B^- \rightarrow K^- \Lambda_c^+ \bar{A}_c$ decays [11] with a data sample of $386 \times 10^6 \bar{B}B$ pairs but the distributions of the intermediate $K\Lambda_c$ systems have not been presented. The full Belle data sample of $(772 \pm 11) \times 10^6 \bar{B}B$ pairs permits an improved study of $B^- \rightarrow K^- \Lambda_c^+ \bar{A}_c$ and a test for the existence of the Ξ_c(2930).

The same $B$ decay mode can be used to study the $\Lambda_c^+ \bar{A}_c$ invariant mass. In this system, Belle has previously observed a charmonium-like state, the $Y(4630)$, in the initial state radiation (ISR) process $e^+e^- \rightarrow \gamma_{\text{ISR}} \Lambda_c^+ \bar{A}_c$ [12] with a measured mass of $[4634^{+2}_{-4}(\text{stat.})^{+8}_{-8}(\text{syst.})]$ MeV/c² and a width of $[92^{+40}_{-24}(\text{stat.})^{+19}_{-19}(\text{syst.})]$ MeV. As this mass is very close to that of the $Y(4660)$ observed by Belle in the ISR process $e^+e^- \rightarrow \gamma_{\text{ISR}} \pi^+ \pi^- \psi$ [13, 14], many theoretical explanations assume they are the same state [15–17]. In Refs. [18, 19], where the $Y(4660)$ is modeled as an $f_0(980)\psi$ bound state, the authors predict that it should have a spin partner—a $f_0(980)\eta_c(2S)$ bound state denoted as the $Y_0^-$—with a mass and width of $(4613 \pm 4)$ MeV/c² and around 30 MeV, respectively, and a large partial width into $\Lambda^+_c \bar{A}_c^-$ [17, 19].

In this Letter, we perform an updated measurement of $B^- \rightarrow K^- \Lambda_c^+ \bar{A}_c$ [20] and observe the $\Xi_c(2930)^0$ signal with a significance of 5.1σ. This analysis is based on the full data sample collected at the $\Upsilon(4S)$ resonance by the Belle detector [21] at the KEKB asymmetric energy electron-positron collider [22]. Simulated signal events with $B$ meson decays are generated using EVTGEN [23], while the inclusive decays are generated via PYTHIA [24]. These events are processed by a detector simulation based on GEANT3 [25]. Inclusive Monte Carlo (MC) samples of $\Upsilon(4S) \rightarrow \bar{B}B$ ($B = B^+$ or $B^0$) and $e^+e^- \rightarrow q\bar{q} (q = u, d, s, c)$ events at $\sqrt{s} = 10.58$ GeV are used to check the backgrounds, corresponding to more than 5 times the integrated luminosity of the data.

We reconstruct the $\Lambda_c^+$ via the $\Lambda_c^+ \rightarrow pK^-\pi^+$, $pK_S^0$, $\Lambda^+\pi^-$, $pK_S^0\pi^+\pi^-$, and $\Lambda^+\pi^+\pi^-$ decay channels. When a $\Lambda_c^+$ and $\bar{A}_c$ are combined to reconstruct a $B$ candidate, at least one is required to have been reconstructed via the $pK^-\pi^-$ or $\bar{p}K^-\pi^+$ decay process. For charged tracks, information from different detector subsystems, including specific ionization in the central drift chamber, time measurements in the time-of-flight scintillation counters and the response of the aerogel threshold Cherenkov counters, is combined to form the likelihood $L_i$ for species $i$, where $i = \pi, K$, or $p$ [26]. Except for the charged tracks from $\Lambda \rightarrow p\pi^-$ and $K_S^0 \rightarrow \pi^+\pi^-$ decays, a track with a likelihood ratio $R_{K_p}^\tau = L_K/(L_K + L_\pi) > 0.6$ is identified as a kaon, while a track with $R_{K_p}^\pi < 0.4$ is treated as a pion [26]. With this selection, the kaon (pion) identification efficiency is about 94% (98%), while 5% (2%) of the kaons (pions) are misidentified as pions (kaons). A track with $R_{p/\bar{p}}^\tau = L_{p/\bar{p}}/(L_{p/\bar{p}} + L_\pi) > 0.6$ and $R_{p/\bar{p}}^\pi = L_{p/\bar{p}}/(L_{p/\bar{p}} + L_K) > 0.6$ is identified as a proton/anti-proton with an efficiency of about 98%; fewer than 1% of the pions/kaons are misidentified as protons/anti-protons.

The $K_S^0$ candidates are reconstructed from pairs of oppositely-charged tracks, treated as pions, and identified by a multivariate analysis with a neural network [27] based on two sets of input variables [28]. Candidate $\Lambda$ baryons are reconstructed in the decay $\Lambda \rightarrow p\pi^-\pi^-$ and selected if the $p\pi^-$ invariant mass is within 5 MeV/c² ($5\sigma$) of the $\Lambda$ nominal mass [1].

We perform a vertex fit to signal $B$ candidates. If
there is more than one B signal candidate in an event, we select the one with the minimum $\chi^2_{\text{vertex}}$ from the vertex fit. We require $\chi^2_{\text{vertex}} < 50$ with a selection efficiency above 96%. As the continuum background level is very low, continuum suppression is not necessary.

The B candidates are identified using the beam-energy constrained mass $M_{bc}$ and the mass difference $\Delta M_B$. The beam-energy constrained mass is defined as $M_{bc} \equiv \sqrt{E_{\text{beam}}^2/c^2 - (\sum \vec{p}_i)^2/c^2}$, where $E_{\text{beam}}$ is the beam energy and $\vec{p}_i$ are the three-momenta of the B-meson decay products, all defined in the center-of-mass system (CMS) of the $e^+e^-$ collision. The mass difference is defined as $\Delta M_B \equiv M_B - m_B$, where $M_B$ is the invariant mass of the B candidate and $m_B$ is the nominal B-meson mass [1].

Figure 1 shows clear evidence of $\Lambda_c^+$ and $\Lambda_c^-$ in the distribution of $M_{\Lambda_c^-}$ versus $M_{\Lambda_c^+}$ (left panel) from the selected $B^\rightarrow K^-\Lambda_c^-\Lambda_c^+$ data candidates in the $B$ signal region of $|\Delta M_B| < 0.018 \text{GeV/c}^2$ and $M_{bc} > 5.27 \text{GeV/c}^2$ ($\sim 3\sigma$), illustrated by the green box in the right panel’s distribution of $\Delta M_B$ versus $M_{bc}$. The $\Lambda_c$ signal region (the central green box in the left panel) is defined as $|M_{\Lambda_c} - m_{\Lambda_c}| < 10 \text{ MeV/c}^2$ ($\sim 2.5\sigma$), where $m_{\Lambda_c}$ is the nominal mass of the $\Lambda_c$ baryon [1]. As the mass resolution of $\Lambda_c$ candidates is almost independent of the $\Lambda_c$ decay mode, according to the signal MC simulation, the same requirement is placed on all $\Lambda_c$ decay modes. The non-$\Lambda_c$ background in the $\Lambda_c$ signal region is estimated as half of the total number of events in the four red sideband regions minus one quarter of the total number of events in the four blue sideband regions of the left panel.

![Figure 1](image1.png)

FIG. 1: Signal-enhanced distribution of $M(\Lambda_c^-)$ versus $M(\Lambda_c^+)$ (left panel) and of $\Delta M_B$ versus $M_{bc}$ (right panel) from the selected $B^\rightarrow K^-\Lambda_c^-\Lambda_c^+$ candidates, summing over all five reconstructed $\Lambda_c$ decay modes. Each panel shows the events falling in the solid green signal region of the other panel. The dashed red and blue boxes in the left panel show the $\Lambda_c$ sideband regions used for the estimation of the non-$\Lambda_c$ background (see text).

To obtain the $B^\rightarrow K^-\Lambda_c^-\Lambda_c^+$ signal yields, we perform an unbinned two-dimensional (2D) simultaneous extended maximum likelihood fit to the $\Delta M_B$ versus $M_{bc}$ distributions for the five reconstructed $\Lambda_c$ decay modes. The model used to fit the $M_{bc}$ distribution is a Gaussian function for the signal shape plus an ARGUS function [29] for the background. The model for the $\Delta M_B$ distribution is the sum of a Gaussian function for the signal plus a first-order polynomial for the background. The Gaussian resolutions are fixed to the values from the fits to the individual MC distributions, and the relative signal yields among the five final states is fixed according to the relative branching fraction between the final states and the detection acceptance and efficiency of the intermediate states.

Figure 2 shows the projections of the fit superimposed on the $\Lambda_c$-signal-enhanced $M_{bc}$ and $\Delta M_B$ distributions, summing over all five reconstructed $\Lambda_c$ decay modes. We observe $153 \pm 14$ signal events with a signal significance above $10\sigma$, and extract the branching fraction of $B(B^\rightarrow K^-\Lambda_c^+\Lambda_c^-) = \left[4.80 \pm 0.43 \text{(stat.)}\right] \times 10^{-4}$.

![Figure 2](image2.png)

FIG. 2: The $\Lambda_c$-signal-enhanced distributions of (a) $M_{bc}$ in the $\Delta M_B$ signal region and (b) $\Delta M_B$ in the $M_{bc}$ signal region for $B^\rightarrow K^-\Lambda_c^+\Lambda_c^-$, combining five exclusive final states. The dots with error bars are data, the solid blue curves are the best-fit projections to the distributions and the dashed magenta lines are the fitted backgrounds.

The Dalitz distribution of the reconstructed $B^\rightarrow K^-\Lambda_c^+\Lambda_c^-$ candidates is shown in Fig. 3. A vertical band enhancement near $M(K^-\Lambda_c^+) \sim 2.93 \text{ GeV/c}^2$ is observed; no signal band is apparent in the $M(\Lambda_c^+\Lambda_c^-)$ horizontal direction nor in the $M(K^-\Lambda_c^-)$ diagonal direction.

![Figure 3](image3.png)

FIG. 3: Dalitz distribution of reconstructed $B^\rightarrow K^-\Lambda_c^+\Lambda_c^-$ candidates in the B signal region. The black dots are data; the shaded region is the MC simulated phase-space distribution.
in Fig. 4. The shaded histogram is from the normalized \( \Lambda^+ \) and \( \bar{\Lambda}^- \) mass sidebands, and the dot-dashed line is the sum of the contributions from normalized \( e^+e^- \rightarrow q\bar{q} \) and \( \Upsilon(4S) \rightarrow BB \) generic MC samples. Since they are consistent, we take the \( \Lambda^+ \) and \( \bar{\Lambda}^- \) mass sidebands to represent the total background, neglecting the small possible contribution of background with real \( \Lambda^+ \) and \( \bar{\Lambda}^- \). A clear \( \Xi_c(2930) \) signal is observed. No structure is seen in the \( \Lambda^+ \) and \( \bar{\Lambda}^- \) mass sidebands, nor in the generic MC samples, nor in the wrong-sign-combination distribution of \( K^- \Lambda^- \).

An unbinned simultaneous extended maximum likelihood fit is performed to the \( K^- \Lambda^+ \) invariant mass spectra for selected \( B^- \) and \( \Lambda^- \)-signal events and the \( \Lambda^+ \) and \( \bar{\Lambda}^- \) mass sidebands. An S-wave Breit-Wigner (BW) function convolved with a Gaussian function with the phase space factor and efficiency curve included (the mass resolution of Gaussian function being fixed to 4.46 MeV/c\(^2\) from the signal MC simulation) is taken as the \( \Xi_c(2930) \) signal shape. Direct three-body \( B \) decays are modeled by the shape corresponding to \( B^- \rightarrow K^-\Lambda^+\bar{\Lambda}^- \) MC-simulated decays distributed uniformly in phase space. A second-order polynomial is used to represent the \( \Lambda^+ \) and \( \bar{\Lambda}^- \) mass-sideband distribution, which is normalized to represent the total background in the signal events in the fit.

The fit results are shown in Fig. 4. The fitted mass and width of the \( \Xi_c(2930) \) are \( M_{\Xi_c(2930)} = [2928.9 \pm 3.0 \text{ (stat.)}] \) MeV/c\(^2\) and \( \Gamma_{\Xi_c(2930)} = [19.5 \pm 8.4 \text{ (stat.)}] \) MeV, where a fit bias of 1.37 MeV/c\(^2\) on the \( \Xi_c(2930) \) mass, determined using MC simulation, has been corrected. The yields of the \( \Xi_c(2930) \) signal and the phase-space contribution are \( N_{\Xi_c} = 61 \pm 16 \) and \( N_{\text{phsp}} = 70 \pm 19 \). The statistical significance of the \( \Xi_c(2930) \) signal is 5.1\( \sigma \), calculated from the difference of the logarithmic likelihoods [30], \(-2\ln(L_0/L_{\text{max}}) = 32.8\), where \( L_0 \) and \( L_{\text{max}} \) are the maximized likelihoods without and with a signal component, respectively, taking into account the difference in the number of degrees of freedom (\( \Delta\text{ndf} = 3 \)). The signal significance remains at 5.1\( \sigma \) when convolving the likelihood with a Gaussian function whose width equals the total systematic uncertainty. Alternative fits to the \( K^-\Lambda^+ \) mass spectra are performed: (a) using a first-order or third-order polynomial as background shape; (b) changing the \( \Xi_c(2930) \) mass resolution by 10\%; and (c) using an energy-dependent BW function as the \( \Xi_c(2930) \) signal shape. The \( \Xi_c(2930) \) signal significance is larger than 5.1\( \sigma \) in all cases.

The product branching fraction of \( B(B^- \rightarrow \Xi_c(2930)\Lambda^-) \) \( \mathcal{B}(\Xi_c(2930) \rightarrow K^-\Lambda^+) = [1.73 \pm 0.45 \text{ (stat.)}] \times 10^{-4} \) is calculated as \( N_{\Xi_c}^{\text{total}}/\left[2N_B\varepsilon_{\Xi_c}^B\mathcal{B}(\Lambda^+ \rightarrow pK^-\pi^+)\right]^2 \), where \( N_{\Xi_c}^{\text{total}} \) is the fitted \( \Xi_c(2930) \) signal yield; \( N_B = N_{Y(4S)}\mathcal{B}(Y(4S) \rightarrow B^+B^-) \) \( N_{Y(4S)} \) is the number of accumulated \( Y(4S) \) events and \( \mathcal{B}(Y(4S) \rightarrow B^+B^-) = 0.514 \pm 0.006 \) [1]; \( \mathcal{B}(\Lambda^+ \rightarrow pK^-\pi^+) = (6.35 \pm 0.33\%) \) is the world-average branching fraction for \( \Lambda^+ \rightarrow pK^-\pi^+ \) [1]; \( \varepsilon_{\Xi_c}^B = \sum_i \varepsilon_i^B \times \Gamma_i/\Gamma(pK^-\pi^+) \) \( (i \text{ is the } \Lambda_i \text{ decay mode index})\), \( \varepsilon_i^B \) is the detection efficiency from MC simulation and \( \Gamma_i \) is the partial decay width of \( \Lambda^+_c \rightarrow pK^-\pi^+, \Lambda^+_S \), \( \Lambda^- \), \( \Lambda^+_S \pi^+\pi^- \), and \( \Lambda^+\pi^-\pi^- \) [1]. Here, \( \mathcal{B}(\Lambda^+_S \rightarrow \pi^+\pi^-) \text{ or } \mathcal{B}(\Lambda \rightarrow p\pi^-) \) is included in \( \Gamma_i \) for the final states with a \( K_S^0 \) or a \( \Lambda \).

The \( M_{\Lambda^+_c\bar{\Lambda}^-} \) spectrum is shown in Fig. 5, in which no clear \( Y_1 \) or \( Y(4660) \) signals is evident. An unbinned extended maximum likelihood fit is applied to the \( \Lambda^+_c\bar{\Lambda}^- \) mass spectrum to extract the signal yields of the \( Y_1 \) and \( Y(4660) \) in \( B \) decays. In the fit, the signal shape of the \( Y_1 \) or \( Y(4660) \) is obtained from MC simulation directly with the input parameters \( M_{Y_1} = 4166 \text{ MeV}/c^2 \) and \( \Gamma_{Y_1} = 30 \text{ MeV for } Y_1 \) [17], and \( M_{Y(4660)} = 4643 \text{ MeV}/c^2 \) and \( \Gamma_{Y(4660)} = 72 \text{ MeV for } Y(4660) \) [1]; a third-order polynomial is used to describe all other contributions. The fit results are shown in Figs. 5(a) and (b) for the \( Y_1 \) and \( Y(4660) \), respectively. From the fits, we have \( 10 \pm 23 \) \( Y_1 \) signal events with a statistical signal significance of 0.7\( \sigma \), and \( 10 \pm 26 \) \( Y(4660) \) signal events.

As the statistical signal significance of each \( Y \) state is less than 3\( \sigma \), 90\% C.L. Bayesian upper limits on \( \mathcal{B}(B^- \rightarrow K^-Y)\mathcal{B}(Y \rightarrow \Lambda^+_c\bar{\Lambda}^-) \) are determined to be \( 1.2 \times 10^{-4} \) and \( 2.0 \times 10^{-4} \) for \( Y_1 \) and \( Y(4660) \), respectively, by solving the equation \( \int_{0}^{\infty} \mathcal{L}(B)dB/ \int_{0}^{\infty} \mathcal{L}(B)dB = 0.9 \), where \( B = n_Y/\left[\varepsilon_{\Xi_c}^B N_{BB}\mathcal{B}(\Lambda^+_c \rightarrow pK^-\pi^+)\right]^2 \) is the assumed product branching fraction; \( \mathcal{L}(B) \) is the corresponding maximized likelihood of the data; \( n_Y \) is the number of \( Y \) signal events; and \( \varepsilon_{\Xi_c}^B = \sum_i \varepsilon_i^B \times \Gamma_i/\Gamma(pK^-\pi^+) \) being the total efficiency from MC simulation for mode \( i \). To take the systematic uncertainty into account, the above likelihood is convolved with a Gaussian function whose width equals the total
systematic uncertainty.

There are several sources of systematic uncertainties in the branching fraction measurements. The detection efficiency relevant (DER) errors include those for tracking efficiency (0.35%/track), particle identification efficiency (1.9%/kaon, 0.9%/pion, 2.4%/proton and 2.0%/anti-proton), as well as Λ (3.0%) and Ξ (1.7%) selection efficiencies. Assuming all the above systematic error sources are independent, we add them in quadrature for each decay mode, yielding 6.0–8.3%, depending on the mode. For the four branching fraction measurements, the final DER errors are summed in quadrature over the five Λ decay modes using weight factors equal to the product of the total efficiency and the Λc partial decay width. We estimate the systematic errors associated with the fitting procedure by changing the order of the background polynomial, the range of the fit, and the values of the masses and widths of the Y and Υ (4660) by ±1σ, and by enlarging the mass resolution by 10%; the deviations from nominal in the fitted results are taken as systematic errors. Uncertainties for $B(Λ_c^+ \to pK^-\pi^+)$ and $Γ_c/Γ(pK^-\pi^+)$ are taken from Ref. [1]. The final errors on the $Λ_c$ partial decay widths are summed in quadrature over the five modes with the detection efficiency as a weighting factor. The world average of $B(Υ(4S) \to B^+B^-)$ is $51.4\pm 0.6$% [1], which corresponds to a systematic uncertainty of 1.2%. The systematic uncertainty on $N_{Υ(4S)}$ is 1.37%. Assuming all sources listed in Table I to be independent, the total systematic uncertainties on the branching fraction measurements are summed in quadrature.

The following systematic uncertainties are considered for the Ξc mass and width. Half of the correction due to the fitting bias on the Ξc mass is taken conservatively as a systematic error. By enlarging the mass resolution by 10%, the difference in the measured Ξc width is 0.7 MeV, which is taken as a systematic error. By changing the background shape, the differences of 0.3 MeV/c² and 0.9 MeV in the measured Ξc mass and width, respectively, are taken as systematic uncertainties.

### Table I: Relative systematic uncertainties (%) in the branching fraction measurements.

<table>
<thead>
<tr>
<th>Branching fraction</th>
<th>DER</th>
<th>Fit</th>
<th>$Λ_c$ decays</th>
<th>$N_{Β+}$</th>
<th>Sum</th>
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<td>13.9</td>
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<td>8.65</td>
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<td>1.82</td>
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<tr>
<td>$B_4$</td>
<td>4.53</td>
<td>23.1</td>
<td>12.8</td>
<td>1.82</td>
<td>26.8</td>
</tr>
</tbody>
</table>

The signal-parametrization systematic uncertainty is estimated by replacing the constant total width of a mass-dependent width of $Γ_1 = Γ^0_1 × \Phi(M_{K^-Λ_c^+})/\Phi(M_{Ξ_c(2930)})$, where $Γ^0_1$ is the width of the resonance, $\Phi(M_{K^-Λ_c^+}) = P/M_{K^-Λ_c^+}$ is the phase space factor for an S-wave two-body system (P is the K–momentum in the K–Λ_c^+ CMS) and $M_{Ξ_c(2930)}$ is the K–Λ_c^+ invariant mass fixed at the Ξc(2930) nominal mass. The differences in the measured Ξc(2930) mass and width are 0.2 MeV/c² and 5.3 MeV, respectively, which are taken as the systematic errors. An alternative fit to the $M(Κ^-Λ_c^+)$ spectra with interference between the Ξc(2930) and the phase-space contribution included gives $M_{Ξ_c(2930)} = (2917.0 ± 5.5)$ MeV/c² and $Γ_{Ξ_c(2930)} = (13.8 ± 6.9)$ MeV; the differences of $–11.9$ MeV/c² and $–5.7$ MeV from the nominal mass and width, respectively, are taken as systematic errors. Assuming all the sources are independent, we add them in quadrature to obtain the total systematic uncertainties on the Ξc(2930) mass and width of $–12.0$ MeV/c² and $–7.9$ MeV, respectively.

In summary, using $(772 ± 11) \times 10^6 BB$ pairs, we perform an updated analysis of $B^– \to K^-Λ_c^+Λ_c^-$. In the $K^-Λ_c^+$ mass spectrum, the charmed baryon state Ξc(2930)⁰ is clearly observed for the first time with a statistical significance greater than 5σ. The measured mass and width are $M_{Ξ_c(2930)} = (2928.9 ± 3.0^{+0.8}_{–12.0})$ MeV/c² and $Γ_{Ξ_c(2930)} = (19.5 ± 8.4^{+5.5}_{–7.9})$ MeV. The branching fraction is $B(B^– \to K^-Λ_c^+Λ_c^-) = (4.80 ± 0.43 ± 0.68) \times 10^{-4}$, which is consistent with the world average value of $(6.9 ± 2.2) × 10^{-4}$ [1] but with much-improved precision. We measure the product branching fraction $B(B^– \to Ξ_c(2930)Λ_c^-)B(Ξ_c(2930) \to K^-Λ_c^-)$ = $(1.73 ± 0.45 ± 0.24) \times 10^{-4}$, where the first error is statistical and the second systematic. There are no significant signals seen in the $Λ_c^+Λ_c^-$ mass spectrum. We place 90% C.L. upper limits for the Y(4660) and its theoretically predicted spin partner $Y_0$ of $B(B^– \to K^-Y(4660))B(Y(4660) \to Λ_c^+Λ_c^-) < 1.2 × 10^{-4}$ and $B(B^– \to K^-Y_0)B(Y_0 \to Λ_c^+Λ_c^-) < 2.0 × 10^{-4}$ [31].

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[20] Inclusion of charge conjugate states is implicit unless otherwise stated.
[31] Considering the possible change of $B(\Lambda_c^+ \to pK^-\pi^+)$ in the future, we provide $B(B^- \to K^- \Lambda_c^+ \Lambda_c^-)B(\Lambda_c^+ \to pK^-\pi^+)^2 = (1.94 \pm 0.13 \pm 0.18) \times 10^{-6}$ and $B(B^- \to \Xi_c(2930)\Lambda_c^-)B(\Xi_c(2930) \to K^- \Lambda_c^+)B(\Lambda_c^+ \to pK^-\pi^+)^2 = (6.97 \pm 1.81 \pm 0.64) \times 10^{-7}$, where the first errors are statistical and the second systematic with the uncertainty on $B(\Lambda_c^+ \to pK^-\pi^+)$ omitted. The 90% C.L. upper limits on $B(B^- \to K^- Y(4660))B(Y(4660) \to \Lambda_c^+ \Lambda_c^-)B(\Lambda_c^+ \to pK^-\pi^+)^2$ and $B(B^- \to K^- Y_0)B(Y_0 \to \Lambda_c^+ \Lambda_c^-)B(\Lambda_c^+ \to pK^-\pi^+)^2$ are $4.8 \times 10^{-7}$ and $8.0 \times 10^{-7}$, respectively.