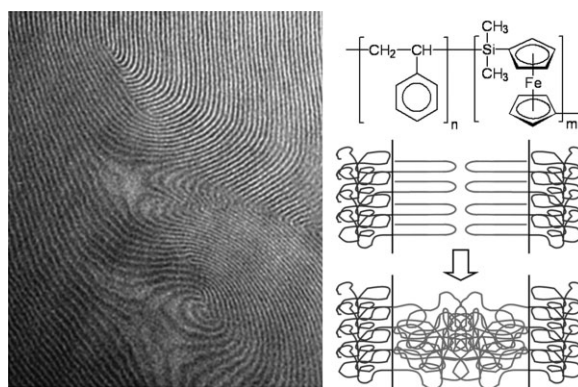


Equilibrium Melting Temperature of Poly(ferrocenyl dimethylsilane) in Homopolymers and Lamellar Diblock Copolymers with Polystyrene^a

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The equilibrium melting temperature, $T_{m,0}$, of poly(ferrocenyl dimethylsilane) (PFDMS) is re-investigated using two PFDMS homopolymers and a PS-*b*-PFDMS copolymer (PS: polystyrene). The melting temperatures of samples crystallized at different crystallization temperatures, T_c , were found to depend in a nonlinear fashion on T_c . Consequently, the Hoffmann–Weeks approach fails in determining $T_{m,0}$. The more reliable Gibbs–Thomson approach results in a $T_{m,0} = (215 \pm 11)^\circ\text{C}$ for the PFDMS homopolymers, and a $T_{m,0} = (179 \pm 5)^\circ\text{C}$ for the PFDMS microphases in the lamellar PS-*b*-PFDMS diblock copolymer.



Introduction

The metal-containing polymer poly(ferrocenyl dimethylsilane) (PFDMS) is currently attracting much attention because of its promising optoelectronic and redox proper-

ties, and because it represents a valuable precursor for ceramics and magnetic materials.^[1–7] Moreover, it counts among the very few metallopolymers accessible by living chain-growth protocols which give access not only to products of defined molar mass and narrow polydispersity but also to a broad variety of copolymer architectures: plenty of PFDMS block copolymers with, e.g., polystyrene (PS),^[8–11] polyisoprene (PI),^[12] poly(dimethyl siloxane) (PDMS),^[8,12,13] poly(ethylene oxide) (PEO),^[14] poly(methyl methacrylate) (PMMA),^[15–21] and poly(2-vinylpyridine)^[22] have been synthesized, and they have been shown to develop fascinating morphologies at the nanometer to micrometer scale.^[2,7,23] Another important feature of PFDMS is its ability to crystallize.^[24] It exhibits a T_g of around 30°C and—depending on molecular mass and thermal history—a melting temperature (T_m) in the range from 110 to 142°C .^[24] Even as a part of block copolymers, the PFDMS crystallinity has already been addressed,^[10,12,13,25] but there is no systematic investigation available so far.

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^a This article is related to the Special Issue “Metal-Containing and Metallo-Supramolecular Polymers and Materials” guest-edited by U. Schubert in *Macromolecular Rapid Communications*.

The crystallization rate of polymers rises with the melting enthalpy for 100% crystalline material at otherwise equal conditions. It amounts to about $26 \text{ J} \cdot \text{g}^{-1}$ only for PFDMS^[26] (compared to, e.g., $277 \text{ J} \cdot \text{g}^{-1}$ for PE^[27] and $197 \text{ J} \cdot \text{g}^{-1}$ for PEO^[28]). PFDMS consequently crystallizes very slowly. Nevertheless, an attempt was made to determine its equilibrium melting temperature, $T_{m,0}$, which was found to be $T_{m,0} = 143 \text{ }^\circ\text{C}$ ^[24] using the linear Hoffmann–Weeks (HW) approach.^[29] The HW formalism is the commonly used procedure when $T_{m,0}$ of semi-crystalline polymers has to be determined, especially because all data required are easily available from conventional differential scanning calorimetry (DSC) measurements: the samples are crystallized isothermally at different temperatures T_c and, subsequently, the melting temperatures, T_m , of the crystallites grown at the respective temperatures T_c are determined. Assuming a linear relationship between T_m and T_c , $T_{m,0}$ is determined by linear extrapolation of the experimentally determined sets of T_m versus T_c pairs to $T_m = T_c$.

The HW procedure is indeed very convenient, but its general validity is still a matter of discussion; careful analysis of various semi-crystalline polymers showed a nonlinear T_m versus T_c behavior relatively frequently, especially when the studies were extended over a broader range of crystallization temperatures.^[24,30] One might speculate that there is in fact not even one semi-crystalline polymer showing a truly linear T_m versus T_c behavior.^[31] The main reason for nonlinearity in the T_m versus T_c relation is that the thickness d_c of crystalline lamellae formed by a crystallizing polymer at a specific crystallization temperature T_c is not strictly inversely proportional to the extent of super-cooling, $\Delta T = T_{m,0} - T_c$, applied during isothermal crystallization.^[30,32] This, however, is a necessary prerequisite for the HW analysis.

Facing the potentially intrinsic nonlinearity of the T_m versus T_c relation, a modified HW Equation was proposed for semi-crystalline polymers which explicitly takes into account the respective thicknesses of crystalline lamellae as an input parameter; this is the Gibbs–Thomson (GT) procedure,^[33] which was originally developed to analyze the size-dependent melting behavior of metallic nanoclusters. However, when taking the thickness of the crystalline lamellae, d_c , as the relevant size in the GT approach, it can be adapted to semi-crystalline polymers as well. Within this model the equilibrium melting temperature is obtained by plotting T_m versus d_c^{-1} for several samples crystallized isothermally at different crystallization temperatures T_c . Linear regression following the Equation $T_m = T_{m,0} - Cd_c^{-1}$, with C being a constant, and extrapolation to $d_c^{-1} = 0$, i.e., to infinite extension of the crystalline lamellae, results in the desired value of $T_{m,0}$. Obviously d_c needs to be determined experimentally as additional information.

Validity of the GT approach has been confirmed for, e.g., PE and PP.^[30,32] Though scientifically more reliable, the GT procedure is not very common so far because d_c has to be determined for several T_c s using small-angle X-ray scattering (SAXS), or at least estimated by TEM, in addition to the DSC studies required anyway to measure the $T_{m,s}$.^[30] Polymers might exist where indeed such efforts are not justified because sufficient linearity is ensured in their T_m versus T_c behavior. However, apparent inconsistencies in the case of PFDMS permit us to assume that this is not correct for this particular metallopolymer. For example, the published value $T_{m,0} = 143 \text{ }^\circ\text{C}$ is neither in harmony with a melting temperature of $145 \text{ }^\circ\text{C}$ observed for PFDMS crystallized under stretching conditions^[26] nor with the melting transition at $143 \text{ }^\circ\text{C}$ found for PS-*b*-PFDMS (PS: polystyrene) diblock copolymers as will be shown later in this paper; by definition, $T_{m,0}$ corresponds to the 100% perfect crystallite of infinite size, which is experimentally inaccessible.^[31,32] As a consequence, the $T_{m,0}$ reported so far might be in error to a more or less extent.

Since $T_{m,0}$ is indeed an important materials parameter,^[31] we decided to spend some effort in its careful re-determination for PFDMS homo- and block copolymers using both the HW and GT approaches. For this purpose, a low-molecular-weight and a high-molecular-weight PFDMS homopolymer sample and an almost symmetrical PS-*b*-PFDMS diblock copolymer were prepared. The latter showed a lamellar micro-morphology below its order-disorder transition (ODT) temperature; based on preliminary SAXS studies carried out in the context of the present work, the PS-*b*-PFDMS diblock copolymer under investigation formed an isotropic monophasic melt above $T_{\text{ODT}} \approx 207 \text{ }^\circ\text{C}$ but microphase-separates at lower temperatures to give a lamellar micromorphology. This behavior is in agreement with the known fact that PS and PFDMS are partially miscible as long as the PFDMS is in its noncrystalline state.^[34] Upon PFDMS crystallization, however, the mixture changes from the weak-segregation into the strong-segregation regime.^[35]

Due to very slow PFDMS crystallization, and in order to realize its crystallization at as high temperatures T_c as possible, the so-called self-seeding technique was applied for the preparation of isothermally crystallized samples.^[32,36] This technique is known to tremendously reduce the time required for nucleation, especially for slowly nucleating polymers at moderate super-coolings. For self-seeding crystallization, a pre-crystallized sample is initially needed which can be prepared by long-term isothermal annealing of the material. Later on, prior to the intended crystallization study, this pre-crystallized sample is heated very carefully, and for a very short period of time only, to a temperature T_{max} (i.e., the self-seeding temperature) that is just slightly above its melting temperature, T_m . At this T_{max} , the crystalline order vanishes to an extent

that the material appears isotropically from the macroscopic point of view; crystalline regions cannot be detected, neither by calorimetry, nor by nuclear magnetic resonance, nor by microscopic techniques.^[37–41] It is still a matter of controversy whether some local chain orientation or even tiny crystallites survive under the applied melting conditions. In any case, these remaining structural entities considerably support nucleation when the material is subsequently cooled down to the respective crystallization temperature. Crystallization supported by the self-seeding technique is the more powerful the smaller the $T_{\max} - T_m$ interval and the shorter the annealing time at T_{\max} .

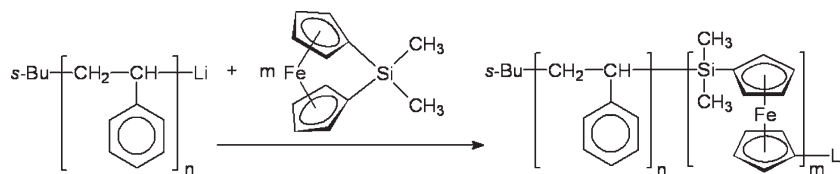
The consequences of self-seeding on the thickness of the crystalline lamellae and the T_m has hardly been addressed so far.^[41] It has been reported for PE and PP,^[41] however, that using this technique one ends up at the same crystal thickness as in the usual procedure of cooling a melt. Throughout, the experimentally determined melting temperatures, T_m , and the thicknesses of the crystalline lamellae, d_c , were found to depend exclusively on the temperature T_c applied during isothermal crystallization.^[42] For PLLA, moreover, the data from the normal and self-seeding crystallizations fall on the same linear T_m versus $1/d_c$ curve.^[43]

Obviously, self-seeding does not influence the T_m versus T_c nor the T_m versus $1/d_c$ behavior. Nevertheless, we carried out additional studies to reconfirm that aspect explicitly for PFDMS. Using the correspondingly crystallized PFDMS samples, finally, the melting behavior was studied using DSC, and the thicknesses of the crystalline PFDMS lamellae were determined using synchrotron SAXS. The sets of data obtained in this manner were taken to calculate $T_{m,0}$ following the HW and the GT approaches, respectively, and thus to compare these two approaches with respect to their ability to yield the equilibrium melting temperature of PFDMS.

Experimental Part

Materials and Syntheses

Dimethylsila[1]ferrocenophane (FS) was prepared as described recently^[16,44] and purified by repeated crystallization from hexane at -60°C in the presence of trioctyl aluminum. Purification of styrene and the solvents was carried out as described elsewhere.^[16] *sec*-BuLi was prepared from *sec*-butyl chloride and lithium dispersion. PFDMS homopolymers having molar masses of $4.9\text{ kg}\cdot\text{mol}^{-1}$ ($P_n \approx 20$; "PFDMS-1") and $60.7\text{ kg}\cdot\text{mol}^{-1}$ ($P_n \approx 250$; "PFDMS-2"), were prepared following published procedures, i.e., by anionic polymerization of dimethylsila[1]ferrocenophane (FS).^[16] The polydispersity was PDI = 1.05 and 1.03, respectively. The PS-*b*-PFDMS diblock copolymer was synthesized by sequential living



■ Scheme 1. Key step of the synthesis of the PS-*b*-PFDMS diblock copolymer.

anionic polymerization in glass reactors equipped with break-seals for the addition of the reagents (Scheme 1).^[8,11,22] First, styrene was polymerized in cyclohexene at room temperature using the required amount of *sec*-BuLi. After completed conversion ($\approx 12\text{ h}$), a solution of FS in cyclohexene was added by smashing the break-seal of the corresponding ampoule. The resulting mixture was stirred for 10 min, thereby allowing the transformation of the styryllithium active centers to the ferrocenyllithium ones. Then, an equal volume of tetrahydrofuran (THF) was added in order to promote the FS propagation from the ferrocenyllithium-based macroanions.^[8–11] After stirring for a further 2 h, the reaction was terminated with degassed methanol. The raw product was isolated by precipitation in hexane. Since it contained PS homopolymer impurities, it was purified by three-fold fractionation; the block copolymer was precipitated from its THF solution by dropwise addition of diethyl ether. The finally remaining traces of PS homopolymer ($\approx 5\%$) were accepted because it has the same molar mass as the PS block in the block copolymer and thus only swells the PS domains slightly. Also, it influences neither the PFDMS crystallization nor the sample's morphology.^[9,45] The resulting diblock copolymer has an overall molar mass (\bar{M}_n) of $24.3\text{ kg}\cdot\text{mol}^{-1}$ (PDI = 1.05), with a PS block of $\bar{M}_n = 12.0\text{ kg}\cdot\text{mol}^{-1}$ and a PFDMS block of $\bar{M}_n = 12.3\text{ kg}\cdot\text{mol}^{-1}$. At room temperature, the volume fraction of PS in the amorphous and crystalline sample is 0.586 and 0.614, respectively, taking a density of $1.05\text{ g}\cdot\text{cm}^{-3}$ for the amorphous PS and densities of 1.294 and $1.455\text{ g}\cdot\text{cm}^{-3}$ for the amorphous and crystalline PFDMS, respectively.^[46] The mass fraction of PFDMS in the sample is 0.494. This composition resulted in a lamellar micro-morphology at temperatures below ODT.

Pre-Crystallization Procedure

In order to prepare the pre-crystallized samples of homopolymers PFDMS-1 and PFDMS-2 needed for subsequent self-seeding crystallization studies, films were cast from methylene chloride solution. After evaporation of the solvent, the films were heated under vacuum (0.1 mbar) to 200°C for 60 min in order to melt all crystallites possibly formed during film preparation, and to remove traces of air and solvent trapped in the samples. Afterward, the samples were cooled down quickly ($20\text{ K}\cdot\text{min}^{-1}$) to 119°C for PFDMS-1 and 125°C for PFDMS-2, which were found to be appropriate temperatures for pre-crystallization. At these temperatures, they were annealed for 60–70 h.

In the case of the PS-*b*-PFDMS diblock copolymer, the powdery "as-synthesized" material was heated in aluminum containers under vacuum to 180°C for 60 min. The molten material was pressed at around 145°C between two parallel steel plates to a film having a thickness of 0.3 mm. The pressed film was put into a DSC pan, placed in a DSC instrument, heated to 180°C by $10\text{ K}\cdot\text{min}^{-1}$,

annealed for 10 min at this temperature, and finally quenched to 127 °C. At this temperature, the material was kept for 64 h for isothermal pre-crystallization.

Self-Seeding Isothermal Crystallization

For isothermal crystallization using the self-seeding technique, the pre-crystallized samples were heated to T_{\max} , i.e., just slightly above the highest observable PFDMS melting transition, at a rate of $10 \text{ K} \cdot \text{min}^{-1}$ in DSC (for PFDMS-1 and PFDMS-2), or $1 \text{ K} \cdot \text{min}^{-1}$ in situ at the beamline (for PS-*b*-PFDMS). T_{\max} was typically 142 °C for PFDMS-1, 145 °C for PFDMS-2, and 147 °C for the PS-*b*-PFDMS diblock copolymer, corresponding to a $\Delta T = T_{\max} - T_m$ ranging between 1 and 3 °C. After annealing at T_{\max} for 1–2 min, the samples were cooled down by $30 \text{ K} \cdot \text{min}^{-1}$ to the respective isothermal crystallization temperature, T_c . Depending on the respective rate of crystallization, the samples were annealed at T_c for 1–10 h. Finally, T_m and d_c of the resulting crystalline materials were determined.

Differential Scanning Calorimetry (DSC)

Differential scanning calorimetry (DSC) measurements were carried out using a Perkin-Elmer DSC 7 instrument under an atmosphere of nitrogen. For determination of the T_m versus T_c behavior, the samples were scanned immediately after completed self-seeding crystallization at a heating rate of $10 \text{ K} \cdot \text{min}^{-1}$ up to a temperature of 200 °C.

Small Angle and Wide Angle Synchrotron X-Ray Scattering (SAXS and WAXS)

SAXS and WAXS measurements were carried out at the beamline A2 in HASYLAB, DESY, Hamburg/Germany at a wavelength $\lambda = 0.15 \text{ nm}$. The samples were sandwiched between two layers of aluminum foil. The sample-to-detector distance for SAXS was around 2.9 m. The accumulation time for each SAXS image and WAXS frame was 30 s. Two-dimensional (2D) SAXS data and 1D WAXS data were recorded simultaneously. The SAXS patterns were calibrated by a rat tail tendon, and the WAXS patterns by poly(ethylene terephthalate). The 2D images were circularly integrated using the Fit2D program,^[47] normalized by the ionization intensity, and the background (i.e., air, aluminum) was subtracted.

For PS-*b*-PFDMS, the samples were crystallized and measured in situ at the beamline. For PFDMS homopolymers, the self-seeding crystallized samples were prepared in DSC prior to the X-ray measurements at DESY. The samples were measured at a heating rate of $1 \text{ K} \cdot \text{min}^{-1}$ from below to above T_c , and the X-ray data were taken at $T_{\text{X-ray}} = T_c$. Because there was a delay of 4–9 d between sample preparation and the X-ray measurements, it was necessary to ensure that the material's structure did not change during that storage at room temperature. In order to verify this, fractions of these stored samples were scanned in a DSC, and the DSC curves thus determined were compared with those obtained from DSC runs carried out with "fresh" samples, i.e., immediately after

crystallization at $T = T_c$. The DSC measurements reliably showed exactly the same traces, indicating that the crystalline morphology did not change during storage at room temperature.

From the SAXS data, the thicknesses of the crystallite lamellae, d_c , were calculated for the PFDMS homopolymers following the correlation analysis developed by Strobl.^[31] In this type of analysis, the structure correlation function $K(z)$ is calculated by cos-Fourier-transformation from the isotropic SAXS pattern $I(q)$,

$$K(z) = C \int_0^{\infty} \cos(qz) q^2 I(q) dq,$$

where C is a constant which reflects the scattering power of the material. If the structure is a regular stack of amorphous and crystalline layers, analysis of the course of that function then yields the values of long period d and lamellar thickness d_c .

Results and Discussion

Conventional versus Self-Seeding Poly(ferrocenyl dimethylsilane) (PFDMS) Crystallization

Poly(ferrocenyl dimethylsilane) (PFDMS) homopolymers crystallize very slowly due to their very low melting enthalpy. Consequently, both samples PFDMS-1 and PFDMS-2 did not crystallize when cooling the melts from 200 °C to room temperature at rates higher than around $30 \text{ K} \cdot \text{min}^{-1}$. If the cooling rate was only $1 \text{ K} \cdot \text{min}^{-1}$, however, they crystallized nicely. Isothermal PFDMS crystallization also proved to be feasible, provided that an appropriate crystallization temperature $80 \text{ °C} < T_c < 130 \text{ °C}$ was selected.

Poly(ferrocenyl dimethylsilane) (PFDMS) crystallization in the PS-*b*-PFDMS diblock copolymer, on the other hand, was found to be even much slower than in the homopolymers; DSC and WAXS studies showed that the PFDMS blocks did not crystallize even if the melt was cooled down from 180 °C at only $0.5 \text{ K} \cdot \text{min}^{-1}$. Vitrification of the PS layers at super-coolings required for efficient PFDMS nucleation and crystal growth might be the reason for the suppressed crystallization observed in the block copolymer. Partial miscibility of PS and amorphous PFDMS at elevated temperatures might additionally hinder proper PFDMS nucleation.^[48] Nevertheless, PFDMS crystallization in the PFDMS-*b*-PS diblock copolymer was proven to be possible if annealing was done at 127 °C for 64 h: two sharp peaks appear in the WAXS patterns at $q = 1.44$ and 1.64 nm^{-1} . They roughly correspond to the 200 and 010 reflections at 1.50 and 1.66 nm^{-1} , respectively, of the melt-crystallized PFDMS homopolymer.^[24,26] Thus, the crystalline morphology of PFDMS in the diblock copolymer is obviously almost the same as in those homopolymers.

Prior to the intended investigation of the equilibrium melting temperature of PFDMS, it was confirmed by suitable experiments that no undesired changes of PFDMS

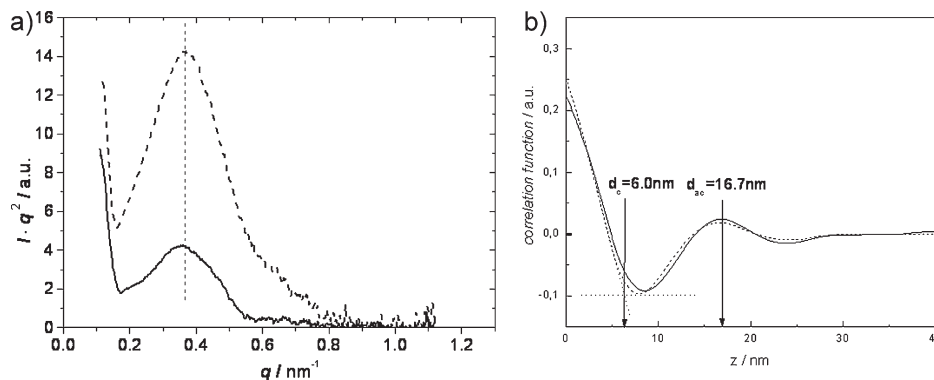


Figure 1. SAXS of two samples of PFDMS-2 isothermally crystallized at $T_c = 120^\circ\text{C}$ (solid line) from the “regular” isotropic melt at 180°C (i.e., without self-seeding), and (dashed line) using the self-seeding procedure as specified in the text; (a) scattering pattern, (b) corresponding correlation function. The solid curve in (a) was magnified 3.4 times for ease of comparison.

crystallization and morphology are induced by application of the self-seeding procedure. To this end, pairs of samples were crystallized at different T_c s, once with and once without self-seeding, and then compared using DSC, WAXS, and SAXS. Figure 1 shows two representative SAXS curves obtained from PFDMS-2 after isothermal crystallization at $T_c = 120^\circ\text{C}$ with and without self-seeding.

Throughout, the SAXS studies proved that identical values of q_{max} are found for samples crystallized at the same T_c with and without self-seeding (Figure 1a). The correlation function analysis^[49] derived from the SAXS curves indicates that the thicknesses of the crystalline PFDMS lamellae are the same for both samples (Figure 1b). Thus, self-seeding does not alter the thicknesses of the crystalline lamellae^a.

Next it was proved that the extension of self-seeding nuclei has no influence on the samples' melting temperatures. For this sake, a series of materials pre-crystallized under identical conditions was heated to different temperatures T_{max} , and annealed there for different periods of time, prior to cooling down to the respective crystallization temperature T_c . After finalized self-seeding crystallization, the materials were scanned in DSC runs to see whether changes can be found in the melting temperature T_m . Figure 2 shows two representative DSC heating runs for two samples of the PS-*b*-PFDMS diblock copolymer which were prepared at different self-seeding conditions. The melting enthalpies were found to be 8.9 and $6.8\text{ J}\cdot\text{g}^{-1}$ for the materials annealed at 145 and 147°C , respectively. The melting temperatures, however, did not differ and hence it is evident that the strength of the self-seeding nuclei does

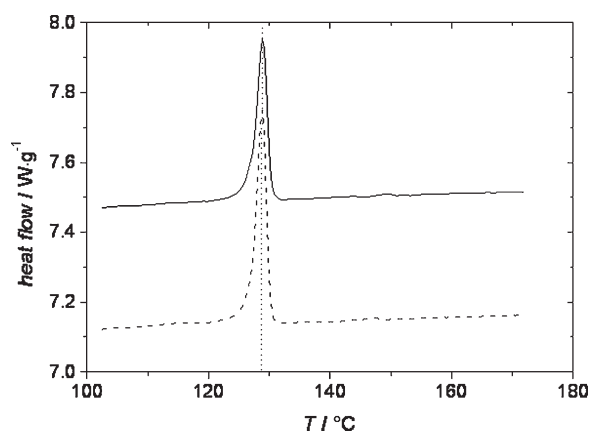


Figure 2. Differential scanning calorimetry (DSC) heating runs at $10\text{ K}\cdot\text{min}^{-1}$ for PS-*b*-PFDMS diblock copolymer samples pre-crystallized at 127°C for 64 h, heated to the self-seeding temperatures $T_{\text{max}} = 147^\circ\text{C}$ (solid line) or $T_{\text{max}} = 145^\circ\text{C}$ (dashed line) prior to quenching at $20\text{ K}\cdot\text{min}^{-1}$ to the isothermal crystallization temperature $T_c = 100^\circ\text{C}$. At T_c , both samples were allowed to crystallize for 180 min.

not induce differences in the materials' behavior. It is only the degree of crystallinity which can differ to some extent depending on the strength of the self-seeding nuclei.

The third aspect to be analyzed was whether the crystalline morphology is affected by the self-seeding procedure and/or the isothermal crystallization temperatures T_c . WAXS studies were carried out, therefore, on samples crystallized at different T_c s, with and without self-seeding. Figure 3 shows three characteristic WAXS patterns obtained for the high-molecular-weight PFDMS-2 crystallized at $T_c = 100, 120,$ and 135°C using an identical self-seeding procedure. The WAXS patterns essentially coincide, allowing the conclusion that no change in the crystalline morphology is detectable in the studied range of temperatures. Thus, we can be sure now that isothermal crystallization at different T_c s with and without self-seeding leads

^a The long period d_{ac} is determined as the location of the first maximum of the correlation function $K(z)$ and a sub-thickness d as the value of z where the tangent at the first minimum of $K(z)$ intersects with the straight line through the first turning point of $K(z)$. Whether d represents the lamellar thickness d_c or d_a can be decided by means of the material's crystallinity as determined by e.g. DSC.

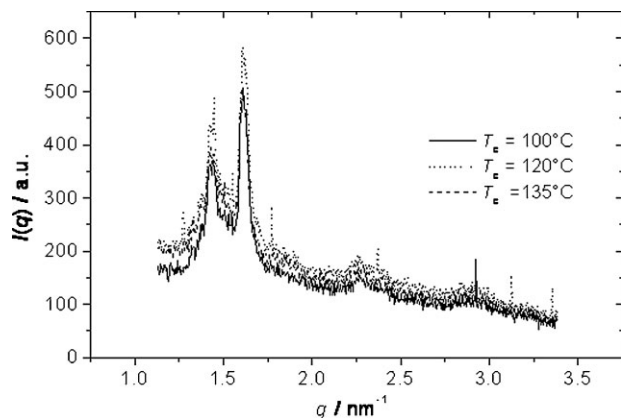


Figure 3. WAXS patterns for PFDMS-2 recorded at different T_c after self-seeding crystallization.

to the same crystal morphology. All in all, the studies proved that application of the self-seeding technique is legitimate in the subsequent investigations planned to determine the $T_{m,0}$ of PFDMS.

Crystal Growth Rate

Isothermal crystallization of the self-seeded PFDMS samples causes an exothermic peak in the DSC thermograms as shown in Figure 4 for PFDMS-2. Such graphs allow complete quantification of crystallization kinetics by integrating the peak areas as a function of crystallization time. For the sake of simplicity, we characterize this crystallization curve by the time t_p of highest crystallization rate, i.e., the time where the exothermic peak maximum appears. These times were plotted in Figure 5 as a function of T_c for

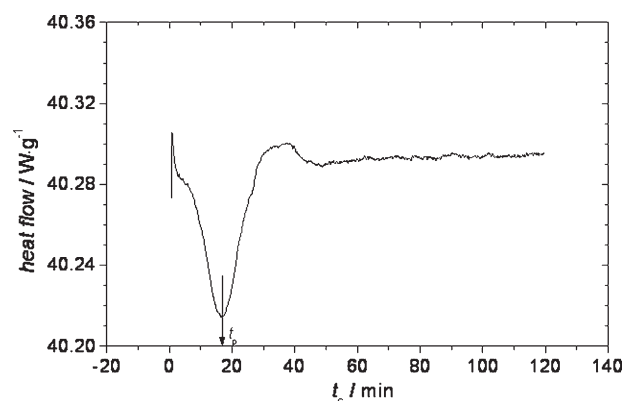


Figure 4. Representative isothermal heat flow as function of crystallization time t_c as measured by DSC for the isothermal self-seeding crystallization of PFDMS-2 at $T_c = 90^\circ\text{C}$. The peak time, t_p , used in the following evaluation is indicated.

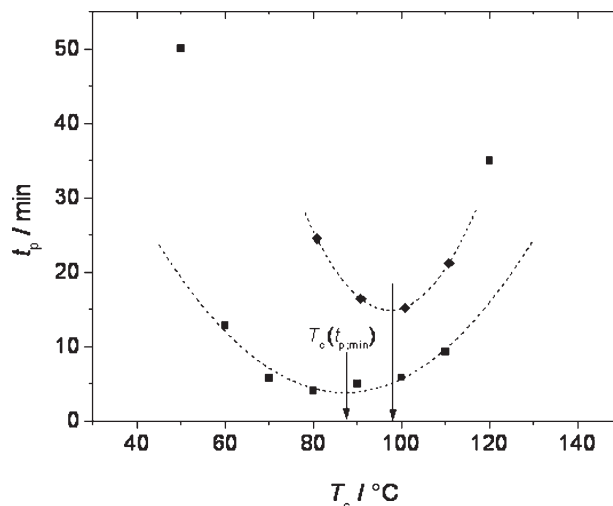


Figure 5. Rate of self-seeding isothermal crystallization as expressed by t_p of the PFDMS homopolymers at different crystallization temperatures, T_c , for PFDMS-1 (■) and PFDMS-2 (◆). The lines through the data points are just a visual guide. The arrows indicate the T_c where t_p has its minimum.

homopolymers PFDMS-1 and PFDMS-2 in the experimentally accessible range $60^\circ\text{C} \leq T_c \leq 120^\circ\text{C}$.

The maximum rates of isothermal crystallization were determined to be at $T_{c,\text{max}} \approx 87^\circ\text{C}$ for PFDMS-1 and at $T_{c,\text{max}} \approx 97^\circ\text{C}$ for PFDMS-2. A higher temperature $T_{c,\text{max}}$ for the higher molecular weight sample is reasonable due to lower chain mobility. For $T_c \leq 60^\circ\text{C}$ and $T_c > 120^\circ\text{C}$, it was not possible to determine reliable values of t_p because of the tremendous broadening and insufficient resolution of the exothermic peaks. For $T_c \leq 60^\circ\text{C}$, frozen chain dynamics might play a role whereas at $T_c > 120^\circ\text{C}$ the low undercooling — and thus the very slow crystal growth — might be relevant.

We note that for the PS-*b*-PFDMS diblock copolymer an analogous estimation of the temperature dependency of the crystallization kinetics was not possible using DSC since self-seeded samples were prepared only for X-ray scattering. Then, the crystal growth rate is very slow, and the samples remained nearly amorphous even under apparently appropriate crystallization conditions, i.e., after the self-seeding approach and for very long crystallization times. For example, the material almost did not crystallize after annealing at $T_c = 110^\circ\text{C}$ for about 4.5 h.

Degree of Crystallinity

As stated above, the crystalline morphology of the PFDMS is the same for all T_c s in the homopolymers and the diblock copolymer. Therefore, we apply for the moment the published melting enthalpy of $\Delta H_{0,\text{PFDMS}} = 26 \text{ J} \cdot \text{g}^{-1}$ for

Table 1. Normalized melting enthalpies for PFDMS-1, PFDMS-2, and the PS-*b*-PFDMS.

Samples	DSC heating rate	Melting enthalpy $\text{J} \cdot \text{g}^{-1}$	Crystallinity	Crystallinity
	$^{\circ}\text{C} \cdot \text{min}^{-1}$		% ^{a)}	% ^{b)}
PFDMS-1	2	19.3	72	49
PFDMS-2	2	14.7	54	37
PS- <i>b</i> -PFDMS	2	26.7 ^{c)}	99	68
PFDMS-1	10	19.7	73	50
PFDMS-2	10	12.7	47	32
PS- <i>b</i> -PFDMS	10	23.5 ^{c)}	87	60

^{a)} $\Delta H_0 = 26 \text{ J} \cdot \text{g}^{-1}$ [26], ^{b)} $\Delta H_0 = 38 \text{ J} \cdot \text{g}^{-1}$ (cf. text); ^{c)}Normalized with respect to the mass fraction of PFDMS.

100% crystallized PFDMS^[26] to estimate the degrees of crystallinity for both sets of materials. The highest determined DSC melting enthalpies were $19.9 \text{ J} \cdot \text{g}^{-1}$ for the low-molecular-weight PFDMS-1 (after annealing at 119°C for 66 h), and $14.6 \text{ J} \cdot \text{g}^{-1}$ for the high-molecular-weight homopolymer PFDMS-2 (annealing at 125°C for 67 h). Obviously, the degrees of crystallinity in the PFDMS homopolymers do not exceed 50–75% (Table 1 and Figure 6). In the PS-*b*-PFDMS block copolymer, on the other hand, the enthalpies determined from the normalized melting peak areas were considerably larger. Values of 23–26 $\text{J} \cdot \text{g}^{-1}$ were observed upon crystallization at 127°C for 64 h, which would correspond to PFDMS crystallinities of 90–100%. This is rather unlikely, however. It is not an easy task to identify the completion of isothermal crystallization. With prolongation of the crystallization time, the melting enthalpy may increase slightly further. The data of Table 1 may consequently be too low, and the DSC degrees

of crystallinity may even be higher than those given in Figure 6.

One might argue that the published $\Delta H_{0,\text{PFDMS}} = 26 \text{ J} \cdot \text{g}^{-1}$ is too low. Therefore, for comparison purposes, the degrees of crystallinity of the same samples were additionally determined by 1D analysis of the SAXS data recorded at different T_c s, following the correlation analysis. Figure 6 shows the calculated degrees of crystallinity for PFDMS-2. It is evident that the crystallinities obtained from SAXS are clearly below those from DSC data. Usually, the inverse situation is found; especially for samples with low crystallinity, the degrees of crystallinity determined from DSC are lower than those determined from SAXS measurements. Hence, the larger values of crystallinity from the DSC data in the case of our studies might indicate that the true value of $\Delta H_{0,\text{PFDMS}}$ is actually higher than $26 \text{ J} \cdot \text{g}^{-1}$. Assuming identical DSC and SAXS crystallinities, $\Delta H_{0,\text{PFDMS}} = (38 \pm 3) \text{ J} \cdot \text{g}^{-1}$ would result. The DSC crystallinities based on that value are given in the last column of Table 1.

Using that ΔH_0 , however, the block copolymer crystallinity still remains above that of the homopolymer. Possibly, $\Delta H_0 = 38 \text{ J} \cdot \text{g}^{-1}$ is not valid for the block copolymer either. Although the crystalline morphology of the PFDMS is almost the same in the homopolymer and in the block copolymer, the ideal melting enthalpies must not necessarily be the same; the amorphous phases of the homopolymer and the block copolymer may have different enthalpic levels. If it is speculated that the crystallinity of the block copolymer must be below that of the homopolymer, then ΔH_0 (block) must be greater than ΔH_0 (homo). This, however, seems rather unlikely. We therefore rely on the result that the crystallinity of the block copolymer is higher than that of the homopolymer. The problem is treated in more detail, and a reliable explanation of the crystallization kinetics and the structural background will be presented in a forthcoming paper.

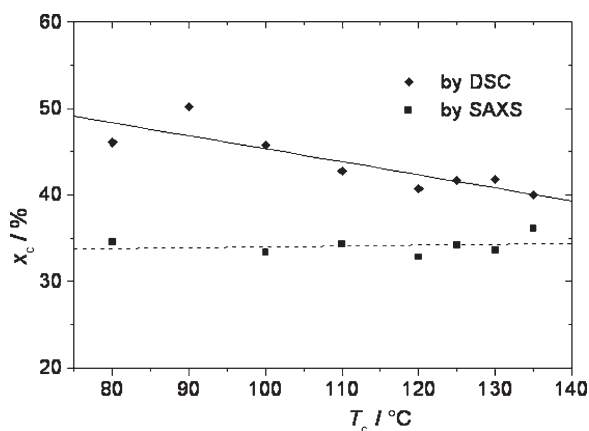


Figure 6. Degrees of crystallinity x_c by DSC (◆◆◆) and SAXS (■ ■ ■) of PFDMS-2 samples crystallized at different T_c s.

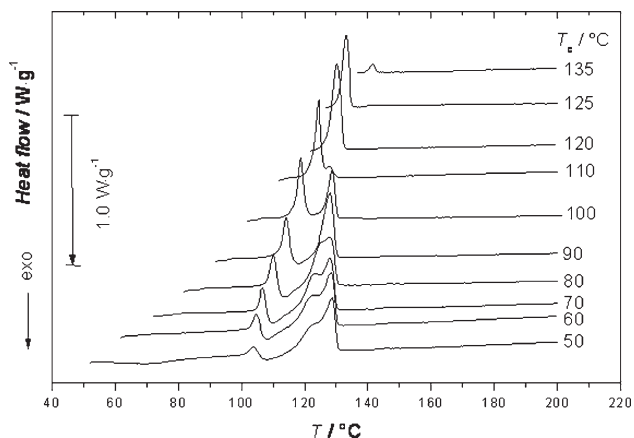


Figure 7. Heat flow versus temperature curves for PFDMS-1, scanned immediately after self-seeding crystallization at different isothermal temperatures. Heating rate: $10 \text{ K} \cdot \text{min}^{-1}$.

The T_m versus T_c Behavior

The DSC heating runs for the homopolymers PFDMS-1 and PFDMS-2 after isothermal self-seeding crystallization at different T_c s are shown in Figure 7 and 8, respectively. Multiple melting peaks were found for both homopolymers when crystallized at lower T_c , whilst for higher T_c only a single melting endotherm can be identified. The lower the T_c , the lower the temperature at which the first endotherm appears. For the low-molecular weight sample PFDMS-1, moreover, an exothermal transition can be seen if isothermal crystallization was performed at $T_c \leq 90^\circ\text{C}$ (Figure 7). This exothermic peak is attributed to re-crystallization during the heating run.^[32] Re-crystallization occurs when rather unstable original crystallites as formed during isothermal annealing at lower T_c s melt very early during the heating run, and can thus re-crystallize

instantaneously again to give more stable crystallites. Later, these new crystallites melt again but only at clearly higher temperatures. In the case of the high-molecular-weight PFDMS-2 (Figure 8), an exothermic re-crystallization peak is not detectable, maybe due to the lower overall chain mobility. This indicates that, as the molecular weight increases, the tendency of the PFDMS chains to re-crystallize during the heating scans decreases. On the other hand, absence of an exothermic peak during heating runs does not necessarily indicate absence of re-crystallization.^[50] Furthermore, occurrence of re-crystallization depends on the heating rate; during WAXS measurements where a heating rate of $1 \text{ K} \cdot \text{min}^{-1}$ was applied, re-crystallization has also been observed for the PFDMS-2 samples at all applied T_c .

For the PS-*b*-PFDMS diblock copolymer, the corresponding DSC curves measured immediately after completed self-seeding crystallization are given in Figure 9. Here, a single melting endotherm can be observed for all samples, independent of the applied T_c . In addition, there is no indication of an exotherm pointing toward re-crystallization during the DSC run. For further evaluation of the DSC traces, the position of the first melting peak is taken as the relevant value of T_m , because it can be reliably assumed that this peak is related to the crystallites formed during isothermal crystallization.

Since the presence of melting peaks at different temperatures as found for homopolymers at lower T_c is related to crystalline lamellae of different thicknesses, the DSC results indicate that the thicknesses of the crystalline lamellae in the diblock copolymer are defined much better than those in the homopolymers. This is a well-known phenomenon.^[51] This observation might be summarized by the statement that the more perfect the lamellar structure in a sample, the lower its tendency to recrystallize.

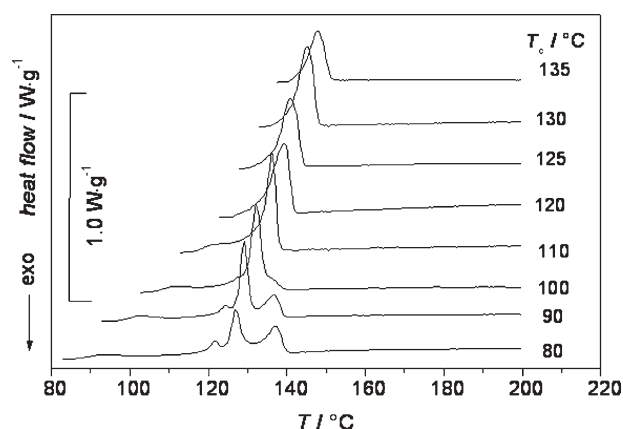


Figure 8. Heat flow versus temperature curves for PFDMS-2 scanned immediately after the self-seeding crystallization at different temperatures. Heating rate: $10 \text{ K} \cdot \text{min}^{-1}$.

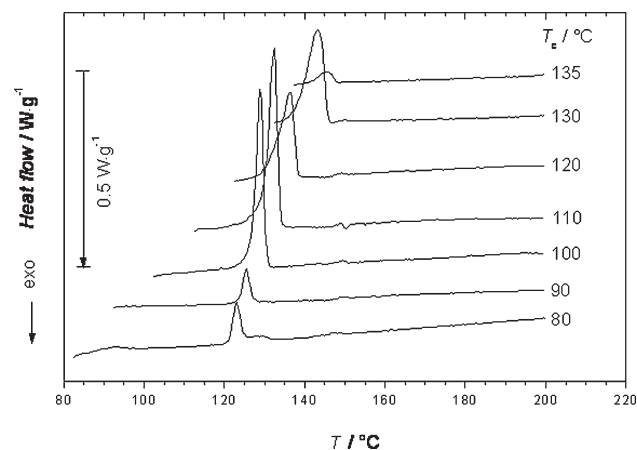


Figure 9. Heat flow versus temperature curves for the PS-*b*-PFDMS diblock copolymer after self-seeding crystallization at different temperatures. Heating rate: $10 \text{ K} \cdot \text{min}^{-1}$.

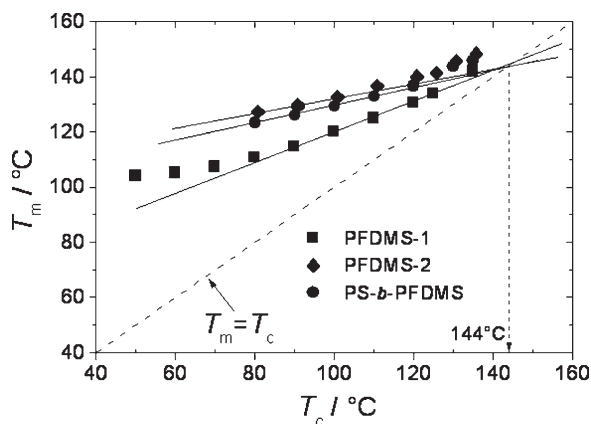


Figure 10. Determination of $T_{m,0}$ following the HW formalism: T_m versus T_c plots for PFDMS-1 (■ ■ ■), PFDMS-2 (◆ ◆ ◆), and PS-*b*-PFDMS (● ● ●).

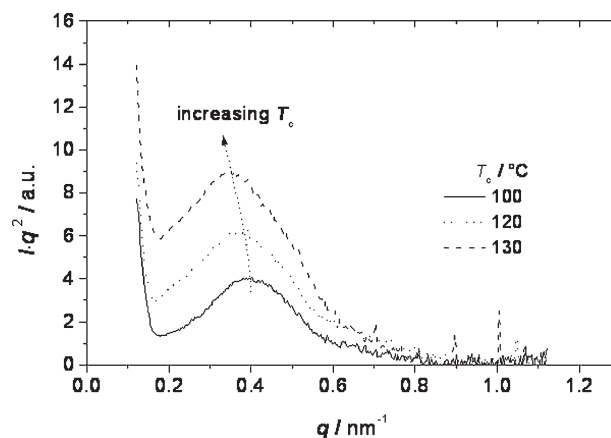


Figure 11. Representative synchrotron SAXS curves obtained for PFDMS-2 samples.

The Equilibrium Melting Temperature $T_{m,0}$

The Hoffmann–Weeks (HW) Procedure

As discussed above, the peak positions of the first melting endotherm observed has been taken as the melting temperature T_m of PFDMS homo and diblock copolymers crystallized at different temperatures T_c . For the $T_{m,0}$ determination following the HW procedure, T_m is plotted as a function of T_c (Figure 10). Obviously, a nonlinear T_m versus T_c relationship is found for all samples. If the data are nevertheless extrapolated linearly, taking into account only the data measured in the $80\text{ °C} \leq T_c \leq 120\text{ °C}$ range, a value of $T_{m,0} = 144\text{ °C}$ is found, which is in excellent agreement with the one reported in the literature.^[24]

It is surprising, however, that the values of $T_{m,0}$ of the homopolymers and the block copolymer coincide. Moreover, we cannot ignore that the T_m versus T_c data deviate substantially from a linear behavior especially for those samples obtained by isothermal crystallization at rather high values of T_c . Also, for PFDMS-2, deviation from a linear behavior is found even if T_c is lower than 80 °C . This nonlinearity may be due to the fact that the thicknesses of the crystallite lamellae are not strictly inversely proportional to the super-cooling ΔT , i.e., $T_{m,0} - T_c$.

The Gibbs–Thomson (GT) Procedure

Since the HW analysis fails due to the obvious nonlinearity of the T_m versus T_c relation, a more reliable value of $T_{m,0}$ of PFDMS was determined using the GT equation, where the lamellae thickness is explicitly taken into account. The required values of the crystallites' lamellae thicknesses, d_c , were obtained from SAXS patterns. Figure 11 shows some representative scattering curves for PFDMS-2. From these scattering data, the thicknesses of the crystallite lamellae, d_c , were calculated for the PFDMS homopolymers following

the correlation analysis. For the PS-*b*-PFDMS diblock copolymer, it must be considered that the diblock copolymer by itself forms a lamellar micro-morphology, consisting of a stack of its amorphous (PS) and crystallizable (PFDMS) part. If the PFDMS chains are oriented perpendicular to the lamellae's surfaces, each semi-crystalline PFDMS layer consists of two layers of PFDMS crystallites (Figure 12), and the correlation analysis yields $2d_c$.

The T_m versus d_c^{-1} pairs thus obtained are plotted in Figure 13 for the PFDMS-2 and the diblock copolymer, taking the T_m s from the DSC scans (Figures 8 and 9) and the d_c values from SAXS. The linear fit yields $T_{m,0} = (215 \pm 11)\text{ °C}$ for PFDMS-2 and $T_{m,0} = (179 \pm 5)\text{ °C}$ for PS-*b*-PFDMS. Different values of $T_{m,0}$ for PFDMS as homopolymer and being part of a block copolymer are reasonable because of the different environments for the PFDMS crystallites. In any case, our findings prove that the $T_{m,0}$ of PFDMS given in the literature^[24] is far too low.

Discussion and Conclusion

Failure of the HW approach, or discrepancies between the results of HW and GT analyses, have already been reported

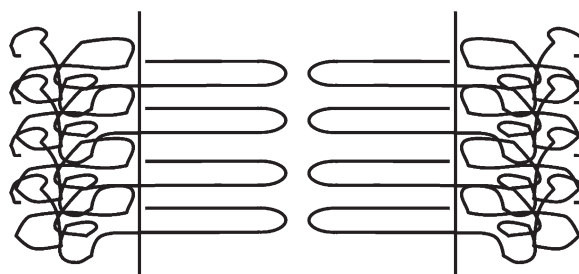


Figure 12. Schematic drawing of assumed structure and orientation of the PFDMS crystallites in the microphase-separated lamellar PS-*b*-PFDMS diblock copolymer.

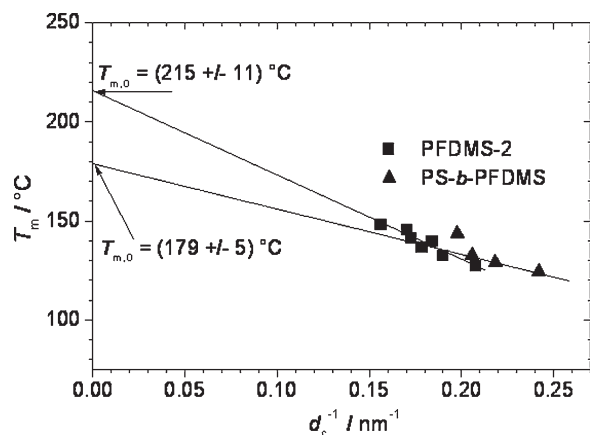


Figure 13. T_m versus d_c^{-1} for PFDMS-2 and for PS-*b*-PFDMS diblock copolymer.

several times.^[52] The HW approach for the determination of the equilibrium melting temperature $T_{m,0}$ presupposes that T_m and T_c differ due to lamellar thickening upon heating but are both proportional to the reverse d^{-1} of the respective lamellar thicknesses. In contrast, the GT approach demands only that $T_{m,0} - T_m = c_1 + c_2/d_c^{-1}$ (where c_1 and c_2 represent otherwise constant thermodynamic parameters). Obviously, the HW presuppositions are sharper and, consequently, more often violated than the GT one. The violation of the HW presuppositions causes a nonlinear dependency of T_m on T_c ,^[53] as found particularly here for PFDMS as a homopolymer and as a component of a block copolymer. The HW analysis of the dependency of T_m on T_c —neglecting the mentioned nonlinearity—confirmed a value of $T_{m,0} = 144$ °C as reported in the literature^[24] and also determined by HW analysis. The GT approach, in contrast, yielded $T_{m,0} = 215$ °C for the homopolymer and $T_{m,0} = 179$ °C for the block copolymer component. For the reasons given above we are convinced that these values are more reliable than the former ones. Moreover, their difference reasonably reflects the different thermodynamic boundary conditions for crystallization in the homopolymer and in the block copolymer.

Another remarkable result of our investigations is that PFDMS as a component of a block copolymer crystallizes to a higher extent than as a homopolymer. This result of DSC experiments holds although comparison with SAXS analyses indicates that the heat of fusion of PFDMS $\Delta H_0 = 26$ J · g⁻¹ reported in the literature^[26] may be too low.

Acknowledgements: The authors would like to thank the *Deutsche Forschungsgemeinschaft (DFG)*, the *Fonds der Chemischen Industrie (FCI)*, and the *Dr. Otto Röhm foundation* for financial support of this work. The helpful discussions with Prof. G. Strobl (Freiburg University), and the technical support provided

by Dr. S. S. Funari (DESY, Hamburg), and Mrs. Y. Xu (DKI) for the X-ray measurements, are gratefully acknowledged.

Received: December 22, 2009; Revised: February 27, 2010;
DOI: 10.1002/macp.200900718

Keywords: block copolymers; Gibbs–Thomson equation; Hoffmann–Weeks equation; melting point; polyferrocenylsilane

- [1] D. A. Foucher, B.-Z. Tang, I. J. Manners, *J. Am. Chem. Soc.* **1992**, *114*, 6246.
- [2] I. Manners, *Synthetic Metal-Containing Polymers*, VCH, Weinheim 2004.
- [3] A. S. Abd-El-Aziz, *Macromol. Rapid Commun.* **2002**, *23*, 995.
- [4] K. Kulbaba, I. Manners, *Macromol. Rapid Commun.* **2001**, *22*, 711.
- [5] I. Korczagin, R. G. H. Lammertink, M. A. Hempenius, S. Golze, G. J. Vancso, *Adv. Polym. Sci.* **2006**, *200*, 91.
- [6] D. E. Herbert, U. F. J. Mayer, I. Manners, *Angew. Chem., Int. Ed.* **2007**, *46*, 5060.
- [7] V. Bellas, M. Rehahn, *Angew. Chem., Int. Ed.* **2007**, *46*, 5082.
- [8] Y. Ni, R. Rulkens, I. J. Manners, *J. Am. Chem. Soc.* **1996**, *118*, 4102.
- [9] R. G. H. Lammertink, M. A. Hempenius, E. L. Thomas, G. J. Vancso, *J. Polym. Sci., Part B: Polym. Phys.* **1999**, *37*, 1009.
- [10] W. Li, N. Sheller, M. D. Foster, D. Balaisishis, I. Manners, B. Annis, J.-S. Lin, *Polymer* **2000**, *41*, 719.
- [11] V. Bellas, M. Rehahn, *Macromol. Rapid Commun.* **2007**, *28*, 1415.
- [12] J. A. Massey, K. Temple, L. Cao, Y. Rharbi, J. Raez, M. A. Winnik, I. J. Manners, *J. Am. Chem. Soc.* **2000**, *122*, 11577.
- [13] J. Raez, I. Manners, M. A. J. Winnik, *J. Am. Chem. Soc.* **2002**, *124*, 10381.
- [14] J. F. Gohy, B. G. G. Lohmeijer, A. Alexeev, X. S. Wang, I. Manners, M. A. Winnik, U. S. Schubert, *Chem. Eur. J.* **2004**, *10*, 4315.
- [15] C. Klöninger, M. Rehahn, *Macromolecules* **2004**, *37*, 8319.
- [16] C. Klöninger, M. Rehahn, *Macromolecules* **2004**, *37*, 1720.
- [17] U. Datta, M. Rehahn, *Macromol. Rapid Commun.* **2004**, *25*, 1615.
- [18] Ch. Klöninger, D. Knecht, M. Rehahn, *Polymer* **2004**, *45*, 8323.
- [19] Ch. Klöninger, M. Rehahn, *Macromol. Chem. Phys.* **2007**, *208*, 833.
- [20] X. Wang, M. A. Winnik, I. Manners, *Macromolecules* **2005**, *38*, 1928.
- [21] X. S. Wang, M. A. Winnik, I. Manners, *Macromol. Rapid Commun.* **2002**, *23*, 210.
- [22] V. Bellas, M. Rehahn, *Macromol. Chem. Phys.* **2009**, *210*, 320.
- [23] I. Manners, *Science* **2001**, *294*, 1664.
- [24] R. G. H. Lammertink, M. A. Hempenius, I. Manners, G. Julius Vancso, *Macromolecules* **1998**, *31*, 795.
- [25] R. G. H. Lammertink, M. A. Hempenius, G. J. Vancso, *Langmuir* **2000**, *16*, 6245.
- [26] V. S. Papkov, M. V. Gerasimov, I. I. Dubovik, S. Sharma, V. V. Dementiev, K. H. Pannell, *Macromolecules* **2000**, *33*, 7107.
- [27] *Polymer Handbook*, J. Brandrup, E. H. Immergut, Eds., 3rd edition, Wiley, New York 1989.

- [28] S. Z. D. Cheng, B. J. Wunderlich, *J. Polym. Sci., Polym. Phys. Ed.* **1986**, *24*, 577.
- [29] J. D. Hoffman, J. J. Weeks, *J. Res. Natl. Bur. Stand.* **1962**, *66A*, 13.
- [30] K. Yamada, M. Hikosaka, A. Toda, S. Yamazaki, K. Tagashira, *Macromolecules* **2003**, *36*, 4802.
- [31] G. Strobl, *The Physics of Polymers*, 2nd edition, Springer, Berlin Heidelberg 1996.
- [32] G. Strobl, *Prog. Polym. Sci.* **2006**, *31*, 398.
- [33] H. Zhou, G. L. Wilkes, *Polymer* **1997**, *38*, 5735.
- [34] H. B. Eitouni, N. P. Balsara, H. Hahn, J. A. Pople, M. A. Hemenius, *Macromolecules* **2002**, *35*, 7765.
- [35] M. D. Whitmore, J. Noolandi, *Macromolecules* **1988**, *21*, 1482.
- [36] D. J. Blundell, A. Keller, A. J. Kovacs, *J. Polym. Sci. B* **1966**, 481.
- [37] A. Maus, E. Hempel, T. Thurn-Albrecht, K. Saalwächter, *Eur. Phys. J.* **2007**, *E 23*, 91.
- [38] A. T. Lorenzo, M. L. Arnal, J. J. Sanchez, A. J. Müller, *J. Polym. Sci., Part B: Polym. Phys.* **2006**, *44*, 1738.
- [39] A. J. Müller, M. L. Arnal, *Prog. Polym. Sci.* **2005**, *30*, 559.
- [40] J. Xu, Y. Ma, W. Hu, M. Rehahn, G. Reiter, *Nat. Mater.* **2009**, *8*, 348.
- [41] G. Strobl, *Eur. Phys. J. E* **2000**, *3*, 165.
- [42] B. Heck, T. Hugel, M. Iijima, E. Sadiku, G. Strobl, *New J. Phys.* **1999**, *1*, 17.
- [43] T.-Y. Cho, G. Strobl, *Polymer* **2006**, *47*, 1036.
- [44] Y. Ni, R. Rulken, I. J. Manners, *J. Am. Chem. Soc.* **1996**, *118*, 4102.
- [45] I. W. Hamley, *The Physics of Block Copolymers*, Oxford University Press, New York 1998.
- [46] J. Bodycomb, D. Yamaguchi, T. Hashimoto, *Macromolecules* **2000**, *33*, 5187.
- [47] A. P. Hammersley, *ESRF Internal Report ESRF98HA01T*, FIT2D V9.129 Reference Manual V3.1 1998.
- [48] H. B. Eitouni, N. P. Balsara, H. Hahn, J. A. Pople, M. A. Hemenius, *Macromolecules* **2002**, *35*, 7765.
- [49] G. R. Strobl, M. Schneider, *J. Polym. Sci., Polym. Phys. Ed.* **1980**, *18*, 1343.
- [50] A. A. Minakov, D. A. Mordvintsev, C. Schick, *Polymer* **2004**, *45*, 3755.
- [51] F. Zhang, B. Stühn, *Colloid Polym. Sci.* **2007**, *285*, 371.
- [52] M. Al-Hussein, G. Strobl, *Macromolecules* **2002**, *35*, 1672.
- [53] H. Marand, J. Xu, S. Srinivas, *Macromolecules* **1998**, *31*, 8219.