

Dipole model analysis of high precision HERA data

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Abstract

We analyse, within a dipole model, the inclusive DIS cross section data, obtained from the combination of the H1 and ZEUS HERA measurements. We show that these high precision data are very well described within the dipole model framework, which is complemented with a valence quark structure functions. We discuss the properties of the gluon density obtained in this way.

1 Introduction

Many investigations have shown that HERA inclusive and diffractive DIS cross sections are very well described by the dipole models [1–3]. Interest in the dipole description emerge from the fact that dipole picture provides a natural description of QCD reaction in the low- x region. Due to the optical theorem, dipole models allow a simultaneous description of many different physics reactions, like inclusive DIS processes, inclusive diffractive processes, exclusive J/ψ , ρ , ϕ production, diffractive jet production, or diffractive and non-diffractive charm production. In the dipole picture, all these processes are determined by the same, universal, gluon density [4–6].

In the era of the LHC, the precise knowledge of gluon density is very important because the QCD-evolved gluon density determines the cross sections of most relevant physics processes, e.g. Higgs production. Any significant deviation of the predicted cross section from their Standard Model value could be a sign of new physics.

The validity of the dipole approach was experimentally established, a decade ago, by a comparison of the dipole predictions with HERA F_2 and diffractive data in the low x region [1], [2]. In the meantime, the precision of data obtained from HERA experiments increased substantially. The H1 and ZEUS experiments have combined their inclusive DIS cross sections which, due to a substantial reduction of systematic measurements errors, led to an increase of precision by about a factor two [7]. In the same way the quality of the inclusive

charm data was substantially improved [8]. Finally, recently, the exclusive J/ψ production was much more precisely measured [9]. All these reactions were used in the past to establish the dipole approach. It is therefore interesting to re-evaluate these reactions because the dipole picture provides a somewhat different approach to the gluon density than the usual pdf approach. In the usual pdf approach the gluon density contributes to F_2 mainly through the evolution of the sea quarks, the direct gluon contribution is only of the order of a few percent. On the other hand, in the dipole models the gluon density is directly connected to the sea quarks. In the pdf scheme the evolution is evaluated in the collinear approximation whereas the dipole approach uses the k_T factorization.

The direct connection between the dipole production and gluon density is particularly clearly seen in the exclusive J/ψ production, which was therefore proposed as a testing ground of the properties of the gluon density [10]. Presently, the exclusive J/ψ production is precisely measured in heavy ion collisions at RHIC and LHC. These measurements combined with their dipole analysis can become a new source of information about the gluonic structures of nuclei [12, 13].

Another important application of the dipole description is the investigation of the gluonic high density states. These can be characterized by the degree by which a dipole is absorbed or multiply scattered in such states. The states with the highest gluon densities are produced today in the high energy heavy ion scattering at RHIC and LHC. This is now a very lively field of saturation investigation [14, 15].

The aim of this paper is to investigate the additional information which is contained in the new, combined HERA data. The most precise data were obtained in the region of higher Q^2 's (Q^2 from 3.5 to $O(10000)$ GeV^2), where the DGLAP evolution is known to describe data very well. Therefore, as discussed below, in this investigation we use the so called BGK dipole model, because it uses the DGLAP evolution scheme.

This paper concentrates first on the *inclusive* DIS measurements in the low x region. Here, the contribution of the valence quarks is small, below 7%, and has therefore been neglected until now. However, the combined H1 and ZEUS HERA data achieve however a precision of about 2%, so the contribution of the valence quarks can no longer be neglected. The present paper addresses the question to what extent the contribution of the valence quark and the dipoles are compatible with each others. To do so we use the HARAFitter framework [16] which allows to treat consistently QCD evolution together with the valence quark and dipoles contributions.

The paper is organized as follows: in Section 2 we recall the main properties of the dipole approach and review various models in order to motivate our choice. In Section 3 we discuss the results of fits and in Section 4 we compare the fits with data. Section 5 contains the summary.

2 Dipole models

In the dipole picture the deep inelastic scattering is viewed as a two stage process; first the virtual photon fluctuates into a dipole, which consists of a quark-antiquark pair (or a $q\bar{q}g$ or $q\bar{q}gg \dots$ system) and in the second stage the dipole interacts with the proton [17] , [18–24]. Dipole denotes a quasi-stable quantum mechanical state, which has a very long life time ($\approx 1/m_p x$) and a size r , which remains unchanged during scattering. The wave function Ψ determines the probability to find a dipole of size r within a photon. This probability depends on the value of external Q^2 and the fraction of the photon momentum carried by the quarks forming the dipole, z . Neglecting the z dependence, in a very rough approximation, $Q^2 \sim 1/r^2$.

The scattering amplitude is a product of the virtual photon wave function, Ψ , with the dipole cross section, σ_{dip} , which determines a probability of the dipole-proton scattering. Thus, within the dipole formulation of the γ^*p scattering

$$\sigma_{T,L}^{\gamma^*p}(x, Q^2) = \int dr^2 \int dz \Psi_{T,L}^*(Q, r, z) \sigma_{\text{dip}}(x, r) \Psi_{T,L}(Q, r, z), \quad (2.1)$$

where T, L denotes the virtual photon polarization and $\sigma_{T,L}^{\gamma^*p}$ the total inclusive DIS cross section.

Several dipole models have been developed to test various aspects of the data. They vary due to different assumption made about the physical behavior of dipole cross sections. In the following we will shortly review them and motivate our the choice of the model used for the present investigation.

2.1 GBW model

The dipole model became an important tool in investigations of deep-inelastic scattering due to the initial observation of Golec-Biernat and Wüesthoff (GBW) [1], that a simple ansatz for the dipole cross section was able to describe simultaneously the total inclusive and diffractive cross sections.

In the GBW model the dipole-proton cross section σ_{dip} is given by

$$\sigma_{\text{dip}}(x, r^2) = \sigma_0 \left(1 - \exp \left[-\frac{r^2}{4R_0^2(x)} \right] \right), \quad (2.2)$$

where r corresponds to the transverse separation between the quark and the antiquark, and R_0^2 is an x dependent scale parameter which has a meaning of saturation radius, $R_0^2(x) = (x/x_0)^{\lambda_{GBW}}$. The free fitted parameters are: the cross-section normalisation, σ_0 , as well as x_0 and λ_{GBW} . In this model saturation is taken into account in the eikonal approximation and the saturation radius is intimately related to the gluon density in the transverse plane, see below. The exponent λ_{GBW} determines the growth of the total and diffractive cross section with decreasing x . For dipole sizes which are large in comparison to the saturation radius,

R_0 , the dipole cross section saturates by approaching a constant value σ_0 , i.e. saturation damps the growth of the gluon density at low x .

The GBW model provided a good description of data from medium Q^2 values (≈ 30 GeV²) down to low Q^2 (≈ 0.1) GeV²). Despite its success and its appealing simplicity the model has some shortcomings; in particular it describes the QCD evolution by a simple x dependence, $\sim (1/x)_{BGW}^\lambda$, i.e the Q^2 dependence of the cross section evolution is solely induced by the saturation effects. Therefore, it does not match with DGLAP QCD evolution, which is known to describe data very well from $Q^2 \approx 4$ GeV² to very large $Q^2 \approx 10000$ GeV².

2.2 BGK model

The evolution ansatz of the GBW model was improved in the model proposed by Bartels, Golec-Biernat and Kowalski, (BGK) [2], by taking into account the DGLAP evolution of the gluon density in an explicit way. The model preserves the GBW eikonal approximation to saturation and thus the dipole cross section is given by

$$\sigma_{\text{dip}}(x, r^2) = \sigma_0 \left(1 - \exp \left[-\frac{\pi^2 r^2 \alpha_s(\mu^2) x g(x, \mu^2)}{3\sigma_0} \right] \right). \quad (2.3)$$

The evolution scale μ^2 is connected to the size of the dipole by $\mu^2 = C/r^2 + \mu_0^2$. This assumption allows to treat consistently the contributions of large dipoles without making the strong coupling constant, $\alpha_s(\mu^2)$, un-physically large.

The gluon density, which is parametrized at the starting scale μ_0^2 , is evolved to larger scales, μ^2 , using LO or NLO DGLAP evolution. We consider here three forms of the gluon density:

- the *soft* ansatz, as used in the original BGK model

$$xg(x, \mu_0^2) = A_g x^{-\lambda_g} (1-x)^{C_g}, \quad (2.4)$$

- the *soft + hard* ansatz

$$xg(x, \mu_0^2) = A_g x^{-\lambda_g} (1-x)^{C_g} (1 + D_g x + E_g x^2), \quad (2.5)$$

- the *soft + negative* gluon

$$xg(x, \mu_0^2) = A_g x^{-\lambda_g} (1-x)^{C_g} - A'_g x^{-\lambda'_g} (1-x)^{C'_g}, \quad (2.6)$$

The free parameters for this model are σ_0 , μ_0^2 and the parameters for gluon A_g , λ_g , C_g or additionally D_g , E_g , or A'_g , λ'_g , C'_g . Their values are obtained by a fit to the data. The fit results were found to be independent on the parameter C , which was therefore fixed as $C = 4$ GeV², in agreement with the original BGK fits.

2.3 IIM model

Although we do not use the IMM (Iancu, Itacura and Mounier) model in this paper we mention it here because it may take better into account the saturation effects than it is the case in the BGK or GBW models. The last models use for saturation the eikonal approximation, whereas the IIM model uses a simplified version of the Balitsky-Kovchegov equation [27]. The explicit formula for σ_{dip} can be found in [3]. We do not use this model because we concentrate here on the higher Q^2 data, which precise description requires an equally precise transition to the DGLAP regime.

2.4 Dipole model with valence quarks

The dipole models are valid in the low- x region where the valence quark contribution is small. Therefore, this contribution was usually neglected which was justified as long as the experimental errors were relatively large. Theoretically, it is very difficult to treat valence quarks inside the dipole framework because, until now, the dipole amplitudes are not well defined in the region of high x . The problem may be solved, in future, by the analytic continuation of the dipole (or BFKL) amplitudes from the low x to the high x region [28]. However, for the purpose of this paper, we propose to take an heuristic approach and just to add the valence quark contribution from the standard pdf's fits to the dipole predictions. In this approach the dipole contribution plays a role of the sea quarks in the standard pdf's. This procedure is justified by the fact that the sea quark contribution disappears at larger x . The HERAFitter project is well suited for this purpose since the dipole model and the valence quarks contributions are a part of the same framework.

3 Results from fits

No	Q_0^2	σ_0	A_g	λ_g	C_g	C	Np	χ^2	χ^2/Np
1	1.1	143.14	1.605	-0.056	5.884	4.0	201	198.17	0.986
3	1.3	123.18	1.589	-0.094	6.937	4.0	201	200.70	0.998
5	1.5	112.44	1.685	-0.109	8.124	4.0	201	202.26	1.006
7	1.7	97.91	1.603	-0.137	8.849	4.0	201	203.55	1.013
9	1.9	90.98	1.624	-0.149	9.696	4.0	201	202.18	1.006

Table 1: BGK fit with valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft gluon*.

In this section we investigate how well the dipole model can describe the new, precise, HERA data which were obtained in the region of $Q^2 > 3.5 \text{ GeV}^2$. Since the quality of data in the region of $Q^2 < 1 \text{ GeV}^2$ was not improved until now we concentrate here on the higher Q^2 region where the valence quark contribution becomes relevant.

3.1 Dipole fits with valence quarks

First, we show that it is possible to combine the dipole and valence quark contributions and obtain a good fit to the data. For the purpose of this investigation we choose the BGK model because it uses the DGLAP evolution. The fits were performed within the HERAFitter framework, i.e.; the QCD evolution is the same as in the standard HERAFitter pdf fits. The results of the BGK fit, with valence quarks, are shown in Table 1. The fit is performed in the low x range, $x < 0.01$, for various μ_0^2 values. The value of μ_0^2 plays a role of the starting scale of the QCD evolution which is usually denoted by Q_0^2 in the pdf fits. N_p denotes the number of measured values of the reduced cross section, σ_r , which were used in the fit. The parameters σ_0 of the dipole model and the parameters for gluon A_g , λ_g , C_g are obtained from the fit at a given value of Q_0^2 (in GeV^2). The value of the parameter C was fixed, as explained above.

The table shows that the BGK model with valence quarks taken from the usual HERAFitter pdf fit, is describing the precise HERA data very well for all Q_0^2 value. The fit quality improve slightly with diminishing Q_0^2 . This could indicate that HERA data in the low range of $Q^2 \sim 3.5 \text{ GeV}^2$, retain some sensitivity to the saturation effects. In the BGK model the saturation effects increases with decreasing Q_0^2 value.

In the Table 2 we show results of the standard HERAPDF fits. They are performed in the same Q^2 range as the dipole fits but in the full x range. The full x range is necessary to fix the contribution of valence quarks.

No	Q_0^2	HF Scheme	χ^2	N_p	χ^2/N_p
1	1.1	RT	604.64	592	1.021
3	1.3	RT	586.33	592	0.990
5	1.5	RT	579.72	592	0.979
7	1.7	RT	576.76	592	0.974
9	1.9	RT	575.08	592	0.971

Table 2: HERAPDF fit for σ_r for H1ZEUS-NC-(e+p), H1ZEUS-NC-(e-p) and H1ZEUS-CC-(e+p), H1ZEUS-CC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 1.0$.

No	Q_0^2	HF Scheme	χ^2	N_p	χ^2/N_p
1	1.1	RT	472.52	550	0.859
3	1.3	RT	469.80	550	0.854
5	1.5	RT	469.06	550	0.853
7	1.7	RT	468.67	550	0.852
9	1.9	RT	468.34	550	0.852

Table 3: HERAPDF fit for σ_r for H1ZEUS-NC-(e+p), H1ZEUS-NC-(e-p) and H1ZEUS-CC-(e+p), H1ZEUS-CC-(e-p) data in the range $Q^2 \geq 8.5$ and $x \leq 1.0$.

No	Q_0^2	σ_0	A_g	λ_g	C_g	C	N_p	χ^2	χ^2/N_p
1	1.1	91.60	2.227	-0.022	9.322	4.0	162	131.78	0.813
3	1.3	83.393	2.047	-0.069	10.019	4.0	162	132.10	0.815
5	1.5	77.121	1.969	-0.098	10.825	4.0	162	132.23	0.816
7	1.7	71.975	1.922	-0.120	11.538	4.0	162	132.88	0.820
9	1.9	69.128	1.897	-0.135	12.175	4.0	162	132.03	0.815

Table 4: Dipole model BGK fit with valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 8.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft gluon*.

Table 2 shows a very good agreement with data of the standard pdf fit. The agreement is similar as in the dipole fits, if corrected for the number of points and the number of free parameter, which is $N_{free} = 10$ for the HARAPDF fit and $N_{free} = 4$ in case of the BGK fit with the *soft* gluon assumption. In difference to the dipole fits, the quality of the HERAPDF fit is deteriorating with decreasing Q_0^2 scale.

Table 3 and 4 show HERAPDF and BGK dipole fits in the higher Q^2 range, $Q^2 > 8.5$ GeV². We see that the quality of fits clearly improves in the higher Q^2 region. In case of HARAPDF fit the χ^2/N_p improves from 0.97 to 0.85 and in case of the BGK fit from ~ 1.0 to 0.82. Moreover, the BGK fits do not show any dependence from the starting scale, Q_0^2 . The HERAPDF fits do still show some slight deterioration with decreasing Q_0^2 but the effect is much smaller than seen in Table 2.

In Fig. 1 we show a comparison the gluon density obtained in the fits with valence quarks and compare it to the gluon density obtained in the HERAPDF fit. We see that the two gluon densities, at NLO, differ at smaller scales but then start to approach each other at higher scales. It is interesting to observe that the convergence of the two gluon densities is much slower in LO, Fig. 2.

3.2 Fits with alternative forms of the gluon density

In this section we investigate whether the more involved forms of the gluon density, eq.(2.5) and eq.(2.6), can improve the data description. In table 5 and 6 we show the fit results for the fits with *soft + hard* gluon of eq.(2.5), in the lower $Q^2 > 3.5$ GeV² and higher $Q^2 > 8.5$ GeV² regions. We observe that the fit quality improves significantly by adding a "hard" component, $D_g x + E_g x$, to a classic soft gluon of eq.(2.4). The value of χ^2 diminishes by about $\Delta\chi^2 \approx 20$ for $Q^2 > 3.5$ GeV² and by about $\Delta\chi^2 = 15$ for $Q^2 > 8.5$ GeV², which is a much larger drop than the increase of the parameter number (just by 2).

In Table 7 we show the fit results for the fits with the *soft + negative* gluon of eq.(2.6). The fit in the lower Q^2 range is not significantly improved by the addition of the negative gluon term. In the higher Q^2 range, $Q^2 > 8.5$ GeV², the fit improves somewhat, although not so clearly as in the "hard" case.

No	Q_0^2	σ_0	A_g	λ_g	C_g	D_g	E_g	χ^2	χ^2/Np
1	1.1	217.09	1.976	-0.012	22.502	-35.364	1339.3	181.34	0.930
2	1.3	181.82	1.847	-0.059	21.597	-25.051	1030.3	180.80	0.927
3	1.5	165.17	1.871	-0.082	24.623	-23.630	1237.7	180.80	0.927
4	1.7	147.12	1.903	-0.099	26.720	-20.584	1310.2	181.70	0.932
5	1.9	132.26	1.948	-0.111	28.211	-18.008	1322.4	180.81	0.927

Table 5: Dipole model BGK fit with valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft + hard gluon*. $Np = 201$ and $C=4.0$ GeV².

No	Q_0^2	σ_0	A_g	λ_g	C_g	D_g	E_g	χ^2	χ^2/Np
1	1.1	254.97	2.524	-0.027	24.857	-46.523	1639.8	117.34	0.752
2	1.3	154.25	2.171	-0.041	13.728	-20.261	340.97	121.79	0.781
3	1.5	292.89	2.358	-0.034	31.168	-50.312	2585.8	115.51	0.740
4	1.7	221.52	2.483	-0.051	34.010	-44.156	2630.6	115.78	0.742
5	1.9	174.46	2.490	-0.070	35.347	-37.706	2499.7	116.18	0.745

Table 6: Dipole model BGK fit with valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 8.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft + hard gluon*. $Np = 162$ and $C = 4.0$ GeV².

Q_0^2	Q^2	σ_0	A_g	λ_g	C_g	A'_g	B'_g	C'_g	χ^2	χ^2/Np
1.9	3.5	115.09	0.874	-0.253	3.669	-0.014	-0.606	25.0	200.49	1.028
1.9	8.5	111.94	0.799	-0.290	3.922	0.020	-0.642	25.0	119.48	0.766

Table 7: Dipole model BGK fit with valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data. NLO fit. RT HF Scheme. *Soft + negative gluon*. $C = 4.0$ GeV², $N_p = 201$ for $Q^2 > 3.5$ GeV² and $N_p = 162$ for $Q^2 > 8.5$ GeV²

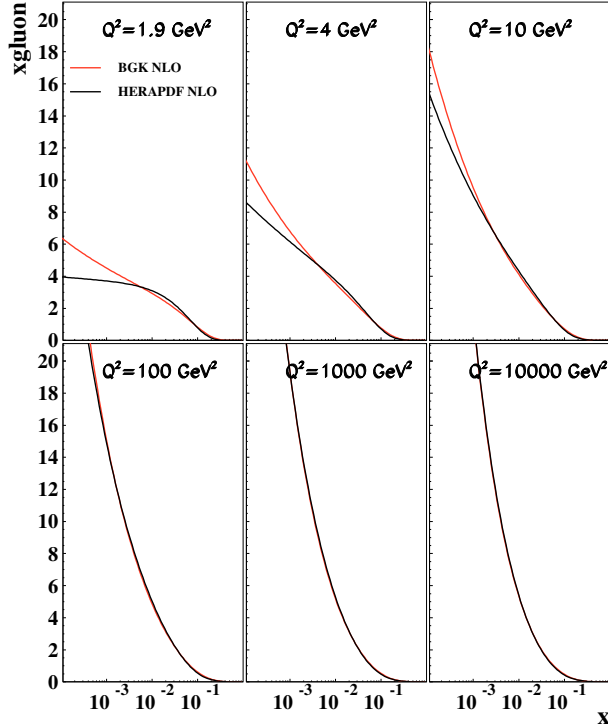


Figure 1: Comparison between the dipole (soft) and HERAPDF gluon in NLO.

3.3 Fits without or with fitted valence quarks

To better understand the meaning of the fits which are using alternative forms of the gluon density we performed also fits without valence quarks, and with valence quarks fitted to data. In Table 8 and Table 9 we show fits performed without valence quarks for the soft and soft+hard forms of the gluon density in the region of $Q^2 > 3.5 \text{ GeV}^2$.

No	Q_0^2	Q^2	σ_0	A_g	λ_g	C_g	χ^2	χ^2/Np
1	1.9	3.5	115.09	2.038	-0.097	4.969	197.83	1.004

Table 8: Dipole model BGK fit without valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft gluon*. $C = 4.0 \text{ GeV}^2$ and $Np = 201$.

The contribution of the valence quarks in the low x region are large enough to be able to determine them in this region only. In Table 10 we show an example of a fit with parameters of the valence quarks fitted to data together with the parameters of the gluon density. The fit is performed for $Q^2 > 3.5 \text{ GeV}^2$, in the low x range, $x < 0.01$. In Table 11 we give, for completeness, the parameters of the valence quarks determined in this way. Note, that the fit with fitted valence quarks is better than the fit with fixed valence quarks of Table 1 and

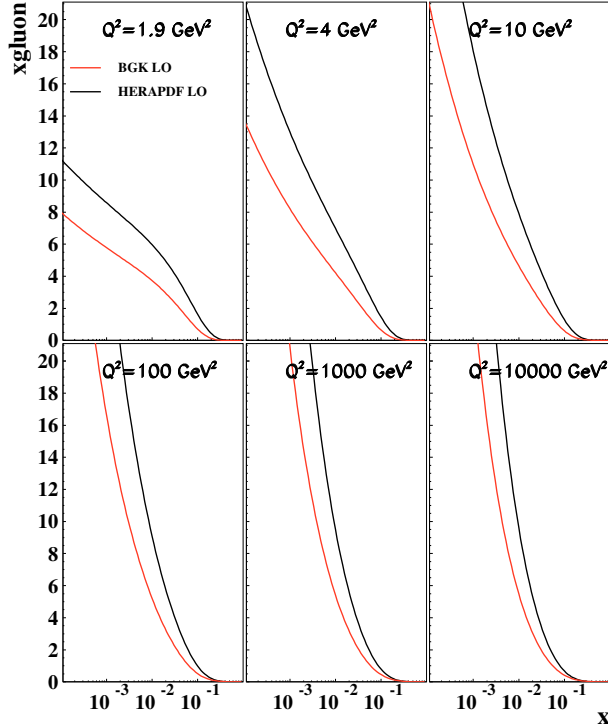


Figure 2: Comparison between the dipole (soft) and HERAPDF gluon in LO.

No	Q_0^2	Q^2	σ_0	A_g	λ_g	C_g	D_g	E_g	χ^2	χ^2/N_p
1	1.9	3.5	119.18	1.970	-0.104	5.001	3.347	-19.340	196.26	1.006

Table 9: Dipole model BGK fit without valence quarks for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft + hard gluon*. $C = 4.0 \text{ GeV}^2$ and $N_p = 201$.

it is also better than the fit without valence quarks, Table 8.

Fig. 3 shows the comparison between the NLO gluon densities determined with the *soft* and *soft + hard assumptions*. The *soft* gluon density is taken from the fit of Table 1. The *soft+hard* gluon density shown on the LHS of Fig. 3 is taken from the fit of Table 5 and was obtained with the fixed valence quark contribution. The RHS of this figure shows the *soft+hard* gluon density obtained from the fit of Table 10. Here, the contribution of valence quarks is fitted to data together with the gluon density. Both fits, of Table 5 and Table 10, have a very similar quality, the form of gluon densities differs, however, at lower scales in the high x region; the one with the fixed valence quarks shows a clear bump around $x \approx 0.1$, the another shows no bump and has a form similar to the *soft* case. In all fits which we performed, the bump in the *soft+hard* gluon density fitted with the fixed valence quarks was always present. independently of the Q^2 cut or the LO or NLO QCD evolution. This bump disappears, however, when the valence quark contribution is fitted. Therefore, we do

not attribute a physical meaning to this bump, especially that it is in the region which is not directly tested by data and it contributes only to the low- x region through the QCD evolution. Nevertheless, its existence emphasizes the necessity of a full fit to the data, i.e. of a fit in which the gluon density is fitted *together* with the valence quarks.

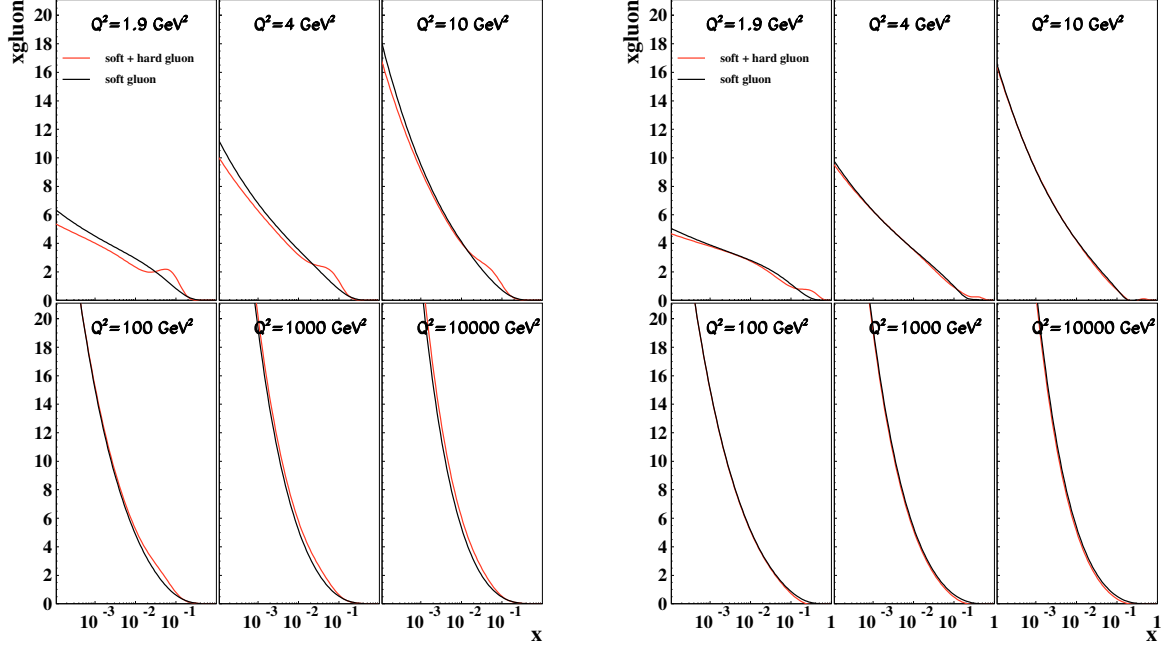


Figure 3: Comparison between the NLO gluon densities determined with the soft and soft + hard assumptions. LHS shows the gluon distribution functions determined with the fixed valence quark contribution. RHS shows the gluon distribution functions determined with the contribution of valence quarks fitted to data in the $x < 0.01$ region.

No	Q_0^2	Q^2	σ_0	A_g	λ_g	C_g	χ^2	χ^2/Np
1	1.9	3.5	88.040	1.766	-0.115	6.747	182.89	0.978

Table 10: Dipole model BGK fit with valence quarks fitted for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft gluon*. $C = 4.0 \text{ GeV}^2$ and $Np = 201$.

4 Comparison with HERA data

In Fig. 4 we show a comparison of the dipole BGK fit with the HERA reduced cross section data. Figure shows an excellent agreement of the fit with data. In Fig. 5 we show a comparison of F_L structure function obtained from the dipole BGK fit with HERA data. In both figures we use the BGK fit of Table 1, with $Q_0^2 = 1.9 \text{ GeV}^2$.

No	A_{uv}	B_{uv}	C_{uv}	E_{uv}	A_{dv}	C_{dv}	C_{Ubar}	A_{Dbar}	B_{Dbar}	C_{Dbar}
1	3.717	0.665	4.652	9.694	2.189	4.291	2.582	0.100	-0.165	2.405

Table 11: Parameters for valence quarks from : dipole model BGK fit with valence quarks fitted for σ_r for H1ZEUS-NC-(e+p) and H1ZEUS-NC-(e-p) data in the range $Q^2 \geq 3.5$ and $x \leq 0.01$. NLO fit. RT HF Scheme. *Soft gluon*. Parameter $cBGK = 4.0$, $N_p = 201$.

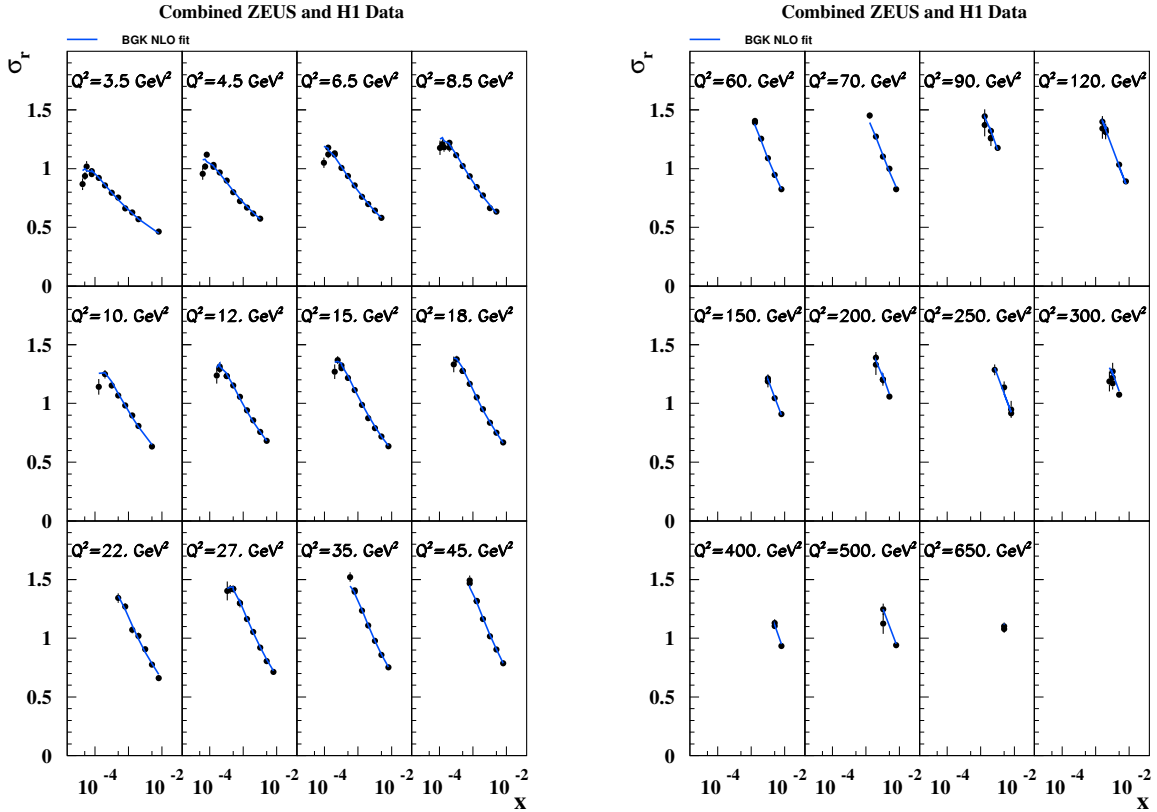


Figure 4: Comparison of the dipole BGK fit of Table 1 with the reduced cross sections of HERA data.

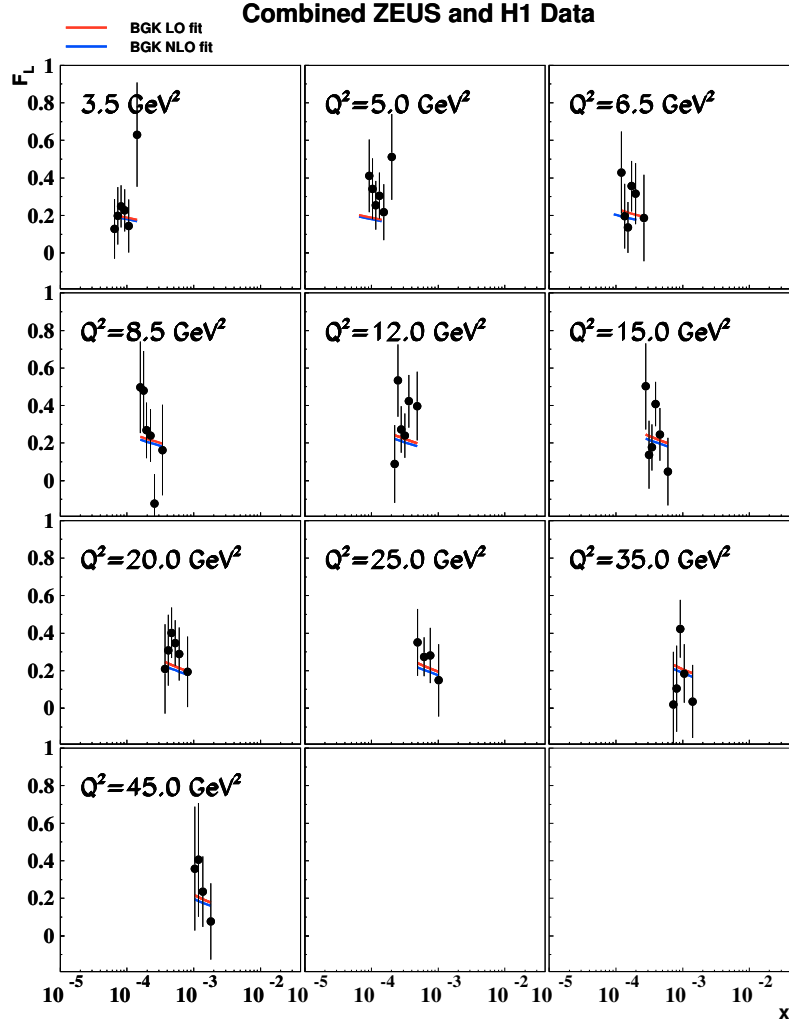


Figure 5: Comparison of F_L structure function obtained from the dipole BGK fit of Table 1 with HERA data.

5 Summary

We have shown that the k_T factorized, DGLAP evolved gluon density, evaluated within the BGK model, describe the combined, precise HERA data in the low- x region, very well. The valence quark contribution added to the dipole model improves the fit significantly. Therefore, for precise dipole evaluations the gluon contribution should be complemented by valence quarks.

The resulting gluon density obtained from fits with fitted valence quarks could be used for the prediction of LHC cross sections, provided that the dipole amplitude, which is now only well defined in the low- x region, can be analytically continued to the high x region [28].

As a byproduct of this investigation we observe that the fits of all dipole and pdf types improve significantly when the Q^2 cut on data is increased from $Q^2 > 3.5$ to $Q^2 > 8.5$ GeV². We have checked this with the dipole model with quarks and without quarks, with various forms of the gluon density, as well as with the standard HERAPDF1.0 fit. The persistence of this effect indicate some shortcomings of the theoretical description; it could be due to the lack of higher order QCD corrections or to saturation effects. We note, that the higher order corrections diminish logarithmically with increasing Q^2 whereas the saturation effects diminish like a power of Q^2 ¹, or faster. In our view, the relatively fast change of χ^2/N_p with the increased Q^2 cut indicates that the effect is due to saturation, at least to large extent. In this way, the increase of precision in HERA data offers a novel testing ground for saturation study in the well measured region above $Q^2 > 3.5$ GeV². The study of this type may become very interesting when, in the near future, the combined HERA I and HERA II data, with yet further increased precision, is published.

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¹The degree of saturation is characterized by the size of the dipole, r_S , which, at a given x , starts to interact multiple times (in about 60% of cases). We recall that the saturation scale, $Q_S^2 = 2/r_S^2$, was determined at HERA as 0.5 GeV² at $x = 10^{-3}$ and as about 1 GeV² at $x = 10^{-4}$ [6]. Therefore, to avoid multiple scattering of dipoles, the Q^2 cut should be by about a factor of 10 higher than the saturation scale.

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