

VERITAS OBSERVATIONS OF THE NOVA IN V407 CYGNI

E. ALIU¹, S. ARCHAMBAULT², T. ARLEN³, T. AUNE⁴, M. BEILICKE⁵, W. BENBOW⁶, A. BOUVIER⁴, S. M. BRADBURY⁷, J. H. BUCKLEY⁵, V. BUGAEV⁵, K. BYRUM⁸, A. CANNON⁹, A. CESARINI¹⁰, L. CIUPIK¹¹, E. COLLINS-HUGHES⁹, M. P. CONNOLLY¹⁰, W. CUI¹², G. DECERPRIT¹³, R. DICKHERBER⁵, C. DUKE¹⁴, J. DUMM¹⁵, V. V. DWARKADAS¹⁶, M. ERRANDO¹, A. FALCONE¹⁷, Q. FENG¹², J. P. FINLEY¹², G. FINNEGAN¹⁸, L. FORTSON¹⁵, A. FURNISS⁴, N. GALANTE⁶, D. GALL¹⁹, S. GODAMBE¹⁸, S. GRIFFIN², J. GRUBE¹¹, G. GYUK¹¹, D. HANNA², J. HOLDER²⁰, H. HUAN²¹, G. HUGHES¹³, T. B. HUMENSKY²², P. KAARET¹⁹, N. KARLSSON¹⁵, M. KERTZMAN²³, Y. KHASSEN⁹, D. KIEDA¹⁸, H. KRAWCZYNSKI⁵, F. KRENNRICH²⁴, M. J. LANG¹⁰, K. LEE⁵, G. MAIER¹³, P. MAJUMDAR³, S. MCARTHUR⁵, A. MCCANN², J. MILLIS²⁵, P. MORIARTY²⁶, R. MUKHERJEE¹, P. D. NUÑEZ¹⁸, R. A. ONG³, M. ORR²⁴, A. N. OTTE²⁷, D. PANDEL²⁸, N. PARK²¹, J. S. PERKINS^{29,30}, M. POHL^{13,31}, H. PROKOPH¹³, J. QUINN⁹, K. RAGAN², L. C. REYES³², P. T. REYNOLDS³³, E. ROACHE⁶, H. J. ROSE⁷, J. RUPPEL^{13,31}, D. B. SAXON²⁰, M. SCHROEDTER⁶, G. H. SEMBROSKI¹², C. SKOLE¹³, A. W. SMITH¹⁸, D. STASZAK², I. TELEZHINSKY^{13,31}, G. TEŠIĆ², M. THEILING¹², S. THIBADEAU⁵, K. TSURUSAKI¹⁹, J. TYLER², A. VARLOTTA¹², S. VINCENT¹⁸, M. VIVIER²⁰, S. P. WAKELY²¹, J. E. WARD⁹, T. C. WEEKES⁶, A. WEINSTEIN²⁴, T. WEISGARBER²¹, R. WELSING¹³, D. A. WILLIAMS⁴, AND B. ZITZER¹²

¹ Department of Physics and Astronomy, Barnard College, Columbia University, NY 10027, USA

² Physics Department, McGill University, Montreal, QC H3A 2T8, Canada

³ Department of Physics and Astronomy, University of California, Los Angeles, CA 90095, USA

⁴ Santa Cruz Institute for Particle Physics and Department of Physics, University of California, Santa Cruz, CA 95064, USA

⁵ Department of Physics, Washington University, St. Louis, MO 63130, USA

⁶ Fred Lawrence Whipple Observatory, Harvard-Smithsonian Center for Astrophysics, Amado, AZ 85645, USA

⁷ School of Physics and Astronomy, University of Leeds, Leeds LS2 9JT, UK

⁸ Argonne National Laboratory, 9700 South Cass Avenue, Argonne, IL 60439, USA

⁹ School of Physics, University College Dublin, Belfield, Dublin 4, Ireland

¹⁰ School of Physics, National University of Ireland Galway, University Road, Galway, Ireland

¹¹ Astronomy Department, Adler Planetarium and Astronomy Museum, Chicago, IL 60605, USA

¹² Department of Physics, Purdue University, West Lafayette, IN 47907, USA

¹³ DESY, Platanenallee 6, 15738 Zeuthen, Germany

¹⁴ Department of Physics, Grinnell College, Grinnell, IA 50112-1690, USA

¹⁵ School of Physics and Astronomy, University of Minnesota, Minneapolis, MN 55455, USA

¹⁶ Department of Astronomy and Astrophysics, University of Chicago, Chicago, IL 60637, USA

¹⁷ Department of Astronomy and Astrophysics, 525 Davey Lab, Pennsylvania State University, University Park, PA 16802, USA

¹⁸ Department of Physics and Astronomy, University of Utah, Salt Lake City, UT 84112, USA

¹⁹ Department of Physics and Astronomy, University of Iowa, Van Allen Hall, Iowa City, IA 52242, USA; daniel-d-gall@uiowa.edu, kazuma-tsurusaki@uiowa.edu

²⁰ Department of Physics and Astronomy and the Bartol Research Institute, University of Delaware, Newark, DE 19716, USA

²¹ Enrico Fermi Institute, University of Chicago, Chicago, IL 60637, USA

²² Physics Department, Columbia University, New York, NY 10027, USA

²³ Department of Physics and Astronomy, DePaul University, Greencastle, IN 46135-0037, USA

²⁴ Department of Physics and Astronomy, Iowa State University, Ames, IA 50011, USA

²⁵ Department of Physics, Anderson University, 1100 East 5th Street, Anderson, IN 46012, USA

²⁶ Department of Life and Physical Sciences, Galway-Mayo Institute of Technology, Dublin Road, Galway, Ireland

²⁷ School of Physics & Center for Relativistic Astrophysics, Georgia Institute of Technology, 837 State Street NW, Atlanta, GA 30332-0430, USA

²⁸ Department of Physics, Grand Valley State University, Allendale, MI 49401, USA

²⁹ CRESST and Astroparticle Physics Laboratory NASA/GSFC, Greenbelt, MD 20771, USA

³⁰ University of Maryland, Baltimore County, 1000 Hilltop Circle, Baltimore, MD 21250, USA

³¹ Institut für Physik und Astronomie, Universität Potsdam, 14476 Potsdam-Golm, Germany

³² Physics Department, California Polytechnic State University, San Luis Obispo, CA 94307, USA

³³ Department of Applied Physics and Instrumentation, Cork Institute of Technology, Bishopstown, Cork, Ireland

Received 2011 December 14; accepted 2012 May 10; published 2012 July 6

ABSTRACT

We report on very high energy ($E > 100$ GeV) gamma-ray observations of V407 Cygni, a symbiotic binary that underwent a nova outburst producing 0.1–10 GeV gamma rays during 2010 March 10–26. Observations were made with the Very Energetic Radiation Imaging Telescope Array System during 2010 March 19–26 at relatively large zenith angles due to the position of V407 Cyg. An improved reconstruction technique for large zenith angle observations is presented and used to analyze the data. We do not detect V407 Cygni and place a differential upper limit on the flux at 1.6 TeV of $2.3 \times 10^{-12} \text{ erg cm}^{-2} \text{ s}^{-1}$ (at the 95% confidence level). When considered jointly with data from *Fermi*-LAT, this result places limits on the acceleration of very high energy particles in the nova.

Key words: gamma rays: general – novae, cataclysmic variables – white dwarfs

Online-only material: color figures

1. INTRODUCTION

In March of 2010, the *Fermi*-LAT Collaboration announced a new GeV transient in the galactic plane, FGL J2102+4542, that was identified as a nova outburst in the symbiotic binary

V407 Cygni (hereafter V407 Cyg). At least seven GeV transients located near the Galactic plane have been discovered by EGRET, *Fermi*-LAT, and *AGILE*. Only two have been identified at other wavelengths: V407 Cyg, which is the first nova to be detected at GeV energies, and J0109+6134, which was likely a background

blazar (Abdo et al. 2010; Vandenbrouke et al. 2010). The physical nature of the other five sources is unknown (Abdo et al. 2010; Chaty 2012; Hays et al. 2010; Sabatini et al. 2010; Tavani et al. 1997; Vandenbrouke et al. 2010), and some of these GeV transients may represent a new class of gamma-ray emitting objects.

The *Fermi*-LAT Collaboration reported variable gamma-ray emission in the 0.1–10 GeV band from FGL J2102+4542 during 2010 March 10–26 (MJD 55265–55281; Cheung et al. 2010). Its flux in gamma rays, binned on a day-to-day basis, peaked 2010 March 13–14 with a flux of 9×10^{-7} photons cm $^{-2}$ s $^{-1}$ above 100 MeV (Abdo et al. 2010). The GeV gamma-ray activity lasted approximately two weeks. The initial report of GeV emission triggered Very Energetic Radiation Imaging Telescope Array System (VERITAS) observations of the object at very high energy (VHE; $E > 100$ GeV) as part of an ongoing campaign to observe transients detected by *Fermi*-LAT.

Using multi-wavelength data, it was determined that the new transient was most likely associated with V407 Cyg, a binary system consisting of a Mira-type pulsating red giant and a white dwarf companion (Abdo et al. 2010). A nova outburst from V407 Cyg was detected in the optical waveband on 2010 March 10 (Nishiyama & Kabashima 2010) with a magnitude of ~ 6.9 , while pre-outburst magnitudes from the previous two years of monitoring ranged between magnitudes 9 and 12 (Abdo et al. 2010). V407 Cyg has been optically monitored for decades and has experienced previous outbursts, but the system had never been observed to be as bright as during the nova (e.g., Munari et al. 1990; Kolotilov et al. 1998). The onset of the optical outburst corresponds to the first significant detection of the source by the *Fermi*-LAT on 2010 March 10.

Novae in red giant/white dwarf systems have been known to produce expanding shocks that can result in X-ray emission (e.g., the recurrent nova RS Oph; Sokoloski et al. 2006; Bode et al. 2006), and indeed, X-ray emission from V407 Cyg was detected after the nova (Abdo et al. 2010; Nelson et al. 2012). Based on the observed X-ray emission from the 2006 nova outburst of RS Oph, before the launch of *Fermi*-LAT, Tatischeff & Hernanz (2007) suggested that particles could be accelerated in novae up to TeV energies, but gamma-ray emission from a nova had never previously been detected. Here, we discuss the VERITAS observations of V407 Cyg and their implications for gamma-ray emission from the nova. We also describe an improved event reconstruction technique for stereo observations by imaging atmospheric Cherenkov telescopes (IACTs) made at large zenith angles (LZA).

2. OBSERVATION AND ANALYSIS

2.1. VERITAS Observations

VERITAS is a ground-based VHE gamma-ray observatory located at the Fred Lawrence Whipple Observatory in southern Arizona. It consists of four IACTs sensitive from approximately 100 GeV to above 30 TeV. Each VERITAS telescope has a 12 m tessellated reflector with a total area of 110 m 2 . Each camera's focal plane contains 499 closed-packed circular photomultiplier tubes, giving a total field of view (FoV) of $3^\circ 5'$. Gamma-rays incident onto the upper atmosphere induce a particle cascade, called an air shower, in which some charged particles have sufficient speed to emit Cherenkov light. The direction and energy of the original gamma ray can be reconstructed from images of the Cherenkov light recorded by the telescopes. When observing at small zenith angles ($< 40^\circ$), the array has an energy

Table 1
Summary of the VERITAS Data Sets Presented in This Work

Target	Period	Useful Duration (min)	Zenith Angle ($^\circ$)
V407 Cyg	2010 Mar 19–26	304	50–66
Crab Nebula	2010 Mar 12–16	203	55–65

resolution of 15% at 1 TeV and an angular resolution of better than 0.1° at 1 TeV (Acciari et al. 2008). For observations at LZA, the energy and angular resolution are degraded and the energy threshold is increased.

VERITAS observed V407 Cyg for several nights after the announcement of the *Fermi*-LAT detection, during days 9–16 of the outburst (2010 March 19–26). The zenith angle of these observations ranged between 50° and 66° . The VERITAS telescopes are regularly operated in a mode called wobble mode, during which the location of the object to be observed is offset from the center of the FoV by 0.5° , allowing for simultaneous source and background measurements (Fomin et al. 1994). The offset direction cycles between north, south, east, and west for sequential observing segments to reduce systematic effects. After filtering the data for contamination due to poor weather or instrumental problems, 304 minutes of live time remained from the original 335 minutes of observations; see Table 1.

To test the improved reconstruction technique discussed in Section 2.2, VERITAS observations of the Crab Nebula were also analyzed. We selected 203 minutes of good time intervals from 17 data segments taken on the Crab Nebula during 2010 March 12–16 (MJD 55267–55271) with similar zenith angles ranging from 55° to 65° . All data were analyzed using the standard analysis package for VERITAS data (Cogan 2007).

2.2. Event Reconstruction

The raw data were calibrated and cleaned, and quality selection criteria based on the number of photomultiplier tubes contained in the images and the position of the image in the camera were applied. The shape and orientation of the gamma-ray images were parameterized by their principal moments (Hillas 1985). In order to produce gamma-ray images of the sky, it is necessary to reconstruct the putative source location for each shower in the camera plane (hereafter “arrival direction”). When imaging showers with multiple IACTs, the arrival direction of a shower is usually found using simple geometric arguments. The major axes of the images produced by a shower in each IACT camera intersect near the location of the arrival direction. The shower arrival direction is calculated by minimizing the perpendicular distance to each image's semi-major axis, weighted by the size of each image. This method, here called the standard method, is effective at small zenith angles. However, at LZA, the major axes of the air shower images from an individual gamma-ray event are generally close to parallel. Thus, the uncertainty of the intersecting point increases, resulting in a loss of angular resolution. Due to this effect, a reconstruction technique that does not depend on the intersection of the axes is desirable for LZA observations.

The displacement method is a direction reconstruction algorithm that is useful for LZA observations (Şentürk et al. 2011). In very general terms, it consists of calculating the arrival direction using the shape and brightness of a given air shower image. More specifically, it relies on the determination of the *disp* parameter, defined as the angular distance from the image centroid to the arrival direction. This method was used by

Table 2
Selection Criteria Used for the VERITAS Analysis

Parameter	Selection Criteria
Image size	>400 digital counts (~ 75 photoelectrons)
Mean scaled width	$0.05 < \text{MSW} < 1.15$
Mean scaled length	$0.05 < \text{MSL} < 1.3$
Height of shower maximum	>7 km
θ	<0.1

Note. For an explanation of these parameters, see Section 2.3.

several experiments in the past (Lessard 2001; Kranich et al. 2003; Domingo-Santamaría et al. 2005), with varying ways of calculating *disp*.

The basis of the displacement method is the relationship of the *disp* parameter to other image parameters (Hofmann et al. 1999). The implementation of the algorithm in VERITAS is as follows: we estimate *disp* as a function of three other image parameters, *size*, *length*, and *width* (Hillas 1985), using Monte Carlo simulated gamma-ray showers. The method results in two different arrival directions, one on each side of a telescope image along the semi-major axis, also known as head–tail ambiguity (Hofmann et al. 1999). This ambiguity is eliminated by choosing the cluster of arrival directions closest to one another, one coming from each image. Finally, the arrival direction is estimated independently for each telescope image and an average weighted by *size* is taken. This method proves to be more powerful than the standard method (Cogan 2007) when reconstructing events with zenith angle larger than 50° . Quantitatively, an improvement of $\sim 30\%$ in detection significance for a source having 1% of the strength of the Crab Nebula has been observed.

2.3. Event Selection

The cosmic-ray background rate for IACTs is typically more than 10^4 times the gamma-ray rate, so it is important to reduce this background while retaining as many gamma-ray events as possible. By exploiting the differences in the development of gamma-ray and cosmic-ray-induced showers, the background due to cosmic rays can be reduced significantly, while still retaining a high fraction of gamma-ray-like events. The background reduction is performed by placing standard selection criteria, optimized using Monte Carlo simulations, and real data from the Crab Nebula on the shower image parameters. The selection criteria for the size of the telescope images, the mean scaled width and mean scaled length parameters (Daum 1997; Krawczynski et al. 2006), the height of maximum Cherenkov emission, and the angular distance from the anticipated source location to the reconstructed arrival direction of each shower (θ) are given in Table 2.

To perform a background subtraction of the surviving cosmic-ray events, an estimation of these background counts is made using the reflected-region background model (Aharonian et al. 2001). Events within an angular distance θ of the anticipated source location are considered ON events. Background measurements (OFF events) are taken from regions of the same size and at the same angular distance from the center of the FoV. For this analysis, a minimum of eight background regions was used. The excess number of events from the anticipated source location is found by subtracting the number of OFF events (scaled by the relative exposure, α) from the ON events. Statistical signifi-

cances are calculated using a modified version of Equation (17) of Li & Ma (1983) to allow for varying number of off-source regions due to nearby star 60 Cygni (Aharonian et al. 2004). More details about VERITAS, the calibration procedure, and the analysis techniques can be found in Acciari et al. (2008).

3. RESULTS

Analysis of the VERITAS data did not show a significant detection at the location of V407 Cyg. The results from both event reconstruction methods were used to calculate upper limits on the flux from V407 Cyg with the method described by Rolke et al. (2005) and the assumption of a Gaussian-distributed background.

The upper limits for V407 Cyg are calculated at the decorrelation energies of 1.8 TeV for the standard method and 1.6 TeV for the displacement method and assume that any emission takes the form of a power law with a photon index of -2.5 . The decorrelation energy is the energy at which the dependence of the upper limit calculation on the assumed photon index is minimized. This energy is found by performing multiple upper limit calculations, with different spectral indices, and determining the region where the resulting upper limit functions intersect. The energy threshold for the observations of V407 Cygni with VERITAS, defined as the maximum of the product of the assumed spectral shape and the effective area, is 1.2 TeV for both methods.

The analysis results for V407 Cyg are presented in Table 3 and Figure 1. In addition, results from observations of the Crab Nebula at similar zenith angles are presented in Table 3 and Figure 2. The efficiency of the displacement method for event reconstruction can be observed in the increase in both gamma-ray rate and significance for the Crab measurements. The increased sensitivity also results in a reduction of the upper limit for V407 Cyg compared to the standard method.

4. DISCUSSION

The GeV detection of V407 Cyg provides evidence for previously unobserved gamma-ray emission from novae in white dwarf/red giant systems. Expanding shock waves have been known to accelerate particles to high energies, and gamma rays are observed from supernova remnants. The discovery by the *Fermi*-LAT team, however, suggests that the same phenomenon occurs in some novae, adding a new class of gamma-ray emitting objects. The lack of a significant detection in the VHE band suggests that either particles were not accelerated to sufficient energies to produce VHE photons during the V407 Cyg outburst or that VHE photons were produced, but then absorbed.

The key to creating gamma rays is the acceleration of sufficiently energetic charged particles. In the case of V407 Cyg, the expanding matter from the nova collides with the stellar wind from the red giant and causes a shock, which accelerates the particles near the shock to relativistic energies. A rough estimation for the maximum energy attainable by first-order *Fermi* acceleration of a particle at a shock can be found, following the discussion of Longair (2011). If B is the magnetic flux density where a shock proceeds and the shock travels with velocity U , the maximum energy of a particle with charge Ze is $E_{\text{max}} = \int ZeBU^2 dt$, where t is the time allowed for particle acceleration. This means that the highest attainable energy is proportional to the magnetic field in which the nova travels, the square of shock speed, and the time for acceleration. The mean magnetic field in the shock can be

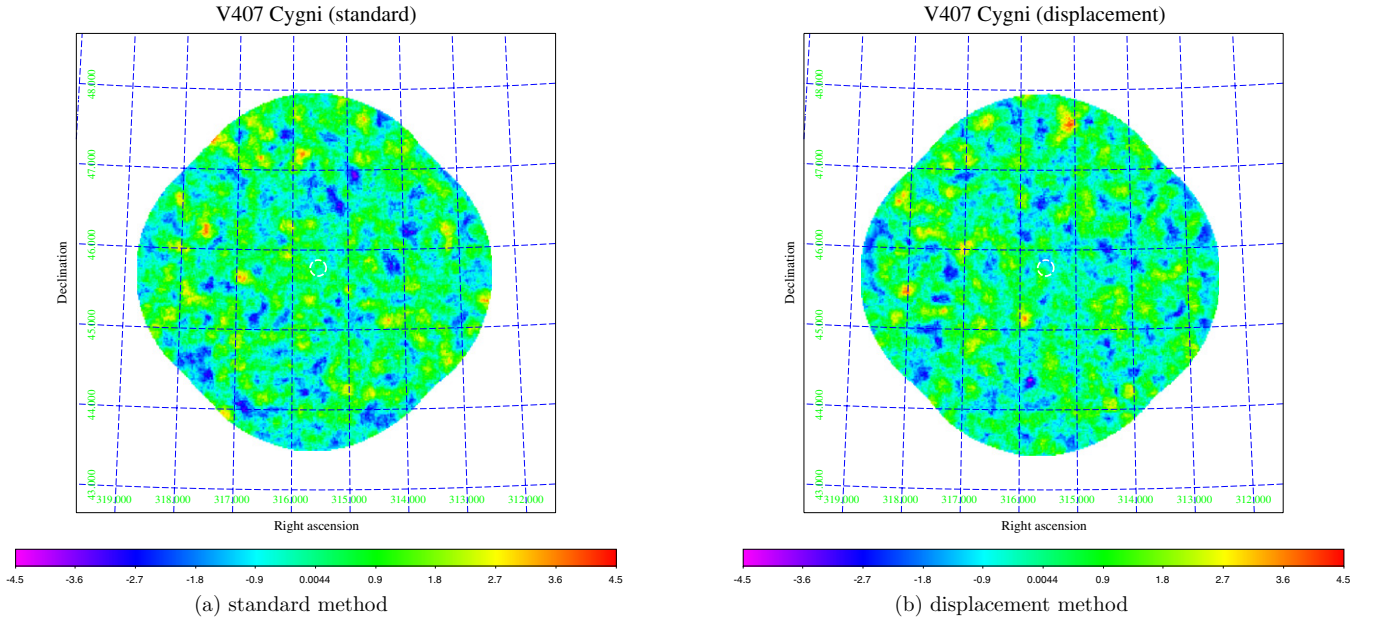


Figure 1. VERITAS significance maps of V407 Cyg reconstructed by the two methods. The white circle at the center indicates the location of the *Fermi*-LAT detection and the size of the source region used for the VERITAS analysis.

(A color version of this figure is available in the online journal.)

Table 3
A Summary of the VERITAS Analysis Results

Reconstruction Method	Standard		Displacement	
	Crab	V407 Cyg	Crab	V407 Cyg
Source				
Exposure (minutes)	203	304	203	304
ON (source) counts	255	91	300	76
OFF (background) counts	744	841	351	630
α (see Section 2.3)	0.111	0.125	0.111	0.125
Significance	15.7σ	0.5σ	25.1σ	1.0σ
Rate (photons minute ⁻¹)	0.97 ± 0.09	0.02 ± 0.04	1.40 ± 0.09	0.03 ± 0.03
Energy threshold (TeV)	1.5	1.2	1.7	1.2
Decorrelation energy (TeV)	...	1.8	...	1.6
Flux upper limit at decorrelation energy (95% c.l.) ($E^2 dF/dE$; erg cm ⁻² s ⁻¹)	...	2.7×10^{-12}	...	2.3×10^{-12}

Notes. Flux upper limits for V407 Cyg from VERITAS observations are calculated assuming a photon power-law index of -2.5 and taken at the decorrelation energy, minimizing the effect of the assumed spectral index. The displacement method produces both a higher significance for the Crab Nebula data and a more sensitive upper limit for V407 Cyg.

estimated as $B = [32\pi\rho(R)kT]^{1/2}$, where $\rho(R)$ is the density of gas molecules with respect to the distance from the center of the red giant and T is the wind temperature (Abdo et al. 2010). We assume a wind temperature of $T = 700$ K, corresponding to the temperature of the dust envelope measured by Munari et al. (1990). Espey & Crowley (2008) were able to directly measure the temperature of the red giant wind in a similar symbiotic system, EG And, and found that it can reach ~ 8000 K near the red giant. Using this wind temperature would increase the estimate for maximum particle energy by a factor of three.

Orlando & Drake (2012) carried out detailed hydrodynamic simulations of the V407 Cyg nova with various gas distribution models and could accurately reproduce the X-ray light curve of V407 Cygni. Their model for the distribution of gas that best reproduced the light curve included what they call a “circumbinary density enhancement,” a region of density exceeding the

typical R^{-2} profile of the stellar wind in the binary system, and had a binary separation of 15.5 AU (Orlando & Drake 2012). For the temporal profile of the nova shock velocity, we used the equation that Munari et al. (2010) found from fits to the broad components of the $H\alpha$ spectra they measured beginning at day 2.3 after the outburst (2010 March 13) and thereafter: $U = 4320 - 5440 \log t + 2635(\log t)^2 - 460(\log t)^3$. For the velocity between day 0 and day 2 of the outburst, we can assume two cases that bound the possible velocity profiles: (1) the nova shell experienced free expansion at a constant velocity before day 2, with the assumption that the mass collected by the nova shell during this period was small (free expansion model) or (2) extrapolate the above equation for the velocity to times before day 2 (extrapolation model). We then find that at the start of VERITAS observations of V407 Cyg (day 9 of the outburst), $E_{\max} \sim 1.4$ TeV for the free expansion model and ~ 3.0 TeV for

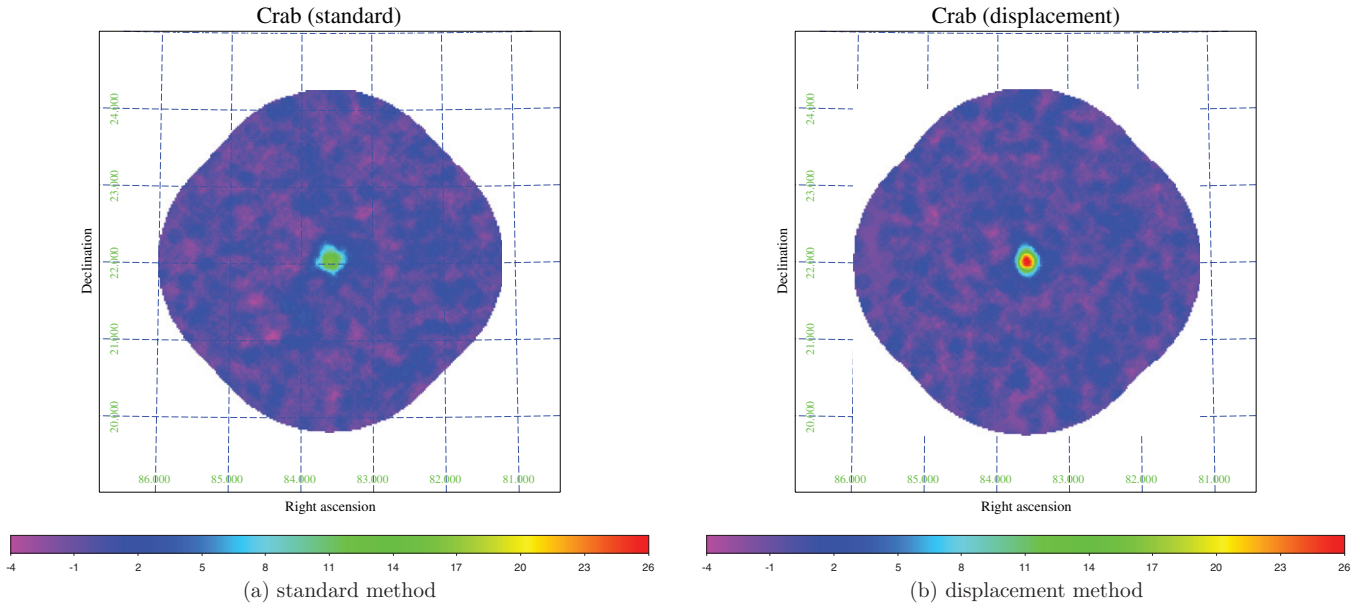


Figure 2. VERITAS significance maps of the Crab Nebula reconstructed by the two methods.

(A color version of this figure is available in the online journal.)

the extrapolation model. This suggests that particles could have been accelerated to TeV energies by the time of the VERITAS observations.

To check the importance of absorption, we calculated the opacity along the photon path for gamma rays generated at the shock front. Electron-positron pair production via photon-photon collision is the dominant interaction (Gould & Schröder 1967). We modeled the red giant spectrum as a black-body with a temperature of 2500 K and a radius of $500 R_{\odot}$, found the photon density as a function of position and energy following Nunez (2011), and used the cross section for the photon-photon collision from Gould & Schröder (1967). We found that the opacity for TeV photons only becomes significant when the TeV emission region is located directly behind the red giant with the system viewed edge-on. Though this case cannot be ruled out, it is statistically unlikely. In addition, if the suggested orbital parameters of Munari et al. (1990) are accurate, the system is unlikely to have been in such an orientation at the time of the nova.

The upper limits placed by VERITAS can put some restrictions on the gamma-ray emission mechanism in V407 Cyg. Two physical models of gamma-ray production at the shock front have been suggested (Abdo et al. 2010). In the hadronic model, gamma rays are produced in the decay of π^0 particles generated by collisions of high-energy protons accelerated in the shock. In the leptonic model, gamma rays are produced via inverse-Compton scattering of infrared photons emitted from the red giant on high-energy electrons accelerated in the shock.

The electron threshold energy for the production of a gamma-ray photon via inverse-Compton scattering off the red giant photons can be estimated as $E > (E_{\gamma}/2)(1 + \sqrt{1 + m_e^2 c^4/E_{\gamma} \epsilon})$, where E is the electron threshold energy, E_{γ} is the gamma-ray energy, and ϵ is the energy of the red giant photons. The electron threshold energy for a 1 TeV gamma-ray scattering off 0.6 eV photons at the peak of the red giant spectrum is 1.1 TeV. Though the above calculation indicates that particles could reach TeV energies if continuously accelerated for the full nine days from the initial outburst to the start of the VERITAS observations, the inverse-Compton cooling time would be significantly less than a

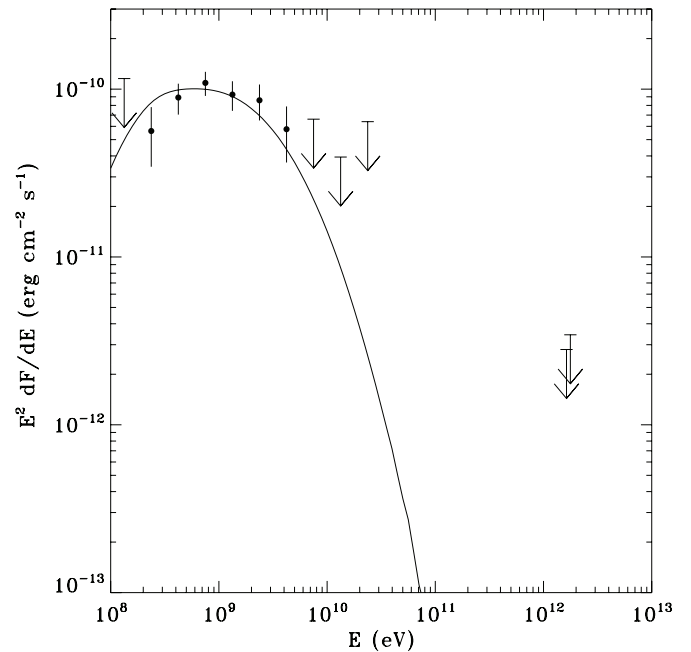


Figure 3. Spectrum of V407 Cyg (FGL J2102+4542) measured by the *Fermi*-LAT (Abdo et al. 2010) and VERITAS upper limits. Vertical bars indicate 1σ statistical error, and arrows indicate 2σ upper limit. The rightmost arrows show the 99% confidence level (3σ) VERITAS upper limit calculated using the displacement method (at 1.6 TeV) and the standard method (at 1.8 TeV) for event reconstruction (see Section 3). The fitting curve was constructed with the method of Kamae et al. (2006) with the parameters ($s_p, \log(E_{cp})$) = (2.15, 1.5) (see Section 4).

day (the time estimated by Abdo et al. 2010 for 5 GeV electrons), meaning electrons that are accelerated in the first few days of the outburst would not likely retain sufficient energy to produce VHE photons by the time of the VERITAS observations. VHE emission near the time of the VERITAS observations would therefore require freshly accelerated particles, however, recently accelerated particles would likely not have enough time to reach TeV energies. If electrons did reach TeV energies, they would be approaching the Klein–Nishina regime, where the cross section

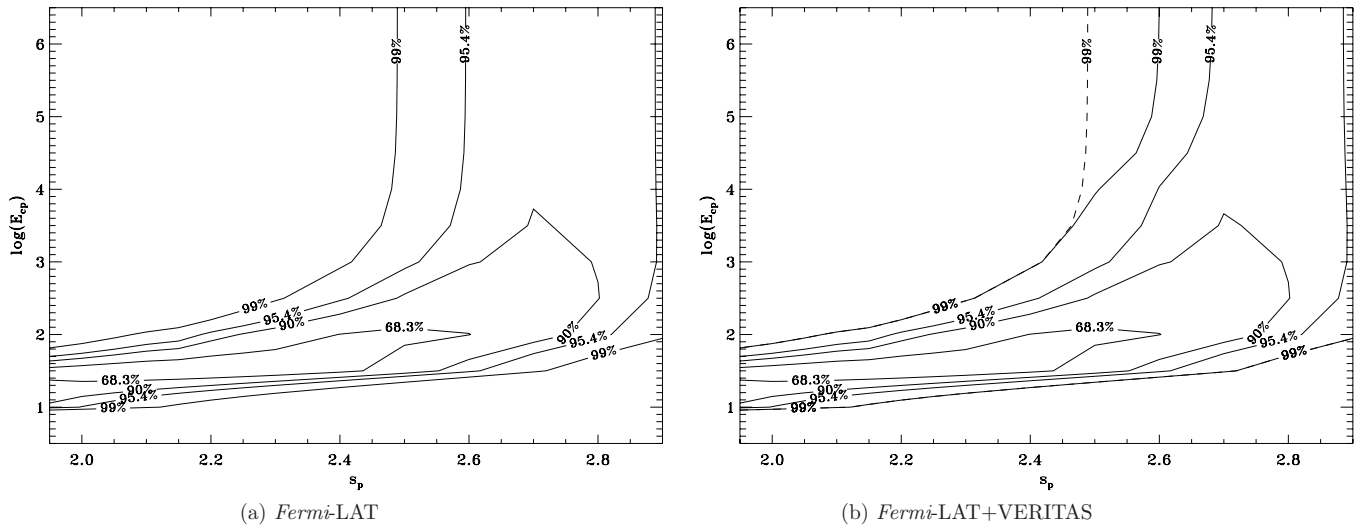


Figure 4. Confidence region maps for the hadronic model of gamma-ray production in V407 Cyg constrained by *Fermi*-LAT data alone (a) and constrained by the *Fermi*-LAT data plus the VERITAS upper limit (b). The x-axis gives the spectral slope (s_p) and the y-axis is the logarithm of the cutoff energy (E_{cp} in GeV). The numerical values correspond to the confidence level of the fit with the given values of the parameters. The 99% contour from *Fermi*-LAT data alone is shown as a dashed line in (b) for comparison.

for inverse-Compton scattering would be reduced, resulting in a longer cooling time. However, electrons with TeV energies would be well above the exponential cutoff, $3.2^{+2.6}_{-0.1}$ GeV, of the electron spectrum in the best-fit leptonic model of Abdo et al. (2010). These two factors imply that the VERITAS upper limits place no new constraints on leptonic models.

For hadronic models, the *Fermi*-LAT data alone provide relatively poor constraints on the extension of the proton spectrum to high energies. In particular, the cutoff energy (E_{cp}) is not well bounded from above if the spectral index is steep. The VERITAS data can be used to improve the constraints on the hadronic model parameters. To do so, we re-fit the hadronic model used by Abdo et al. (2010) to the *Fermi*-LAT points with the addition of the VERITAS upper limit. Figure 3 shows the *Fermi*-LAT data (Abdo et al. 2010) and the VERITAS flux upper limit compared to the best fitting hadronic model. The gamma-ray spectrum is calculated via the method of Kamae et al. (2006), assuming a cosmic proton spectrum of the following form: $N_p = N_{p,0}(W_p + m_p c^2)^{-s_p} e^{-W_p/E_{cp}}$ (protons GeV^{-1}), where (W_p , E_{cp} , and m_p are kinetic energy, cutoff energy, and mass of the proton and s_p is the power-law index).

Figure 4 shows a confidence region map for the parameters of the hadronic model using both the *Fermi*-LAT and VERITAS data. The gamma-ray spectrum was modeled as described above and fitted to the *Fermi*-LAT data by varying s_p and E_{cp} . The spectrum was then compared to the VERITAS upper limit, and a χ^2 value for the VERITAS data point alone was calculated. Specifically, we calculated the model flux in the VERITAS energy band and compared this to the flux upper limit determined via the displacement method. This χ^2 value was then added to the χ^2 calculated for the *Fermi*-LAT data. The confidence levels were then calculated for the two parameters of interest, s_p and E_{cp} . As can be seen from Figure 4, the VERITAS observations place greater restrictions on the model proton spectral index for high cutoff energies. The 90% confidence limits are $E_{cp} \lesssim 5$ TeV (comparable to $E_{\text{max}} = 3$ TeV, calculated above for the extrapolation model) and $E_{cp} \gtrsim 0.01$ TeV (much lower than E_{max}). It is possible that the peak energy of the particles produced by the shock could be reduced if the magnetic field is weaker than estimated above. Nelson et al. (2012)

argue for a larger binary separation of 20–25 AU, based on the presence of lithium burning in the Mira, and evidence that the white dwarf in the system is massive. Using this larger separation distance would lead to weaker magnetic fields. The limits placed by the VERITAS observations are near the threshold for the observations, so it is also possible that simply not enough particles were accelerated to high enough energies to produce a significant detection by VERITAS.

We thank Pierre Jean of the *Fermi*-LAT team for providing the contour data from the *Fermi*-LAT results and useful discussion. This research is supported by grants from the U.S. Department of Energy Office of Science, the U.S. National Science Foundation, and the Smithsonian Institution; by NSERC in Canada, by Science Foundation Ireland (SFI 10/RFP/AST2748); and by STFC in the UK. We acknowledge the excellent work of the technical support staff at the Fred Lawrence Whipple Observatory and at the collaborating institutions in the construction and operation of the instrument.

REFERENCES

- Abdo, A. A., Ackermann, M., Ajello, M., et al. 2010, *Science*, **329**, 817
- Acciari, V. A., Beilicke, M., Blaylock, G., et al. (VERITAS Collaboration) 2008, *ApJ*, **679**, 1427
- Aharonian, F., Akhperjanian, A., Barrio, J., et al. (HEGRA Collaboration) 2001, *A&A*, **370**, 112
- Aharonian, F., Akhperjanian, A., Beilicke, M., et al. 2004, *A&A*, **421**, 529
- Bode, M., O’Brien, T. J., Osborne, J. P., et al. 2006, *ApJ*, **652**, 629
- Chaty, S. 2012, in Proc. of the Twelfth Marcel Grossmann Meeting on General Relativity, ed. T. Damour, R. T. Jantzen, & R. Ruffini (Singapore: World Scientific), 1001
- Cheung, C. C., Donato, D., Wallace, E., et al. 2010, ATel. #2487
- Cogan, P. 2007, in Proc. 30th International Cosmic Ray Conference, Mérida, Vol. 3, ed. R. Caballero, J. C. D’Olivo, G. Medina-Tanco, L. Nellen, F. A. Sánchez, & J. F. Valdés-Galicia (Mexico City, Mexico: Universidad Nacional Autónoma de México), 1385
- Daum, A., Hermann, G., Hess, M., et al. 1997, *Astropart. Phys.*, **8**, 1
- Domingo-Santamaría, E., et al. 2005, in Proc. of the 29th Int. Cosmic Ray Conf., ed. B. Sripathi Acharya, S. Gupta, P. Jagadeesan et al. (Mumbai: Tata Institute of Fundamental Research), 363
- Espey, B. R., & Crowley, C. 2008, in ASP Conf. Ser. 401, RS Ophiuchi (2006) and the Recurrent Nova Phenomenon, ed. A. Evans, M. F. Bode, T. J. O’Brien, & M. J. Darnley (San Francisco, CA: ASP), 166

- Fomin, V. P., Stepanian, A. A., Lamb, R. C., et al. 1994, *Astropart. Phys.*, **2**, 137
- Gould, R. J., & Schröder, G. P. 1966, *Phys. Rev. Lett.*, **16**, 252
- Hays, E. A., & Fermi LAT Collaboration 2010, *BAAS*, **41**, 670
- Hillas, A. M. 1985, in Proc. of the 19th Int. Cosmic Ray Conf. (La Jolla), **3**, 445 (<http://catalog.hathitrust.org/Record/006239301>)
- Hofmann, W., Jung, I., Konopelko, A., et al. 1999, *Astropart. Phys.*, **122**, 135
- Kamae, T., Karlsson, N., Mizuno, T., Abe, T., & Koi, T. 2006, *ApJ*, **647**, 692
- Kolotilov, E. A., Munari, U., Popova, A. A., et al. 1998, *Astron. Lett.*, **24**, 451
- Kranich, D., Stark, L. S., & for the HEGRA Collaboration 2003, in Proc. of the 28th Int. Cosmic Ray Conf., ed. T. Kajita, Y. Asaoka, A. Kawachi, Y. Matsubara, & M. Sasaki (Tsukuba), 3023
- Krawczynski, H., Carter-Lewis, D. A., Duke, C., et al. 2006, *Astropart. Phys.*, **25**, 380
- Lessard, R. W., Buckley, J. H., Connaughton, V., & Le Bohec, S. 2001, *Astropart. Phys.*, **15**, 1
- Li, T.-P., & Ma, Y.-Q. 1983, *ApJ*, **272**, L317
- Longair, M. S. 2011, *High Energy Astrophysics* (3rd ed.; Cambridge: Cambridge Univ. Press)
- Munari, U., Joshi, V. H., Ashok, N. M., et al. 2011, *MNRAS*, **410**, L52
- Munari, U., Margoni, R., & Stagni, R. 1990, *MNRAS*, **242**, 653
- Nelson, T., Donato, D., Mukai, K., Sokoloski, J., & Chomiuk, L. 2012, *ApJ*, **748**, 43
- Nishiyama, K., & Kabashima, F. 2010, International Astronomical Union Central Bureau for Astronomical Telegrams, reported by H. Maehara, No. 2199
- Ñunez, P. D., LeBohec, S., & Vincent, S. 2011, *ApJ*, **731**, 105
- Orlando, S., & Drake, J. J. 2012, *MNRAS*, **419**, 3
- Rolke, W. A., López, A. M., & Conrad, J. 2005, *Nucl. Instrum. Methods Phys. Res. A*, **551**, 493
- Sabatini, S., Tavani, M., Striani, E., et al. 2010, *ApJ*, **712**, L10
- Şentürk, G. D., et al. 2011, in Proc. of the 32nd Int. Cosmic Ray Conf. (Beijing), **9**, 126 (<http://www.ihep.ac.cn/english/conference/icrc2011/paper/proc/v9.pdf>)
- Sokoloski, J. L., Luna, G. J. M., Mukai, K., & Kenyon, S. J. 2006, *Nature*, **442**, 276
- Tatischeff, V., & Hernanz, M. 2007, *ApJ*, **663**, L101
- Tavani, M., Mukherjee, R., Mattox, J. R., et al. 1997, *ApJ*, **479**, L109
- Vandenbrouke, J., Buehler, R., Ajello, M., et al. 2010, *ApJ*, **718**, L166