Structural and magnetic properties of Cu substituted manganites studied by EXAFS and dc magnetization measurements

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**A B S T R A C T**

We report extended X-ray absorption fine structure (EXAFS) measurements at the Mn K-edge and dc magnetic measurements performed on (La\textsubscript{1−x}Ca\textsubscript{x})\textsubscript{3}Mn\textsubscript{1−y}Cu\textsubscript{y}O\textsubscript{3} samples (x=0.37 and 0.75 and y=0.03 or 0.08), thus comparing the effects determined by the partial Mn substitution with Jahn–Teller Cu\textsuperscript{2+} ions with those induced by the already studied analogous substitution of Mn with non-Jahn-Teller Cr\textsuperscript{3+}.

Hence, in the present paper we complete our study on the influence of chemical Mn site substitutions on the nanoscale phase inhomogeneity characteristic of these strongly correlated electron systems. We observe that the Cu\textsuperscript{2+} substitution determines two very different effects in the two compositions belonging from the opposite sides of the phase diagram. In the x=0.37 case it weakens the ferromagnetic ordering thus favouring the antiferromagnetic charge ordered phase component, while in the x=0.75 one it does not significantly modify the charge ordering already present in the system, differently from what previously observed in presence of non-Jahn-Teller Cr\textsuperscript{3+} substitution.

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1. Introduction

Manganese perovskites (La\textsubscript{1−x}Ca\textsubscript{x})MnO\textsubscript{3} present a strong interplay between charge, lattice and spin degrees of freedom. Their phase diagram is characterized, in the range 0.20 < x < 0.50, by a transition from a paramagnetic insulating phase to a ferromagnetic metallic (FM) one upon cooling [1]. Such a coupling between a magnetic phase transition and a conductive one is the basis for the so called negative colossal magnetoresistance effect (CMR) also higher than 90% in presence of a magnetic field

A significant role seems also to be played by a strong spin-lattice coupling displayed by the presence around the Curie temperature $T_C$ of magnetic polarons of about 12 Å, namely local magnetically ordered regions in an insulating matrix where the charge carrier is trapped like in a polaronic state [4]. In this last framework, experimental evidences of an inhomogeneous phase above and below $T_C$ where nanometric metallic and insulating domains coexist and percolate are reported in literature [4–7]. These research developments on manganites nanoscopic phase texture also point out the occurrence of further physical properties like non-volatile magnetoresistive memory at low magnetic fields and at specific chemical compositions which could allow to continue to miniaturize transistors and electronic components useful for computer engineering and spintronics applications [8]. Therefore the debate about the homogeneity of the low temperature phase and its outcomes at all Ca content is still open.

The present main research topic on these compounds is the identification of the real nature of the insulating phase component for 0.20 < x < 0.50 at low T in the phase separation framework; it could be a CO phase with a very low coherence length and therefore not detectable by lots of techniques, becoming more extended and then noticeable for x > 0.50. In this general scenario we can reasonably assume that the partial substitution of the Mn magnetic ion could...
contribute to an enhancement of one phase component over the other [9–11]. This result should be obtained due to the strong coupling between CO and Jahn–Teller lattice distortion (induced by the presence of the Mn$^{3+}$ ion or of its Jahn–Teller substitutes), as already evidenced by our EXAFS measurements on (La$_{1-x}$Ca$_x$)$_2$MnO$_3$ with $x > 0.50$ [12]. Another relevant effect on the conduction bandwidth and then on the magnetic exchange interaction along the partially substituted Mn$^{3+}$–O$^2$–M paths can be determined by the different ionic radius of the $B$ site element in the perovskite $ABO_3$ structure and by the 3d electronic configuration of the substituting ions (note the following ionic radii in octahedral coordination: Mn$^{3+} = 0.645$ Å; Mn$^{4+} = 0.53$ Å; Cr$^{3+} = 0.615$ Å; Ni$^{2+} = 0.69$ Å; Cu$^{2+} = 0.73$ Å).

In a previous paper [13] we pointed out how Mn-substituting Cr$^{3+}$ or Ni$^{2+}$ favour strengthening of the metallic-ferromagnetic phase component for $x > 0.50$ since they are not associated with a Jahn–Teller distortion and act as random impurities. In this earlier case the resulting disordered Jahn–Teller distortion plays the role of a site-random or effective magnetic field weakening the Jahn–Teller CO phase [9,10].

In the present paper we report the EXAFS study of local order and structure around the manganese ion on Cu$^{2+}$ substituted samples (La$_{0.63}$Ca$_{0.37}$)$_2$Mn$_{0.92}$Cu$_{0.08}$O$_3$, (La$_{0.25}$Ca$_{0.75}$)$_2$Mn$_{0.92}$Cu$_{0.08}$O$_3$ and (La$_{0.25}$Ca$_{0.75}$)$_2$Mn$_{0.97}$Cu$_{0.03}$O$_3$. Substituting copper ions, in contrast to Cr$^{3+}$ or Ni$^{2+}$, should be able to strengthen the CO phase component since they are associated to a lattice Jahn–Teller distortion (Cu$^{2+}$ is characterized by a $d_{3z^2}$ Jahn–Teller electronic configuration), thus completing the Mn ion site substitution scenario. We expect to observe a bigger effect in the (La$_{0.63}$Ca$_{0.37}$)$_2$Mn$_{0.92}$Cu$_{0.08}$O$_3$ sample, considering that the pristine compound (La$_{0.63}$Ca$_{0.37}$)$_2$MnO$_3$ is metallic and ferromagnetic below 270 K, while on the other side of the phase diagram (La$_{0.25}$Ca$_{0.75}$)$_2$Mn$_{0.97}$Cu$_{0.03}$O$_3$ exhibits CO below $T_C = 225$ K.

All the samples were characterized by dc magnetization measurements and X-ray powder diffraction revealing an orthorhombic $Pnma$ crystalline structure, and thereafter extended X-ray absorption fine structure (EXAFS) spectroscopy measurements were performed between 20 and 300 K at the Mn K absorption edge. Furthermore the same samples were also investigated by neutron powder diffraction (NPD), scanning electron microscopy (SEM), transmission electron microscopy (TEM) and anelastic spectroscopy in order to investigate both the additional lattice disorder and the Mn average valence variation induced by Cu$^{2+}$ doping.

2. Experimental method and data analysis

The studied samples were pellets of polycrystalline (La$_{0.63}$Ca$_{0.37}$)$_2$Mn$_{0.92}$Cu$_{0.08}$O$_3$, (La$_{0.63}$Ca$_{0.37}$)$_2$Mn$_{0.97}$Cu$_{0.03}$O$_3$ and (La$_{0.25}$Ca$_{0.75}$)$_2$Mn$_{0.92}$Cu$_{0.08}$O$_3$ powders dispersed in a cellulose matrix optimising the jump at the absorption edge. Polycrystalline (La$_{1-x}$Ca$_x$)$_2$Mn$_{1-x}$Cu$_x$O$_3$ samples were prepared by means of a solid state reaction: binary oxides (CaO 99.95% ALDRICH; La$_2$O$_3$ 99.99% A.E.S.A., Mn$_2$O$_3$ 99.99% A.E.S.A., CuO 99.98% A.E.S.A); all the samples were finally sieved (80 mesh), then mixed in stoichiometric amounts and finally reacted at high temperature in air. Four thermal treatments with intermediate grinding were carried out, the first at 1523 K for 15 h and the remaining at 1603 K for 18 h.

EXAFS data have been recorded in transmission geometry at the beamline E4 of the HASYLAB synchrotron radiation source (Hamburg) using a double-crystal Si(1 1 1) monochromator and a helium flux cryostat. The EXAFS oscillations were extracted from the experimental data using standard procedures and were normalized using the Lengeler–Eschenberger method (Fig. 1) [14].

The $k^2$ weighted $\chi(k)$ data were Fourier transformed (FFT) in the $k$ range between 2.68 and 15.48 A$^{-1}$ (Fig. 2). The obtained Fourier transform shows many correlating peaks up to $R = 8$ Å. The analysis was restricted to the first Mn–O coordination shell by inverse Fourier transforming the first FFT peak in the $R$ range between 0.98 and 2.26 Å. The so-obtained filtered data were fitted using the standard single scattering EXAFS formula and FEFF8.2 theoretical amplitude and phase functions [15]. Starting from the orthorhombic unit cell model typical of Jahn–Teller distorted manganites the FEFF code computes overlapping muffin tin spherically symmetric potentials and their corresponding phase-shifts thus generating the different Mn–O path contributions in order to fit the data.

We have performed two alternative fitting procedures: at first using only one Mn–O first shell with coordination number $N_{Mn-O} = 6$; afterwards using two sub-shells, one corresponding to a Mn–O coordination number $N_{Mn-O} = 4$ and the other corresponding to $N_{Mn-O} = 2$. In particular, we considered four longer distances $R_{Mn-O} = 2.10$ Å and two shorter $R_{Mn-O} = 1.93$ Å at 300 K for the $x = 0.75$ Cu-substituted samples [12,13], vice versa four shorter and two longer bond lengths for the $x = 0.37$ case. These starting models are the same already reported in previous papers for the corresponding un-substituted compounds [12,13].

The two sub-shells fit has been possible thanks to the extension of the data in the $k$ range, the window chosen in the FFT space, the good signal to noise ratio, which enabled a sufficient number of independent fitting parameters to be used. The fitting parameters were the Mn–O mean square relative displacements $\sigma^2$ and the interatomic bond lengths $R_{Mn-O}$ both in the one shell and in the two shells analysis. The coordination numbers were fixed to the previously cited values.

In order to correctly compare the parameters behaviour as a function of temperature for the substituted and for the corresponding pristine samples, all the
fitting criteria in the data analysis performed on the samples were identical to those reported in previous papers [12,13].

The dc magnetic measurements were performed by using a commercial magnetometer (MPMS by Quantum Design) on the same sample powders.

3. Results and discussion

We present a comparison of the local structure parameters around the Mn ions and magnetic properties, obtained by the EXAFS and dc magnetization studies respectively, in the Cu-substituted and pristine compounds for the cited Ca concentrations thus characterizing also on a local scale the Mn substitution effects on polarons and CO formation.

Zero field cooled (ZFC) and FC susceptibility measurements performed on both pristine and substituted samples, with $x = 0.37$ and $x = 0.75$ are shown in Figs. 3 and 4, respectively.

In the (La0.63Ca0.37)Mn0.92Cu0.08O3 sample the Mn/Cu substitution effect is extremely marked: copper ions induce a strong distortion of the surrounding lattice both due to their Jahn–Teller nature and to their greater dimensions respect to the Mn ones. These distortions frustrate the double exchange magnetic interactions between the neighboring Mn ions thus reducing the $e_g$ electron conduction bandwidth and favouring the Mn ions superexchange interaction associated to the antiferromagnetic CO phase component (see the $m(T)$ behaviour in Fig. 3) [16]. In addition to local distortions, Cu$^{2+}$ substitution induces a re-equilibration of the $[\text{Mn}^{3+}]/[\text{Mn}^{4+}]$ ratio determining a charge transfer effect. For example the 8% Cu$^{2+}$ content increases the average Mn valence from 3.37 in (La0.63Ca0.37)MnO3 to 3.49 in (La0.63Ca0.37)(Mn0.92Cu0.08)O3 approaching the boundary between the FM and CO long-range phases [1].

The described magnetic outcome is much more evident in (La0.63Ca0.37)(Mn0.92Cu0.08)O3, where we start from a FM macroscopic phase before substitution, while it is less marked in the (La0.25Ca0.75)(Mn0.92Cu0.08)O3 with respect to the sudden decrease evidenced in the (La0.63Ca0.37)MnO3 sample [17,18] and to a Debye–Wall effect. This striking effect is strictly connected to an enhanced CO phase, going towards the Jahn–Teller lattice distortion is favoured in presence of Jahn–Teller substituting ions like Cu$^{2+}$.

We describe now the behaviour of the structural EXAFS parameters, firstly focusing our attention on the Mn–O single shell results. The best-fit parameters values are reported in Table 1 for the (La0.63Ca0.37)Mn0.92Cu0.08O3 and (La0.25Ca0.75)Mn0.92Cu0.08O3 samples at $T = 300$ K.

In Fig. 5 the local order parameter $\sigma^2$ of the MnO6 coordination shell is reported as a function of temperature for the substituted sample (La0.63Ca0.37)Mn0.92Cu0.08O3 (empty squares) and for the pristine (La0.63Ca0.37)MnO3 and (La0.25Ca0.75)(Mn0.92)O3 samples (empty triangles and full circles) which respectively belong to the long-range CO and FM respectively.

The EXAFS results before and after the Cu substitution show a clear enhancement of the local disorder as a function of $T$ in (La0.63Ca0.37)(Mn0.92Cu0.08)O3 with respect to the sudden decrease evidenced in the (La0.63Ca0.37)MnO3 sample [17,18] and to a Debye-like trend (continuous line). This striking effect is strictly connected to an enhanced CO phase, going towards the $\sigma^2$ behaviour of the charge ordered (La0.25Ca0.75)(Mn0.92)O3 sample [12]. We point out how the effect induced by the Jahn–Teller Cu$^{2+}$ substituting ion is extremely different from that induced on the same sample com-

### Table 1

<table>
<thead>
<tr>
<th>Fit sub-shell</th>
<th>$(\text{La}<em>{0.63}\text{Ca}</em>{0.37})[\text{Mn}<em>{0.92}\text{Cu}</em>{0.08}]_{\text{O}_3}$</th>
<th>$(\text{La}<em>{0.25}\text{Ca}</em>{0.75})<em>{\text{Mn}</em>{0.92}\text{Cu}<em>{0.08}}</em>{\text{O}_3}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Mn–O</td>
<td>4</td>
<td>$(0.85(1), 0.85(1))$</td>
</tr>
<tr>
<td>Mn–O</td>
<td>2</td>
<td>$(2.03(1), 2.03(1))$</td>
</tr>
<tr>
<td>Mn–O</td>
<td>6</td>
<td>$(1.98(1), 1.98(1))$</td>
</tr>
<tr>
<td>Mn–O single-shell</td>
<td></td>
<td>$(0.85(1), 2.03(1))$</td>
</tr>
</tbody>
</table>
position by non-Jahn–Teller Cr$^{3+}$ and Ni$^{2+}$, where we revealed (see Ref. [13]) only a progressive broadening and lowering of the PM–FM transition determined by the fact that the citric species do not participate in double exchange. Moreover, we have also to consider that the cited charge transfer effect induced by the Mn substitution with a divalent ion is not sufficient to determine the huge effect observed both in magnetic and structural properties only in the Jahn–Teller Cu$^{2+}$ case and not in the Ni$^{2+}$ one [13]. Therefore the local Jahn–Teller distortion coupled to the CO phase component and the greater internal chemical pressure, both influenced by the Cu$^{2+}$ substitution, appears to be determinant to interpret the obtained results.

In (La0.63Ca0.37)(Mn0.92Cu0.08)O3 the $\sigma^2$ parameter follows a temperature dependent Debye-like behaviour $\sigma_D^2$ at high temperatures ($T \geq 225$ K) with a Debye temperature $\theta_D$ around 500 K, plus a temperature independent static contribution $\sigma_S^2 = 0.00183$ Å$^2$ due to the presence of a Jahn–Teller distorted environment. The high temperature Debye–Waller factor shown in Fig. 5 has been fitted considering $\sigma^2(T) = \sigma_D^2 + \sigma_S^2$, as already explained in previous papers [12,13], where the temperature dependent dynamic contribution $\sigma_D^2$ is expressed as an infinite series whose first and most important terms are [19]:

$$\sigma_D^2 = \frac{3h^2}{M \theta_D^2} \left[ \frac{1}{4} + \frac{T}{\theta_D} \int_0^{\theta_D/T} dx \frac{x^4}{e^{x-1}} \right]$$

(1)

where $M$ is the mass of the diffuser atom. The fitting parameters were the static contribution $\sigma_S^2$ and the Debye temperature $\theta_D$, found to be about 510 K, in accordance with previous experiments [12,13]. Otherwise below $T_C$ the $\sigma^2(T)$ behaviour is characteristic of a CO-like phase characterized by a breaking of the high temperature crystal symmetries [12,13].

For what concerns the (La0.25Ca0.75)(Mn0.92Cu0.08)O3 and (La0.25Ca0.75)(Mn0.92Cr0.08)O3 samples we evidence in Fig. 5, after comparison with the charge ordered (La0.25Ca0.75)MnO$_3$, the local order parameter $\sigma^2$ substantially remains charge ordering-like on a long-range scale, in agreement with the magnetization results. Clearly, in the Jahn–Teller Cu$^{2+}$ substitution case the induced random impurity effect is not sufficient to reduce CO at a local level as occurs with non-Jahn–Teller Ni$^{2+}$ or Cr$^{3+}$ ions in the $x > 0.50$ Ca-content range [13].

Moreover in all the samples we suggest the two Mn–O interatomic distances as an alternative model, which anyway does not exclude the single distance one [12]. Therefore, as published in previous papers [12,13], we distinguish a bimodal distribution of Mn–O interatomic bond lengths, with a fit residual about 35% lower at all temperatures than that obtained by the one shell analysis, obtaining an improvement of the fit at each temperature respect to the single shell model (see Ref. [13]). In this framework, the quality of the fit with one or two Mn–O distances was checked by means of the residual $\rho = \sum k_i^2 |x_{\text{exp}}(k_i) - x_{\text{theo}}(k_i)|^2 / \sum k_i^2 |x_{\text{exp}}(k_i)|^2$, where $x_{\text{exp}}(k_i)$ and $x_{\text{theo}}(k_i)$ are the experimental and the calculated EXAFS signals at the wave vector $k_i$. The Mn–O interatomic distances thus obtained for the sample (La0.25Ca0.75)(Mn0.92Cu0.08)O3 are reported in Fig. 7 and can be compared with the chromium doped (La0.25Ca0.75)(Mn0.92Cr0.08)O3 and un-doped (La0.25Ca0.75)MnO$_3$ cases presented in Ref. [13]. While the partial substitution of Mn with Cr reduced the difference between longer and shorter Mn–O distances respect to the pristine compound [13], denoting a weakening of the Jahn–Teller distortion and of the CO phase coherence, the John–Teller Cu substitution determines an increasing of the longer Mn–O distance and therefore accentuates the divergence between the two different sub-shell bond lengths (see also Table 1). This result confirms the presence of a local Jahn–Teller distorted and insulating phase component also.
In summary, we reveal a strong correlation between the local order and structure studied by EXAFS and the magnetization behaviour and therefore we confirm a relation between charge localization (connected to local disorder), and magnetic degrees of freedom. We emphasize that a local probe like EXAFS provides information on systems that are only apparently homogeneous revealing, as in the present case, the influence of the Mn partial substitution on peculiar local scale features. Moreover in the present case, further EXAFS measurements at the Cu K-edge would be interesting. This should be one way to probe the Cu valence and to determine quantitatively the extent of distortions about the copper site.

4. Conclusions

We measured the EXAFS spectra and dc magnetic susceptibility of \((\La_{0.63} \Ca_{0.37})_3(\Mn_{0.92} \Cu_{0.08})_3\) and \((\La_{0.25} \Ca_{0.75})_3(\Mn_{0.92} \Cu_{0.08})_3\) manganites, between 20 and 300 K. We found that the Cu substitution weakens in a relevant and outstanding way the FM long-range phase in the compound with \(x < 0.50\) while the CO phase in the second class of samples \(x > 0.50\) results to be only slightly conditioned. Moreover, the local structure parameters still confirm in all the studied Cu-substituted samples the presence of the same anomalies characteristic of the CO pristine compounds, like the Debye–Waller factor plateau below \(T_{CO}\) and the Jahn–Teller splitting of the Mn–O distances in the MnO6 octahedra.

The results we have obtained verify the natural tendency to phase separation in manganites and the possibility in the Cu substitution case, as already happened after Cr and Ni substitutions [13], to tune the relative weight of one phase component (CO) respect to the other (FM) in this characteristic nanometric phase percolation scenario [9,13].

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