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PARTON DISTRIBUTIONS AND α_S FOR THE LHC

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We report on recent determinations of NNLO parton distributions and of $\alpha_s(M_Z)$ based on the world deep-inelastic data, supplemented by collider data. Some applications are discussed for semi-inclusive processes at the LHC.

1 Introduction

The precise knowledge of the parton distribution functions (PDFs) and the strong coupling constant $\alpha_s(M_Z^2)$ is of central importance for the description of any high energy reaction, in which hadrons are involved. In this way both the search for new particles, including the Higgs boson(s), and the precision description of the associated scattering cross sections at the LHC do crucially depend on these quantities. This also applies for detailed investigations of other heavy systems, which can now be explored in greater detail, as $W^\pm$, Z^0 -, jet- and top-quark production. On the other hand, both the world precision data on deep-inelastic scattering (DIS) and precise enough hadron collider data allow to fit the parton distribution functions and $\alpha_s(M_Z^2)$ in a correlated way and to test QCD. The corresponding analyses have to be based on the most accurate theoretical descriptions available. In case of deep-inelastic scattering the massless contributions to the structure functions are known to NNLO [1], while the heavy quark contributions are presently available to NLO [2]. At the 3-loop level in the large- Q^2 region a series of Mellin moments for the massive operator matrix elements has been calculated in [3]. All logarithmic contributions for the charm quark contributions to the structure function $F_2(x, Q^2)$ were calculated in [4]. Recently, an interpolation between the threshold and high Q^2 -region for the heavy flavor Wilson coefficients has been given in [5], also based on the moments in Ref. [3] and the known small- x contributions.

Currently there are results from six NNLO PDF-analyses available, four of which were published, with most recent results given in [6–11]. Due to the strong correlations in the fit the non-perturbative parameters of the PDFs and $\alpha_s(M_Z^2)$ have to be determined together. A detailed account of the systematic errors of the different experiments is necessary. In part of the phase space target mass effects [12] and higher twist terms have to be accounted for. For a recent survey on the theory of deep-inelastic scattering, see e.g. [13]. Since not all data sets are of the same quality it may be helpful to refer to those of the highest quality first and to carefully

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add other data sets then, rather than fitting a wide host of data of quite different quality. Proceeding in this way, our focus is narrower but also sharper [14]. In the following we give a brief survey on the status of the twist-2 parton distribution functions, on $\alpha_s(M_Z^2)$, and on hard scattering cross sections at the LHC, which start to develop sensitivity on the PDFs.

2 Parton Distribution Functions

During the last years various updates on the PDFs at NNLO (and at NLO) have been given by different groups [6–11]. Here it is mandatory to refer to the combined HERA data [15], not yet included by all of the groups. In Fig. 1 we compare several PDFs at NNLO with those being obtained in the recent analysis [8].

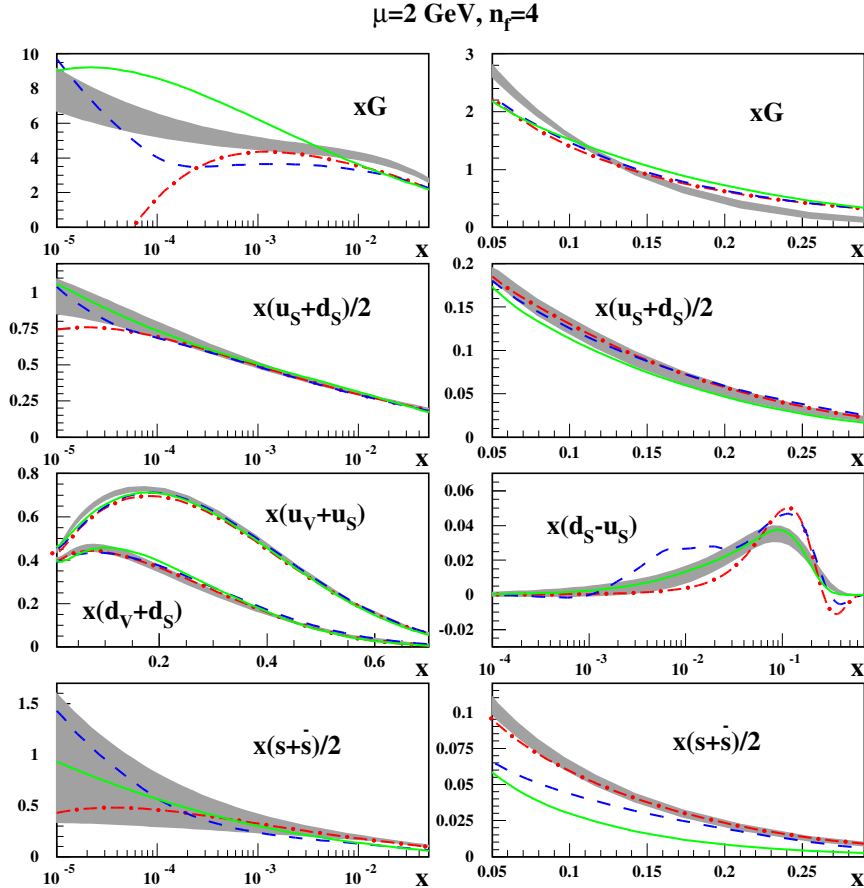


Figure 1: The 1σ band for the 4-flavor NNLO ABM11 PDFs at the scale of $\mu = 2 \text{ GeV}$ versus x obtained in [8] (shaded area) compared with the ones obtained by other groups (solid lines: JR09 [6], dashed dots: MSTW [7], dashes: NN21 [9]); from Ref. [8] ©(2012) by the American Physical Society.

	a	b	γ_1	γ_2	γ_3	A
u_v	0.712 ± 0.081	3.637 ± 0.138	0.593 ± 0.774	-3.607 ± 0.762	3.718 ± 1.148	
d_v	0.741 ± 0.157	5.123 ± 0.394	1.122 ± 1.232	-2.984 ± 1.077		
u_s	-0.363 ± 0.035	7.861 ± 0.433	4.339 ± 1.790		0.0280 ± 0.0036	0.0808 ± 0.0122
Δ	0.70 ± 0.28	11.75 ± 1.97	-2.57 ± 3.12			0.316 ± 0.385
s	-0.240 ± 0.055	7.98 ± 0.65				0.085 ± 0.017
g	-0.170 ± 0.012	10.71 ± 1.43	4.00 ± 4.21			

Table 1: The parameters of the PDFs in and their 1σ errors obtained in the scheme with $n_f = 3$ flavors; from [8], ©(2012) by the American Physical Society.

In [8] the shapes of the PDFs were parameterized by $xf(x) = Ax^{a+P(x)}(1-x)^b$, $P(x) = x(\gamma_1 + \gamma_2 x + \gamma_3 x^2)$, resp. $P_{us}(x) = (1 + \gamma_3 \ln(x))(1 + \gamma_1 x + \gamma_2 x^2)$. The fit parameters of the PDFs obtained are listed in Tab. 1.

Also the full error correlation matrices are presented in Ref. [8]. While the valence and u, d -sea quark distributions do widely agree for the $+$ -combinations, there are still significant differences in the gluon distribution at the starting scale of $\mu = 2$ GeV between different groups. In case of MSTW [7] the gluon distribution takes negative values at low x . The measurement of $F_L(x, Q^2)$ can partly clarify the situation since it is carried out at lower scales of Q^2 . In Fig. 2 predictions from four PDF-sets are compared. MSTW [7] yields lower values in the small- x region, while NN21 [17] gives high values. ABM11 [8] and JR09 [6] describe these data very well. To draw a final conclusion, more accurate data would be needed, however. They are expected to come from experiments at a future electron-ion collider [18] or at LHeC [19]. There are some differences visible in the $\bar{d} - \bar{u}$ distributions and the strange quark distribution. The present and upcoming LHC data on $W^\pm, Z^0$ boson production and the off-resonance Drell-Yan (DY) process may help to improve the situation here. The presently worst known distribution is that of the s -quark, while the c - and b -quark distributions (in some scheme to be specified) are driven by the gluon distribution mainly.

In the analysis [8] we included data down to scales of $Q^2 = 2.5$ GeV². In the lower W^2 -region contributions of twist-4 operators are found, which were param-

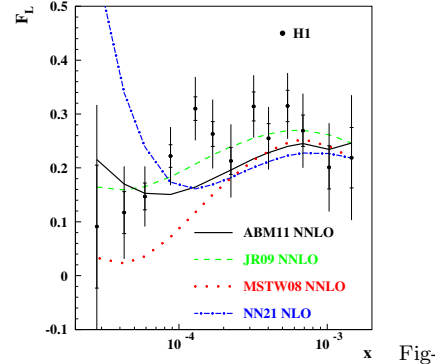


Figure 2: The data on F_L versus x obtained by the H1 collaboration [16] confronted with the 3-flavor scheme NNLO predictions based on the different PDFs (solid line: ABM11 [8], dashes: JR09 [6], dots: MSTW [7]). The NLO predictions based on the 3-flavor NN21 PDFs [17] are given for comparison (dashed dots). The value of Q^2 for the data points and the curves in the plot rises with x in the range of 1.5 to 45 GeV²; from Ref. [8], ©(2012) by the American Physical Society.

eterized by a phenomenological shape $H(x)/Q^2$ additively. In Tab. 2 the corresponding fit results are given.

	$H_2^p(x)/\text{GeV}^2$	$H_2^{ns}(x)/\text{GeV}^2$	$H_T^p(x)/\text{GeV}^2$
$x = 0.1$	-0.036 ± 0.012	-0.034 ± 0.023	-0.091 ± 0.017
$x = 0.3$	-0.016 ± 0.008	0.006 ± 0.017	-0.061 ± 0.012
$x = 0.5$	0.026 ± 0.007	-0.0020 ± 0.0094	0.0276 ± 0.0081
$x = 0.7$	0.053 ± 0.005	-0.029 ± 0.006	0.031 ± 0.006
$x = 0.9$	0.0071 ± 0.0026	0.0009 ± 0.0041	0.0002 ± 0.0015

Table 2: The parameters of the twist-4 contribution to the DIS structure functions for the fit to NNLO accuracy in QCD; from [8], ©(2012) by the American Physical Society.

3 The Strong Coupling Constant $\alpha_s(M_Z^2)$

The strong coupling constant $\alpha_s(M_Z^2)$ is measured together with the parameters of the PDFs, the heavy quark masses m_c and m_b , and the higher twist parameters within the analysis [8]. The present experimental accuracies of $O(1\%)$ require NNLO corrections, since at NLO the scale uncertainties amount to $O(5\%)$, cf. [20]. In Tab. 3 we compare the response of the data sets from BCDMS [21], NMC [22], SLAC [23], HERA [15] and the Drell-Yan data [24] at NLO and NNLO and to the NLO values given in earlier experimental analyses. Within the experimental errors an entirely consistent picture arises. Moreover, the new results agree with those of Refs. [6, 25–28].

Experiment	$\alpha_s(M_Z)$		
	NLO _{exp}	NLO	NNLO
BCDMS	0.1111 ± 0.0018	0.1150 ± 0.0012	0.1084 ± 0.0013
NMC	$0.117 \pm_{-0.016}^{+0.011}$	0.1182 ± 0.0007	0.1152 ± 0.0007
SLAC		0.1173 ± 0.0003	0.1128 ± 0.0003
HERA comb.		0.1174 ± 0.0003	0.1126 ± 0.0002
DY		0.108 ± 0.010	0.101 ± 0.025
ABM11		0.1180 ± 0.0012	0.1134 ± 0.0011

Table 3: Comparison of the values of $\alpha_s(M_Z)$ obtained by BCDMS and NMC at NLO with the individual results of the fit in the present analysis at NLO and NNLO for the HERA data the NMC data the BCDMS data the SLAC data and the DY data; from [8], ©(2012) by the American Physical Society.

As very well-known, both JR [29] and ABM [8, 30], along with other groups, carried out systematic fits including both jet data from the Tevatron and in [8] also from LHC^a.

^aContrary statements, as given in [31, 32], are incorrect; see Ref. [8] for all details. It is very problematic to call present fits including jet data and the world DIS data NNLO analyses, since the corresponding jet scattering cross sections are available in NLO *only*. In Ref. [8] we have accounted for threshold resummation contributions beyond NLO, however.

Experiment	$\alpha_s(M_Z)$		
	NLO_{exp}	NLO	NNLO*
D0 1 jet	$0.1161^{+0.0041}_{-0.0048}$	0.1190 ± 0.0011	0.1149 ± 0.0012
D0 2 jet		0.1174 ± 0.0009	0.1145 ± 0.0009
CDF 1 jet (cone)		0.1181 ± 0.0009	0.1134 ± 0.0009
CDF 1 jet ($k_\perp$)		0.1181 ± 0.0010	0.1143 ± 0.0009
ABM11		0.1180 ± 0.0012	0.1134 ± 0.0011

Table 4. Comparison of the values of $\alpha_s(M_Z)$ obtained by D0 with the ones based on including individual data sets of Tevatron jet data into the analysis at NLO. The NNLO* fit refers to the NNLO analysis of the DIS and DY data together with the NLO and soft gluon resummation corrections (next-to-leading logarithmic accuracy) for the 1 jet inclusive data.

In Tab. 4 we summarize the results including individual sets of Tevatron jet data [33]. Both the results in NLO and NNLO are 1σ compatible with our values on $\alpha_s(M_Z^2)$ obtained without these data. In case we include the ATLAS jet data [34] into the analysis along with the DIS and Drell-Yan data we obtain $\alpha_s(M_Z^2) = 0.1141 \pm 0.0008$, demanding $p_\perp^{\text{jet}} > 100$ GeV. Here the jet cross section was described at NLO adding threshold resummation, see [8] for details.

The ATLAS and CMS jet data span a wider kinematic range than those of Tevatron and will allow very soon even more accurate measurements. At the moment NLO QCD analyses are carried out with first results being obtained in [35, 36]. The values of $\alpha_s(M_Z^2)$ come out rather low. Including the scale uncertainties the following NLO values are obtained for the 3/2 jet ratio at CMS

$$\alpha_s(M_Z^2) = 0.1143 \pm 0.0064 \quad [0.1191 \pm 0.0006] \quad [37] \quad \text{NNPDF21} \quad (1)$$

$$\alpha_s(M_Z^2) = 0.1130 \pm 0.0080 \quad [0.180 \text{ (favored value)}] \quad [38] \quad \text{CT10} \quad (2)$$

$$\alpha_s(M_Z^2) = 0.1135 \pm 0.0096 \quad [0.1202^{+0.0012}_{-0.0015}] \quad [39] \quad \text{MSTW08}, \quad (3)$$

varying the corresponding PDF-sets given for a wider range in $\alpha_s(M_Z^2)$ [36]. Note that the minima in χ^2 found by the fitting groups are reached at much higher values of $\alpha_s(M_Z^2)$ (quoted in parenthesis). In the near future the scale uncertainties in (1–3) will significantly diminish upon the arrival of the NNLO corrections. A comparable NLO value has been reported using ATLAS jet data [35] $\alpha_s(M_Z^2) = 0.1151 \pm 0.0050$ (exp.) $^{+0.0080}_{-0.0073}$ (th.) Recently the jet energy-scale error has been improved by CMS [36], leading to a significant reduction of the experimental error. Moreover, the gluonic NNLO corrections for jet-production are now available [40], with the other NNLO terms to be expected later this year. The gluonic NNLO K -factor is positive. As shown in Fig. 2 of Ref. [40] the scale dependence for $\mu = \mu_F = \mu_R$ behaves flat over a wide range of scales. It is expected that also the error due to scale variation will turn out to be very small. Therefore, a definite answer on $\alpha_s(M_Z^2)$ at NNLO from both ATLAS and CMS alone using jet measurement is imminent. It should be mentioned that in this determination both $\alpha_s(M_Z^2)$ and the gluon distribution shall be fitted due to the strong correlations of both quantities,

which also will yield an independent input on $xG(x, \mu^2)$ based on LHC data only and will allow for interesting comparisons with the gluon distributions determined by other means.

Data Set	ABM11	BBG	NN21	MSTW
BCDMS	0.1048 ± 0.0013	0.1126 ± 0.0007	0.1158 ± 0.0015	0.1101 ± 0.0094
NMC	0.1152 ± 0.0007	0.1153 ± 0.0039	0.1150 ± 0.0020	0.1216 ± 0.0074
SLAC	0.1128 ± 0.0003	0.1158 ± 0.0034	> 0.124	$\begin{cases} 0.1140 \pm 0.0060 \text{ ep} \\ 0.1220 \pm 0.0060 \text{ ed} \end{cases}$
HERA	0.1126 ± 0.0002		$\begin{cases} 0.1199 \pm 0.0019 \\ 0.1231 \pm 0.0030 \end{cases}$	0.1208 ± 0.0058
DY	0.101 ± 0.025	—	—	0.1136 ± 0.0100
	0.1134 ± 0.0011	0.1134 ± 0.0020	0.1173 ± 0.0007	0.1171 ± 0.0014

Table 5: Summary of recent NNLO QCD analyses of the DIS world data, supplemented by related measurements using other processes; from [8], ©(2012) by the American Physical Society.

Let us compare the results of the NNLO fits ABM11 [8], BBG [26], NN21 [41] and MSTW08 [39] on $\alpha_s(M_Z^2)$ w.r.t. their individual response to the deep-inelastic and Drell-Yan data sets used, cf. Tab. 5. While the $\alpha_s(M_Z^2)$ values obtained in the ABM11 and BBG analyses lead to 1σ consistent values between the final value of $\alpha_s(M_Z^2) = 0.1134$ and those of the individual data sets, for NN21 and MSTW08 variations are being observed. NN21 agrees with ABM11 and BBG for the NMC data but reports a higher value for BCDMS, with the converse for MSTW08. Both NN21 and MSTW08 come up with a larger value for the HERA data, however, MSTW08 does not include the combined HERA data yet. For the SLAC data MSTW08 agrees with ABM11 and BBG for the ep data but obtains a much larger value for the ed data, while NN21 obtains $\alpha_s(M_Z^2) > 0.124$ in case of the SLAC data. Because of the large errors, the Tevatron Drell-Yan data are not very decisive on the value of $\alpha_s(M_Z^2)$. Despite the similar final values of $\alpha_s(M_Z^2)$ obtained by NN21 and MSTW08, those for the individual data sets vary significantly.

We outlined in [42] that a consistent F_L -treatment for the NMC data and the BCDMS-data, cf. [26], is necessary and will lead to a change of the value of $\alpha_s(M_Z^2)$. Furthermore, we analyzed the sensitivity on kinematic cuts applied to remove higher twist effects. In the flavor non-singlet case this can be achieved cutting for $W^2 > 12.5 \text{ GeV}^2$, cf. [26]. As we have shown in Section 2, there are also higher twist contributions in the lower x -region. They can be removed applying a cut $Q^2 > 10 \text{ GeV}^2$, which, however, is not applied by NN21 and MSTW08. We performed a fit ignoring the higher twist terms and allowed for the range of data down to values of $Q^2 > 2.5 \text{ GeV}^2$, [8]. Here we obtain $\alpha_s(M_Z^2) = 0.1191 \pm 0.0016$, very close to the values found by NN21 and MSTW08.

In Tab. 6 a general overview on the values of $\alpha_s(M_Z^2)$ at NNLO is given, with a few determinations effectively at $N^3\text{LO}$ in valence analyses [26,27], and the hadronic Z -decay [47]. The BBG, BB, GRS, ABKM, JR, and ABM11 analyses find lower values of $\alpha_s(M_Z^2)$ with errors at the 1–2% level, while NN21 and MSTW08 find larger values with comparable accuracy to the former ones, as has been discussed

before. We also add a yet unpublished NNLO value from CT10 quoting a significantly larger error. In the analysis of thrust in e^+e^- data two groups find low values, also with errors at the 1% level. Higher values of $\alpha_s(M_Z^2)$ are found for the e^+e^- 3-jet rate, the hadronic Z -decay, and τ -decay. $\alpha_s(M_Z^2)$ has also been determined in different lattice simulations to high accuracy. The N³LO values for $\alpha_s(M_Z^2)$ in the valence analyses [26, 27] yield slightly larger values than at NNLO, but are fully consistent with these values within errors.

	$\alpha_s(M_Z)$	
BBG	$0.1134^{+0.0019}_{-0.0021}$	valence analysis, NNLO [25]
BB	0.1132 ± 0.0022	valence analysis, NNLO [26]
GRS	0.112	valence analysis, NNLO [27]
ABKM	0.1135 ± 0.0014	HQ: FFNS $n_f = 3$ [24]
ABKM	0.1129 ± 0.0014	HQ: BSMN-approach [24]
JR	0.1124 ± 0.0020	dynamical approach [6]
JR	0.1158 ± 0.0035	standard fit [6]
ABM11	0.1134 ± 0.0011	[8]
MSTW	0.1171 ± 0.0014	[39]
NN21	0.1173 ± 0.0007	[41]
CT10	0.118 ± 0.005	[11]
Abbate et al.	$0.1135 \pm 0.0011 \pm 0.0006$	e^+e^- thrust [43]
Gehrmann et al.	$0.1131^{+0.0028}_{-0.0022}$	e^+e^- thrust [44]
3 jet rate	0.1175 ± 0.0025	Dissertori et al. 2009 [45]
Z-decay	0.1189 ± 0.0026	BCK 2008/12 (N ³ LO) [46,47]
τ decay	0.1212 ± 0.0019	BCK 2008 [46]
τ decay	0.1204 ± 0.0016	Pich 2011 [48]
τ decay	0.1191 ± 0.0022	Boito et al. 2012 [49]
lattice	0.1205 ± 0.0010	PACS-CS 2009 (2+1 fl.) [50]
lattice	0.1184 ± 0.0006	HPQCD 2010 [51]
lattice	0.1200 ± 0.0014	ETM 2012 (2+1+1 fl.) [52]
lattice	0.1156 ± 0.0022	Brambilla et al. 2012 (2+1 fl.) [53]
BBG	$0.1141^{+0.0020}_{-0.0022}$	valence analysis, N ³ LO(*) [25]
BB	0.1137 ± 0.0022	valence analysis, N ³ LO(*) [26]
world average	0.1184 ± 0.0007	[54] (2009)
	0.1183 ± 0.0010	[48] (2011)

Table 6. Summary of recent NNLO QCD analyses of the DIS world data, supplemented by related measurements using other processes; from [8], [©](2012) by the American Physical Society.

Finally we would like to mention, that the present DIS world data together with the $F_2^{c\bar{c}}(x, Q^2)$ data, are competitive in the determination of the charm quark mass. In a recent analysis [55] the value of $m_c(m_c) = 1.24 \pm 0.03$ (exp.) $^{+0.04}_{-0.00}$ (th.) is

obtained at NNLO. The analysis in Ref. [55] is presently the only one, in which all known higher order heavy flavor corrections to deep-inelastic scattering have been considered. Let us note, that there is still a wide spectrum in the way heavy flavor corrections are accounted for in [6–11] w.r.t. the present knowledge on the side of theory.

4 Hard Scattering Cross Sections at the LHC

The presently obtained precision predictions on the PDFs at NNLO may be used to predict different scattering cross sections being measured at the LHC. This applies in particular to those cross sections, resp. cross section ratios, which can be measured at high precision, cf. [56]. Here observables of central importance are the $W^\pm/Z^0$ cross section ratios. In Fig. 3 we show results obtained by ATLAS [57].

Due to the ratios considered the luminosity errors do widely cancel and the present experimental accuracy nearly compares to the size of the PDF-errors. Later this year even more precise results are expected and the analyses are still ongoing. One should also notice, that the corresponding CMS bands [58] are shifted by about 1σ at present. For LHCb similar analyses have been performed, cf. [59]. PDF predictions not hitting the yellow bands in Fig. 3 cannot be considered^b as outliers [32]. After all, precision data like this will have an impact on the PDF fits very soon as has happened in improving PDF-sets so often in the past, see e.g. [7] and references therein.

Let us now consider more differential distributions in case of $W^\pm$ and Z^0 production. In Ref. [60] we investigated the rapidity distributions for $W^\pm, Z^0$ -production for ATLAS and LHCb and the $e^\pm$ -asymmetry for $W^\pm$ production in CMS. This requires detailed Monte Carlo studies and comparisons based on the codes DYNLO

^bWe thank M. Klein for a corresponding clarification.

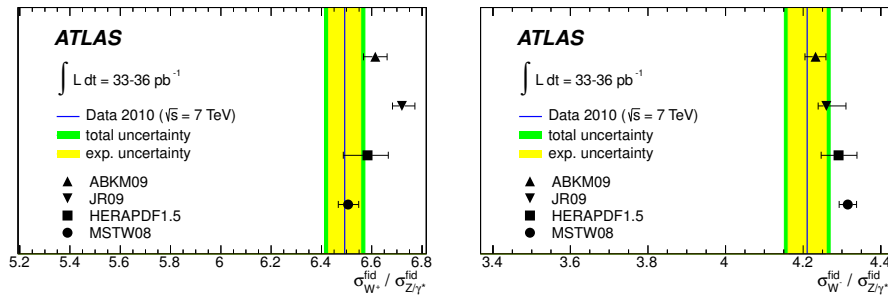


Figure 3: Measured and predicted fiducial cross section ratios, $\sigma(W^+)/\sigma(Z^0)$ (left) and $\sigma(W^-)/\sigma(Z^0)$ (right). The experimental uncertainty (inner yellow band) includes the experimental systematic errors. The total uncertainty (outer green band) includes the statistical uncertainty and the small contribution from the acceptance correction. The uncertainties of the ABKM [25], JR [6] and MSTW08 [7] predictions are given by the PDF uncertainties considered to correspond to 68 % CL and their correlations are derived from the eigenvector sets. The results for HERAPDF comprise all three sources of uncertainty of that set; from [57], ©(2012) by the American Physical Society.

1.3 [61] and FEWZ 3.1 [62]. As an example, we show the comparison of the ABM11 prediction for the LHCb data in Fig. 4. Tab. 7 summarizes the fit results for the predictions based on the ABM11-distributions for ATLAS, CMS and LHCb. χ^2/NDP values of 1.3–1.07 are obtained. Here the NNLO standard value of $\alpha_s(M_Z^2)$ of ABM11 has been used. Contrary to this, a recent benchmarking of LHC cross sections in Table 6 of [31] reported values of χ^2/NDP in the range of 1.602–1.923 at NNLO for comparisons assuming $\alpha_s(M_Z^2) = 0.117$ here, which are not confirmed in [60]. For the MSTW08 distributions we find χ^2/NDP of 26.2/9 for the Z -production at LHCb. The benchmarking of this channel has not been performed in [31].

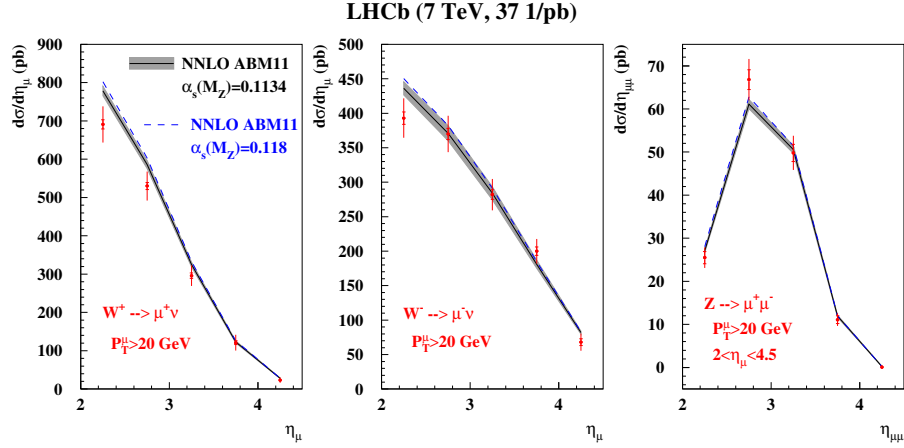


Figure 4: The LHCb data of Ref. [59] on the rapidity distribution of the charged leptons μ^+ (left panel), μ^- (central panel), and Z -boson (right panel) in comparison with the NNLO predictions computed with the ABM11 PDFs taking into account the uncertainties due to PDFs (shaded area). The dashed curves display the ABM11 predictions obtained with the value of $\alpha_s(M_Z) = 0.118$; from [60], ©(2013) by SISSA.

Experiment	ATLAS [57]	CMS [66]	LHCb [59]	LHCb [67]
Final states	$W^+ \rightarrow l^+ \nu$ $W^- \rightarrow l^- \nu$ $Z \rightarrow l^+ l^-$	$W^+ \rightarrow e^+ \nu$ $W^- \rightarrow e^- \nu$	$W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$	$Z \rightarrow e^+ e^-$
Luminosity (1/pb)	35	840	37	940
NDP	30	11	10	9
χ^2	34.7(7.7)	11.8(4.7)	13.0(4.5)	11.5(4.2)

Table 7. The value of χ^2 obtained for different samples of the Drell-Yan LHC data with the NNLO ABM11 PDFs. The figures in parenthesis give one standard deviation of χ^2 equal to $\sqrt{2NDP}$; from [60] ©(2013) SISSA.

At present the inclusive $t\bar{t}$ -production cross section is measured by ATLAS and CMS with

$$\sigma(t\bar{t}) = 186 \pm 13 \text{ (stat.)} \pm 20 \text{ (syst.)} \pm 7 \text{ (lum.) } pb \text{ ATLAS, [63]} \quad (4)$$

$$\sigma(t\bar{t}) = 161.9 \pm 2.5 \text{ (stat.)} \pm 5.1 \text{ (syst.)} \pm 3.6 \text{ (lum.) } pb \text{ CMS [64]} \quad (5)$$

$$\sigma(t\bar{t}) = 139 \pm 10 \text{ (stat.)} \pm 26 \text{ (syst.)} \pm 3 \text{ (lum.) } pb \text{ CMS, jets [65],} \quad (6)$$

for $m_t = 172.5$ GeV. The current difference between both experiments has still to be understood and reconciled. Furthermore, the calculation of the inclusive $t\bar{t}$ -production cross section at NNLO has to be completed. The scattering cross section exhibits a very strong dependence on the top quark mass m_t . In the ABM11 analysis a value of $\sigma(t\bar{t}) = 130.4^{+2.9}_{-7.2} \text{ (scale)} \pm 5.9 \text{ (PDF)}^c$ is obtained for $m_t = 173$ GeV using HATHOR [68], with a variation from $\sigma(t\bar{t}) = 167.9 \dots 122.6 pb$ for m_t in the range [165, 175 GeV]. On the other hand, JR, MSTW08 and NN21 give values in the range of $159.3 - 167.7 pb$ for $m_t = 173$ GeV. In this comparison, the ABM11 value comes out lower. With a variation to somewhat lower values of m_t it is also getting closer to the higher values of CMS and ATLAS and consistent with the lower value by CMS. In the future a final LHC analysis will deliver a unique value for $\sigma(t\bar{t})$ with smaller errors and a determination of $m_t^{\overline{\text{MS}}}$ being fully consistent with the quantum field-theoretic description of the scattering cross section.

The inclusive Higgs-boson production cross section is dominated by the gluonic channel and depends on both the gluon distribution function and the strong coupling constant as $\propto \alpha_s^2(M_H^2) xG(x, M_H^2) \otimes xG(x, M_H^2)$, with M_H the Higgs boson mass and $\otimes$ denoting the Mellin convolution. Due to this the errors both in α_s and xG enter the cross section twice and a precise prediction requires that both quantities have to be understood very accurately. Detailed tables of the production cross sections have been given in Refs. [8, 56, 69, 70] and in Ref. [71]. After first evidence for a Higgs-like particle has been found in the mass range ~ 125 GeV [72, 73], the search is ongoing and further detailed results are expected to be presented for the $\sqrt{s} = 7$ and 8 TeV runs at LHC very soon.

5 Conclusions

We reported on recent results of NNLO determinations of the parton distribution functions and of $\alpha_s(M_Z^2)$, based on the world precision data on deep-inelastic scattering and sensitive data from hadron colliders as Tevatron and the LHC, comparing the results obtained by different fitting groups. There are still differences being observed. In case of the gluon distribution the upcoming precision measurement of the jet cross sections at LHC may deliver a new impact. In various analyses low values of $\alpha_s(M_Z^2)$ are obtained in NNLO, for which we quote the value of ABM11 with $\alpha_s(M_Z^2) = 0.1134 \pm 0.0011$. We have given reasons why the NNLO values for $\alpha_s(M_Z^2)$ are coming out larger for MSTW08 and NNPDF and think that the current difference is understood. The correct inclusion of systematic error correlations and higher twist effects are important. The inclusion of the ATLAS jet data

^c In [32] the cross section for ABM11 [8] is quoted by 7 pb too low.

into the present analysis leads to $\alpha_s(M_Z^2) = 0.1141 \pm 0.0008$. The gluon distribution being obtained is softer than that in case of fitting to the Tevatron jet data. A rather consistent picture has been obtained for the inclusive $W^\pm$ and Z^0 production cross section ratios by ATLAS, CMS and LHCb, which show some sensitivity to the quark and anti-quark densities. For the inclusive $t\bar{t}$ -cross section there are still differences between the measurements at ATLAS and CMS. The different fitting groups also report different results here. However, the scattering cross sections do strongly depend on $\alpha_s(M_Z^2)$, $xG(x, Q^2)$ and the value of m_t assumed in the data analysis. Furthermore, the NNLO corrections are not yet complete for this process.

The gluon density will soon be constrained by the present LHC data on two and three jet production. These data will also yield an improved value of $\alpha_s(M_Z^2)$. There is a need of continuous benchmarking between groups using correct statistical methods. An effective rescaling of experimental errors, even with individual factors changing for different data sets, is neither necessary nor can it lead to an objective result. It is evident that highly precise and systematically well-controlled data sets have to be given preference over data with a poorer systematics in global analyses to determine the parton distribution functions and $\alpha_s(M_Z^2)$.

The so-called ‘pdf4lhc recommendation’ [74] unfortunately left out two out of three NNLO analyses since 2009. Moreover, the effect of the combined HERA data has not been considered in the fits chosen there. On the other hand, all NNLO analyses are available in terms of parameterizations in LHAPDF and most of them run very fast, which is important for experimental use. We therefore recommend that the experimental groups *freely* use all NNLO PDFs within their analyses together with the correlation matrices and the right value of $\alpha_s(M_Z^2)$. Since more and more differential distributions are investigated with a need of sophisticated systematic analyses by the LHC experiments, external determinations of theory errors related to the available PDFs become readily impossible. They have to be determined by the experiments to deliver unbiased precision results. The benchmarking between the different fitting groups has to continue to finally provide a thoroughly agreeing picture both on the PDFs and on $\alpha_s(M_Z^2)$ from DIS and precise hard scattering cross sections at high energies.

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Note added. After completion of the present paper a preprint with details on the CT10 analysis appeared [75]. Here $\alpha_s(M_Z)$ is dealt with as a pilot parameter at NNLO and not fitted. As known from other analyses, for a choice of $\alpha_s(M_Z) \in [0.112, 0.127]$ very different values of $\chi_{\min}^2$ are obtained, cf. e.g. Figs. 4.16–4.18 of Ref. [8].

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