

Broadband terahertz generation with a stair-step echelon

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Abstract: A method to overcome limitations of conventional broadband terahertz generation techniques is presented. A stair-step echelon allows for the creation of superior tilted-pulse-fronts to yield larger frequencies and bandwidths, energy conversion efficiencies exceeding 5%.

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1. Introduction

Ultrafast terahertz (THz) transients are of great interest for linear and nonlinear THz spectroscopy and compact particle acceleration [1]. Tilted-pulse-front (TPF) [2] THz generation in lithium niobate (LN) has become ubiquitous due to its compatibility with easily accessible high-power 800 nm/1 μ m laser technology. However, the use of diffraction grating (DG) based TPFs (DG-TPFs) present formidable imaging challenges, particularly for large pump bandwidths and beam sizes. Furthermore, as the optical pump spectrum red-shifts and broadens due to repeated down-conversion to THz frequencies (referred to as cascading), a spatio-temporal break-up of the pump pulse results [3]. This limits the energy conversion efficiency and produces THz pulses with spatio-temporal distortions. In order to address this issue, terahertz generation in organic materials and semiconductors have been proposed and demonstrated [4-5]. However, they rely on the development of pump lasers at wavelengths $> 1.3 \mu$ m, which maybe more challenging in relation to well-developed 1 μ m laser technology. We describe new mechanisms to eliminate these limitations by engineering the transverse-momentum distribution of the pump pulse via a stair-step echelon [6] to produce TPFs which do not contain the same distortions as DG-TPFs. This constitutes one of two approaches [7] presented by us towards scaling single-cycle pulse energies from the current mJ level to the 10 mJ level.

2. Results

The structure (Fig.1a) produces a discrete optical pulse front composed of many beamlets. Each beamlet propagating in LN at the optical group velocity c/n_g produces a THz wavelet that propagates at phase velocity c/n_{THz} in a direction given by the Cherenkov angle $\gamma = \cos^{-1}(n_{THz}/n_g)$ where $\gamma = 63^\circ$ for LN. The THz wavelets superpose constructively if the time delay Δt and spatial offset Δx between successive optical beamlets are related by $\Delta t = \Delta x \tan \gamma n_g / c$ (Fig.1c), yielding a single-cycle THz plane wave with uniform spatial properties (Fig.1d).

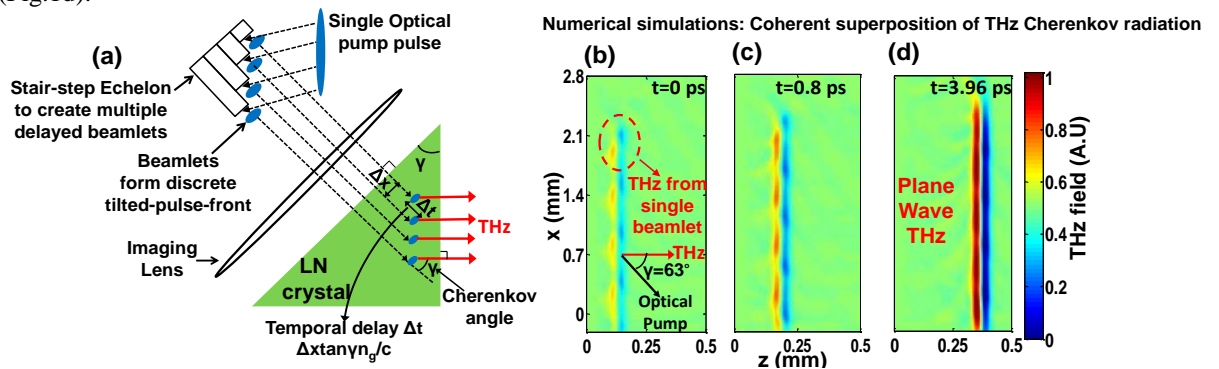


Fig.1(a) Schematic illustration of echelon depicting THz generation by relatively delayed optical beamlets. (b)-(d) Numerical simulations of echelon : Individual THz wavelets superpose coherently when the temporal delay Δt and spatial offset Δx between beamlets are related by $\Delta t = \Delta x \tan \gamma n_g / c$, where $\gamma = 63^\circ$ is the Cherenkov angle for LN

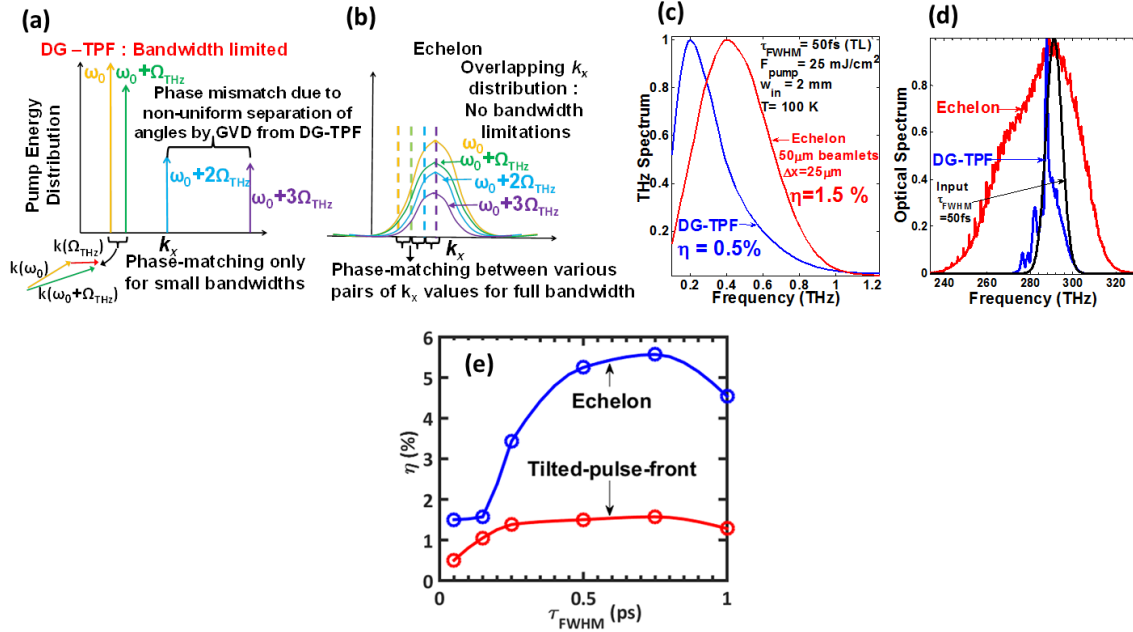


Fig.2.(a) Inherent bandwidth limitation due to group velocity dispersion(GVD) produced by DG-TPFs. Phase-matching conditions are satisfied only for small bandwidths. (b) Echelons contain overlapping transverse momentum (k_x) distributions, which solves the bandwidth limitation. (c) Higher conversion efficiencies, broader THz spectra predicted for echelons. (d) Spectral broadening due to cascading in echelons is larger than with DG-TPFs.(e) Echelons produce conversion efficiencies much larger than conventional TPFs.

In DG-TPFs, each frequency component of the pump is mapped to a different angle (or transverse momentum k_x) due to angular dispersion (Fig.2a). However the variation of angle with frequency is non-uniform due to group velocity dispersion caused by angular dispersion. Consequently, phase-mismatch occurs for large pump bandwidths or due to cascading (Fig.2a). For the echelon, when beamlets are sufficiently sparse, all frequencies have an overlapping distribution in k_x (2b), which circumvents bandwidth limitations. In Fig.2c, THz spectra resulting from DG-TPFs and echelons are compared theoretically. The simulations incorporate spatio-temporal distortions, cascading, absorption and dispersion [7]. For the conditions in Fig.2c, significantly larger conversion efficiencies with broader spectra and larger THz frequencies are observed for the echelon. This is consistent with trends from experiments [3]. The pump spectral broadening due to cascading is also larger (Fig.2d).

3. Conclusion

The use of echelons for THz generation may result in significantly higher conversion efficiencies of several percent ($\eta = 5$ % for 0.5 ps pump pulses is predicted for echelons (Fig. 2e), superior beam properties, and broader THz spectra by eliminating limitations posed by DG-TPFs. The method also alleviates some experimental challenges, particularly for large beam sizes. In comparison to our proposed aperiodically poled structures [7], the advantage is the ease of availability of large aperture crystals while the disadvantage is the relatively lesser flexibility in terahertz pulse shaping.

4. References

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