



# Probing the very high energy $\gamma$ -ray spectral curvature in the blazar PG 1553+113 with the MAGIC telescopes

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## ABSTRACT

PG 1553+113 is a very high energy (VHE,  $E > 100$  GeV)  $\gamma$ -ray emitter classified as a BL Lac object. Its redshift is constrained by intergalactic absorption lines in the range  $0.4 < z < 0.58$ . The MAGIC telescopes have monitored the source's activity since 2005. In early 2012, PG 1553+113 was found in a high state, and later, in April of the same year, the source reached its highest VHE flux state detected so far. Simultaneous observations carried out in X-rays during 2012 April show similar flaring behaviour. In contrast, the  $\gamma$ -ray flux at  $E < 100$  GeV observed by *Fermi*-LAT is compatible with steady emission. In this paper, a detailed study of the flaring state is presented. The VHE spectrum shows clear curvature, being well fitted either by a power law with an exponential cut-off or by a log-parabola. A simple power-law fit hypothesis for the observed shape of the PG 1553+113 VHE  $\gamma$ -ray spectrum is rejected with a high significance (fit probability  $P = 2.6 \times 10^{-6}$ ). The observed curvature is compatible with the extragalactic background light (EBL) imprint predicted by current generation EBL models assuming a redshift  $z \sim 0.4$ . New constraints on the redshift are derived from the VHE spectrum. These constraints are compatible with previous limits and suggest that the source is most likely located around the optical lower limit,  $z = 0.4$ , based on the detection of Ly $\alpha$  absorption. Finally, we find that the synchrotron self-Compton model gives a satisfactory description of the observed multiwavelength spectral energy distribution during the flare.

**Key words:** galaxies: active – BL Lacertae objects: general – BL Lacertae objects: individual: PG1553+113 – cosmic background radiation – gamma-rays: galaxies – gamma-rays: general.

## 1 INTRODUCTION

PG 1553+113 is a blazar found as part of the Palomar-Green Catalogue of UV-excess Stellar Objects (Green, Schimdt & Liebert 1986). Its J2000 coordinates are RA  $15^{\text{h}}55^{\text{m}}43^{\text{s}}.0$ , Dec.  $+11^{\circ}11'24''.4$  (Beasley et al. 2002). It was classified as a BL Lac object due to its featureless optical spectrum (Miller & Green 1983) and significant optical variability (Miller et al. 1988). As occurs in most BL Lac objects, the featureless optical spectrum prevents a spectroscopic redshift measurement. However, several limits have been provided based on indirect measurements (e.g. Sbarufatti, Treves & Falomo 2005; Sbarufatti et al. 2006). The most recent redshift lower limit was estimated assuming that the host galaxy can be used as a standard candle. For absolute  $R$ -band magnitudes  $MR = -22.5$  and  $-22.9$ , Shaw et al. (2013) obtains the limits  $z > 0.24$  and  $0.31$ , respectively. Previously, a more stringent redshift lower limit of  $z > 0.4$  was set by Danforth et al. (2010) based on the detection of intervening Ly $\alpha$  absorbers. This estimation will be used throughout the paper. Danforth et al. (2010) also set a redshift upper limit of  $z < 0.58$  based on the non-detection of any Ly $\beta$  absorbers at  $z > 0.4$ .

The very high energy (VHE)  $\gamma$ -ray emission from PG 1553+113 was discovered independently and almost simultaneously by HESS (Aharonian et al. 2006b) and MAGIC (Albert et al. 2007) in 2005. The integral flux recorded by MAGIC at the time of the discovery was  $F = (10.0 \pm 0.2_{\text{stat}}) \times 10^{-11} \text{ cm}^{-2} \text{ s}^{-1}$  above 120 GeV, and the differential energy spectrum was well described by a power

law with a spectral index  $\Gamma \sim 4$ , compatible with the detection by HESS. The source has been monitored with the MAGIC telescopes since 2005. The results from the 2005–2009 observation campaigns can be found in Aleksić et al. (2012b). Modest flux variability of a factor of  $\sim 2.6$  on a yearly time-scale has been detected at  $E > 150$  GeV, with an integral flux lying in the range from  $1.4$  to  $3.7 \times 10^{-11} \text{ cm}^{-2} \text{ s}^{-1}$ . The observed energy spectra were well fitted by power laws with photon indices in the range  $\Gamma \sim 3.6$ – $4.3$  and compatible within uncertainties.

Extragalactic VHE  $\gamma$ -rays can be absorbed on the way to the Earth via electron–positron pair production when interacting with optical-UV background photons from the extragalactic background light (EBL; Gould & Sreder 1967; Stecker, De Jager & Salamon 1992). The EBL is mainly composed of diffuse optical light emitted by stars and partially reprocessed by dust in the IR, redshifted by the expansion of the Universe (Hauser & Dwek 2001). The uncertainty on its spectral energy distribution (SED) and evolution through the history of the Universe still ranges from 20 to 50 per cent at wavelengths 0.4 and 40  $\mu\text{m}$ , respectively. This uncertainty is mainly due to difficulties in direct measurements.

During the past few years, several different approaches have been developed to model the EBL (e.g. Franceschini, Rodighiero & Vaccari 2008; Finke, Razzaque & Dermeret 2010; Kneiske & Dole 2010; Domínguez et al. 2011; Gilmore et al. 2012; Scully, Malkin & Stecker 2014) and despite the different techniques adopted the resulting EBL models show an overall agreement, differing only marginally.

The  $\gamma$ -ray absorption depends significantly on the energy of the VHE photon, the redshift-dependent SED of the EBL, and the distance to the source. The observed flux ( $F_{\text{obs}}$ ) can be expressed as

$$F_{\text{obs}}(E) = F_{\text{int}}(E)e^{-\tau(E,z)}, \quad (1)$$

where  $F_{\text{int}}$  denotes the intrinsic flux emitted by the source and  $\tau$  the EBL optical depth as a function of the energy and redshift.

The EBL imprint on the VHE  $\gamma$ -ray spectrum can be used to set upper limits on the redshift of the source. This is done by assuming a particular EBL model and a criterion on the intrinsic spectrum, such as a maximum hardness for the reconstructed intrinsic spectrum or the absence of a spectral break with a pile-up at VHE in the reconstructed spectrum.

Different authors have used this  $\gamma$ -ray attenuation technique for PG 1553+113, leading to the following limits:  $z < 0.74$  (Aharonian et al. 2006a),  $z < 0.42$  (Mazin & Goebel 2007),  $z < 0.66$  (Prandini et al. 2010),  $z \leq 0.62$  (Aliu et al. 2015),  $z = 0.49 \pm 0.04$  (Abramowski et al. 2015),  $z < 0.53$  (Biteau & Williams 2015). Limits on the EBL absorption can be estimated independently from EBL models using the VHE spectrum and the redshift of the source under the assumption that the emission of the source can be properly described by a synchrotron self-Compton (SSC) model (Mankuzhiyil, Persic & Tavecchio 2010). This method has previously been used on PG 1553+113 to derive constraints on the  $\gamma$ -ray horizon (Domínguez et al. 2013). Alternatively, the EBL density relative to that of a given model can be evaluated through the joint analysis of the HE or VHE observations of many extragalactic sources, making relatively small assumptions on the shape of the intrinsic spectra (Ackermann et al. 2012b; Abramowski et al. 2013a; Biteau & Williams 2015).

PG 1553+113 was detected in the high-energy (HE,  $E > 100$  MeV)  $\gamma$ -ray band by the Large Area Telescope (LAT) on board the *Fermi* Gamma-ray Space Telescope (Abdo et al. 2010). The energy spectrum for the period 2008 August–2010 August can be well fitted by a power law with spectral index  $\Gamma = 1.67 \pm 0.02_{\text{stat}}$  and  $F(E > 100 \text{ MeV}) = (6.5 \pm 0.6_{\text{stat}}) \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1}$  and its variability index is 93.5 (Nolan et al. 2012). Since the variability index is  $>41.6$ , the source is variable on a monthly time-scale at  $>99$  per cent confidence probability. No flaring activity has been claimed for PG 1553+113 in the HE band to date. Remarkably, a periodic factor of 2 flux enhancement on a monthly time-scale over six years of *Fermi* operation has been recently found by Acero et al. (in preparation).

An extensive multiwavelength (MWL) observation campaign on PG 1553+113 was carried out from 2012 February to June, focused on the characterization of its SED as well as the variability of the source emission at different frequencies. Observations from VHE  $\gamma$ -rays to radio were performed: VHE band by MAGIC, HE band by *Fermi*-LAT, X-rays by *Swift*-X-ray Telescope (XRT), UV-optical observations by *Swift*-Ultraviolet Optical Telescope (UVOT), IR by Rapid Eye Mount (REM) and radio by Metsähovi and Owens Valley Radio Observatory (OVRO). In this paper, the study of the flux variability in the VHE, HE and X-ray bands is presented. The study on the VHE spectrum is focused on the April flare state of the source. A detailed study on the long-term MWL campaign will be presented in a forthcoming paper.

The paper is organized as follows: Section 2 describes the MWL data analysis. In Section 3, the results are presented. The light curves from MAGIC, *Fermi*-LAT and *Swift*-XRT are shown in Section 3.1 while a detailed analysis on the observed VHE energy spectrum is presented in Section 3.2. The intrinsic VHE  $\gamma$ -ray spectrum together with a discussion on EBL imprint and redshift constraints can be

found in Section 4. The SED observed during the flaring state and the theoretical interpretation is described in Section 5. The conclusions can be found in Section 6.

## 2 OBSERVATIONS

### 2.1 VHE $\gamma$ -ray observations with MAGIC

The VHE  $\gamma$ -ray observations were performed by the MAGIC telescopes. The MAGIC system consists of two 17 m-diameter Imaging Atmospheric Cherenkov Telescopes (IACT) located on the Roque de los Muchachos, Canary Island of La Palma ( $28^\circ 46' \text{N}$ ,  $17^\circ 53' \text{W}$ ), at a height of 2200 m above sea level. The system reaches a sensitivity of  $(0.76 \pm 0.03)$  per cent of the Crab nebula flux for  $E > 290$  GeV in 50 h of observations (Aleksić et al. 2012a).

PG 1553+113 was observed with the MAGIC telescopes from 2012 February 26 (MJD 55983) to April 26 (MJD 56043). The data sample after quality cuts consists of 18.3 h in the zenith angle range  $17^\circ$ – $34^\circ$ . The observations were performed in wobble mode (Fomin et al. 1994), with the source located  $0.4$  from the centre of the field of view. The analysis of the data has been performed using the standard MAGIC analysis chain (Moralejo et al. 2009; Lombardi et al. 2011). The energy threshold of the analysis is approximately 70 GeV.

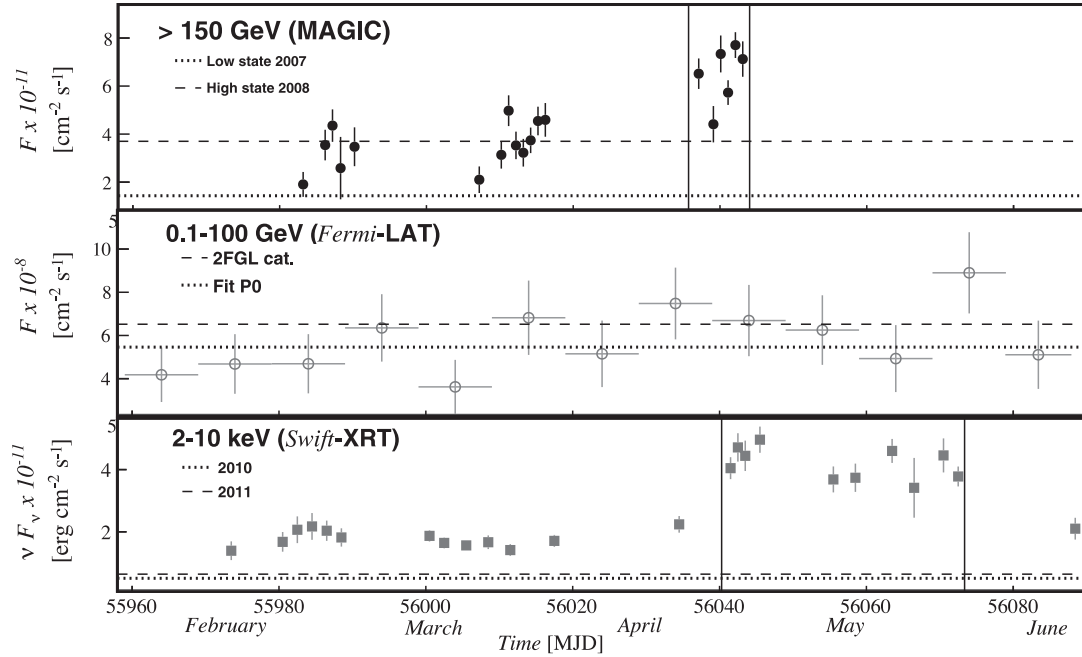
The source was detected with a high statistical significance ( $>70$  standard deviations,  $\sigma$ ) during the time period 2012 February–April. The emission is compatible with a point-like source at the position of PG 1553+113. The mean  $\gamma$ -rate during the flare period is  $4.35 \pm 0.04 \text{ min}^{-1}$  for  $E > 70$  GeV.

### 2.2 HE $\gamma$ -rays observations from *Fermi*-LAT

The *Fermi*-LAT is a pair-conversion telescope operating from 20 MeV to  $> 300$  GeV. Further details about the *Fermi*-LAT can be found in Atwood et al. (2009). The LAT data reported in this paper were collected from 2012 February 2 (MJD 55959) to June 10 (MJD 56088). During this period, the *Fermi* observatory operated almost entirely in survey mode. The analysis was performed with the SCIENTETOOLS software package version v9r32p5. The LAT data were extracted within a  $10^\circ$  region of interest centred at the location of PG 1553+113. Only events belonging to the ‘Source’ class were used. The time intervals when the rocking angle of the LAT was greater than  $52^\circ$  were rejected. In addition, a cut on the zenith angle ( $<100^\circ$ ) was applied to reduce contamination from the Earth limb  $\gamma$ -rays, which are produced by cosmic rays interacting with the upper atmosphere. The spectral analysis was performed with the instrument response functions P7REP\_SOURCE\_V15 using an unbinned maximum-likelihood method implemented in the Science tool *gtlike*. Isotropic (iso\_source\_v05.txt) and Galactic diffuse emission (gll\_iem\_v05\_rev1.fit) components were used to model the background<sup>1</sup> (Ackermann et al. 2012a). The normalizations of both components in the background model were allowed to vary freely during the spectral fitting.

We evaluated the significance of the  $\gamma$ -ray signal from the source by means of the maximum-likelihood test statistic  $TS = 2 (\log L_1 - \log L_0)$ , where  $L$  is the likelihood of the data given the model with ( $L_1$ ) or without ( $L_0$ ) a point source at the position of PG 1553+113 (e.g. Mattox et al. 1996). The model of the region of interest used in *gtlike* includes all point sources from the second *Fermi*-LAT

<sup>1</sup> <http://fermi.gsfc.nasa.gov/ssc/data/access/lat/BackgroundModels.html>



**Figure 1.** Light curve of PG 1553+113 during the 2012 observation campaign. Upper panel: nightly light curve in VHE  $\gamma$ -rays observed with the MAGIC telescopes for  $E > 150$  GeV. For comparison, the flux during the low state of the source in 2007 and the high state in 2008 measured by MAGIC (Aleksić et al. 2012b) is shown by the dotted and dashed lines, respectively. Middle panel: *Fermi*-LAT light curve computed with a 10-d binning for energies between 0.1 and 100 GeV; the dashed grey lines show a zero-order polynomial fit and the mean flux of the source from 2FGL catalogue (Nolan et al. 2012). Lower panel: flux in the X-ray band in one night bins observed by *Swift*-XRT from 2 to 10 keV is represented in grey squares. For comparison, previous measurements from 2010 and 2011 with *Swift*-XRT are plotted as dashed lines. The vertical lines denote the flare intervals observed in the VHE band and X-rays. No hints of flare are observed in the HE band.

catalogue (2FGL; Nolan et al. 2012) as well as from a preliminary third *Fermi*-LAT catalogue from four years of survey observations (Acero et al. 2015) that fall within  $15^\circ$  radius around the source. The spectra of these sources were parametrized by power-law functions, except for 2FGL J1504.3.1+1023, 2FGL J1553.5+1255 and 2FGL J1608.5+1029, for which we used a log-parabola (LP) as in the 2FGL catalogue. A first maximum-likelihood analysis was performed to remove from the model sources having  $TS < 10$  and/or predicted number of counts based on the fitted model  $N_{\text{pred}} < 1$ . A second maximum-likelihood analysis was performed on the updated source model. In the fitting procedure, the normalization factors and the photon indices of the sources lying within  $10^\circ$  of PG 1553+113 were left as free parameters. For the sources located between  $10^\circ$  and  $15^\circ$ , we kept the normalization and the photon index fixed to the values from the 2FGL catalogue. Integrating over the period from 2012 February 2 to June 10 (MJD 55959–56088) the fit with a power-law model in the 0.1–100 GeV energy range results in a  $TS = 908$ , with an integrated average flux of  $(5.7 \pm 0.7_{\text{stat}}) \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1}$  at the decorrelation energy of 2239 MeV and a photon index of  $\Gamma = (1.59 \pm 0.05_{\text{stat}})$  for PG 1553+113. Using an LP,  $dN/dE \propto (E/E_0)^{-\alpha-\beta \log(E/E_0)}$ , the fit yielded for the same period a  $TS = 910$ , with an average flux of  $(4.1 \pm 0.8_{\text{stat}}) \times 10^{-8} \text{ ph cm}^{-2} \text{ s}^{-1}$ , a spectral slope of  $\alpha = 1.49 \pm 0.08_{\text{stat}}$  at the reference energy  $E_0 = 2239$  MeV, and a curvature parameter  $\beta = 0.08 \pm 0.04_{\text{stat}}$ . We used a likelihood ratio test (LRT) to check a PL model (null hypothesis) against an LP model (alternative hypothesis). These values may be compared, following Nolan et al. (2012), by defining the curvature test statistic  $TS_{\text{curve}} = (TS_{\text{LP}} - TS_{\text{PL}})$ . The LRT results in a  $TS_{\text{curve}} = 2$ , corresponding to a  $\sim 1.4 \sigma$  difference. We, therefore, conclude, that no significant curvature was observed in the LAT spectrum of PG 1553+113 during 2012 February–June due to the

poor photon statistics. Similar results were obtained when considering only photons with  $E > 1$  GeV. Above 10 GeV, the analysis is strongly affected by the lack of statistics. No variability during the period was neither found. However, the *Fermi*-LAT spectrum of the source shows curvature when considering a longer integration time interval and we accumulate more photons at the highest *Fermi*-LAT energies. In the 3FGL catalogue (Acero et al. 2015), the spectrum of PG1553+113 is described by an LP.

The  $\gamma$ -ray light curve using 10-d time bins and a PL model is reported in the middle panel of Fig. 1. For each time bin, the spectral shape of PG 1553+113 and of all the sources within  $10^\circ$  of it were fixed to the value obtained over the whole period.

The systematic uncertainty in the flux is dominated by the systematic uncertainty in the effective area (Ackermann et al. 2012a). The systematic uncertainty on the effective area amounts to 10 per cent at 100 MeV, decreasing linearly with the logarithm of energy to 5 per cent between 316 MeV and 10 GeV, and increasing linearly with the logarithm of energy up to 10 per cent at 100 GeV.<sup>2</sup>

### 2.3 X-rays and optical-UV observations from *Swift*

*Swift* target of opportunity observations (Gehrels et al. 2004) of PG 1553+113 were triggered by an increase of the flux emission observed in the VHE band by the MAGIC telescopes (Cortina et al. 2012a,b). The *Swift* observations (XRT and UVOT) were performed in 2012 from February 16 (MJD 55973) to June 24 (MJD 56102). Previous observations in 2010 and 2011 have also been used for comparison purposes. The data taken with XRT on board *Swift* were

<sup>2</sup> [http://fermi.gsfc.nasa.gov/ssc/data/analysis/LAT\\_caveats.html](http://fermi.gsfc.nasa.gov/ssc/data/analysis/LAT_caveats.html)



processed with standard procedures (`xrtpipeline` v0.12.6), filtering, and screening criteria by using the `HEASOFT`<sup>3</sup> package (v6.12). The data were collected both in photon counting (PC) and windowed timing (WT) mode, and XRT event grades 0–12 and 0–2 for the PC and WT events were selected, respectively (Burrows et al. 2005). Source events in WT mode were extracted from a circular region with a radius of 20 pixels (1 pixel  $\sim 2.36$  arcsec), while background events were extracted from a circular region with the same radius away from the source region. Observations in PC mode showed an average count rate of  $>0.5$  counts  $\text{s}^{-1}$ , thus requiring pile-up correction. We extracted the source events from an annular region with an inner radius of 5 pixels (estimated by means of the PSF-fitting technique) and an outer radius of 30 pixels. We extracted background events within an annular region centred on the source with radii 70 and 120 pixels. Ancillary response files were generated with `xrtmkarf`, and account for different extraction regions, vignetting and point spread function (PSF) corrections. We used the most recent spectral redistribution matrices in the calibration data base maintained by HEASARC. We fit the spectrum with an absorbed LP (`logpar` in `XSPEC`; e.g. Massaro et al. 2004) using the photoelectric absorption model `tbabs` (Wilms, Allen & McCray 2000), with a neutral hydrogen column density fixed to its Galactic value ( $3.65 \times 10^{20} \text{ cm}^{-2}$ ; Kalberla et al. 2005).

During the *Swift* pointings, the UVOT instrument observed PG 1553+113 in all its optical (*v*, *b* and *u*) and UV (*w1*, *m2* and *w2*) photometric bands (Poole et al. 2008; Breeveld et al. 2010). We analysed the data using the `uvotsource` task included in the `HEASOFT` package. Source counts were extracted from a circular region of 5 arcsec radius centred on the source, while background counts were derived from a circular region of 10 arcsec radius in the source neighbourhood. Conversion of magnitudes into dereddened flux densities was obtained by using the  $E(B - V)$  value of 0.046 from Schlafly & Finkbeiner (2011), the extinction laws by Cardelli, Clayton & Mathis (1989) and the magnitude-flux calibrations by Bessell, Castelli & Plez (1998).

## 2.4 Infrared observations from REM

PG 1553+113 was observed in the IR regime by the REM telescope from 2012 February 12 (MJD 55969) to July 30 (MJD 56138). The REM (Zerbi et al. 2001; Covino et al. 2004) is a robotic telescope located at the European Southern Observatory (ESO) Cerro La Silla (Chile). It has a Ritchey–Chrétien configuration with a 60-cm *f*/2.2 primary and an overall *f*/8 focal ratio in a fast moving alt-azimuth mount providing two stable Nasmyth focal stations. At one of the two foci, the telescope simultaneously feeds, by means of a dichroic beamsplitter, two cameras: REMIR for the near-infrared band (NIR; Conconi et al. 2004) and REM Optical Slitless Spectrograph (ROSS, Tosti et al. 2004) for the optical band. The cameras both have a field of view of  $10 \text{ arcmin} \times 10 \text{ arcmin}$  and imaging capabilities with the usual NIR (*z*, *J*, *H* and *K*) and Johnson–Cousins *VR*I filters. The REM software system (Covino et al. 2004) is able to manage complex observational strategies in a fully autonomous way. All raw optical/NIR frames obtained with REM telescopes were reduced following standard procedures, i.e. dark frames obtained with the same exposure time are subtracted, sky flat-fields are applied and multiple dithered images are combined to derive sky frames. Multiple scientific frames are then combined to derive the final scientific

images. Instrumental magnitudes were obtained via aperture photometry and absolute calibration has been performed by 2MASS objects in the field. The flux was corrected for Galactic reddening and extinction making use of Schlafly & Finkbeiner (2011) extinction maps.

## 2.5 Radio observations from Metsähovi and OVRO

PG 1553+113 was observed by the Metsähovi 13.7-m radio telescope at 37 GHz during the MWL campaign from 2012 February 19 (MJD 55976) to March 24 (MJD 56010). The measurements were made with a 1 GHz-band dual beam receiver centred at 37 GHz. The observations are ON–ON observations, alternating the source and the sky in each feed horn. A detailed description of the observation and analysis methods can be found in Teräsranta et al. (1998). The detection limit (defined as  $S/N \geq 4$ ) of the telescope is of the order of 0.2 Jy under optimal weather conditions. Given the fact that the typical flux density of PG 1553+113 is close to this limit, the source was significantly detected only on 2012 March 8 (MJD 55994) with a flux  $F = (0.20 \pm 0.05_{\text{stat}})$  Jy.

The source is also monitored at 15 GHz using the 40-m telescope of the OVRO as a part of a larger monitoring program where a sample of  $\sim 1700$  sources are observed twice a week (Richards et al. 2011). The telescope is equipped with dual-beamed off-axis optics and a cooled receiver installed at the prime focus. The two sky beams are Dicke switched using the off-source beam as a reference, and the source is alternated between the two beams in an ON–ON fashion to remove atmospheric and ground contamination. Calibration is referenced to 3C 286 for which the flux density of 3.44 Jy at 15 GHz is assumed (Baars et al. 1977). The systematic uncertainty is about 5 per cent in the flux density scale. Details on the observations, calibration and analysis are given in Richards et al. (2011).

## 3 RESULTS

In this section, a detailed analysis of the  $\gamma$ -ray and X-ray data is presented.

### 3.1 Flux variability

The light curves at VHE  $\gamma$ -rays, HE  $\gamma$ -rays and X-rays are shown in Fig. 1. For the VHE and X-rays bands, a nightly time-scale is used, while for the HE band we have used a 10-d binning. Clear variability is detected in both VHE and X-ray bands. The hypothesis of a constant flux can be rejected with high confidence level,  $P = 1.4 \times 10^{-21}$  ( $\chi^2/\text{ndf} = 143.5/18$ ) in VHE  $\gamma$ -rays and  $P = 1.7 \times 10^{-50}$  ( $\chi^2/\text{ndf} = 302.1/23$ ) in X-rays. The HE flux is compatible with a constant flux of  $F = (5.5 \pm 0.4_{\text{stat}}) \times 10^{-8} \text{ cm}^{-2} \text{ s}^{-1}$  for energies 0.1–100 GeV with a fit probability of  $P = 0.6$  ( $\chi^2/\text{ndf} = 10.7/12$ ). Note that the HE light curve is dominated by the emission at  $E < 10$  GeV, accounting for 95 per cent of the photons.

In the VHE band, two states can be differentiated according to the source flux. In 2012 February–March the average source flux was at a level of  $F(E > 150 \text{ GeV}) = (3.40 \pm 0.15_{\text{stat}}) \times 10^{-11} \text{ cm}^{-2} \text{ s}^{-1}$ , corresponding to  $\sim 11$  per cent of the Crab nebula flux measured by MAGIC (Albert et al. 2008). In 2012 April, the source reached a flux above 150 GeV of  $(7.7 \pm 0.5_{\text{stat}}) \times 10^{-11} \text{ cm}^{-2} \text{ s}^{-1}$ , approximately 24 per cent of the Crab nebula flux. Past MAGIC integral flux measurements above 150 GeV lie in the range between 4 per cent (2007 observations) and 11 per cent (2008 observations) of the Crab nebula flux as reported in Aleksić et al. (2012b). Therefore, we can

<sup>3</sup> <http://heasarc.nasa.gov/heasoft/>

conclude that in 2012 February–March the source was at a level comparable with a previously observed high state in 2008 (dashed line in the upper panel of Fig. 1). In 2012 April, instead, it reached the highest flux observed to date, about six times larger than the low state observed in 2007 and around twice that in February–March of the same year. According to the flux level, we divided the data into two samples: MJD 55983–MJD 56016 (high state) and MJD 56037–MJD 56043 (flare). The probability of a constant fit for both periods independently are low,  $P = 3.2 \times 10^{-3}$  and  $5.1 \times 10^{-3}$ , respectively. The shortest variability time-scale observed is of the order of 1 d. No intranight variability was detected. During the flare, the VHE flux approximately doubled with respect to the high state. The source was also observed by HESS and VERITAS during this high state, and the results of these studies were reported during the publication of this manuscript (Abramowski et al. 2015; Aliu et al. 2015).

The X-ray flux observed in 2012 February–March in the 2–10 keV band is compatible with a constant fit ( $\chi^2/\text{ndf} = 12.7/12$ ,  $P = 0.4$ ), with a mean flux  $(1.71 \pm 0.06) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$ . In 2012 April–May, the source was in a flare state compatible with a constant flux of  $(4.20 \pm 0.14) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$  ( $\chi^2/\text{ndf} = 10.9/9$ ,  $P = 0.3$ ). Later in 2012 June, the source flux decreased to a level compatible with the flux measured during February–March (high state). For comparison, the flux measured in this band from previous observations during 2010 and 2011 was  $(0.59 \pm 0.07_{\text{stat}}) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$ ,  $(0.49 \pm 0.07_{\text{stat}}) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$  and  $(0.64 \pm 0.13_{\text{stat}}) \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$  measured during MJD 55198, MJD 55232 and MJD 55781, respectively. We can conclude that the X-ray flux doubled during the observation campaign. During the flare state, the source reached a level of  $\sim 7$ –10 times the quiescent flux of the source measured during 2010 and 2011.

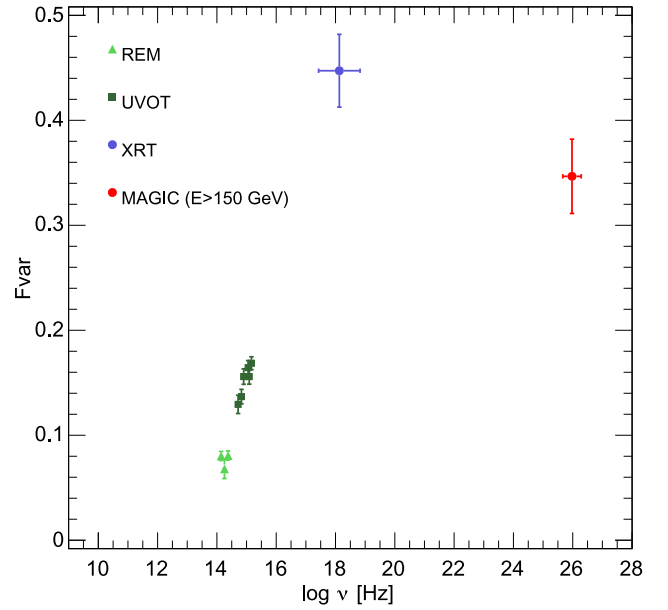
Due to the lack of strictly simultaneous observations, it is difficult to perform an accurate comparison of the VHE  $\gamma$ - and X-ray variability properties. However, the flux evolution in both wavelengths suggests a correlation between the two bands. The SED can, in fact, be properly described in the framework of a one-zone SSC model, pointing to a common origin of the emission in both energy bands as will be discussed in Section 6. Moreover, to quantify the intrinsic variability amplitude, the fractional variability was calculated at each measured frequency. The exception is the observations at 37 GHz since the Metsähovi observations resulted in one single detection. The fractional variability amplitude  $F_{\text{var}}$  is defined as (Vaughan et al. 2003):

$$F_{\text{var}} = \sqrt{\frac{S^2 - \langle \sigma_{\text{err}}^2 \rangle}{\langle F_{\gamma} \rangle^2}}, \quad (2)$$

where  $\langle F_{\gamma} \rangle$  represents the average photon flux,  $S$  the standard deviation of the  $N$  flux measurements and  $\langle \sigma_{\text{err}}^2 \rangle$  the mean squared error.  $F_{\text{var}}$  is estimated for each frequency bin independently. The uncertainty on  $F_{\text{var}}$  is calculated following the prescription from Poutanen, Zdziarski & Ibragimov (2008) as described in Aleksić et al. (2015):

$$\Delta F_{\text{var}} = \sqrt{F_{\text{var}}^2 + \text{err}(\sigma_{\text{NXS}}^2)} - F_{\text{var}}, \quad (3)$$

where  $\sigma_{\text{NXS}}^2$  is given by equation 11 in Vaughan et al. (2003). The study was done only for the period in which we have the full MWL coverage (MJD 55983–56043). Daily fluxes have been used for all the frequencies except for the LAT, for which a 10-d bins have been used in order to detect the source in the individual bins. The fractional variability as a function of the frequency is shown in Fig. 2



**Figure 2.** Fractional variability as a function of the frequency measured by different instruments.

for those bands with positive excess variance ( $S^2$  larger than  $\sigma_{\text{err}}^2$ ). We obtained negative excess variance for the radio (OVRO) and HE  $\gamma$ -ray (*Fermi*-LAT) bands, resulting in  $F_{\text{var}}^2 = -0.003$  and  $F_{\text{var}}^2 = -0.02$ , respectively. Such negative excess variance is interpreted as absence of variability either because there was no variability or because the instruments were not sensitive enough to detect it. Fig. 2 shows clearly that the strongest intrinsic variability is observed in the X-rays (*Swift*/XRT) and the VHE (MAGIC) bands.

### 3.2 X-ray to VHE spectral analysis

In X-rays, PG 1553+113 showed a spectral curvature that can be well described with a log-parabolic shape (see e.g. Perlman et al. 2005; Tramacere et al. 2007). During the 2012 observation campaign, the spectral index and curvature parameter varied in the range 2.0–2.4 and 0.2–0.6, respectively (Table 1). Only for the observations performed on 2012 February 25 and 27 (MJD 55982 and 55984), no significant curvature seems to be present from the fit. No obvious connection was observed between the flux level and the curvature of the X-ray spectra.

As reported in Section 2.2, no significant curvature was observed in the LAT spectrum of PG 1553+113 during 2012 February–June.

In this paper, only the VHE  $\gamma$ -ray spectrum during the 2012 April flare (MJD 56037–56043) is presented, as mentioned in Section 1. The VHE  $\gamma$ -ray spectra observed by MAGIC in 2012 February–March will be presented in a forthcoming paper.

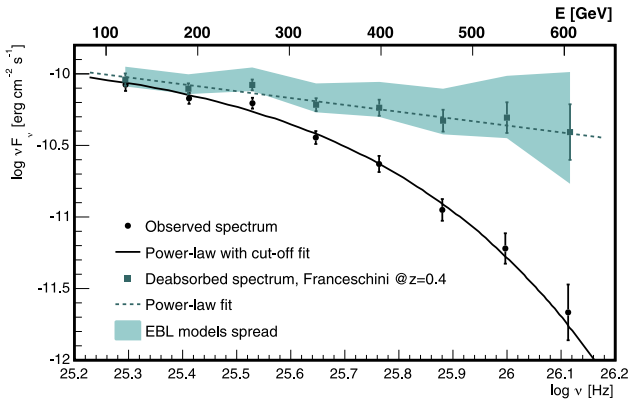
The VHE spectrum during the flare is represented by black circles in Fig. 3. The differential VHE  $\gamma$ -ray spectral points can be found in Table 2. They are corrected for instrumental effects by using the Schmelling unfolding algorithm (Albert et al. 2007).

The observed spectrum shows curvature, and a simple power-law fit can be discarded with a confidence level of  $4.7\sigma$  ( $P = 2.6 \times 10^{-6}$ ,  $\chi^2/\text{ndf} = 36.1/6$ ). The differential spectrum can be well fitted by a power law with an exponential cut-off with a probability of  $P = 0.7$  ( $\chi^2/\text{ndf} = 2.8/5$ ) in the energy range from  $\sim 70$  to 620 GeV:

$$\frac{dF}{dE} = f_0 \left( \frac{E}{200 \text{ GeV}} \right)^{-\Gamma} e^{-E/E_c}, \quad (4)$$

**Table 1.** Log and fitting results of *Swift*-XRT observations of PG 1553+113 using an LP model with a H I column density fixed to the Galactic value in the direction of the source. Fluxes are unabsorbed.

Date (MJD)	Date (UT)	Net exposure time (s)	$\alpha$	$\beta$	Flux 2.0–10 keV ( $10^{-12}$ erg cm $^{-2}$ s $^{-1}$ )	$\chi^2_{\text{red}}$ (d.o.f.)
55973	2012-02-16	2145	$2.13 \pm 0.11$	$0.31 \pm 0.19$	$1.39 \pm 0.31$	1.018 (36)
55980	2012-02-23	2015	$1.96 \pm 0.10$	$0.27 \pm 0.15$	$1.68 \pm 0.32$	0.902 (47)
55982	2012-02-25	1948	$2.40 \pm 0.10$	$0.19 \pm 0.17$	$2.07 \pm 0.43$	0.867 (39)
55984	2012-02-27	2035	$2.18 \pm 0.10$	$0.06 \pm 0.05$	$2.17 \pm 0.43$	1.129 (46)
55986	2012-02-29	1923	$2.11 \pm 0.10$	$0.26 \pm 0.16$	$2.04 \pm 0.33$	1.098 (42)
55988	2012-03-02	2165	$2.18 \pm 0.09$	$0.45 \pm 0.16$	$1.82 \pm 0.29$	0.999 (54)
56000	2012-03-14	1956	$2.14 \pm 0.06$	$0.40 \pm 0.12$	$1.87 \pm 0.18$	0.930 (120)
56002	2012-03-16	2190	$2.22 \pm 0.05$	$0.37 \pm 0.10$	$1.64 \pm 0.17$	1.198 (141)
56005	2012-03-19	2023	$2.24 \pm 0.05$	$0.49 \pm 0.12$	$1.56 \pm 0.15$	1.019 (142)
56008	2012-03-22	1657	$2.25 \pm 0.06$	$0.44 \pm 0.12$	$1.67 \pm 0.22$	1.208 (125)
56011	2012-03-25	1775	$2.27 \pm 0.06$	$0.41 \pm 0.13$	$1.41 \pm 0.19$	1.053 (112)
56017	2012-03-31	2019	$2.30 \pm 0.05$	$0.29 \pm 0.11$	$1.71 \pm 0.19$	1.082 (136)
56034	2012-04-17	912	$2.20 \pm 0.07$	$0.42 \pm 0.14$	$2.24 \pm 0.27$	1.201 (87)
56041	2012-04-24	1062	$2.12 \pm 0.05$	$0.30 \pm 0.11$	$4.05 \pm 0.35$	1.038 (136)
56042	2012-04-25	977	$1.97 \pm 0.06$	$0.43 \pm 0.10$	$4.71 \pm 0.46$	1.099 (136)
56043	2012-04-26	996	$2.17 \pm 0.06$	$0.23 \pm 0.12$	$4.44 \pm 0.49$	1.078 (102)
56045	2012-04-28	999	$2.07 \pm 0.05$	$0.42 \pm 0.09$	$4.96 \pm 0.42$	0.964 (159)
56055	2012-05-08	529	$2.10 \pm 0.07$	$0.58 \pm 0.14$	$3.68 \pm 0.42$	0.863 (83)
56058	2012-05-11	1184	$2.15 \pm 0.06$	$0.25 \pm 0.12$	$3.74 \pm 0.45$	1.017 (108)
56063	2012-05-16	1098	$2.11 \pm 0.05$	$0.32 \pm 0.10$	$4.60 \pm 0.39$	1.165 (155)
56070	2012-05-23	1256	$2.09 \pm 0.07$	$0.36 \pm 0.13$	$4.46 \pm 0.54$	0.901 (96)
56072	2012-05-25	946	$2.11 \pm 0.06$	$0.37 \pm 0.11$	$3.78 \pm 0.32$	1.078 (122)
56088	2012-06-10	955	$2.19 \pm 0.08$	$0.29 \pm 0.16$	$2.07 \pm 0.35$	1.061 (70)
56091	2012-06-13	976	$2.04 \pm 0.07$	$0.52 \pm 0.15$	$2.41 \pm 0.28$	0.965 (89)
56094	2012-06-16	602	$2.14 \pm 0.09$	$0.30 \pm 0.19$	$2.45 \pm 0.45$	1.253 (59)
56095	2012-06-17	983	$2.05 \pm 0.07$	$0.42 \pm 0.14$	$2.49 \pm 0.30$	0.916 (87)
56102	2012-06-24	556	$2.00 \pm 0.09$	$0.61 \pm 0.19$	$2.66 \pm 0.42$	1.116 (60)



**Figure 3.** SED of PG 1553+113 as measured by MAGIC during the flare state of 2012 April. The observed SED is shown as black circles, and the black solid line represents the best fit to a power law with an exponential cut-off. The absorption-corrected spectrum assuming  $z = 0.4$  and using the EBL model by Franceschini et al. (2008) is shown by the green squares; the dashed green line is the best-fitting power law. The green shaded area accounts for the uncertainties derived by the use of different EBL models.

with a normalization constant of  $f_0 = (3.2 \pm 1.4_{\text{stat}} \pm 0.7_{\text{sys}}) \times 10^{-9} \text{ cm}^{-2} \text{ s}^{-1} \text{ TeV}^{-1}$ , a photon index of  $\Gamma = (1.87 \pm 0.37_{\text{stat}} \pm 0.15_{\text{sys}})$  and  $E_c = (110 \pm 24_{\text{stat}} \pm 19_{\text{sys}}) \text{ GeV}$ . A full description of the systematics uncertainties for the MAGIC data analysis can be found in Aleksić et al. (2012a).

The VHE  $\gamma$ -ray differential energy flux can be also well described by an LP:

$$\frac{dF}{dE} = f_0 \left( \frac{E}{200 \text{ GeV}} \right)^{-a-b \log \frac{E}{200 \text{ GeV}}}, \quad (5)$$

where the parameters are given by a flux normalization constant at 200 GeV of  $f_0 = (5.12 \pm 0.27_{\text{stat}} \pm 1.18_{\text{sys}}) \times 10^{-10} \text{ cm}^{-2} \text{ s}^{-1} \text{ TeV}^{-1}$ ,  $a = (3.83 \pm 0.10_{\text{stat}})$  and  $b = (2.09 \pm 0.41_{\text{stat}})$ , the systematic uncertainty on the spectral index is estimated to be  $\pm 0.15$ . The goodness of the fit is given by  $\chi^2/\text{ndf} = 1.8/5$  with a probability  $P = 0.9$ . An LRT shows that for the observed VHE differential spectrum a power law with an exponential cut-off and an LP models are preferred with respect to a simple power law with significances of  $5.8\sigma$  and  $5.9\sigma$ , respectively.

## 4 THE INTRINSIC VHE $\gamma$ -RAY SPECTRUM AND THE ROLE OF THE EBL

### 4.1 Origin of the curvature

The VHE  $\gamma$ -ray spectrum is attenuated by the EBL, as described by equation (1). The optical depth ( $\tau$ ) depends on the redshift of the VHE emitter and the energy of the  $\gamma$ -ray. In order to reconstruct the intrinsic spectrum emitted by a blazar, the redshift and the assumption of an EBL model is required. In the case of PG 1553+113, the uncertainty on the redshift prevents a precise estimation of the intrinsic spectrum. We adopt the optical lower limit from Danforth et al. (2010),  $z = 0.4$ , to study the EBL absorption effect in the observed spectrum, represented in Fig. 3.

**Table 2.** VHE differential energy spectra observed during the 2012 flare. First column represents the energy interval, the second the energy centre of each bin, the second the measured flux after unfolding and the last column is the flux uncertainty.

Energy bin (GeV)	Energy (GeV)	Flux ( $\text{TeV}^{-1} \times \text{cm}^{-2} \times \text{s}^{-1}$ )	Flux uncertainty ( $\text{TeV}^{-1} \times \text{cm}^{-2} \times \text{s}^{-1}$ )
71.2–93.4	81.5	$7.90 \times 10^{-9}$	$0.83 \times 10^{-9}$
93.4–91.1	106.8	$3.69 \times 10^{-9}$	$0.34 \times 10^{-9}$
91.1–160.5	139.9	$1.99 \times 10^{-9}$	$0.18 \times 10^{-9}$
160.5–210.5	183.2	$6.69 \times 10^{-10}$	$0.73 \times 10^{-10}$
210.5–275.9	239.9	$2.55 \times 10^{-10}$	$0.35 \times 10^{-10}$
275.9–361.8	314.0	$7.10 \times 10^{-11}$	$1.36 \times 10^{-11}$
361.8–474.3	410.6	$2.23 \times 10^{-11}$	$0.62 \times 10^{-11}$
474.3–621.9	536.6	$4.68 \times 10^{-12}$	$2.65 \times 10^{-12}$

The curvature measured in the observed VHE spectrum can have different origins: intrinsic electron spectrum curvature, intrinsic self-absorption, Klein–Nishina suppression and/or EBL absorption. The first hypothesis regarding the energy distribution is not likely in the framework of the SED modelling as discussed in Section 5. The assumption of the robust lower limit given by Danforth et al. (2010) allows us to test the possible contribution of intrinsic effect and EBL attenuation.

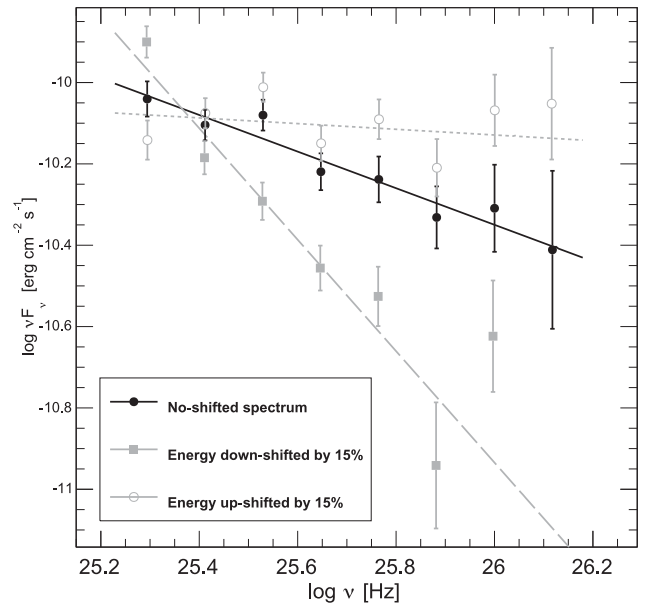
Two possible scenarios can be envisioned considering the possible intrinsic absorption due to pair production within the source. If the  $\gamma$ -ray emission is produced within the broad line region (BLR) populated with optical-UV photons, a softening of the spectrum around tens of GeV would be expected (e.g. Liu & Bai 2006; Reimer 2007; Tavecchio & Mazin 2009). This is typically the case for flat spectrum radio quasars showing strong optical emission lines, although usually weak for BL Lacs. In the far dissipation scenario (e.g. Sikora, Moderski & Madejski 2008), where the emission of  $\gamma$ -rays is assumed to be outside of BLR, the seed photons would come from the IR torus producing a softening in the spectrum at energies typically higher than 1 TeV. None of these scenarios predict intrinsic absorption between 70 and 620 GeV, especially from BL Lac objects with weak BLR emission.

The high flux of the source reached during the flare state allowed a high precision measurement of its spectrum. In addition, the spectrum extends to lower energies than previous measurements performed during lower flux states (Aleksić et al. 2012b). Despite the quality of the data and the high state of the source, no significant  $\gamma$ -ray emission was detected above 620 GeV, in agreement with previous measurements and with the  $\gamma$ -ray absorption expected by the state-of-the-art EBL models given the redshift limits. According to present generation of EBL models (Franceschini et al. 2008; Kneiske & Dole 2010; Domínguez et al. 2011; Gilmore et al. 2012), the observations during the flare reach an optical depth of  $\tau \sim 3$ , which corresponds to  $\sim 95$  per cent photon absorption.

While the observed spectrum shows clear curvature, we find that the spectrum corrected by the EBL effect assuming  $z = 0.4$  can be well described by a simple power law:

$$\frac{dF}{dE} = f_0 \left( \frac{E}{200 \text{ GeV}} \right)^{-\Gamma}, \quad (6)$$

whose parameters using the Franceschini et al. (2008) EBL model are given by a normalization flux at 200 GeV  $f_0 = (9.7 \pm 0.4_{\text{stat}} \pm 2.2_{\text{sys}}) \times 10^{-10} \text{ cm}^{-2} \text{ s}^{-1} \text{ TeV}^{-1}$  and a photon index of  $\Gamma = (2.45 \pm 0.08_{\text{stat}} \pm 0.15_{\text{sys}})$ . The probability of the fit is  $P = 0.9$  ( $\chi^2/\text{ndf} = 2.2/6$ ). The EBL-corrected spectrum

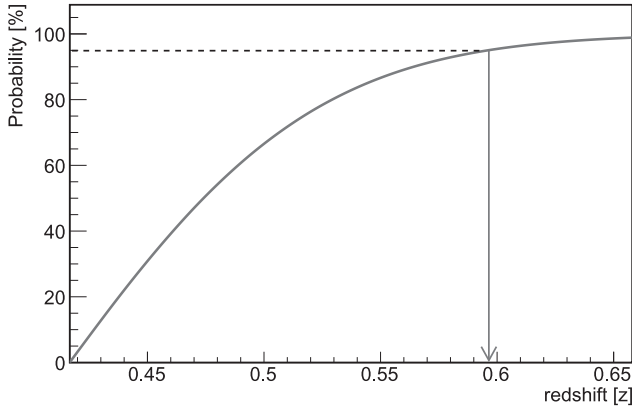


**Figure 4.** MAGIC SED EBL-corrected with Franceschini et al. (2008) model by assuming  $z = 0.4$ . The no-shifted spectrum is represented by the black circles. The solid grey squares show the spectrum considering a shift to lower energies by 15 per cent, and the grey open circles represent the spectrum accounting for a shift to higher energies by 15 per cent.

is shown as green squares in Fig. 3, while the green shaded area represents the uncertainty when assuming different EBL models (Franceschini et al. 2008; Kneiske & Dole 2010; Domínguez et al. 2011; Gilmore et al. 2012).

We tested for a possible shift of 15 per cent in the energy scale due to the uncertainty in the energy measurement (Aleksić et al. 2012a). This was done by performing an eventwise shift in the data while leaving the Monte Carlo (MC) simulations (which are used to determine the energy of each event) unchanged. This simulates a data/MC mismatch, which could occur for numerous reasons including imperfect atmospheric conditions. Both energy-shifted spectra (see Fig. 4), towards lower and higher, are compatible with a power-law fit with  $\chi^2/\text{ndf} = 6.7/5$  and  $\chi^2/\text{ndf} = 8.1/6$ , respectively. The shift to lower energies results in a steepening of the intrinsic spectrum (EBL-corrected according to Franceschini et al. 2008 model assuming  $z = 0.4$ ) with a spectral index  $\Gamma = 3.37 \pm 0.12_{\text{stat}}$ , while the shift to higher energies results in an intrinsic VHE  $\gamma$ -ray spectral index of  $\Gamma = 2.07 \pm 0.08_{\text{stat}}$ .





**Figure 5.** Constraint on the redshift of PG 1553+113 with the LRT when comparing the hypothesis of a power-law fit and a positive curved power law (concave upwards) fit for different distances of the source, using the EBL model from Franceschini et al. (2008).

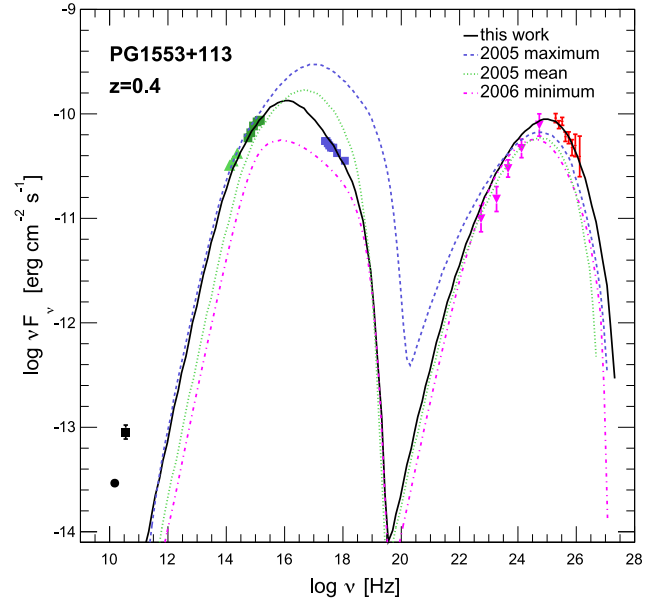
The fact that the EBL-corrected VHE spectrum, assuming as a redshift the robust optical lower limit (Danforth et al. 2010), is compatible with a simple power law suggests that the curvature measured in the observed spectrum is very likely due to the interaction of the VHE photons with the EBL.

#### 4.2 Redshift estimates

An upper limit on the redshift of PG 1553+113 was estimated by excluding the presence of a pile-up at high energies in the intrinsic VHE  $\gamma$ -ray spectrum. This approach is based on the underlying assumption that HE and VHE spectra connect smoothly and form the second peak of the SED. If the peak is located between 10 and 100 GeV, as usually observed in TeV blazars, then a break between the HE and VHE spectra is expected, with VHE spectral slope softer than the HE slope. In the extreme case that the peak is located at higher energies, i.e. at some TeV, we expect that the HE and VHE spectra connect smoothly and exhibit the same spectral slope. A harder slope at VHE than at HE would imply the presence of an additional component in the SED, which is in general not expected, as discussed in Abramowski et al. (2013b). Therefore, the redshift at which the two slopes equal after the correction for EBL absorption can be considered as an upper limit on the source distance under the assumption that there is no additional component. To find the upper limit, an LRT is performed. This test is used to evaluate the hypothesis of evidence of a break in the intrinsic spectrum, as proposed in Mazin & Goebel (2007). The hypothesis of a simple power-law fitting the EBL-corrected spectrum is compared with that of a curved power law, which can fit better the possible pile-up. For the PG 1553+113 data used in this work, the resulting probability is plotted in Fig. 5. Above a redshift  $\sim 0.42$ , a curved fit with positive curvature, which describes the pile-up, start to describe better the data than a simple power law. At redshift  $z \geq 0.60$ , a curved fit with positive curvature is preferred to a simple power-law fit at the 95 per cent confidence level. Therefore, the assumption that there is no spectral pile-up at high energies gives an upper limit on the source redshift of  $z < 0.60$  within a 95 per cent confidence level.

#### 5 SPECTRAL ENERGY DISTRIBUTION

Fig. 6 shows the quasi-simultaneous SED observed during the flare state on 2012 April from  $\gamma$ -rays to radio. The VHE band



**Figure 6.** SED of PG 1553+113 during the 2012 April flare state modelled with the one-zone SSC model of Maraschi & Tavecchio (2003). From high to low energies: the EBL-corrected MAGIC spectra using Franceschini et al. (2008) assuming  $z = 0.4$  (red dots, see the text), the *Fermi*-LAT data from MJD 55959–56088 (pink triangles), *Swift*-XRT (purple squares) and *Swift*-UVOT (green squares) data from MJD 56045 (good representation of the X-ray and optical-UV state during the VHE flare), IR data from REM (green triangles) from MJD 56047, Metsähovi (black square) from MJD 55994 (single detection) and mean radio observation in the period MJD 55959–56088 from OVRO (black circle). For comparison, the SSC models for previous source states (Aleksić et al. 2012b) have been plotted in coloured dashed lines.

is represented by the MAGIC observations during from MJD 56037–56043. The HE differential energy spectrum was derived using *Fermi*-LAT data, which covers the time interval from MJD 56030–56088. Because no variability was detected at HE by the LAT, we used a longer time interval to improve the statistics. The X-ray spectrum shown in Fig. 6 represents the data collected by *Swift*-XRT on MJD 56045. The optical-UV data is the *Swift*-UVOT observation from the same day. The IR flux is estimated from REM telescope observations made on MJD 56046. At 37 GHz, the single detection of the source by Metsähovi on MJD 55994 is shown. The radio flux at 15 GHz measured by OVRO is compatible with a steady emission; therefore, for the SED shown in Fig. 6, the mean flux from the period MJD 56037–56043 has been used.

The SED of PG 1553+113 data during the flare state has been modelled by using a one-zone SSC model (Maraschi & Tavecchio 2003). The emitting region is assumed to be spherical and populated by relativistic electrons. The electron spectrum is assumed to be a smoothed broken power law as a function of the energy (electron Lorentz factor) between  $\gamma_{\min}$  and  $\gamma_{\max}$  and break at  $\gamma_b$ :

$$N(\gamma) = K\gamma^{-n1} \left(1 + \frac{\gamma}{\gamma_b}\right)^{n1-n2}, \quad (7)$$

where  $K$  is the normalization factor, and  $n1$  and  $n2$  the spectral indices before and after the break. The region is filled with a tangled magnetic field and moves out of the jet with a given bulk Lorentz factor ( $\Gamma$ ). The observable effect of bulk Lorentz factor depends on the viewing angle of the jet, which is taken into account in the Doppler factor ( $\delta$ ) used for the SED modelling. According

**Table 3.** One-zone SSC model parameters of the SED fit during the flare state on 2012. The following quantities are reported: the minimum, break, and maximum Lorentz factors and the low and high energy slope of the electron energy distribution, the magnetic field intensity, the electron density, the radius of the emitting region and the Doppler factor. The derived luminosities are shown in Table 4.

Model	$\gamma_{\min}$ ( $10^3$ )	$\gamma_b$ ( $10^4$ )	$\gamma_{\max}$ ( $10^5$ )	$n_1$	$n_2$	$B$ (G)	$K$ ( $\text{cm}^{-3}$ )	$R$ ( $10^{16}\text{cm}$ )	$\delta$
This work	3.7	3.6	8.0	1.60	3.83	0.045	19.5	6.0	40
Maximum <sup>a</sup>	1.0	3.0	5.2	2.00	3.75	0.800	$3.8 \times 10^3$	1.0	35
Minimum <sup>a</sup>	5.0	1.3	4.1	2.00	3.55	0.200	$25.0 \times 10^3$	1.0	35
Mean <sup>a</sup>	1.5	3.2	2.2	2.00	4.00	0.500	$5.4 \times 10^3$	1.0	35

Notes. The models marked as <sup>a</sup> correspond to previous activity states of the source (see Aleksić et al. 2012b) and are shown for comparison.

**Table 4.** Luminosities derived from the one-zone SSC model of the SED during the flare state on 2012. The following quantities are reported: the kinetic energy of the electrons, (cold) protons (assuming one proton per emitting electron), and magnetic field, and the frequency of the synchrotron peak. The model parameters are shown in Table 3.

Model	$L_{\text{kin}(e)}$ ( $10^{45} \text{ erg s}^{-1}$ )	$L_{\text{kin}(p)}$ ( $10^{44} \text{ erg s}^{-1}$ )	$L_B$ ( $10^{43} \text{ erg s}^{-1}$ )	$\log 10(\nu_{\text{syn}})$
This work	2.18	1.49	5.83	16.1
Maximum <sup>a</sup>	0.52	0.6	39.2	17.0
Minimum <sup>a</sup>	0.52	0.6	2.5	15.9
Mean <sup>a</sup>	0.52	0.6	15.3	16.7

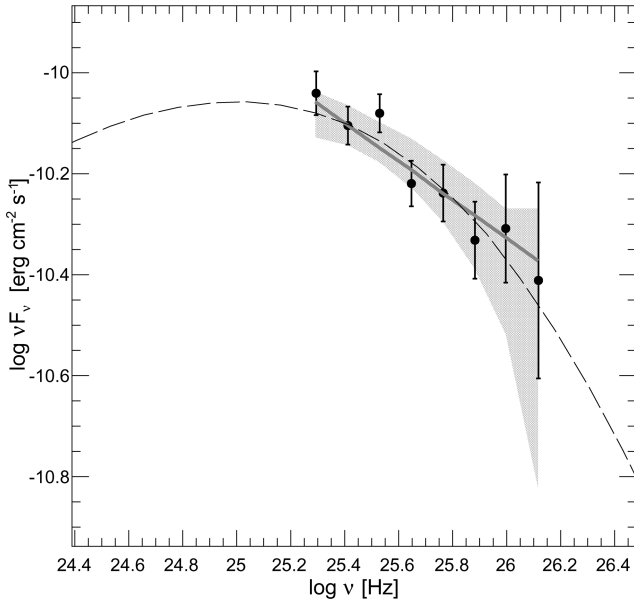
Notes. The models marked as <sup>a</sup> correspond to previous activity states of the source (see Aleksić et al. 2012b) and are shown for comparison.

to the SSC model, the electrons emit synchrotron radiation due to their interaction with the magnetic field creating a low-energy photon field, which can in turn interact with the same population of electrons via inverse Compton (IC), producing the high energy emission. The synchrotron component considered is self-absorbed below  $10^{11}$  Hz and thus cannot reproduce the radio emission. This emission is likely from the superposition of multiple self-absorbed jet components (Königl 1981).

The parameters used for the modelling as well as those of SSC models reproducing previous observations (Aleksić et al. 2012b) of the source in different states, for comparison purpose, can be found in Table 3. During the strong flare in 2012, the magnetic field and the electron population normalization are significantly smaller than in previous states of the source, while the emitting region size is six times larger. However, as given by the causality relation  $R < (c t \delta)/(1 + z)$ , the allowed flux variability time-scale is  $\sim 19$  h (assuming  $z = 0.4$ ), which is compatible with the variability detected in the source as shown in Fig. 1. The IC energy peak moved to higher energies with respect to previous observations and more energetic particles were involved, requiring a larger  $\gamma_{\max}$  parameter for the modelling as shown in Table 3. This could point to different origins of the high states of the source. The derived luminosities from the SSC modelling are shown in Table 4. The electron and (cold) proton luminosities are higher than previous high states. For the luminosities calculation, one proton per emitting electron was assumed. It is worth noting that the MWL data of the previous states of the source (Aleksić et al. 2012b) used for comparison were not simultaneous. Moreover, due to the degeneracy of the model parameters the best model is not unique and other parameters could also reproduce the SED. Therefore, strong conclusions cannot be derived from the comparison with previous modelling of the source.

As shown in Fig. 6, the increasing part of both SED bumps shows less variability when compared with the decreasing part. This fact is also in agreement with the light curve discussion on Section 3.1: while X-rays and VHE  $\gamma$ -rays show an increase of the flux in 2012, the emission in the HE band is compatible with a constant flux. The high variability found in X-rays and VHE  $\gamma$ -rays suggests that the flaring activity of this source is driven by the most energetic electrons. Moreover, as discussed previously, the SSC model gives a lower magnetic field with respect to previous models, which implies a longer synchrotron cooling time-scale. This is in agreement with the displacement of the synchrotron peak to higher frequencies, as well as with the higher variability in the high energy component of both peaks.

As shown in Fig. 6, the IC peak of the SED is close to the VHE band. Therefore, curvature would be expected in the intrinsic VHE SED due to the distribution of the relativistic electrons, within the one-zone SSC framework (as mentioned in Section 4). To test if our observations are sensitive enough to detect the expected intrinsic SED curvature, we simulate the MAGIC response assuming the intrinsic emission given by the best SSC modelling of the MWL data shown in Fig. 6. We simulate intrinsic VHE SEDs assuming the same frequency binning and relative errors as in the observed VHE spectral points (only statistical uncertainties have been taken into account). The result of 10 000 realizations are shown in Fig. 7, and are represented by the grey shaded area. Despite the simulated SEDs having, by construction, an evident curvature, 99.2 per cent of the realizations are well described by a simple power law. To be conservative, we require a  $p$ -value of the individual fits  $> 0.9973$ , which allows us to exclude spectral curvature in the simulated spectra at a  $3\sigma$  confidence level. The mean probability of a simple power-law fit is  $P = 0.44 \pm 0.28$  with a mean photon index of  $2.38 \pm 0.10$ . We therefore conclude that the sensitivity of our VHE



**Figure 7.** PG 1553+113 VHE SED. The grey shaded area represents the simulated MAGIC response assuming the intrinsic emission given by the best-fitting SSC model to the MWL data shown in Fig. 6. The grey solid line represents the mean power-law fit of the 10 000 realizations of the toy-MC. The black circles denote the PG 1553+113 VHE spectrum EBL-corrected with Franceschini et al. (2008) model assuming  $z = 0.4$ . The dashed black line represents the best-fitting SSC model from Fig. 6.

measurements do not allow the detection of an intrinsic curvature in the SSC framework and the EBL model from Franceschini et al. (2008).

## 6 CONCLUSIONS

In this paper, we have presented the highest flux state ever detected from the blazar PG 1553+113 in VHE  $\gamma$ -rays. The flare was detected at VHE by the MAGIC telescopes and monitored in HE  $\gamma$ -rays by *Fermi*-LAT, in X-rays by *Swift*-XRT, in optical-UV by *Swift*-UVOT, in infrared by REM, and in radio by Metsähovi and OVRO. While clear variability has been found in both the VHE and X-ray bands, the HE  $\gamma$ -ray flux is compatible with constant emission.

The observed VHE spectrum shows clear curvature and a simple power-law fit is rejected with a confidence level of  $4.7\sigma$ . It is well fitted by a power law with an exponential cut-off or an LP. This curvature most likely originates from EBL absorption if the distance to the source is between the redshift limits measured by Danforth et al. (2010,  $0.4 < z < 0.58$ ). If the real redshift of this source is higher than these limits, the effect would be a hardening of the spectrum or the (unexpected) presence of a pile-up in the intrinsic spectrum, which would denote that either the EBL models predict an overestimated EBL level there is a second emission component at high energies or that more exotic physics needs to be invoked as axion-like particles (e.g. De Angelis, Roncadelli & Mansutti 2007; Hooper & Serpico 2007; Sánchez-Conde et al. 2009) or the effect of Lorentz invariance violation (e.g. Jacob & Piran 2008).

A redshift upper limit of  $z < 0.60$  at 95 per cent C.L. has been derived using the  $\chi^2$  ratio test (Mazin & Goebel 2007). A quasi-simultaneous SED has been compiled for the flare episode in 2012 April. It can be well modelled by a one-zone SSC model. The comparison with previous flux states of the source reveals that the higher frequency part of each SED bump shows higher variability

than the lower frequency part. This fact points to a scenario where the most energetic electrons play a leading role during the flare episodes of the source.

A detailed study of the MWL behaviour and evolution of the SED will be published in a forthcoming paper.

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