

Multiple energy x-ray holography: Incident-radiation polarization effects

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Multiple energy x-ray holography (MEXH) measures both phase and amplitude information for x rays scattered from an incident reference beam, from which three-dimensional atomic images can be directly reconstructed. The angular distribution of the x-ray scattering is highly dependent on the polarization direction via the Thomson scattering cross section. We consider here the effect of incident x-ray polarization on images of Fe atoms reconstructed from theoretical and experimental MEXH data for α -Fe₂O₃(001) (hematite). We also illustrate such polarization effects theoretically in the enhancement of specific atomic structural information of ideal Fe trimers, and a Ge δ -layer buried in Si(001), where the use of different polarization modes and experimental geometries is found to strongly influence atomic images. [S0163-1829(97)04224-0]

I. INTRODUCTION

Gabor proposed holography as a means of surpassing the current limitations of electron microscopy, specifically by experimentally recording the amplitude and phase of scattered wave fronts relative to a direct unscattered reference wave.¹ Later, Szöke observed that the phases of atomically scattered photoelectron or fluorescent x-ray wave fronts can be referenced to the direct or unscattered wave front.² This assumes that the electron or x-ray emission is from a highly localized core excitation. Thus, structural information on the atomic environment surrounding each emitting atom is uniquely encoded in the holographic photoelectron or fluorescence intensities, as detected in the far field outside the sample. Algorithms to retrieve direct three-dimensional atomic images from these holograms have subsequently been developed³⁻⁷ and successfully implemented in application to e.g., photoelectron diffraction data.⁸⁻¹²

These holographic principles have also been demonstrated theoretically for x rays,^{13,14} and more recently experimentally implemented in two distinct forms: single energy *x-ray fluorescence holography* (XFH),¹⁵ and *multiple energy x-ray holography* (MEXH).^{16,17} MEXH is the time-reversed reciprocal method of XFH, and the relationship between these two methods has been discussed in detail in an earlier study.¹⁸ MEXH is found to be more versatile in that aberrations and twin-image overlaps can be suppressed by reconstructing images from holograms covering a range of different energies,^{3,14} but it is also more strongly influenced in both positive and negative ways by the polarization of the incident radiation, and this is the topic of the present paper.

While preliminary MEXH images of Fe atoms in hematite (α -Fe₂O₃(001)) have been reported in a previous study,¹⁶ we here present additional experimental and theoretical atomic images in order to illustrate the effects of the polarization of the incident radiation used, which can enhance or suppress the images at specific atomic locations. These effects are also theoretically demonstrated through the imaging of simple Fe trimers and a Ge δ -layer buried in Si(001), a dilute species-specific structure that is ideally suited for MEXH studies. The limitations and future prospects of this method, specifically as related to polarization effects, are also discussed.

II. IMAGE RECONSTRUCTION OF MEXH DATA

As illustrated in Fig. 1(a), raw MEXH $I(\mathbf{k})$ intensities are measured by detecting the net fluorescence from a sample that is illuminated with coherent radiation incident over different directions and energies. The fluorescing emitter detects the overall radiation field strength resulting from a superposition of wave fronts elastically scattered by

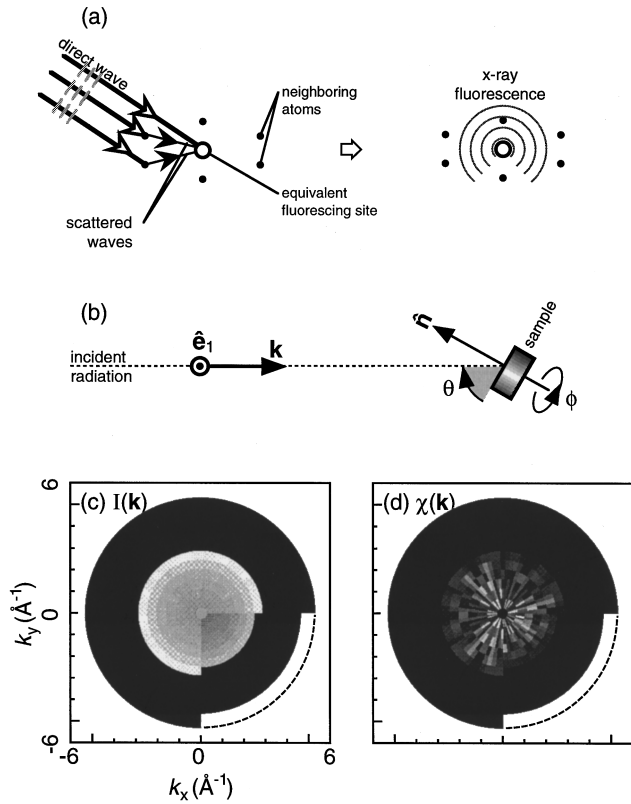


FIG. 1. (a) *Multiple energy x-ray holography*. A coherent far-field excitation x-ray illuminates and photoexcites an emitter, and part of it scatters off of atoms neighboring the emitter in doing so. The emitting atom senses the interference between the direct wave front, and wave fronts elastically scattered by the neighboring atoms. The net photoexcitation is then detected by a stationary, large solid-angle fluorescence detector. Moving the far-field source over a large solid-angle range builds up a holographic interference pattern. (b) Orientation of the sample (where \hat{n} is the surface normal ([001]) with respect to the horizontal (\hat{e}_1) polarization vector of the incident radiation \mathbf{k} . The polar rotation axis $\hat{\theta}$ is parallel to \hat{e}_1 . (c) Schematic \mathbf{k} -space representation of the raw measured MEXH $I(\mathbf{k})$ intensity data set for Fe $K\alpha$ fluorescence from α -Fe₂O₃(001) as excited by horizontally polarized radiation at three energies of $k = 4.561, 4.891, 5.220 \text{ \AA}^{-1}$ ($E = 9.00, 9.65, 10.30 \text{ keV}$). (d) The normalized MEXH $\chi(\mathbf{k})$ data set.

neighboring atoms and the direct, unscattered wave front (i.e., the direct or reference wave). The amount of fluorescence is thus determined by the interference of the scattered and direct wave fronts at the emitting atomic site.^{16–18} Normalized $\chi(\mathbf{k})$ intensities are obtained via $\chi(\mathbf{k}) = [I(\mathbf{k}) - I_0(\mathbf{k})] / \sqrt{I_0(\mathbf{k})}$, where $I(\mathbf{k})$ is the raw measured fluorescence, and $I_0(\mathbf{k})$ is the fluorescence due specifically to the unscattered, direct wave-front excitation of the emitter. These normalized $\chi(\mathbf{k})$ intensities in \mathbf{k} space can be considered an \mathbf{r} -space convolution of the object field $u(\mathbf{r})$ of the scattering atoms surrounding the emitter at the origin, with the convolution kernel describing all of the physics of the x-ray scattering.¹⁹ This object field can in principle be recovered as an image intensity $U(\mathbf{r}')$, via a \mathbf{k} -space deconvolution of the $\chi(\mathbf{k})$ data set that is in its simplest “optical” form given by

$$U(\mathbf{r}') \equiv \iiint_K d^3\mathbf{k} \cdot e^{-i(\mathbf{k} \cdot \mathbf{r}' - k\mathbf{r}')} \chi(\mathbf{k}). \quad (1)$$

This multiple energy imaging algorithm, and its positive aspects in suppressing twin images and other sources of image aberrations was first pointed out by Barton,^{3(a),3(b)} and subsequently proposed in a slightly different form by Tong, Huang, and Wei.^{3(c)} The kernel $e^{-i(\mathbf{k} \cdot \mathbf{r}' - k\mathbf{r}')}$ implies that the scattered x rays are isotropic in amplitude, and that the scattering phase shift is zero, or at least constant, both of which conditions are reasonably well met for the scattering of *unpolarized* x rays.¹⁴ However, since MEXH is most easily done on synchrotron-radiation beamlines whose most intense radiation has very strong linear polarization, we explore here the consequences of the highly anisotropic scattering amplitudes which result largely due to Thomson scattering.

III. THEORETICAL AND EXPERIMENTAL IMAGES OF α -Fe₂O₃(001)

We now show results of applying the reconstruction algorithm of Eq. (1) to experimental and theoretical α -Fe₂O₃(001) (hematite) MEXH $\chi(\mathbf{k})$ data sets, with the radiation polarization being explicitly included in the calculations. The experimental MEXH data was measured by Gog and co-workers at beamline X-14A of the National Synchrotron Light Source at the Brookhaven National Laboratory.¹⁶ Fe $K\alpha$ fluorescence was collected for horizontally polarized incident-radiation spanning three energies from $k = 4.561$ to 5.220 \AA^{-1} ($E = 9.00$ – 10.30 keV) and incident on the α -Fe₂O₃(001) sample over a polar range of $60^\circ \leq \theta \leq 90^\circ$. The experimental geometry is illustrated in Fig. 1(b), with the relationship between the horizontal polarization and the rotation axes of the sample being indicated. The polarization here is parallel to the polar rotation axis ($\hat{\theta}$) of the sample, with the azimuthal rotation axis ($\hat{\phi}$) being parallel to the surface normal. These data points were measured at $\delta k = 0.329 \text{ \AA}^{-1}$ ($\delta E = 650 \text{ eV}$) energy intervals, and at $(\delta\theta, \delta\phi) = (5^\circ, 5^\circ)$ angular intervals, making a total of 435 unique measurements in a symmetry-reduced 1/3rd of the total solid angle above the sample. Figure 1(c) shows the raw measured $I(\mathbf{k})$ intensities in \mathbf{k} space, as viewed down along [001]. Data points in the fourth quadrant have been cut away to reveal the $k = 4.561 \text{ \AA}^{-1}$ $I(\mathbf{k})$ intensities. The dark bands at the perimeter indicate the locations in \mathbf{k} space on these isoenergy surfaces where data was not collected. Note that the weak atomic scattering of x rays renders the anisotropy of the raw $I(\mathbf{k})$ data ($\Delta I/I_0 \approx 0.5\%$) barely discernible with this gray scale; and thus much more demanding statistical accuracy is required in x-ray holographic measurements than in comparable atomic electron holography measurements (where $\Delta I/I_0 \approx 30\%$).²⁰ However, the more ideal and weaker nature of the atomic scattering of x rays also produces atomic images that are relatively free of image aberrations, artifacts, and positions shifts compared to atomic electron holography images.^{14,21–23}

Due to the limited energy range of this $I(\mathbf{k})$ data set, a separate $I_0(\mathbf{k})$ was determined for each of the three different hologram energies via a high-pass filter,²⁴ thereby including in $I_0(\mathbf{k})$ the reference wave as well as any low-angular-

frequency effects in the initial excitation of the x-ray and the subsequent absorption as the fluorescence exits the sample, via

$$I_0(\mathbf{k}) = \frac{\iint_{\tilde{\mathbf{k}}} I(\tilde{\mathbf{k}}) e^{-[(\mathbf{k}-\tilde{\mathbf{k}})/\sigma]^2} \sin \tilde{\theta} d\tilde{\theta} d\tilde{\phi}}{\iint_{\tilde{\mathbf{k}}} e^{-[(\mathbf{k}-\tilde{\mathbf{k}})/\sigma]^2} \sin \tilde{\theta} d\tilde{\theta} d\tilde{\phi}}. \quad (2)$$

All of these effects need to be removed from the raw $I(\mathbf{k})$ intensities before applying the imaging algorithm. Figure 1(d) shows the normalized $\chi(\mathbf{k})$ obtained by this method from the raw $I(k)$ intensities of Fig. 1(c), with holographic modulations more visible.

For comparison, a single-scattering model^{14,25} was used to calculate a theoretical MEXH $\chi(\mathbf{k})$ data set from an ideal α -Fe₂O₃(001) cluster containing only 384 Fe atoms occupying the two inequivalent Fe emitter sites appropriate to this crystal structure. The O atoms were omitted, as they will be negligibly strong in the resulting reconstructed images due to their relatively smaller scattering power.¹⁶ The incident radiation in this model calculation is polarized horizontally with respect to the $\hat{\theta}$ and $\hat{\phi}$ rotations performed on this cluster, just as for the measurement of the experimental $I(\mathbf{k})$ data set discussed above [cf. Fig. 1(b)]. As the sample was rotated over the $60^\circ \leq \theta \leq 90^\circ$ range of the experimental data, the polarization (and the resultant Thomson cross section) rotated also with respect to the crystal structure, but with an average overall effect still of having much stronger scattering for atoms above and below a given emitter along the surface normal or [001] direction.

Figure 2 schematically shows the \mathbf{r} -space regions in the vertical (120) and horizontal (002) planes that would have enhanced/suppressed image intensities due to the effect of using horizontally polarized incident radiation, relative to the image intensities obtained from an unpolarized incident radiation MEXH data set. The image intensity in Fig. 2 is thus the ratio $U_{\text{horizontal}}(\mathbf{r}')/U_{\text{unpolarized}}(\mathbf{r}')$, with white representing a value of unity, and gray, values of < 1 . As can be seen in Fig. 2, the use of horizontally polarized incident radiation would most strongly suppress atomic images in the horizontal (001) plane containing the emitter, with this suppression being less in horizontal planes farther above and below the emitter plane. Images along the [001] axis would be least suppressed by using horizontally polarized incident radiation. Thus, using horizontally polarized radiation in the geometry of Fig. 1(b) would in fact be disadvantageous in the study of horizontal planar structures such as those prevalent in α -Fe₂O₃(001), which is comprised of closely stacked horizontal Fe bilayers.

We now show the reconstructed atomic images obtained from applying Eq. (1) to the experimental (Fig. 3) and theoretical (Fig. 4) single-scattering model MEXH $\chi(\mathbf{k})$ data sets for α -Fe₂O₃(001) in (a) the vertical (120) plane, and (b) the horizontal (002) plane that is 6.89 Å below the emitter. These images are the all superpositions of images for both of the inequivalent Fe emitter sites, with these emitter sites indicated by a common dashed square at the origin, and the relative positions of the Fe scatterers indicated by circles. Fe scatterers just above or below the (002) Fe bilayer are indi-

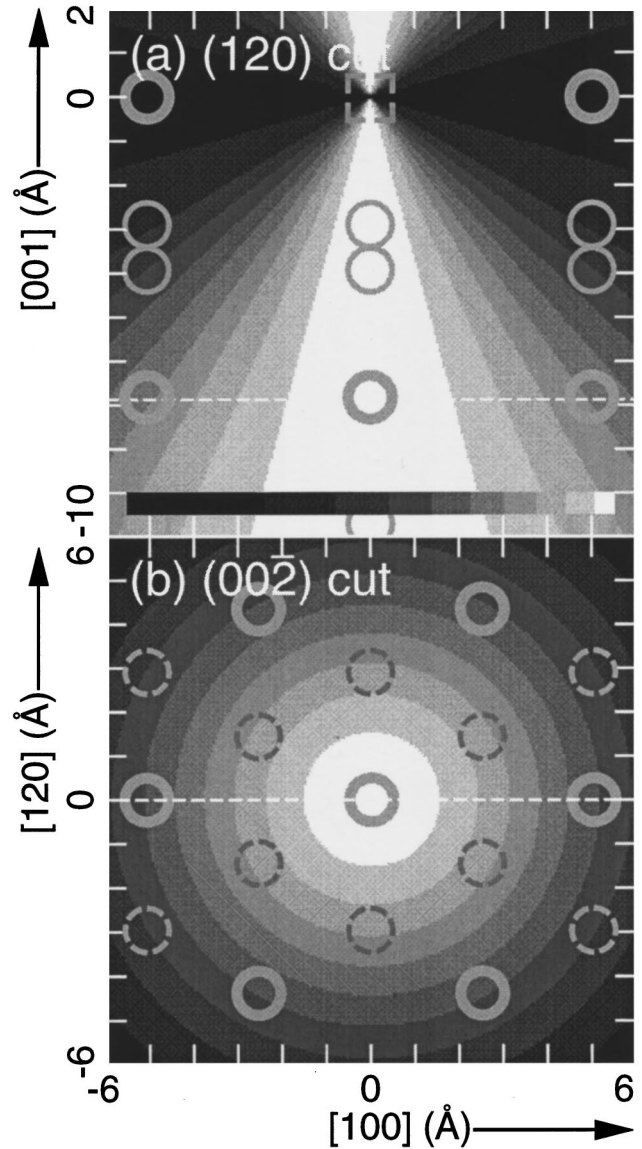


FIG. 2. Theoretically calculated ratio of image intensities obtained using horizontally polarized incident radiation and image intensities obtained using unpolarized incident radiation on α -Fe₂O₃(001), in (a) the vertical (120) plane, and (b) the horizontal (002) plane of 6.89 Å below the emitter. The superposition of both inequivalent Fe emitter sites is indicated by a dashed square, and Fe scatterers are indicated by circles. Fe scatterers in the bilayer just above or below this plane are indicated by dashed circles, and Fe scatterers in relative positions common to both inequivalent Fe emitters are indicated by bold circles. Axes are marked off in 1-Å units.

cated by dashed circles, and Fe scatterers in relative positions common to both inequivalent Fe emitters are indicated by bold circles. The atomic image resolutions expected from $\chi(\mathbf{k})$ data sets of this energy and angular range²⁴ are $\delta x = \delta y \approx 0.6$ Å in the horizontal ([001] and [120]) directions, and $\delta z \approx 2.5$ Å in the vertical ([001]) direction, and these resolutions are indicated by white error bars in the figures. The experimental and theoretical images are in excellent agreement. In particular, both experiment and theory show that three of the Fe atoms from the neighboring upper bilayer intrude into the (002) image plane. This is due to the limited

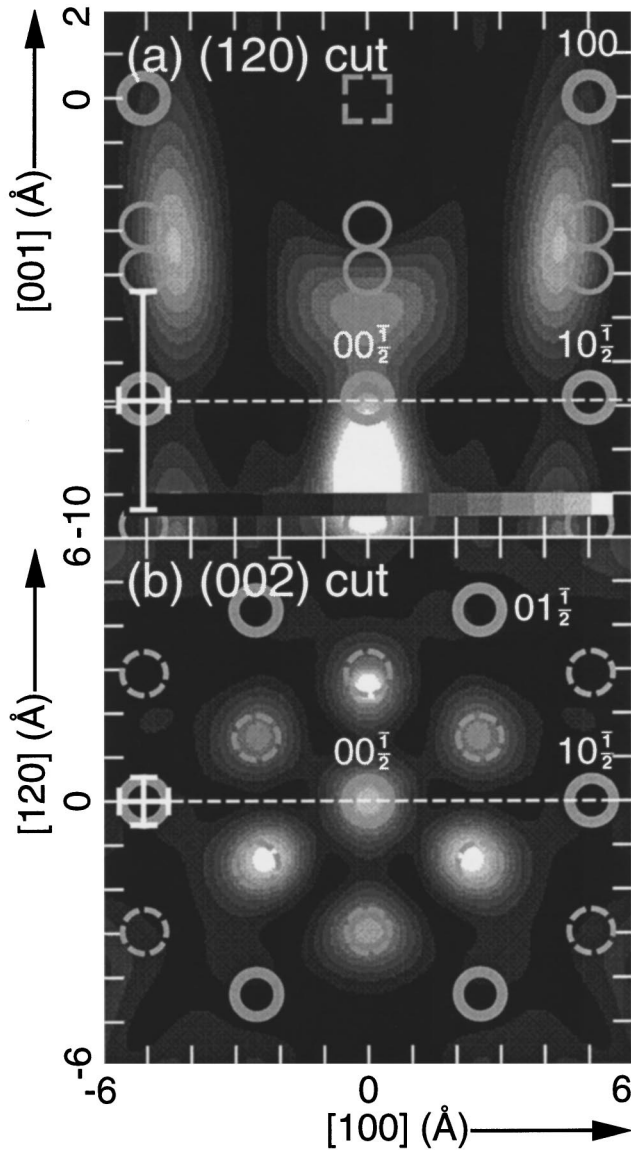


FIG. 3. Reconstructed α -Fe₂O₃(001) atomic images obtained from an experimental MEXH data set measured with horizontally polarized incident radiation, in (a) the vertical (120) plane, and (b) the horizontal (002) plane. The image resolution expected from the energy and angular ranges of this $\chi(\mathbf{k})$ data set are indicated by the white error bars.

energy and angular range of the $\chi(\mathbf{k})$ data set in \mathbf{k} space, which results in atomic images much less resolved in the vertical [001] direction.²⁴ Also, as expected from the suppression of horizontal image intensities that arises from the use of horizontally polarized incident radiation, the Fe atomic images in the (001) plane are not visible in Figs. 3 and 4(a), and the (002) in-plane atoms ($10\frac{1}{2}$, $01\frac{1}{2}$, etc.) are also not visible in Figs. 3 and 4(b), due to their horizontal distance from the [001] axis. The only Fe atomic images that have not been appreciably suppressed lie more nearly along the vertical [001] direction above and below the emitter [cf. Figs. 2(a), 3(a), 4(a)]. In spite of the image suppression due to the polarization of the incident radiation, and the resolution loss due to the limited energy and angular range of these $\chi(\mathbf{k})$ data sets, the resulting images are found to exhibit a

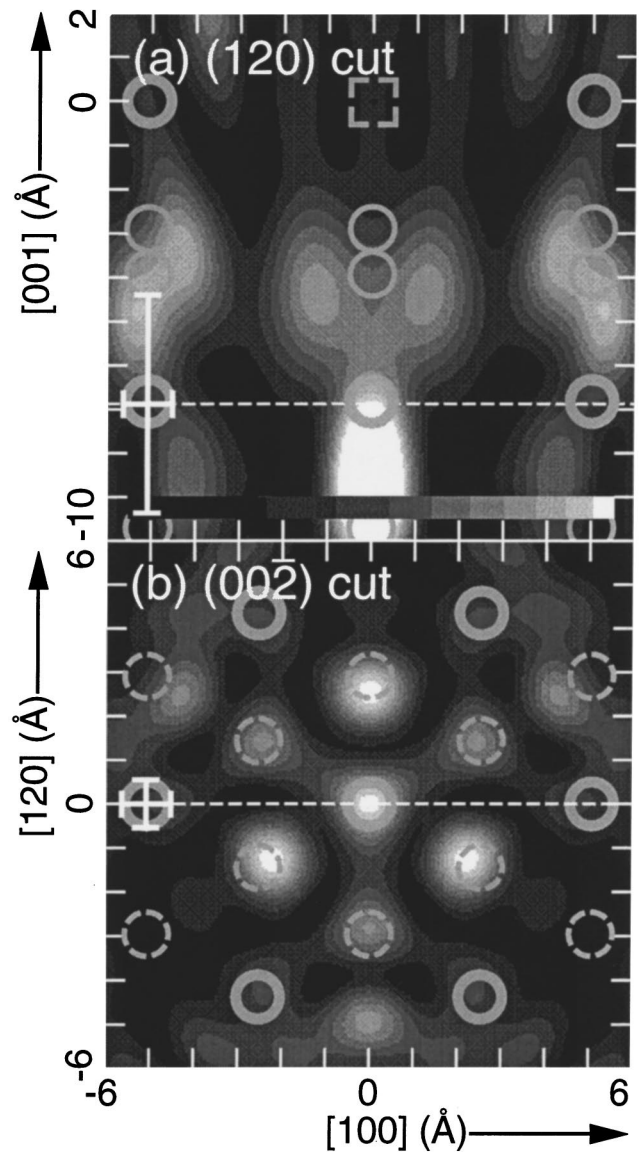


FIG. 4. As in Fig. 3, but for images obtained from a theoretical horizontally polarized MEXH data set.

suppression of real-twin image overlaps, as they are reconstructed from a multiple energy $\chi(\mathbf{k})$ data set.^{3,14,16–18}

IV. EFFECT OF INCIDENT POLARIZATION ON MEXH HOLOGRAPHIC INTENSITIES

In order to best image horizontal and vertical atomic positions simultaneously, utilizing unpolarized incident radiation would be ideal. However, because synchrotron radiation is the most practical experimental MEXH incident-radiation source due to its energy tunability, high brightness, and potentially high-energy resolution (e.g., via undulators and/or crystal monochromators), linearly polarized incident radiation must be considered. We thus now discuss in more detail the effect of incident radiation polarization on the creation of MEXH $\chi(\mathbf{k})$ intensities for two model systems.

Figure 5 shows the Fe atomic scattering factor magnitudes $|f_{\text{Fe}}(\Theta)|$ for x rays at $k = 5.220 \text{ \AA}^{-1}$ ($E = 10.30 \text{ keV}$), where Θ is the angle between the incident (\mathbf{k}) and scattered

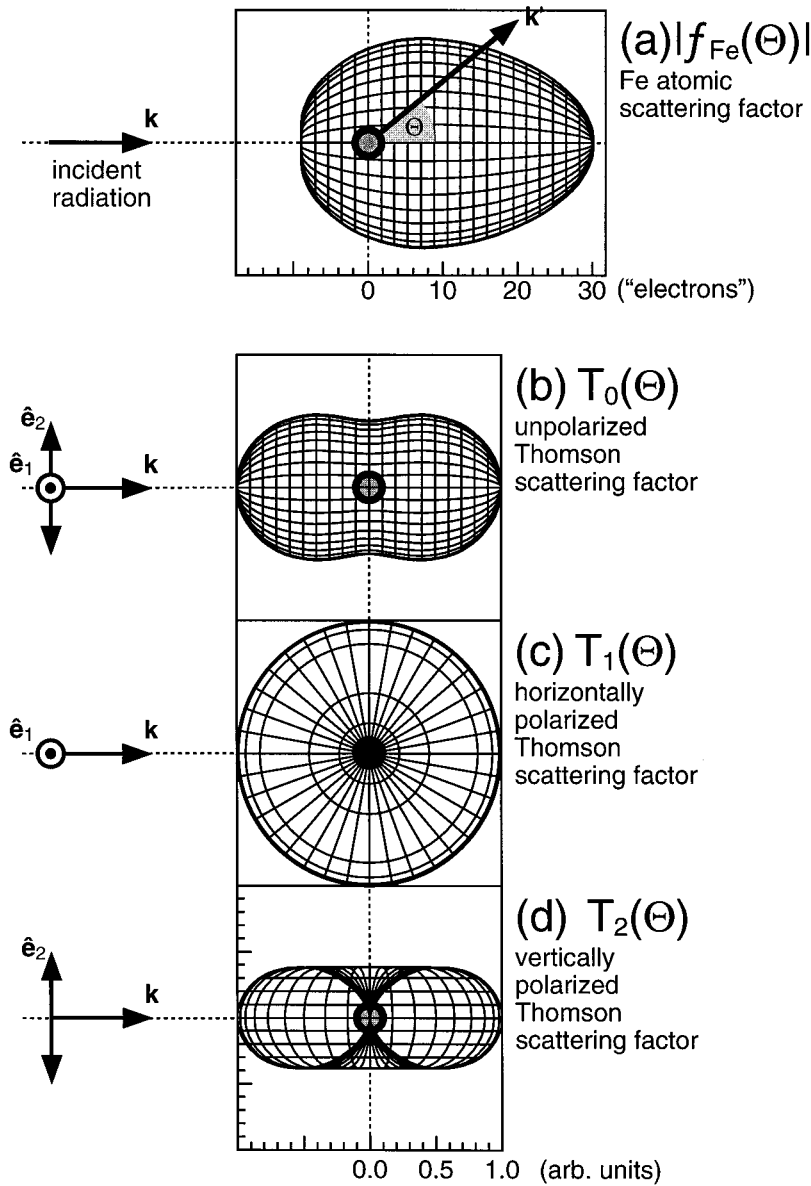


FIG. 5. Scattering factor magnitudes for $k = 5.220 \text{ \AA}^{-1}$ ($E = 10.30 \text{ keV}$) x-rays incident on atomic Fe. Θ is the angle between the incident (\mathbf{k}) and scattered (\mathbf{k}') wave vectors. (a) Fe atomic scattering factor, in relative units of electrons. (b)–(d) Thomson electron scattering factors, in absolute units of $r_e \equiv 2.818 \times 10^{-5} \text{ \AA}$, for (b) unpolarized, (c) vertically polarized, and (d) horizontally polarized incident x rays.

(\mathbf{k}') wave fronts.²⁵ The raw Fe scattering factor magnitude, as given in relative electron-scattering units [Fig. 5(a)], is converted to absolute units by multiplying by the Thomson scattering factor for a single electron interacting with either unpolarized or linearly polarized light, as the case may be. The general form of the well-known Thomson scattering factor is $\sin \Theta_{\hat{\mathbf{e}}}^{\mathbf{k}'}$, where $\Theta_{\hat{\mathbf{e}}}^{\mathbf{k}'}$ is the angle between the polarization vector of the incident radiation $\hat{\mathbf{e}}$, and the direction \mathbf{k}' of the scattered radiation. Figures 5(b)–5(d) show the Thomson scattering factors for unpolarized, horizontally polarized, and vertically polarized light, respectively. The “peanut”-shaped unpolarized electron-scattering factor [Fig. 5(b)] has an azimuthal symmetry about the incident wave-vector axis, while this symmetry is broken for the “bagel”-shaped linearly polarized scattering factors [Figs. 5(c)–(d)], which for horizontal and vertical polarization show preferential scattering in the vertical and horizontal directions, respectively.

The effect of polarized incident radiation in MEXH can be demonstrated by considering the ideal $\chi(\mathbf{k})$ holograms created by simple [001] and [100] linear Fe trimers oriented along the vertical or horizontal direction, respectively, as

shown to the left of Figs. 6–8. In these calculations, the trimers were rotated with respect to the incident-radiation polarization so as to simulate a real experimental situation with a movable sample and a fixed beamline. Figures 6–8 show the expected holographic intensities in the full upper 2π hemisphere above these Fe trimers, as viewed down along the [001] directions, for a hologram energy of $k = 5.220 \text{ \AA}^{-1}$ ($E = 10.30 \text{ keV}$). The holographic interference fringes in the case of unpolarized incident radiation are visible as azimuthal bands for the vertical [001] Fe trimer in Fig. 6(a), and as vertical bands centered along the [100] direction for the horizontal [100] Fe trimer in Fig. 6(b). Figures 9–11 show the atomic image intensities reconstructed via Eq. (1) along these trimers. The full width at half maximum of these image peaks are close to the resolutions expected from the energy and angular range of these $\chi(\mathbf{k})$ data sets ($dx \approx 0.31 \text{ \AA}$; $dz \approx 0.61 \text{ \AA}$).²⁴

Figure 7 shows the MEXH $\chi(\mathbf{k})$ holograms from these same Fe trimers, but in the case of incident-radiation polarized horizontally with respect to a stationary synchrotron source, where the polarization vector $\hat{\mathbf{e}}$ is parallel to the $\hat{\theta}$

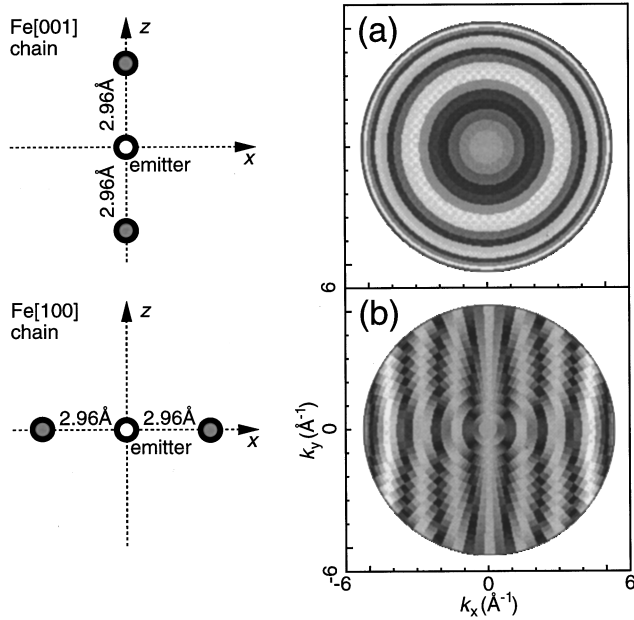


FIG. 6. Normalized holographic intensities $\chi(\mathbf{k})$ as calculated for Fe $K\alpha$ fluorescence excited by $k=5.220 \text{ \AA}^{-1}$ ($E=10.30 \text{ keV}$) unpolarized x rays incident on two different Fe trimers, viewed down along $[001]$. (a) $[001]$ Fe trimer $\chi(\mathbf{k})$. (b) $[100]$ Fe trimer $\chi(\mathbf{k})$.

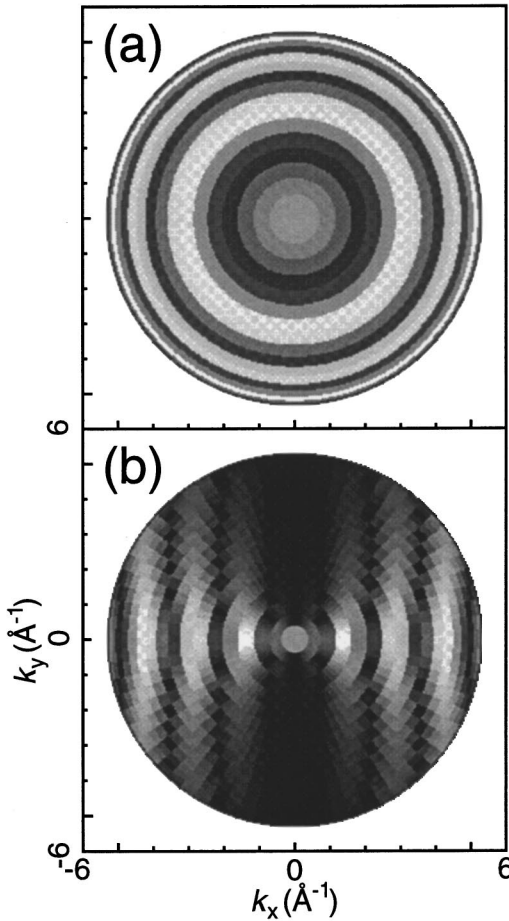


FIG. 7. As in Fig. 6, but for the case of horizontally polarized incident radiation.

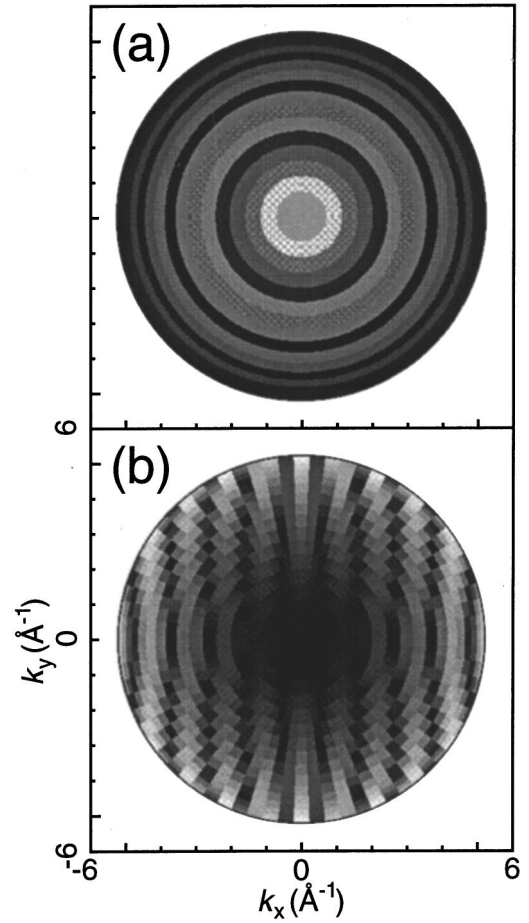


FIG. 8. As in Fig. 6, but for the case of vertically polarized incident radiation.

rotation axis, as shown in Fig. 1(b). Due to the azimuthal symmetry of the horizontal polarization with respect to the $[001]$ Fe trimer for all incident-radiation angles, the resulting hologram in Fig. 7(a) is virtually identical (although relatively less intense, due to the loss of one polarization mode) to the unpolarized incident-radiation hologram in Fig. 6(a). However, the $\chi(\mathbf{k})$ hologram intensities for the $[100]$ Fe trimer in the case of horizontally polarized light are suppressed in directions perpendicular to the $[100]$ trimer axis [Fig. 7(b)]. There, the Thomson scattering factor selects against scattering between the trimer atoms, such that there is much reduced holographic information in these regions. This effect is seen in the reconstructed images of Fig. 10, where the $[001]$ Fe trimer image peaks [Fig. 10(a)] are now approximately four times higher in intensity than the $[100]$ Fe trimer image peaks [Fig. 10(b)]. Thus horizontal polarization emphasizes atomic images along the vertical azimuthal axis, while suppressing horizontal planar atomic images, as noted before for the more complex $\alpha\text{-Fe}_2\text{O}_3(001)$ case.

Figure 8 shows the MEXH $\chi(\mathbf{k})$ holograms from the $[001]$ and $[100]$ Fe trimers, in the case of incident radiation polarized vertically with respect to a stationary synchrotron source, where the polarization vector $\hat{\mathbf{e}}_2$ is now always perpendicular with the $\hat{\theta}$ rotation axis. Again, due to the azimuthal symmetry of the vertical polarization with respect to the $[001]$ Fe trimer for all incident radiation angles, the resulting hologram in Fig. 9(a) is nearly identical in the upper

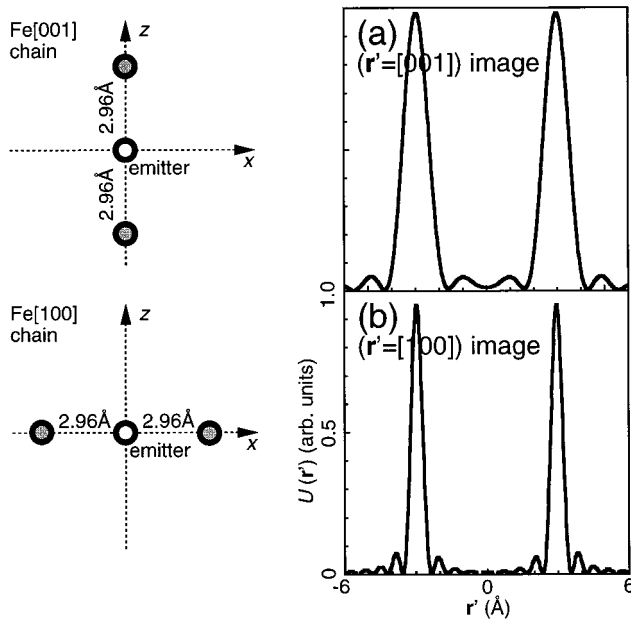


FIG. 9. Reconstructed atomic images obtained from the theoretical unpolarized incident-radiation MEXH $\chi(\mathbf{k})$ data sets of Fig. 6, for (a) the [001] Fe trimer, and (b) the [100] Fe trimer. Image intensities are scaled in arbitrary units relative to each other. The distances r' along the trimer axes (whether [001] or [100]) are marked off in 1-Å units.

polar regions to the unpolarized incident radiation hologram in Fig. 7(a), but is weaker in intensity for lower polar angles, where the polarized Thomson scattering factor selects against scattering between the trimer atoms, leading to reduced holographic information in these regions. For the [100] Fe trimer hologram, the polarized Thomson scattering factor suppresses holographic intensities in the upper polar regions parallel to the [100] trimer orientation. This effect is seen in the reconstructed images of Fig. 11, where the [001] Fe trimer image peaks [Fig. 11(a)] are now approximately half the intensity of the [100] Fe trimer image peaks [Fig. 11(b)]. Thus for horizontally polarized incident radiation, holographic intensities in specific regions are suppressed for both the [001] and [100] trimers, but the net effect for images reconstructed from the full upper 2π hemisphere is to suppress peak intensities along the vertical [001] axis more than in the horizontal plane.

Therefore compared to unpolarized incident radiation, linearly polarized incident radiation reduces the amount of spatial information corresponding to specific atomic sites available in MEXH $\chi(\mathbf{k})$ intensities. Horizontally polarized incident radiation suppresses holographic and reconstructed image intensities for atoms in the horizontal (001) plane of the emitter, while emphasizing intensities for atoms in the vertical [001] axis through the emitter. For low polar angles, vertically polarized incident radiation suppresses low takeoff angle $\chi(\mathbf{k})$ intensities from atoms in the vertical [001] axis through the emitter, while emphasizing $\chi(\mathbf{k})$ intensities from atoms in the horizontal plane of the emitter; these effects are reversed when higher polar angles are considered for vertically polarized incident radiation. When taken over the entire full upper 2π hemisphere of $\chi(\mathbf{k})$ intensities, reconstructed image intensities for atoms along the vertical [001] axis are

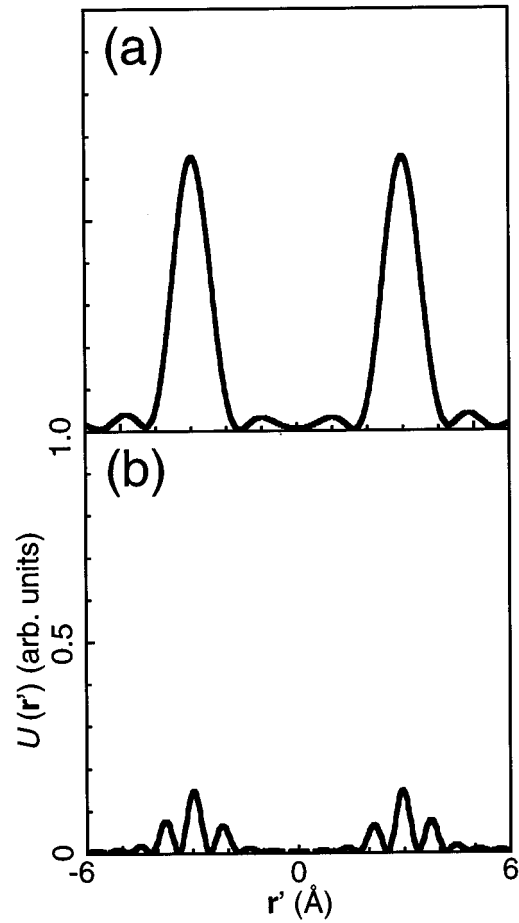


FIG. 10. As in Fig. 9, but for images reconstructed from the horizontally polarized incident-radiation MEXH $\chi(\mathbf{k})$ data of Fig. 7.

suppressed more, relative to atoms in the horizontal plane of the emitter.

V. EFFECT OF INCIDENT POLARIZATION ON MEXH ATOMIC IMAGES

In order to illustrate the possible utility of linearly polarized incident radiation in obtaining MEXH atomic images, we consider a general class of structures where it should be advantageous to utilize horizontally polarized incident radiation, and for which vertical structural information can be much more desirable than horizontal structural information: surface atomic layers and buried epitaxial atomic layers. A specific case theoretically illustrated here is a Ge δ -layer buried in Si(001). The Ge atoms in the δ layer are here assumed to lie in horizontal epitaxial sites with respect to the surrounding Si(001), such that structural information in the horizontal plane of a Ge emitter is relatively unimportant compared to the strained vertical distances between the Ge δ -layer atoms and their Si neighbors above and below them.²⁶ Thus using horizontally polarized incident radiation to record a MEXH Ge $K\alpha$ data set for this system or others like it may prove to be advantageous.

MEXH $\chi(\mathbf{k})$ intensities were calculated for unpolarized, horizontally polarized, and vertically polarized radiation incident on an ideal Ge δ -layer buried in a Si(001) cluster containing a total of 96 atoms, where for simplicity no ver-

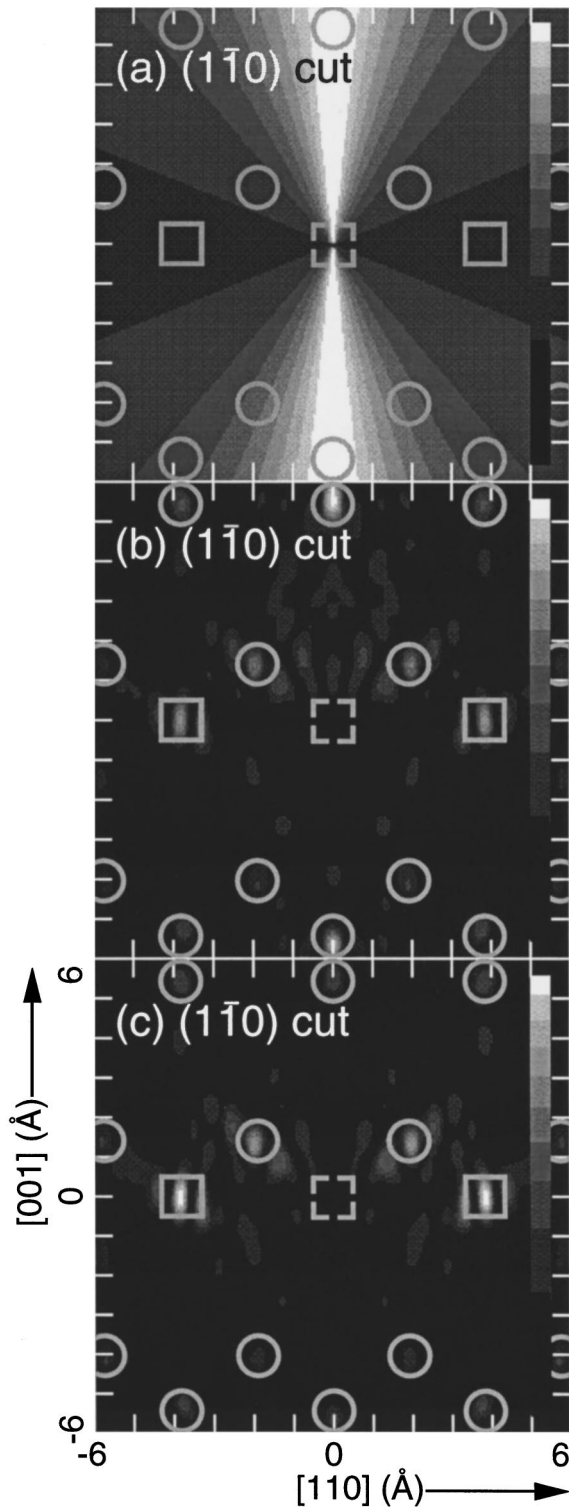


FIG. 13. (a) Ratio of image intensities obtained using horizontally polarized incident radiation to image intensities obtained using unpolarized incident radiation, in the vertical $(\bar{1}10)$ plane. (b) Reconstructed $\text{Si}(001)/\text{Ge-}\delta/\text{Si}(001)$ atomic images obtained from a theoretical horizontally polarized MEXH data set, in the vertical $(\bar{1}10)$ plane. (c) As in (b), but corrected for the effects of the polarization of the incident radiation.

images in the outlying planes of Fig. 12 are strongly suppressed in Fig. 14(b), both due to their distance from the Ge emitter, and the vertical polarization of the incident radiation.

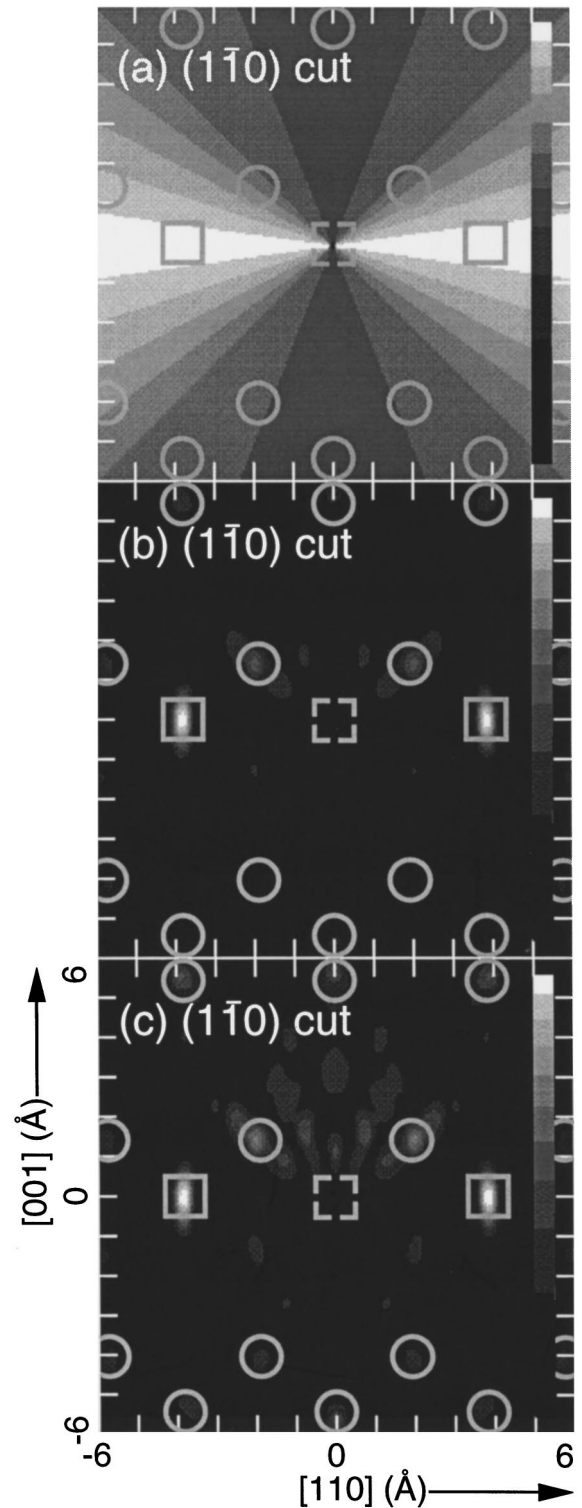


FIG. 14. As in Fig. 13, but for the case of vertically polarized incident radiation.

tion. Thus, recording MEXH intensities with vertically polarized incident radiation would not be most ideal for this particular system, but would be ideal for the imaging of other structures where atomic images in the horizontal (001) plane would be of more interest.

Also, note, the presence of faint image aberrations in Figs. 13(b) and 14(b) near the emitter, as compared to Fig. 12. These aberrations are due to there being less holographic

information overall (i.e., fewer visible hologram fringes) in the MEXH $\chi(\mathbf{k})$ data sets obtained with polarized incident radiation, relative to the set obtained with unpolarized incident radiation (cf. Figs. 6–11).

The relative intensities of the atomic sites that have been preferentially/detrimentally imaged due to the polarization of the incident radiation can be restored. The use of a so-called SWIFT (scattered-wave included fourier transform) correction to the reconstruction algorithm of Eq. (1) is one method for accomplishing this, given by^{4(a)}

$$U(\mathbf{r}') \equiv \iiint_K d^3\mathbf{k} \cdot \frac{e^{-i(\mathbf{k} \cdot \mathbf{r}' - k r')}}{f_{\text{Fe}}(\Theta_{\mathbf{k}}^{\mathbf{r}'}) \sin(\Theta_{\mathbf{e}}^{\mathbf{r}'})} \chi(\mathbf{k}). \quad (3)$$

Here, one is simply dividing by the strength of the scattered wave associated with an atom along \mathbf{r}' . However, doing this type of correction becomes problematic due to the nodes of the Thomson scattering factor $\sin(\Theta_{\mathbf{e}}^{\mathbf{r}'})$, such that the integrand in Eq. (3) will have divisions by zero. An *ad hoc* solution would then be to limit the SWIFT correction to the angular regions where the Thomson scattering factor is appreciably nonzero.^{27,28} However, we also note that the relative intensities of atomic sites enhanced/diminished by polarization effects can be restored simply by dividing the images by the appropriate enhancement ratios $R_{\text{enh}} = U_{\text{polarized}}(\mathbf{r}')/U_{\text{unpolarized}}(\mathbf{r}')$. Determining this ratio does not require the direct experimental measurement of $\chi(\mathbf{k})$ intensities obtained using polarized and unpolarized incident radiation, as we find from an analysis of exact calculations for various geometries that it can be approximated using

$$R_{\text{enh}} = \frac{U_{\text{polarized}}(\mathbf{r}')}{U_{\text{unpolarized}}(\mathbf{r}')} \approx \frac{\iiint_K d^3\mathbf{k} \cdot \sin(\Theta_{\mathbf{e}}^{\mathbf{k}'=\mathbf{r}'})}{\iiint_K d^3\mathbf{k} \cdot \frac{1}{2} \sqrt{1 + \cos^2(\Theta_{\mathbf{k}}^{\mathbf{k}'=\mathbf{r}'})}}. \quad (4)$$

The enhancement ratio is thus approximated by the relative contribution to an atomic image at \mathbf{r}' due to the \mathbf{k} -space domains in $\chi(\mathbf{k})$ that are selectively enhanced by the polarized Thomson scattering factor. Figures 13(c) and 14(c) show the reconstructed atomic images of Figs. 13(b) and 14(b) that have been corrected by dividing out the appropriate approximate polarization (horizontal and vertical, respectively) enhancement factor of Eq. (4). The relative intensities between all of the atomic sites in the corrected images of Figs. 13(c) and 14(c) are now well restored to uniform values, in comparison to the unpolarized atomic images of Fig. 12. Note that the atomic sites that have been detrimentally

affected by the use of polarized incident radiation are discernible in both uncorrected [Figs. 13(b) and 14(b)] and corrected images [Figs. 13(c) and 14(c)]; merely the relative intensities between all atomic sites have been equalized in the corrected polarized images. As expected, the faint image aberrations in Figs. 13(b) and 14(b) persist in the corrected images of Figs. 13(c) and 14(c), due to the overall loss of signal to noise in polarized incident-radiation $\chi(\mathbf{k})$ data sets, compared to the image reconstructed from the unpolarized incident-radiation $\chi(\mathbf{k})$ data set (Fig. 12).

VI. CONCLUSIONS

Multiple energy x-ray holography holds much promise in the imaging of local atomic structure surrounding a specific emitter species of interest, as demonstrated here in the first experimental data obtained with it, and for the model system of buried atomic δ layers. However, there are strong effects in these atomic images due to the Thomson scattering cross section when the linearly polarized radiation is used for exciting the fluorescence signal. Utilization of horizontally or vertically polarized incident radiation can emphasize vertical or horizontal atomic structures, respectively. Approximate procedures for correcting for this nonuniformity in image strengths are also discussed. If full three-dimensional atomic structural information is desired from a specific sample, then either unpolarized or circularly polarized radiation should be used for excitation, or the sample should be mounted such that it can be illuminated by either linear polarization mode relative to the polar and azimuthal rotations being carried out on it. MEXH intensities that would be measured for the case of unpolarized incident radiation also can be determined from the two individual polarized MEXH data sets as obtained over most of the solid angle above the sample.

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