# Towards precise predictions for Higgs-boson production in the MSSM

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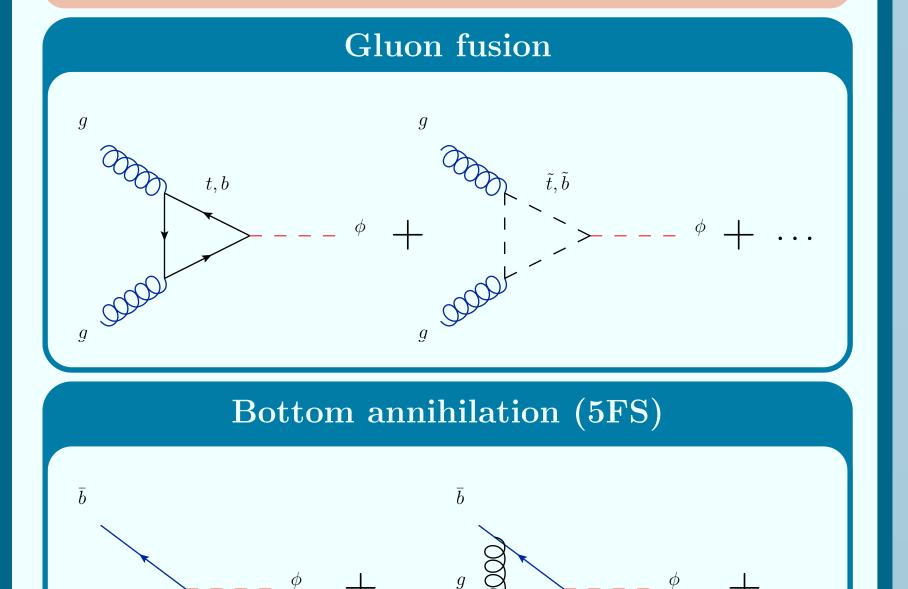
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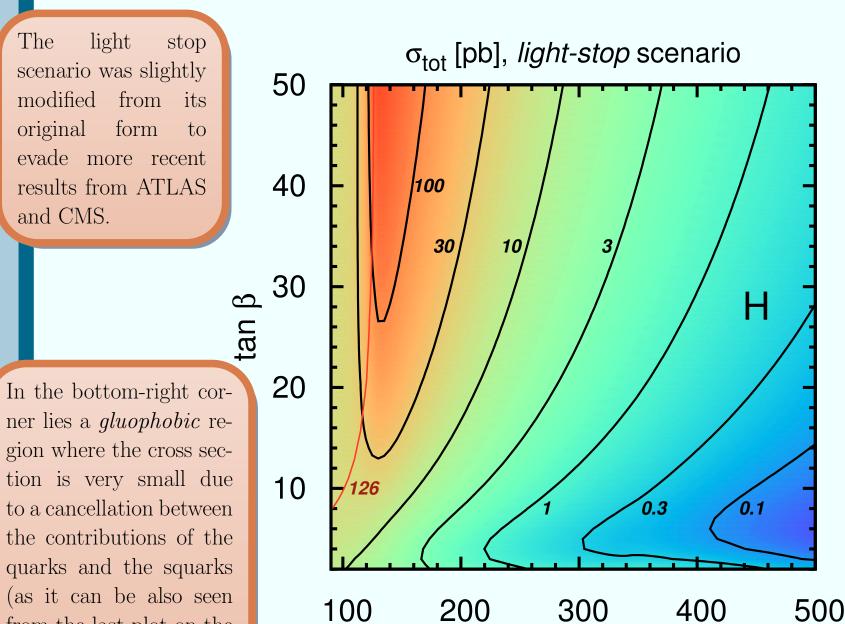
#### Aim of the study

Now that we are entering the era of precision physics at the LHC, it is important to have a thorough understanding of the theoretical uncertainties that affect our predictions. Our work follows this line of research and focuses on the case of the production of the neutral Higgs bosons of the MSSM in gluon fusion and bottom annihilation[1].

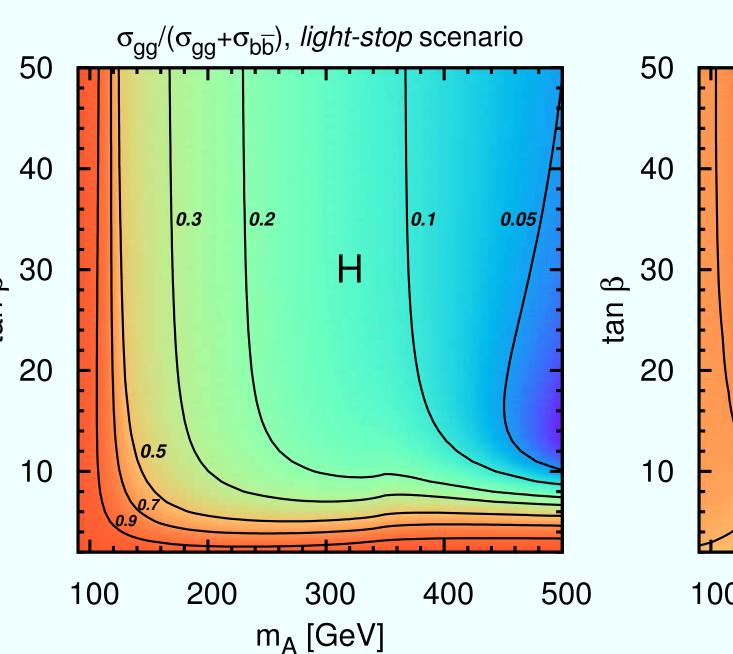


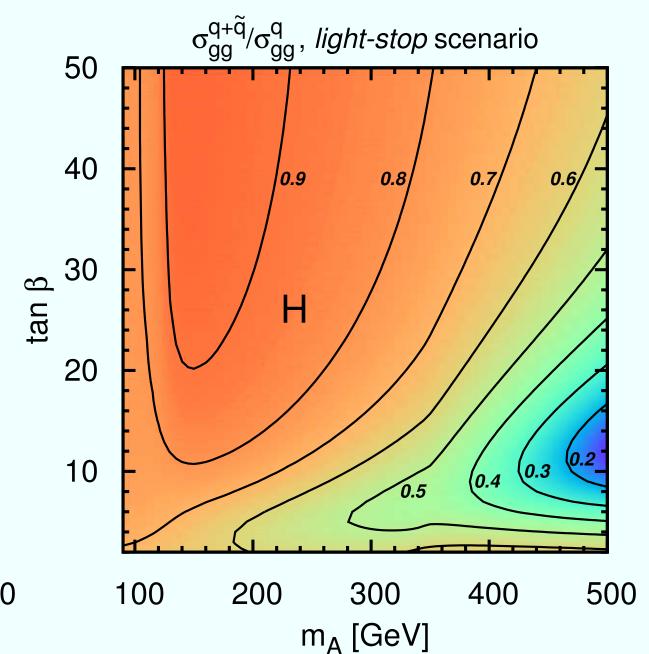
#### Outline of the analysis

We have studied the production of  $\phi = h, H, A$  using the code SusHi in the MSSM benchmark scenarios proposed by Carena et al (2013). Gluon fusion in SusHi includes: the exact top and bottom contribution up to NLO, the top results in the Heavy Quark Effective Field Theory (HQEFT) at NNLO, SUSY-QCD corrections from the stops and sbottom sectors at NLO and the available stops contribution at NNLO. EW contributions are included in a factorized form at NLO. The bottom Yukawa is rescaled to include  $\tan \beta$ -enhanced corrections ( $\Delta_b$ ). Bottom annihilation is included at NNLO in the 5FS (SUSY-QCD corrections are included in a rescaled bottom Yukawa coupling). The results shown in the poster are for **H** production in the so called light stop scenario ( $m_{\tilde{t}_1} = 324 \text{ GeV}$ ,  $m_{\tilde{t}_2} = 672 \text{ GeV}$ ). In this box, from left to right, we show the total cross section, the relative importance of the gluon fusion process and the impact of the SUSY-QCD corrections in the  $\tan \beta - m_A$  plane. In the rest of the poster we present the main uncertainties that affect this prediction for the total rate of H production.

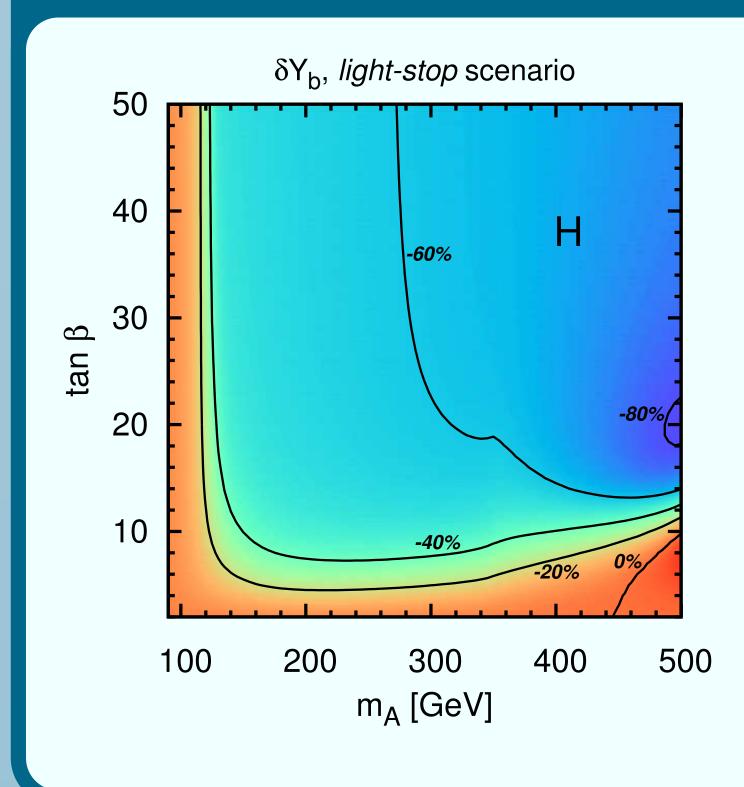


m<sub>A</sub> [GeV]





#### Bottom Yukawa renormalization scheme in gluon fusion

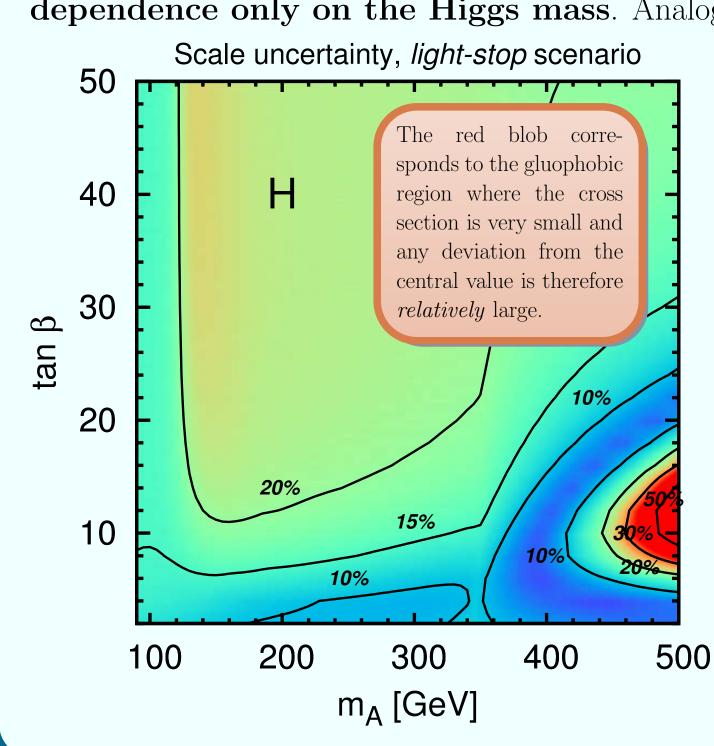


The bottom quark contribution to the gluon fusion process is subleading in the SM. In the MSSM, on the other hand, it can be the dominant one. Traditionally the bottom Yukawa has been extracted from the pole mass  $M_b$ . However **there are** no theoretically justifications to prefer this choice over the one of extracting the Yukawa from the  $\overline{MS}$  mass,  $m_b(\mu_r)$ . On the left we show the relative difference of the cross section for H when  $Y_b \sim m_b(m_H/2)$  with respect to the canonical one obtained with  $Y_b \sim M_b$ . We see that that when the bottom contribution is the dominant one we have deviations up to 60% - 80%. While such a large uncertainty could appear as worrisome, we should remember that in these regions the dominant production process is bottom annihilation, as it can be seen from the middle plot of the top-right box. Analogous considerations are valid for the other Higgses.

from the last plot on the

## Scale uncertainty

The scale variation uncertainty is evaluated with the standard prescription (a rescaling factor r=2 is used and we opt not to include  $\mu_r$  and  $\mu_f$  combinations that could introduce potentially large logarithmic deviation from the central value). Here we show the relative scale-uncertainty band width for H production in gluon fusion (left) and bottom annihilation (right). In the case of gluon fusion, we observe that the uncertainty follows the pattern determined by the available accuracy of the dominant contribution (e.g. where the bottom, known only at NLO, is dominant, the uncertainty is larger with respect to where the top contribution is dominant). In the case of bottom annihilation, being the SUSY-QCD corrections only in the rescaled bottom Yukawa, we see a dependence only on the Higgs mass. Analogous considerations are valid for h and A production.

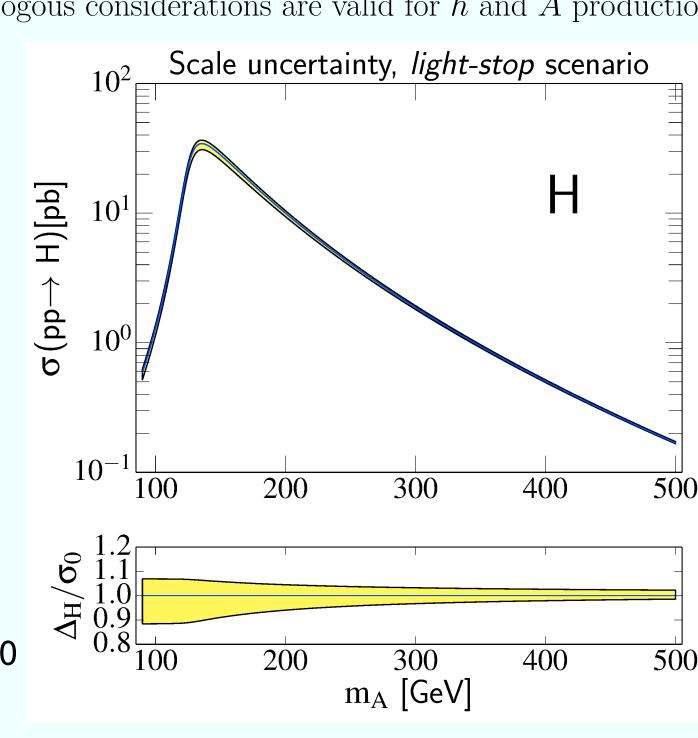


Other sources of uncertainty

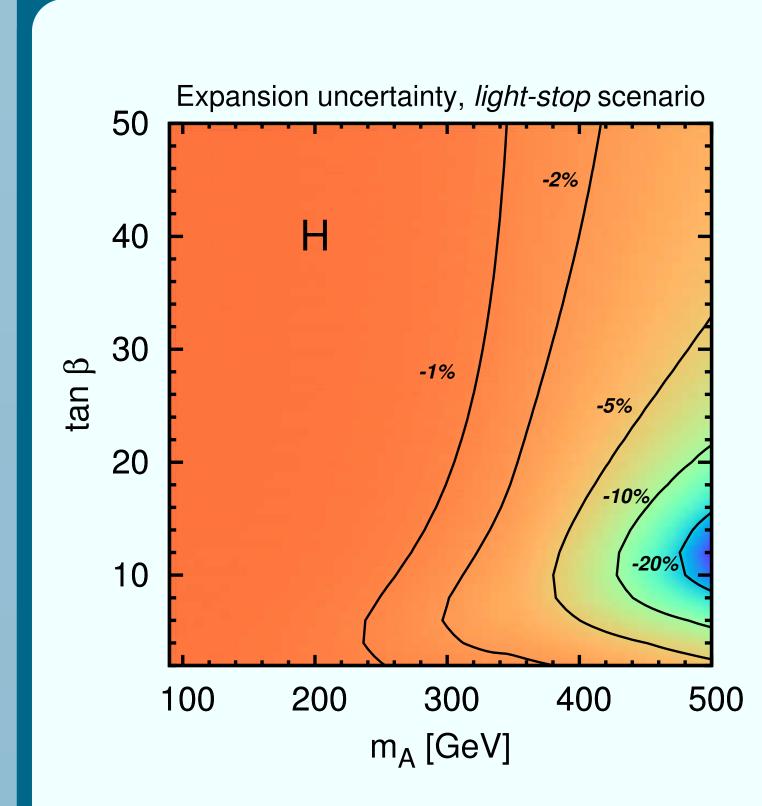
 $\triangleright$  PDF+ $\alpha_s$  uncertainty.

➤ The bottom mass dependence in the PDF.

There are also other sources of uncertainty that we have considered in our study:



### SUSY-QCD expansion uncertainty



Not all SUSY-QCD corrections are available at NLO in an exact form. For example, NLO 2-loop diagrams with quarks, squarks and a gluino are included in SusHi in the form of an expansion in the inverse powers of the SUSY-particle masses.

We estimate the error due to this approximation by rescaling the NLO result for stops and sbottoms by the test factor  $t \equiv \mathcal{A}_{\tilde{q}}^{1l}/\mathcal{A}_{\tilde{q}}^{1l,exp}$ . In the plot we show the relative difference of the

cross section for H when computed with the test factor with respect to its canonical value. We see that, aside from the gluophobic region, the impact in the  $\tan \beta$ - $m_A$  plane is rather modest. For the pseudoscalar uncertainties have a similar pattern but they are smaller in magnitude due to the fact that the squarks appear at NLO and not at LO.

# to the fact that the Acknowledgments

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> Uncertainty due to our incomplete knowledge of NNLO SUSY-QCD corrections.

 $\triangleright$  Estimation of the theoretical uncertainties coming from the  $\Delta_b$  corrections.

## References

[1] E. Bagnaschi, et al. "Towards precise predictions for Higgs-boson production in the MSSM." *JHEP*, 06:167, 2014. doi:10.1007/JHEP06(2014)167.

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