

Search for low-mass Higgs states @ *BABAR*

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Several types of new-physics models predict the existence of light dark matter candidates and low-mass Higgs states. Previous *BABAR* searches for invisible light-Higgs decays have excluded large regions of model parameter space. We present searches for a dark-sector Higgs produced in association with a dark gauge boson and searches for a light Higgs in $\Upsilon(nS)$ decays.

1 Search for Low-Mass Dark Matter at *BABAR*

We have now an overwhelming astrophysical evidence of dark matter. To explain this experimental evidence many theoretical models [1] introducing new dark forces mediated by new gauge boson have been introduced. These models propose Weakly Interacting Massive Particles (WIMPs) that can annihilate into pairs of dark bosons, which subsequently annihilate to lepton pairs (protons are kinematically forbidden). One of this model [2, 3] introduce a new dark sector that couples to the SM with a dark boson (i.e. the dark photon A') through a small kinetic mixing term. Thanks to their large luminosities and low-background environment the *B*-factories offer an ideal place to probe for MeV-GeV dark matter, complementing searches from LHC. The dark boson mass is generated via the Higgs mechanism, adding a dark Higgs boson (h') to the theory. The dark photon and the Higgs bosons could have a comparable mass (GeV-scale). A very minimal scenario has a single dark photon and a single dark Higgs boson. In the *BABAR* analysis [4] we make use of the *Higgsstrahlung* process $e^+e^- \rightarrow A'^* \rightarrow A'h', h' \rightarrow A'A'$ using 521 fb^{-1} . The signal is either fully reconstructed into lepton or pion pair (exclusive mode), or partially reconstructed (inclusive mode). Only two of the three hidden photons are reconstructed in the latter case, and the four-momentum of the third one is identified to that of the recoiling system.

In these searches no significant signal is observed. Using uniform priors in the cross section upper limits on the $e^+e^- \rightarrow A'^* \rightarrow A'h', h' \rightarrow A'A'$ cross section are obtained as a function of the hidden Higgs and hidden photon masses. These limits on the cross section are translated into 90 % upper limit on $\alpha_D \epsilon^2$ shown in Fig. 1, where $\alpha_D = g_D^2/4\pi$, g_D is the dark sector gauge coupling, ϵ is the mixing strength. Values down to 10^{-10} - 10^{-8} are excluded for a large range of hidden photon and hidden Higgs masses, assuming prompt decay. Assuming $\alpha_D = \alpha \simeq 1/137$, limits on the mixing strength in the range 10^{-4} - 10^{-3} are derived as shown in Fig. 2, these limits are an order of magnitude smaller than the current experimental bounds extracted from direct photon production in this mass range.

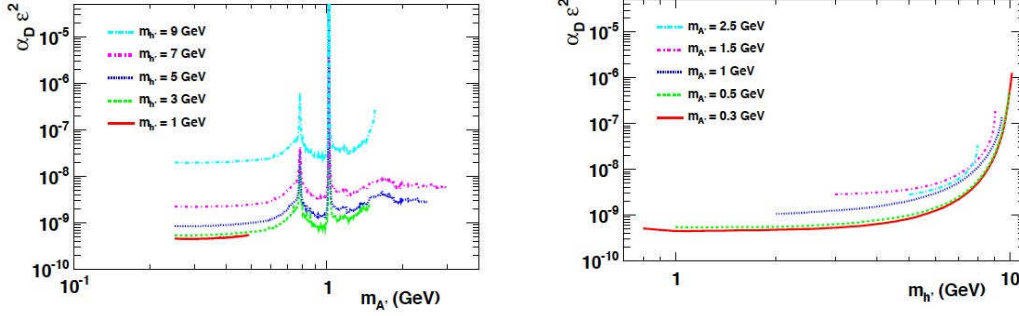


Figure 1: Limits on $\alpha_D \epsilon^2$ for the Dark Photon (left) and Dark Higgs (right).

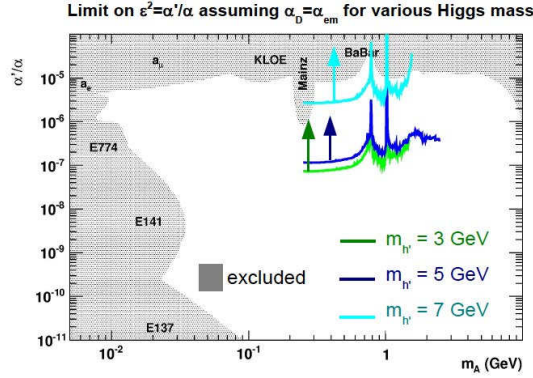


Figure 2: Limit on the mixing strength, ϵ^2 , for various Higgs masses.

2 Search for light Higgs @ *BABAR*

In recent years a number of theoretical models [5] predicted for the existence of a light CP-odd Higgs boson A^0 related to the Next-to-Minimal Supersymmetric Model (NMSSM). Direct searches constrain the mass of $m_{A^0} < 2m_b$, where m_b is the b quark and the decay of $A^0 \rightarrow b\bar{b}$ is forbidden. Of particular interest is to search the lightest CP-odd Higgs boson in Υ decays such as $\Upsilon(nS) \rightarrow \gamma A^0$, where $A^0 \rightarrow SM$ particles. In these $\Upsilon(nS)$ transitions low mass Dark Matter Candidate (χ) can be also directly produced (i.e. $\Upsilon(nS) \rightarrow \chi\chi$). The large data sets available at *BABAR* experiment with more than $500 fb^{-1}$ of data at the $\Upsilon(4S)$, $\Upsilon(3S)$, $\Upsilon(2S)$ resonances and just below the $\Upsilon(4S)$ resonance allow us to place stringent constraints on such theoretical models.

We searched for light Higgs in *BABAR* using two-body-radiative decay on Υ states. The key experimental signature is the monochromatic photon in the CM frame $E_\gamma^* = \frac{m_\Upsilon - m_{A^0}^2}{2m_\Upsilon}$. With this analysis technique we searched for the following transitions:

- $\Upsilon(2S, 3S) \rightarrow \gamma A^0$, $A^0 \rightarrow \mu^+ \mu^-$ [6]

- $\Upsilon(2S, 3S) \rightarrow \gamma A^0$, $A^0 \rightarrow \tau^+ \tau^-$ [7]
- $\Upsilon(2S, 3S) \rightarrow \gamma A^0$, $A^0 \rightarrow \text{hadrons}$ [8]

In all the searches we did not find any evidence of signals but upper limits versus hypothesis mass have been extracted the results are shown in Fig. 3 for the $A^0 \rightarrow \mu^+ \mu^-$, in Fig. 4 for the $A^0 \rightarrow \tau^+ \tau^-$ (where τ is reconstructed both in $\tau \rightarrow e \nu \bar{\nu}$ and $\tau \rightarrow \mu \nu \bar{\nu}$) and in Fig. 5 for the $A^0 \rightarrow \text{hadrons}$.

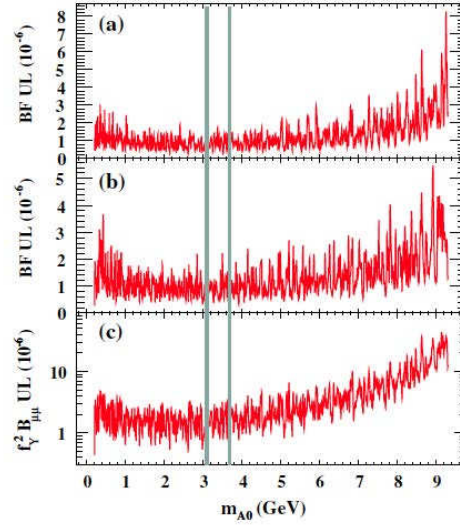


Figure 3: 90 % C.L. upper limits on (a) $\mathcal{B}(\Upsilon(2S) \rightarrow \gamma A^0) \times \mathcal{B}_{\mu^+ \mu^-}$, (b) $\mathcal{B}(\Upsilon(3S) \rightarrow \gamma A^0) \times \mathcal{B}_{\mu^+ \mu^-}$ and (c) effective coupling $f_{\Upsilon}^2 \times \mathcal{B}_{\mu^+ \mu^-}$ as a function of m_{A^0} . The shaded areas show the regions around the J/ψ and the $\psi(2S)$ resonances excluded from the search.

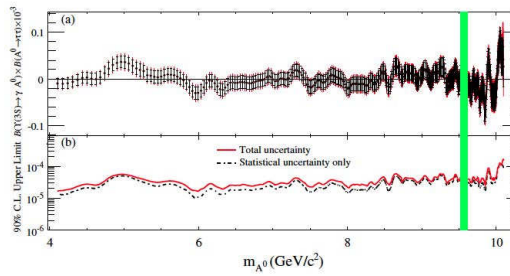


Figure 4: (a) Product branching fractions as a function of the Higgs mass. (b) The corresponding 90% C.L. upper limits on the product of the branching fractions versus the Higgs mass values.

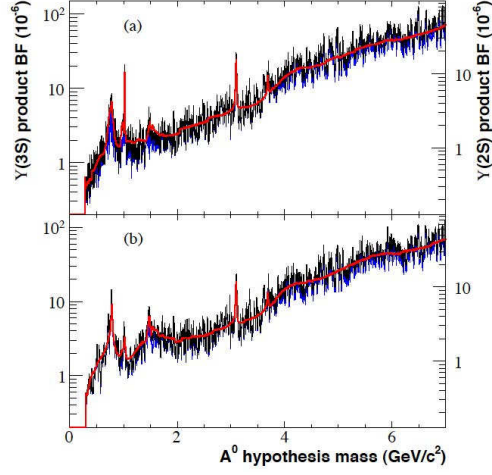


Figure 5: 90 % C.L. upper limits on product branching fractions (BF) (left axis) $\mathcal{B}(\Upsilon(3S) \rightarrow \gamma A^0) \mathcal{B}(A^0 \rightarrow \text{hadrons})$ and (right axis) $\mathcal{B}(\Upsilon(2S) \rightarrow \gamma A^0) \mathcal{B}(A^0 \rightarrow \text{hadrons})$ for (a) CP-all analysis and (b) CP-odd analysis.

3 Conclusion

We have searched for evidence of dark sector candidates and evidence of CP-odd light Higgs in the $\Upsilon(2S)$ and $\Upsilon(3S)$ data sample at *BABAR*. We did not observe any significant signal but more stringent limit have been set on space parameters of NP model.

References

- [1] B. Batell, M. Pospelov and A. Ritz, Phys. Rev. D **79**, 115008 (2009)
- [2] N. Arkani-Hamed, D. P. Finkbeiner, T. R. Slatyer and N. Weiner, Phys. Rev. D **79**, 015014 (2009)
- [3] R. Essig, P. Schuster and N. Toro, Phys. Rev. D **80**, 015003 (2009)
- [4] arXiv:1202.1313 [hep-ex]
- [5] R. Dermisek, J. F. Gunion and B. McElrath, Phys. Rev. D **76**, 051105 (2007)
- [6] B. Aubert *et al.*, Phys. Rev. Lett. **103**, 081803 (2009)
- [7] B. Aubert *et al.*, Phys. Rev. Lett. **103**, 181801 (2009)
- [8] J. P. Lees *et al.*, Phys. Rev. Lett. **107**, 221803 (2011)